

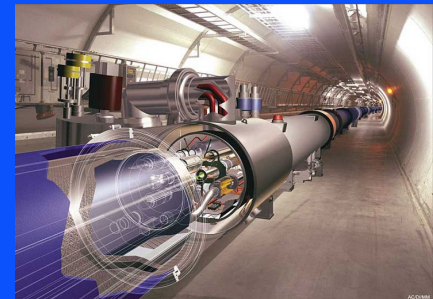


**Superconductivity & Particle AcceleratorS 2024**  
**Polish Academy of Sciences**

**PREDICTING LIFETIME OF IRRADIATED  
METASTABLE MATERIALS AT EXTREMELY LOW  
TEMPERATURES**

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Laboratory of Extremely Low Temperatures  
Cracow University of Technology**





## Outline:

1. Motivation: radiation sources at cryogenic temperatures
2. Strain induced fcc-bcc phase transformation
3. Radiation induced damage
4. Radiation induced hardening
5. Conclusions

Research Centre –  
Laboratory of Extremely Low Temperatures (RC-LENT)

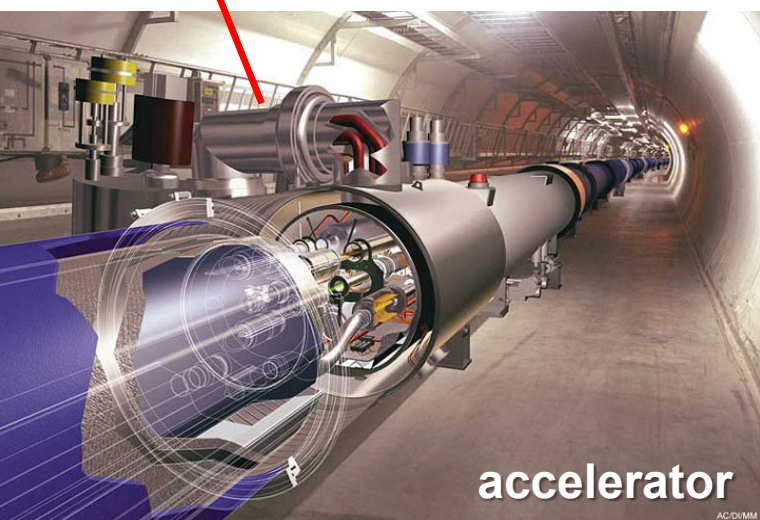
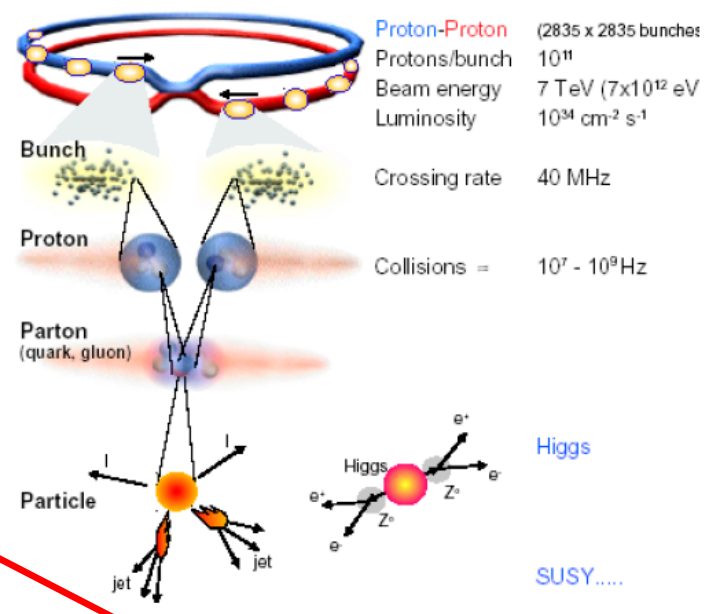
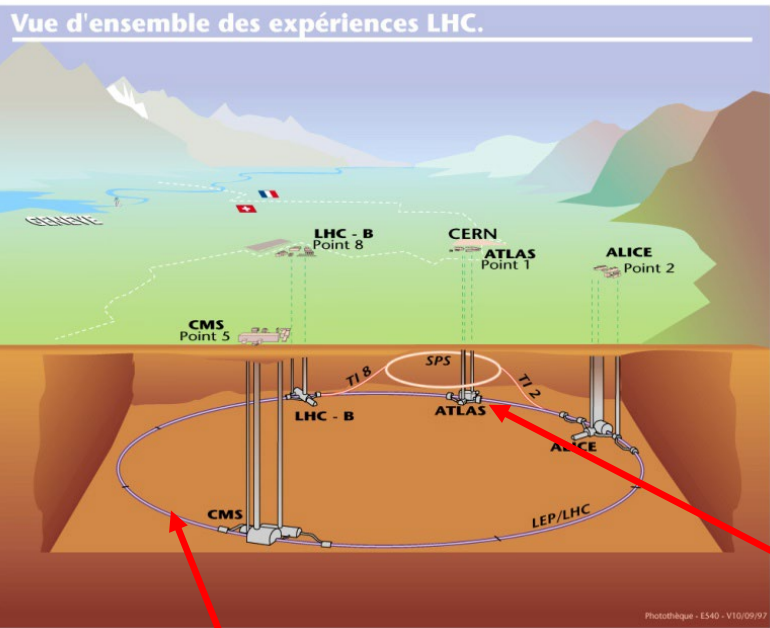




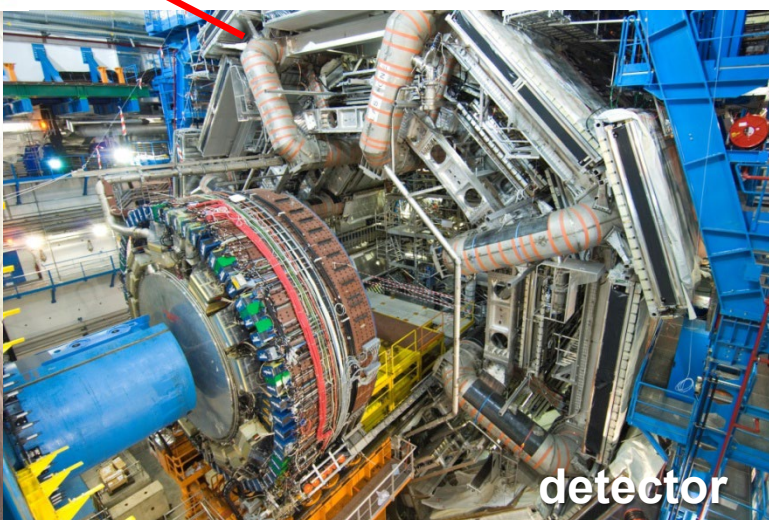
# Motivation: CERN Large Hadron Collider *active phenomena at cryogenic temperatures*

**LHC** is the largest scientific instrument in the world based on the principle of super-conductivity!

**LHC** operates in super-fluid helium at 1.9 K



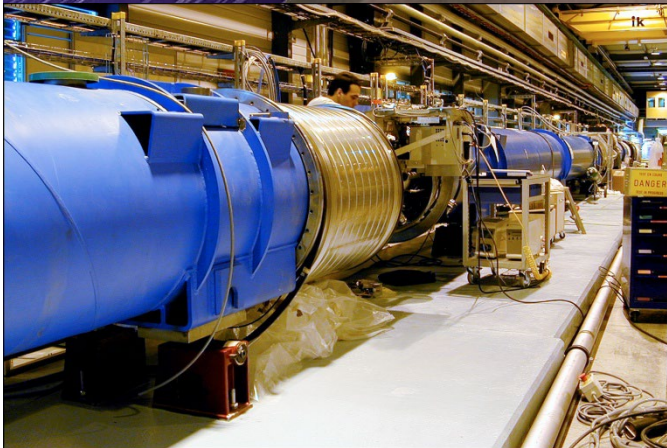
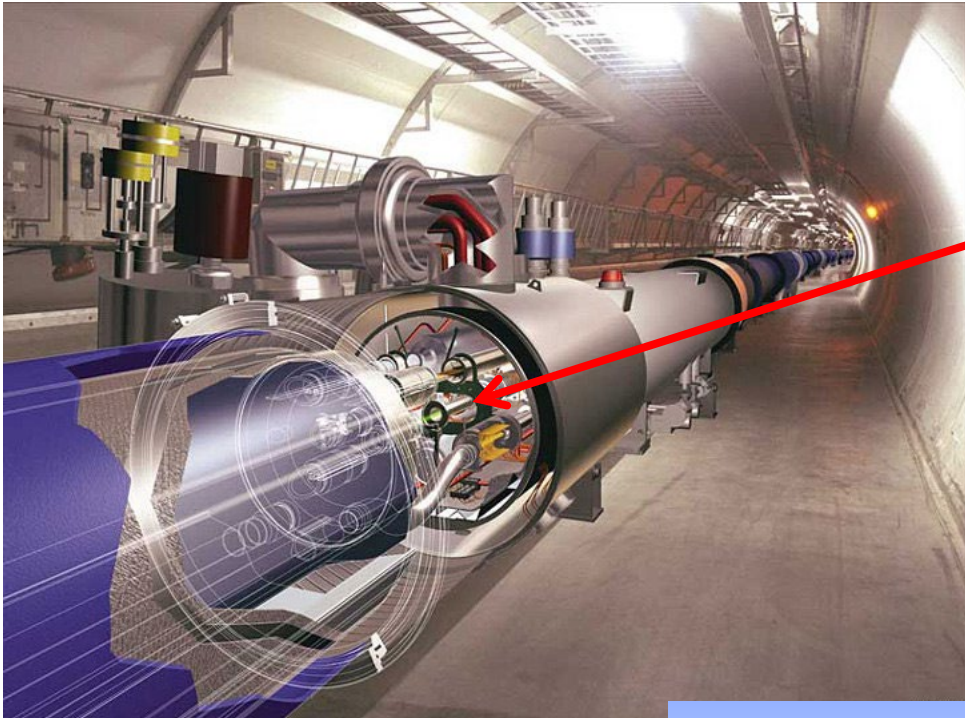
accelerator



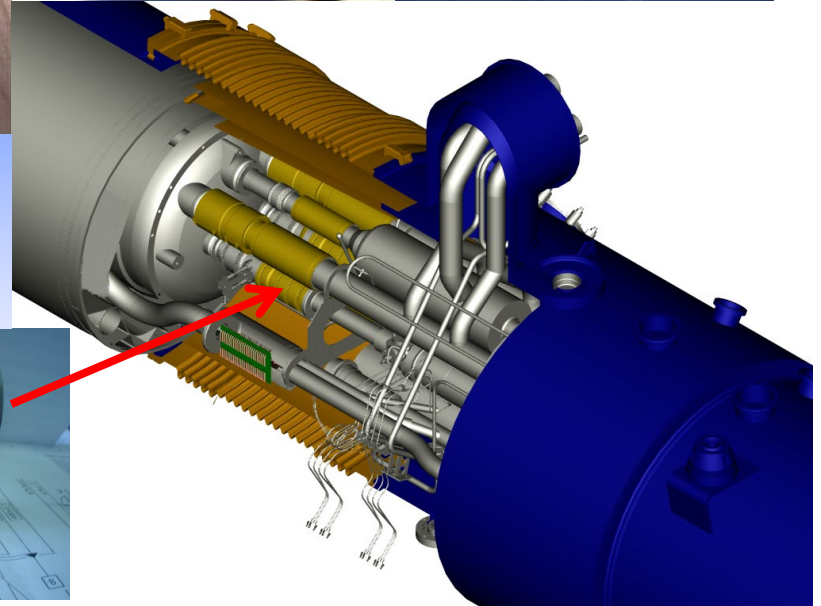
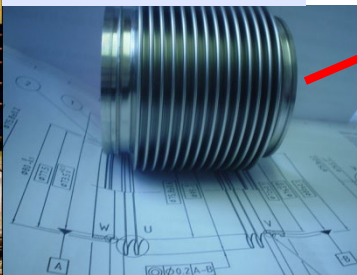
detector



# Motivation: CERN Large Hadron Collider *expansive phenomena at cryogenic temperatures*



**20000  
expansion  
bellows**





## Materials used at low temperatures: LHC, FCC, ITER

### Modern materials for extremely low temperatures (0-5 K):

#### I Metals:

Al, Ag, Be, Cu, Sn, Ta, Ti, Zr, etc



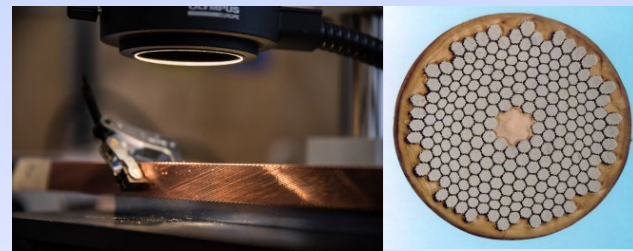
#### II Alloys:

304L, 316L, 316LN, 316Ti, P506, JK2LB, etc.



#### III Composites:

NbTi, NbAl, Nb<sub>3</sub>Sn, MgB<sub>2</sub>, Bi2212, etc.





# General program of research

*Coupled dissipative phenomena at cryogenic temperatures*

## I constitutive models of single phenomena

Strain induced phenomena at extremely low temperatures

Done

Intermittent plastic flow  
**IPF**

Plastic strain induced  
**Fcc-Bcc** phase transformation

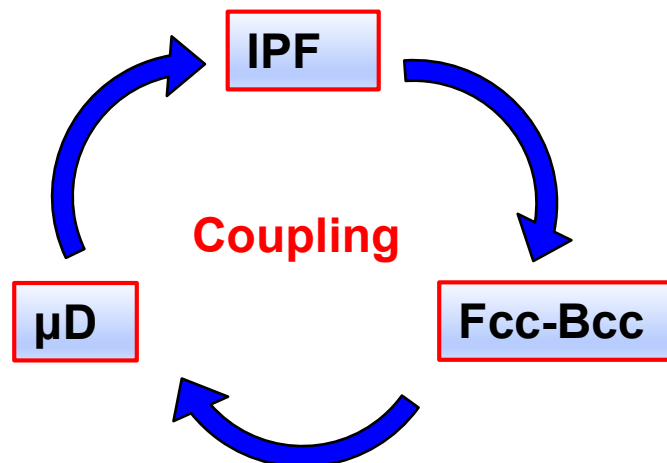
Nucleation and evolution of  
micro-damage  **$\mu$ D**

Fracture at low temperatures **LTF**

Creep at low temperatures **LTC**

In progress

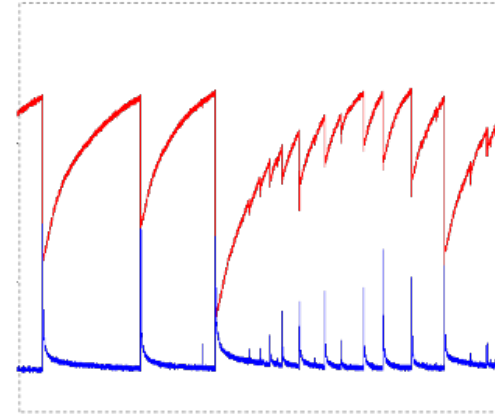
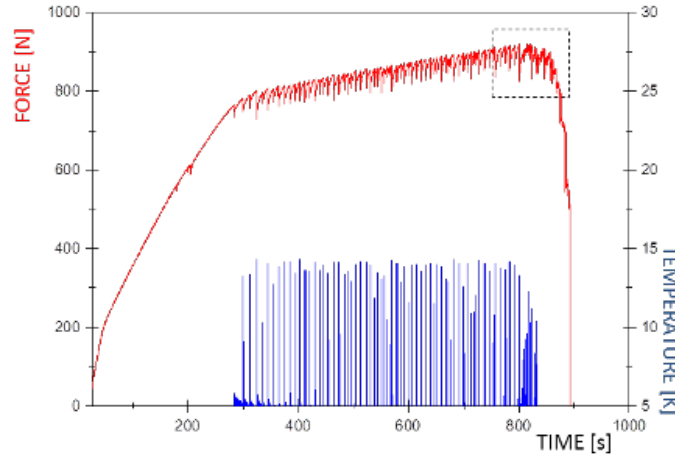
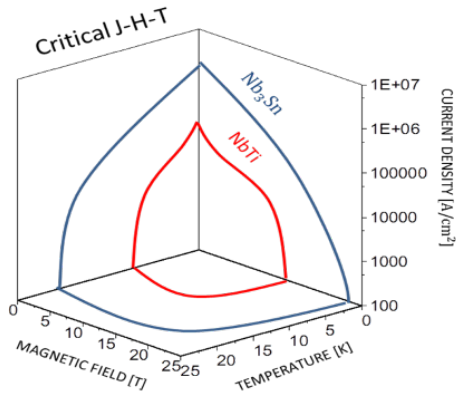
## II coupling between the phenomena



In progress



# Problem 8: IPF in the intermetallic composites (superconductors)



The critical surface for **NbTi** and **Nb<sub>3</sub>Sn** filaments

The intermittent plastic flow in NbTi specimen in LHe (4.2 K)

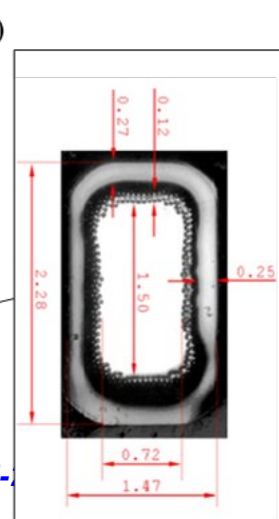
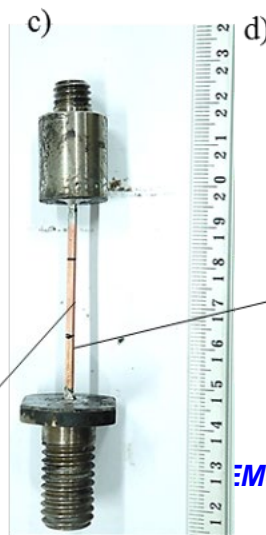
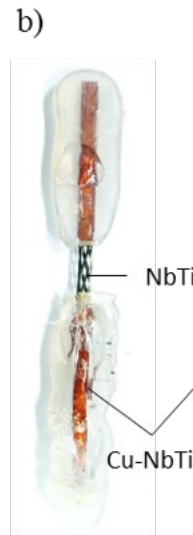
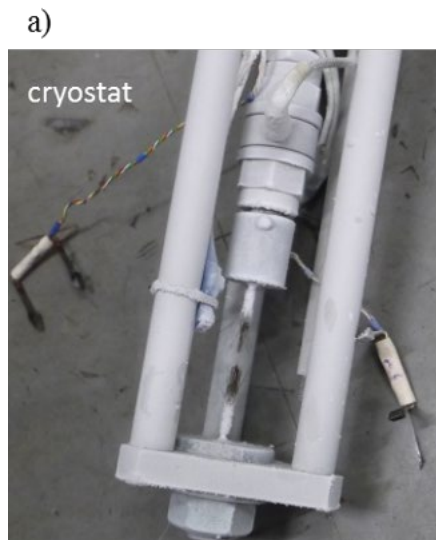
Superconducting multifilament composite based on NbTi wires:

a) the specimen in the test set-up (cryostat insert),

b) NbTi filaments extracted from the copper matrix,

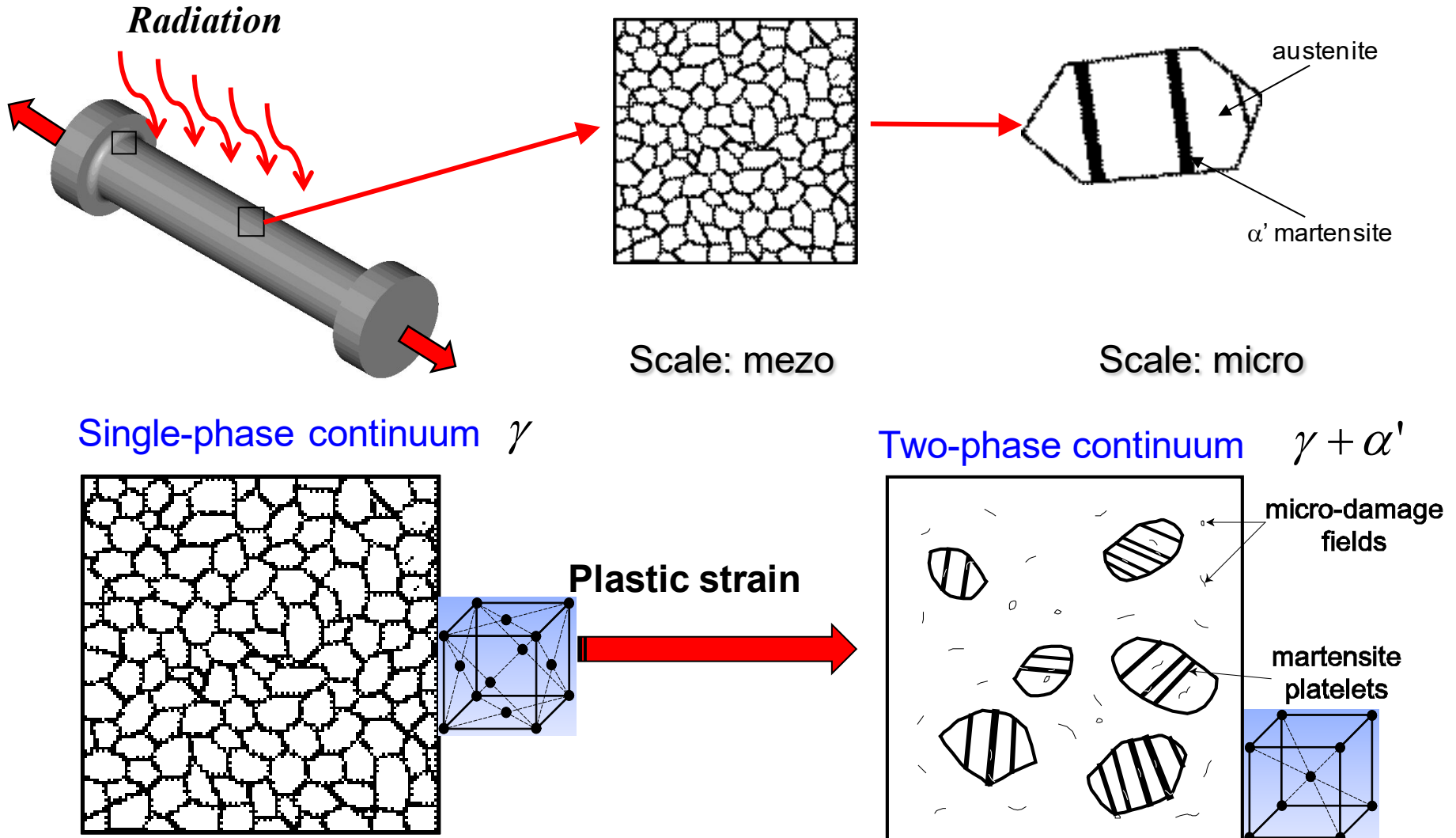
c) the NbTi specimen ready for testing,

d) the cross section of the NbTi specimen.





# Coupled field problems: radiation versus phase transformation Pres

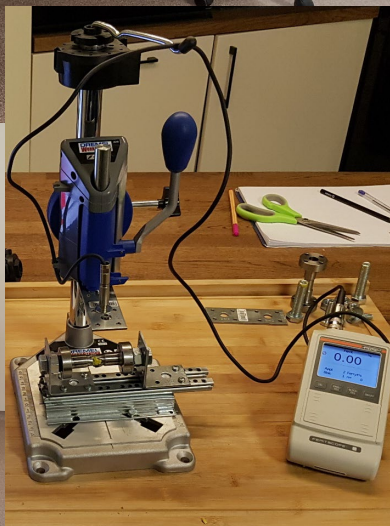






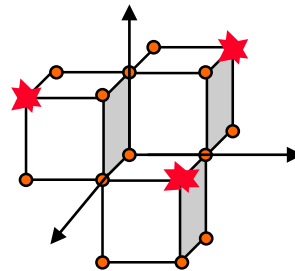
# Laboratory of Extremely Low Temperatures **RC-LENT**

*ogenic temperatures*





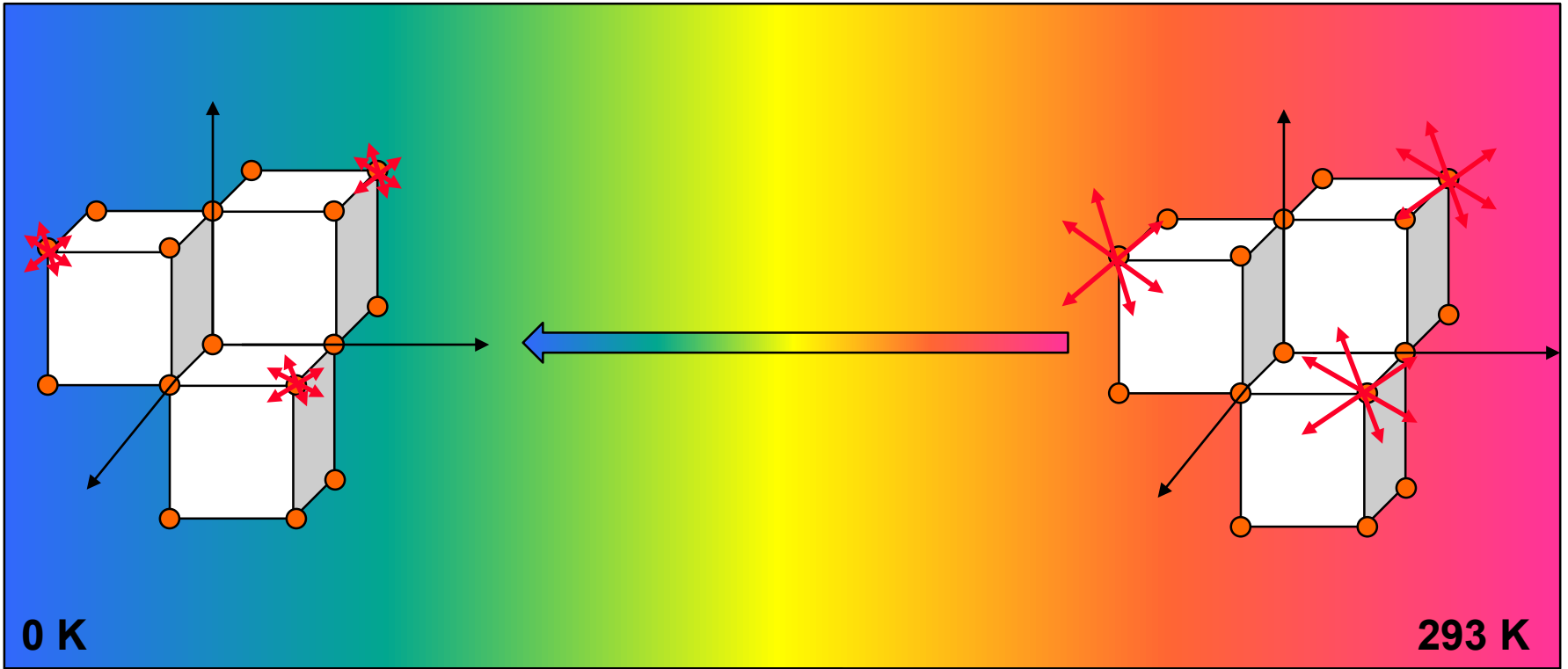
# Thermodynamics in the proximity of absolute zero





# Thermodynamics of lattice at low temperatures

phenomena at cryogenic temperatures



Temperature = measure of lattice excitation

$$\lim_{T \rightarrow 0} S(T, V, N) = 0$$

III-rd law of thermodynamics (Nernst; for perfect crystals)



Hamilton operator of lattice vibrations:

$$H = H_0 + \sum_{\alpha, \underline{k}} \hbar \omega_{\alpha}(\underline{k}) a_{\alpha}^{\dagger}(\underline{k}) a_{\alpha}(\underline{k})$$

Number of phonons operator and its eigenvalues:

$$a_{\alpha}^{\dagger}(\underline{k}) a_{\alpha}(\underline{k}) \longrightarrow N_{\alpha}(\underline{k})$$

Energy of lattice vibrations:

$$E = E_0 + \sum_{\alpha, \underline{k}} \hbar \omega_{\alpha}(\underline{k}) N_{\alpha}(\underline{k})$$

Phonon energy of lattice:

$$\Delta E = \sum_{\alpha, \underline{k}} \hbar \omega_{\alpha}(\underline{k}) f_0(\omega_{\alpha}) \Rightarrow \int_0^{\omega_{\max}} f_0(\omega) \nu(\omega) \hbar \omega d\omega$$

$$E_{ph} \sim N \left( \frac{T}{\Theta} \right)^3 T$$



# Thermodynamics of lattice at low temperatures

phenomena at cryogenic temperatures

Specific heat at constant volume:

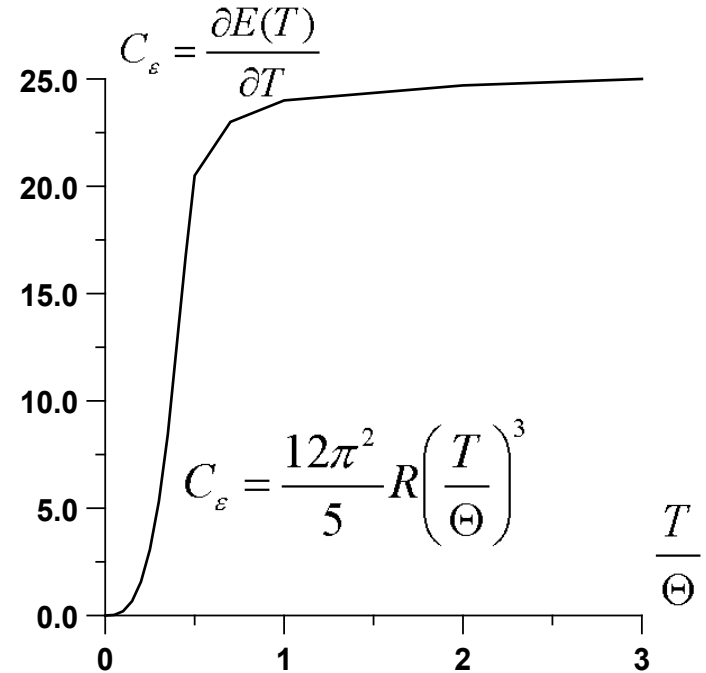
$$(C_{ph})_V = \left( \frac{\partial E_{ph}}{\partial T} \right)_V \sim N \left( \frac{T}{\Theta} \right)^3 \quad (C_{el})_V \sim \frac{T}{T_F}$$

$$C_V = (C_{el})_V + (C_{ph})_V \sim aT + bT^3$$

$$\lim_{T \rightarrow 0} C_V = 0$$

$$dQ = mC_V dT \quad \frac{dT}{dQ} = \frac{1}{mC_V}$$

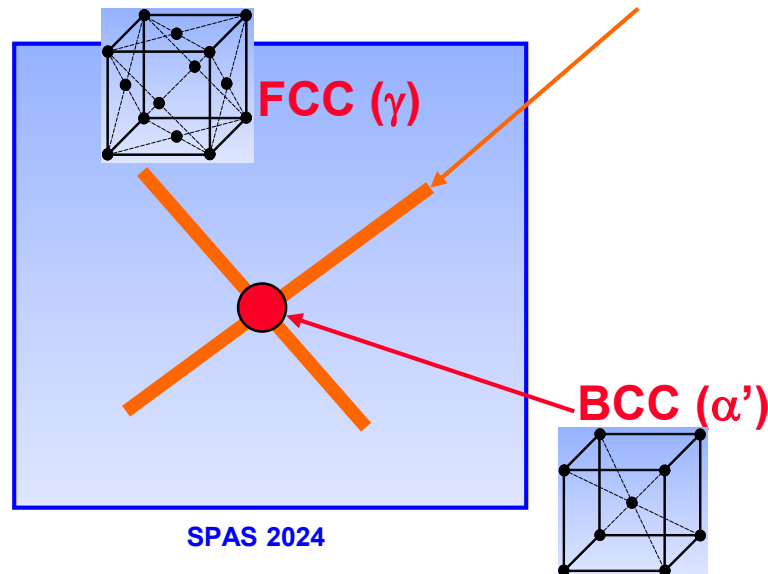
$$T \rightarrow 0 \Rightarrow \frac{dT}{dQ} \rightarrow \infty$$



Quantum  $\Delta Q$  causes large  $\Delta T \rightarrow$  “thermodynamic instability”



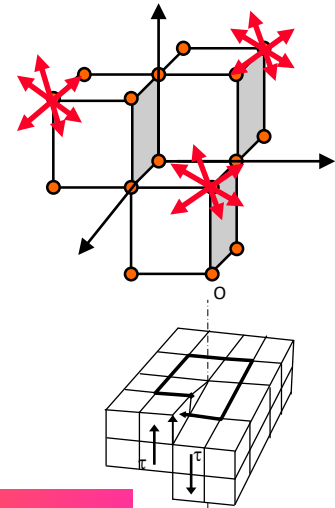
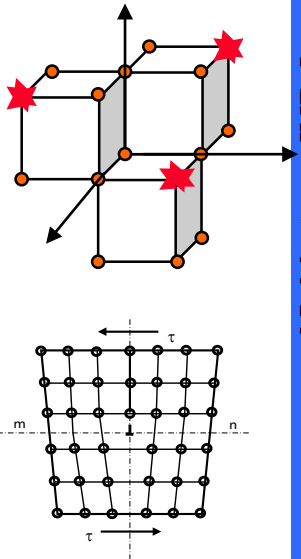
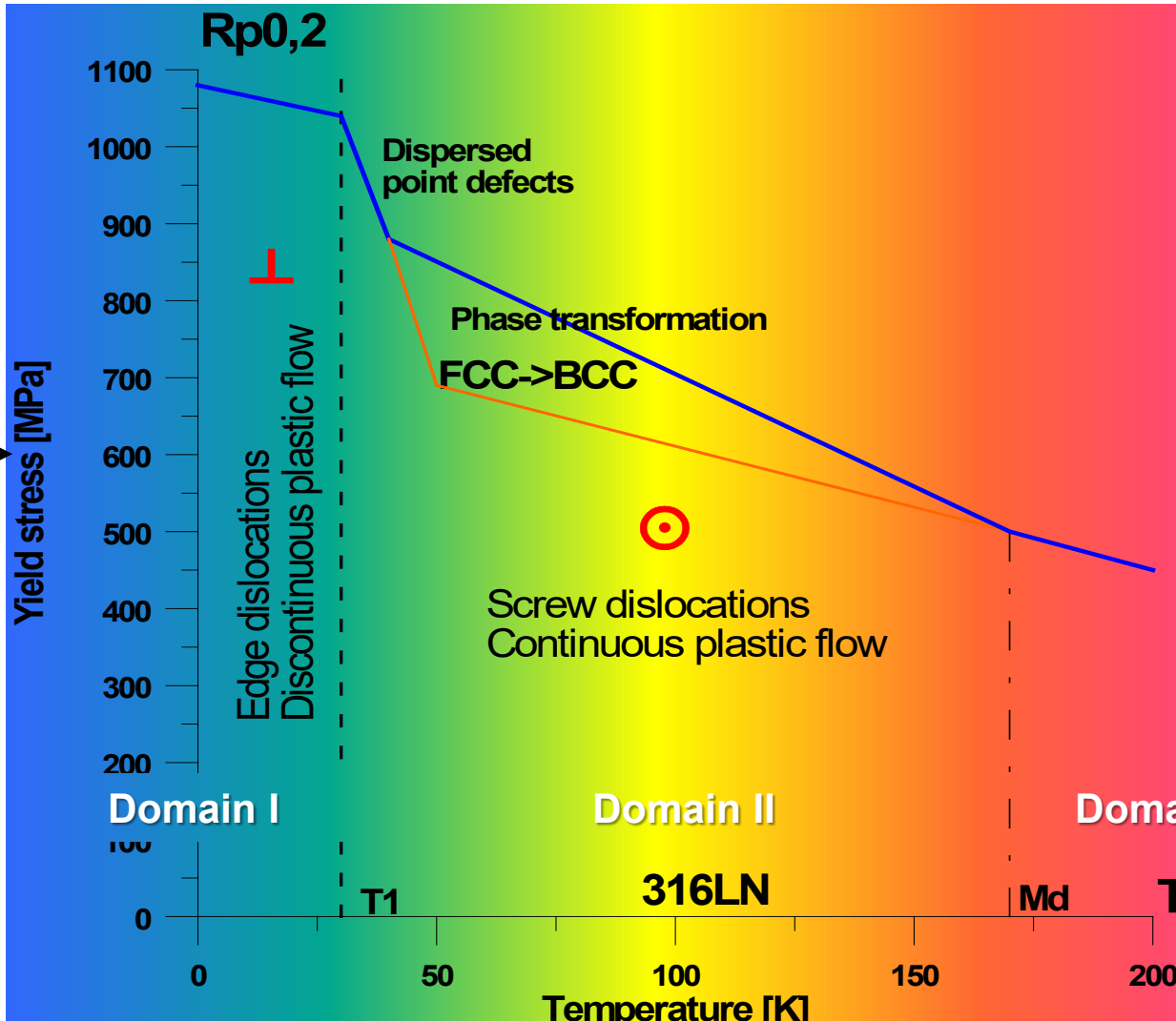
# Strain induced fcc-bcc phase transformation





# Mechanisms of plastic flow at cryogenic temperatures

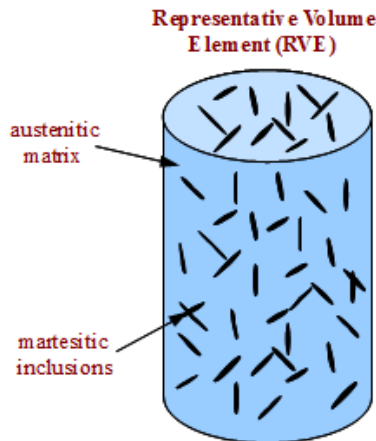
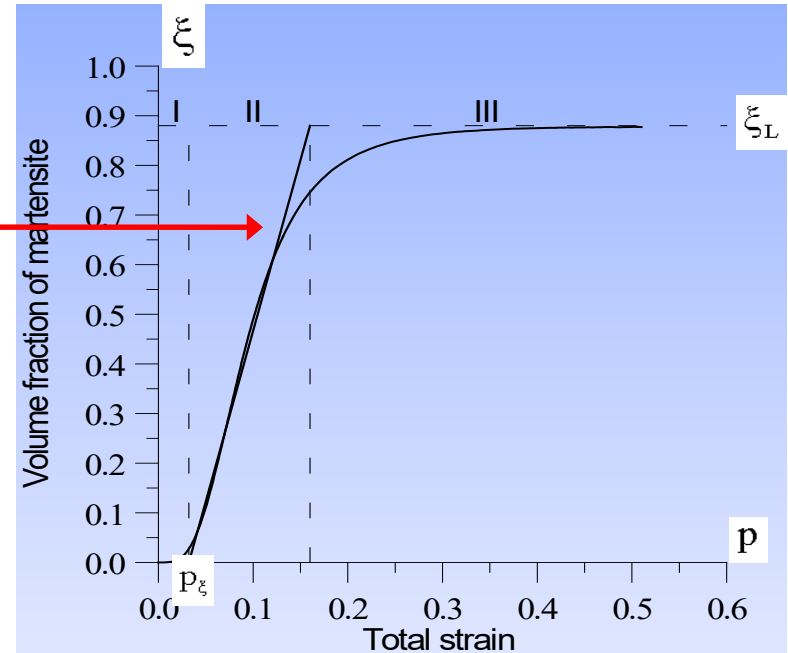
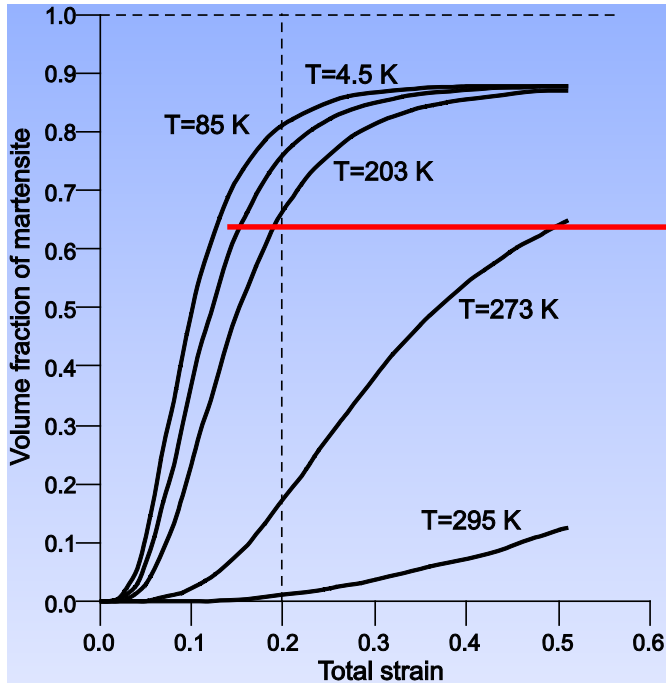
...nic temperatures





# Kinetics of Fcc-Bcc phase transformation

phenomena at cryogenic temperatures



$$\xi = \frac{dV_{\xi}}{dV} ; 0 \leq \xi \leq 1$$

$$\dot{\xi} = A(T, \dot{\underline{\underline{\epsilon}}}^p, \underline{\underline{\sigma}}) \dot{p} H((p - p_{\xi})(\xi_L - \xi))$$

$\xi$  – volume fraction of  $\alpha'$  phase





# Micromechanics: transformation strain

...ative phenomena at cryogenic temperatures

$$\underline{\underline{\varepsilon}}^{bs} = \frac{1}{V} \int_V \underline{\underline{\varepsilon}}_{\mu}^{bs} dV$$

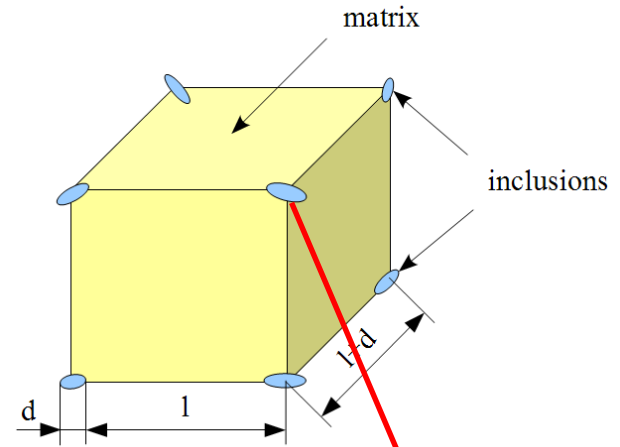
$\approx 0$

$$\underline{\underline{\varepsilon}}^{bs} = \frac{V_\gamma}{V} \frac{1}{V_\gamma} \int_{V_\gamma} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV + \frac{V_\alpha}{V} \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV$$

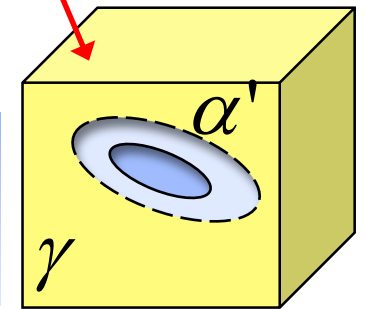
$$\underline{\underline{\varepsilon}}^{bs} = \frac{V_\alpha}{V} \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV = \xi \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV = \xi \langle \underline{\underline{\varepsilon}}_{\mu}^{bs} \rangle$$

$$\underline{\underline{\varepsilon}}_{\mu}^{bs} = \begin{pmatrix} 0 & 0 & \frac{\gamma}{2} \\ 0 & 0 & 0 \\ \frac{\gamma}{2} & 0 & \Delta v \end{pmatrix}_{(\bar{x}, \bar{y}, \bar{z})}$$

$$\langle \underline{\underline{\varepsilon}}_{\mu}^{bs} \rangle = \frac{1}{3} \Delta v \underline{\underline{I}}$$



**Type Eshelby  
ellipsoidal  
inclusion**



$$\underline{\underline{\sigma}} = \underline{\underline{E}} : \left( \underline{\underline{\varepsilon}} - \underline{\underline{\varepsilon}}^p - \underline{\underline{\varepsilon}}^{th} - \xi \underline{\underline{\varepsilon}}^{bs} \right)$$

$$\underline{\underline{\varepsilon}}^{bs} = \xi \frac{1}{3} \Delta v \underline{\underline{I}}$$



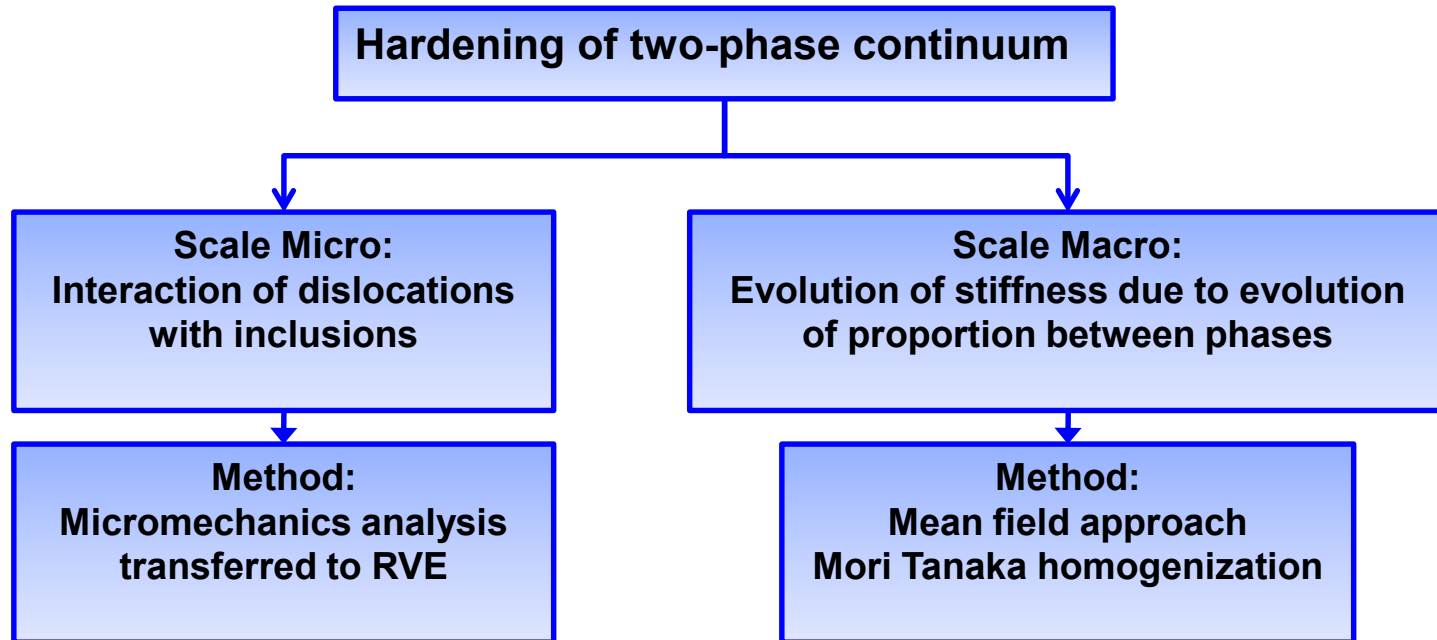
# Constitutive description of two-phase continuum at cryogenic temperatures

**Yield condition:**  $f_c(\underline{\underline{\sigma}}, \underline{\underline{X}}, R) = J_2(\underline{\underline{\sigma}} - \underline{\underline{X}}) - \sigma_y - R = 0$       $J_2(\underline{\underline{\sigma}} - \underline{\underline{X}}) = \sqrt{\frac{3}{2}(\underline{\underline{s}} - \underline{\underline{X}}) : (\underline{\underline{s}} - \underline{\underline{X}})}$

**Mixed hardening** depending on the phase transformation parameter:

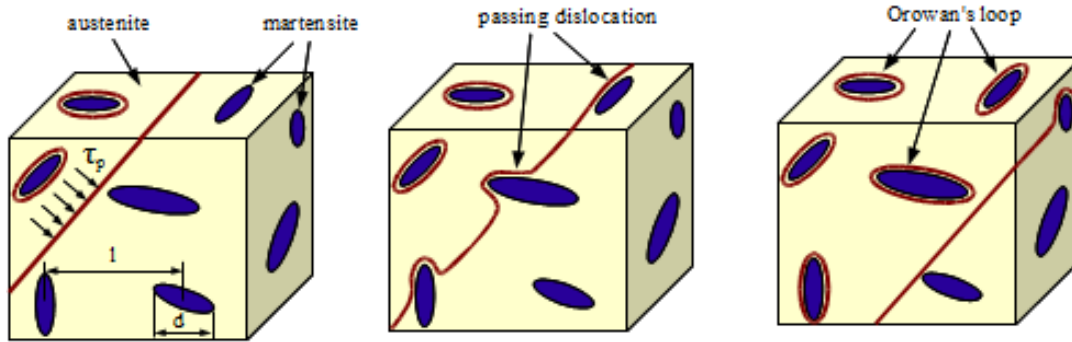
$$d\underline{\underline{X}} = d\underline{\underline{X}}_a + d\underline{\underline{X}}_{a+m} = \frac{2}{3} C_X(\xi) d\underline{\underline{\varepsilon}}^p$$

$$dR = C_R(\xi) dp$$





# Scale Micro: interaction of dislocations with inclusions at low temperatures

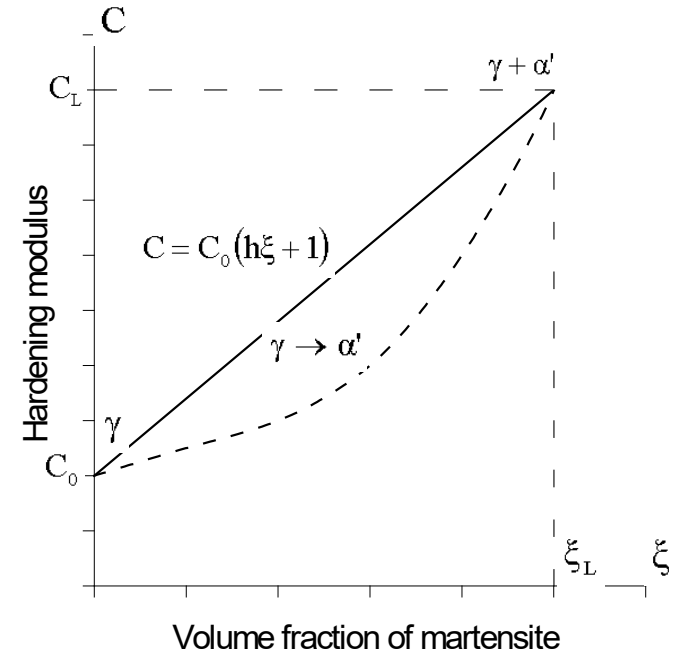


$$dX_a = \frac{2}{3} C_0 d \varepsilon^p$$

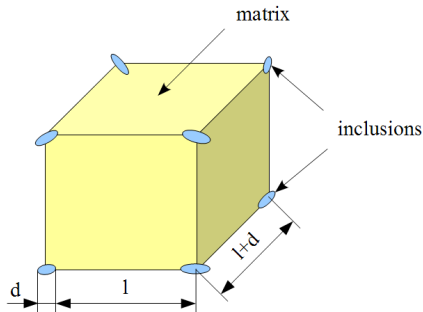
**Initial hardening of matrix (austenite)**

$$dX_a = \frac{2}{3} C_0 \phi(\xi) d \varepsilon^p$$

**Hardening of matrix containing inclusions**



## Micromechanics analysis



$$\tau_p = \frac{Gb}{d} \left( \frac{6\xi_0}{\pi} \right)^{\frac{1}{3}} \left( 1 + \frac{\xi - \xi_0}{3\xi_0} \right)$$



$$\phi(\xi) = 1 + h\xi ; \quad 0 \leq \xi \leq 1$$

$$C = C_0 \phi(\xi)$$

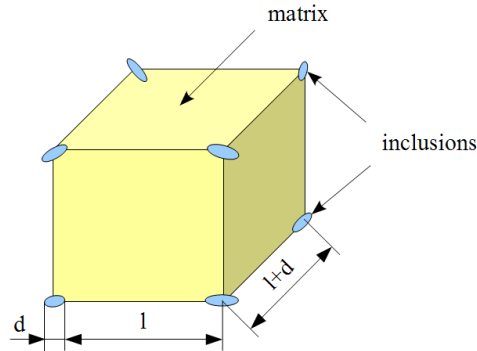


# Scale Macro: evolution of proportion between phases ogenic temperatures

Elastic-plastic matrix:

$$\underline{\underline{\Delta\sigma}}_a = \underline{\underline{E}}_{ta} : \underline{\underline{\Delta\varepsilon}}$$

$$\underline{\underline{E}}_{ta} = 3k_a \underline{\underline{J}} + 2\mu_a \underline{\underline{K}} - 2\mu_a \frac{\underline{\underline{n}} \otimes \underline{\underline{n}}}{1 + \frac{C(\xi)}{3\mu_a}}$$



Elastic inclusions:

$$\underline{\underline{\Delta\sigma}}_m = \underline{\underline{E}}_m : \underline{\underline{\Delta\varepsilon}}$$

$$\underline{\underline{E}}_m = 3k_m \underline{\underline{J}} + 2\mu_m \underline{\underline{K}}$$

$$\mu_m = \frac{E}{2(1+\nu)} \quad k_m = \frac{E}{3(1-2\nu)}$$



„Linearization”: extraction of isotropic part of tangent stiffness operator

$$\underline{\underline{E}}_{ta} = 3k_{ta} \underline{\underline{J}} + 2\mu_{ta} \underline{\underline{K}}$$

$$\mu_{ta} = \frac{E_t}{2(1+\nu)} \quad k_{ta} = \frac{E_t}{3(1-2\nu)}$$

$$\underline{\underline{E}}_t = \frac{EC}{E+C}$$



Homogenization:

$$\underline{\underline{\Delta\sigma}} = \underline{\underline{E}}_H : \underline{\underline{\Delta\varepsilon}}$$



# Constitutive description of two-phase continuum at cryogenic temperatures

Kinematic hardening

$$\Delta X_{\equiv a+m} = \Delta \sigma_{\equiv a+m}$$

$$dX_{\equiv a+m} = \frac{2}{3} \beta C_{a+m}(\xi) d\varepsilon^p$$

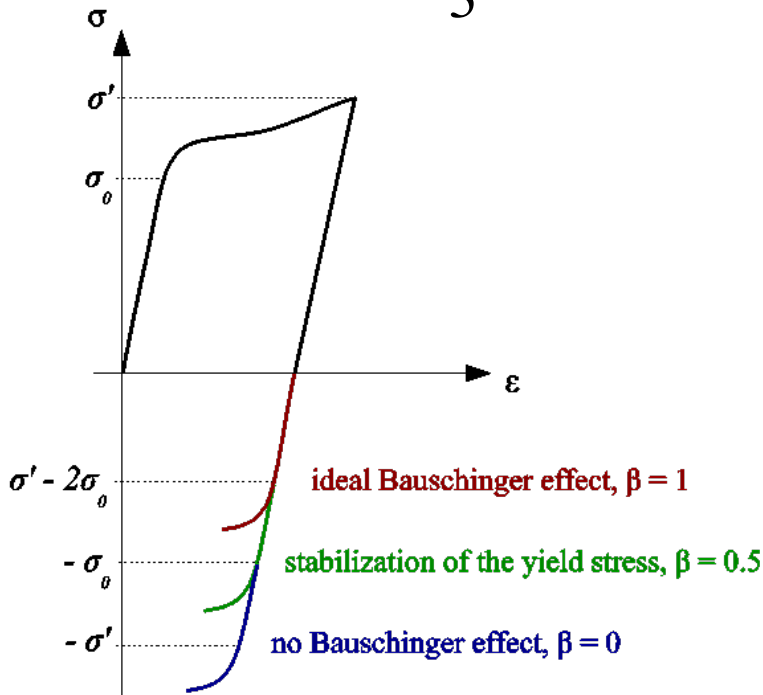
$$\beta = \frac{\sigma' + \sigma'^-}{2(\sigma' - \sigma_0)}$$

Parametization: *Życzkowski, 1981*

Isotropic hardening

$$\Delta R = \|\Delta \sigma_{a+m}\|$$

$$dR = (1 - \beta) C_{a+m}(\xi) dp$$



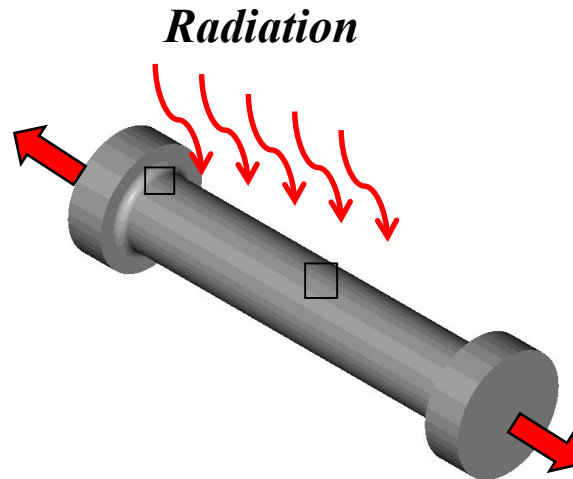
## Evolution laws of hardening parameters

$$dR = C_R(\xi) dp = (1 - \beta) C_{a+m}(\xi) dp$$

$$dX_{\equiv a+m} = \frac{2}{3} C_X(\xi) d\varepsilon^p = \frac{2}{3} \beta C_{a+m}(\xi) d\varepsilon^p$$



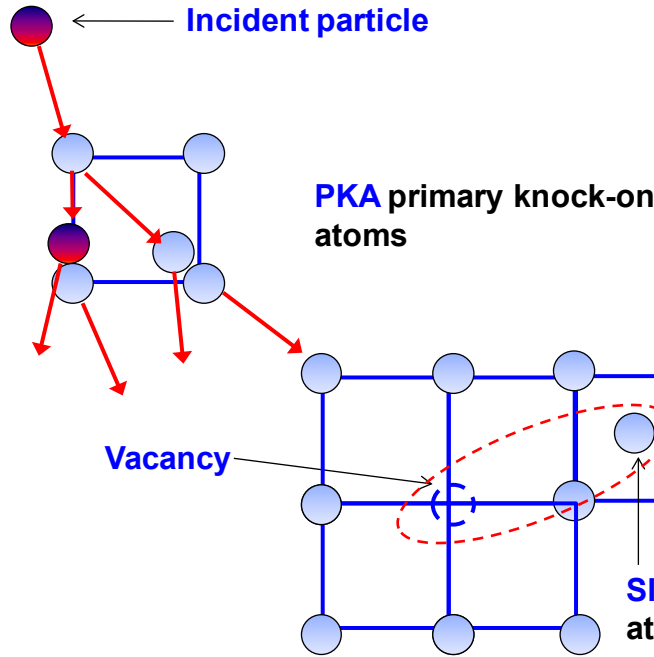
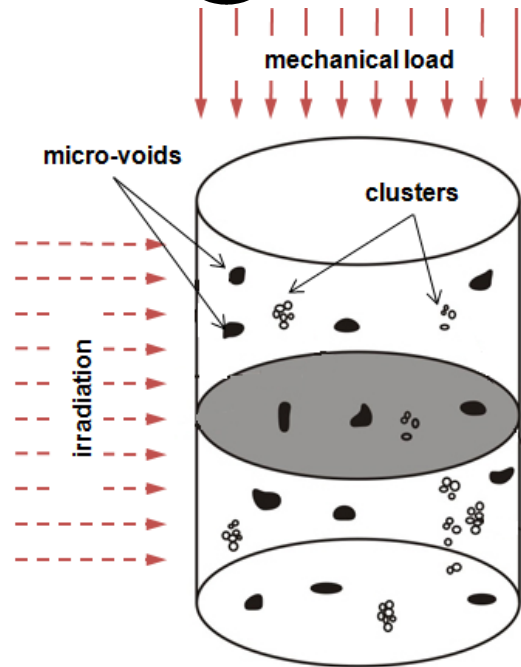
# Radiation induced damage





# Radiation induced micro-damage

*coupled dissipative phenomena at cryogenic temperatures*



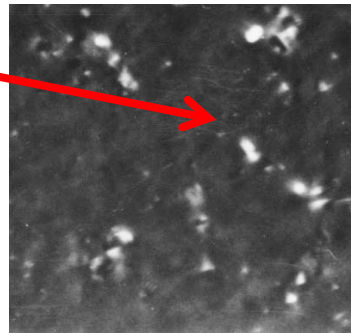
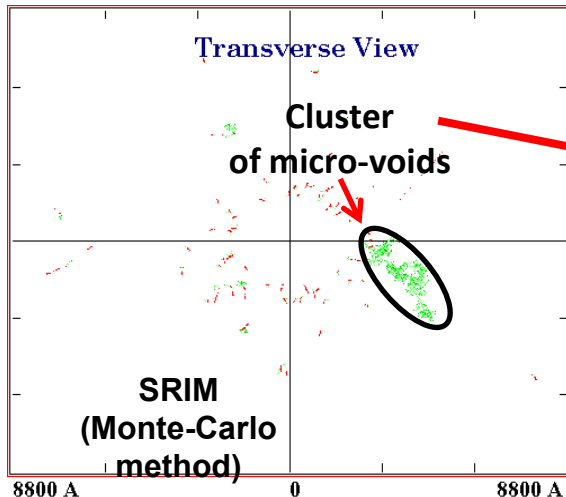
Displacement cascade and formation of Frenkel pairs

Kinchin & Pease, 1955

Norgett, Robinson, Torrens, 1975

$$N_{NRT} = \frac{0.8E_{dam}}{2\bar{E}_d}$$

$$E_{dam} = NIEL$$



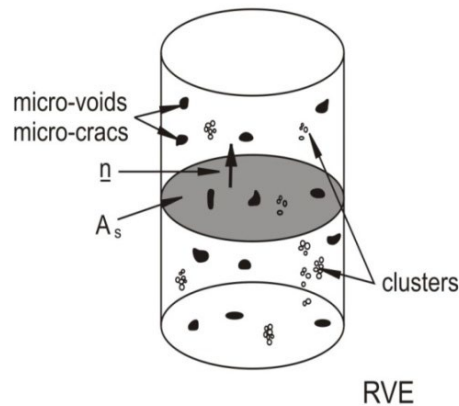
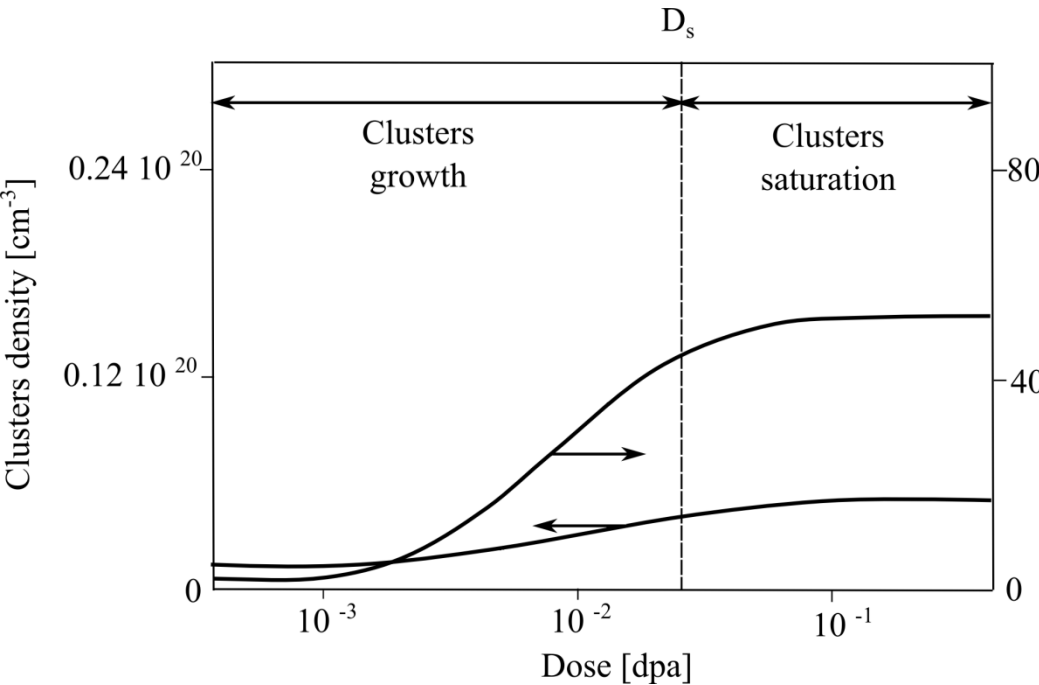
## Defects due to irradiation:

1. SFT – stacking fault tetrahedron
2. Faulted or perfect dislocation loops
3. Voids – 3D vacancy clusters
4. Cavities – 3D vacancy clusters with impurities (He)



# Lattice defects after irradiation

Coupled dissipative phenomena at cryogenic temperatures



$q_c$  – number of clusters per unit volume  
 $r_c$  – average radius of clusters

$$D = \frac{dS_D}{dS} ; 0 \leq D \leq 1$$

$$\xi = \frac{dV_D}{dV} ; 0 \leq \xi \leq 1$$

$dpa$

**Physics**

$$q_c = \begin{cases} C_{qI} (dpa)^{n_{qI}} & \text{for } dpa < D_s \\ C_{qII} (dpa)^{n_{qII}} & \text{for } dpa \geq D_s \end{cases}$$

$$r_c = \begin{cases} C_r (dpa)^{n_r} & \text{for } dpa < D_s \\ r_{cr} & \text{for } dpa \geq D_s \end{cases}$$

$$q_A = \left( \sqrt[3]{q_V} \right)^2 = q_c^{2/3}$$

$$D_{r0} = q_A \pi r_{c0}^2$$

**Mechanics**

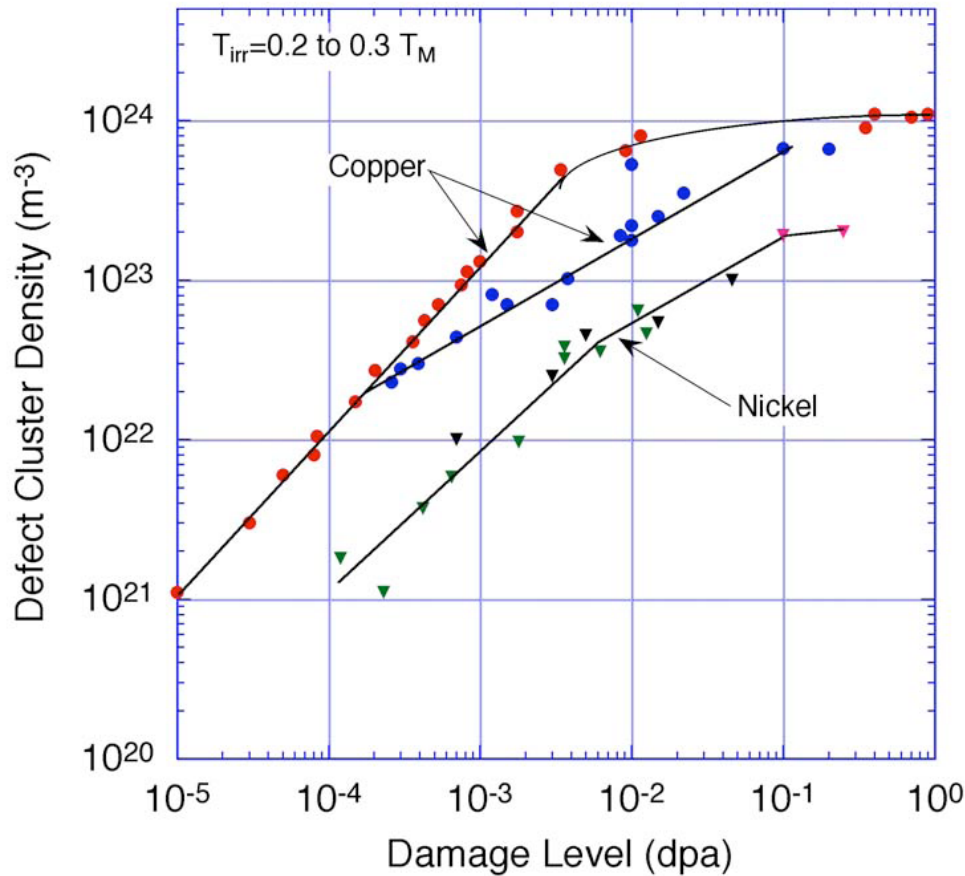




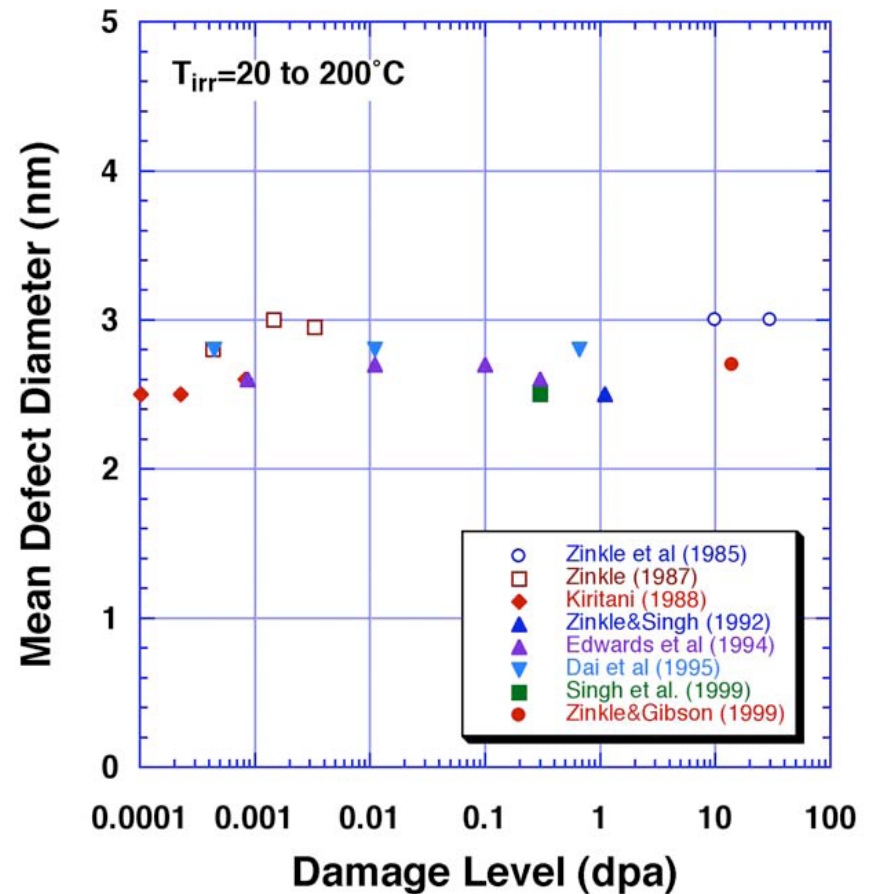
# Irradiated metals and alloys: Nickel and Copper

...ena at cryogenic temperatures

### COMPARISON OF DEFECT CLUSTER ACCUMULATION IN NEUTRON-IRRADIATED NICKEL AND COPPER



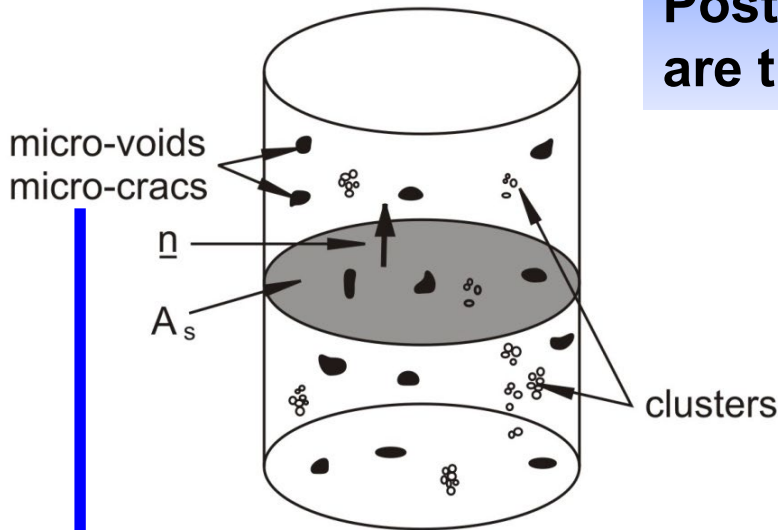
### Measured Average Image Width of Defect Clusters in Neutron and Ion-Irradiated Copper



**Source:** S.J. Zinkle „Microstructure evolution in irradiated metals and alloys: fundamental aspects”, Italy, 2004.



# Radiation and mechanical damage: additive formulation at high temperatures



RVE

**Postulate: both micro-damage components are treated in additive way**

$$D_r = D_{r0} + \int_0^{\hat{p}} dD_{rm}$$

**radiation induced damage**

$$\underline{\underline{D}}_r = \frac{D_r}{3} \underline{\underline{I}}$$

$\underline{\underline{I}}$  - identity tensor

**isotropic**

$$d\underline{\underline{D}}_m = \underline{\underline{C}} \underline{\underline{Y}} \underline{\underline{C}}^T dp$$

**anisotropic**

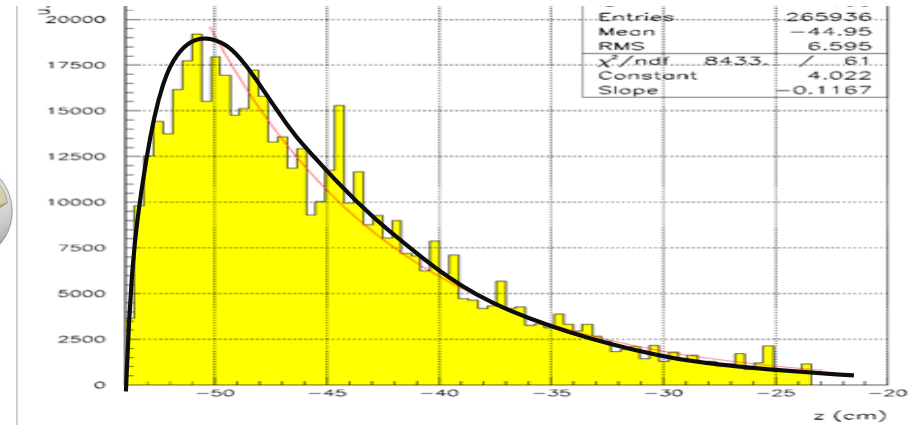
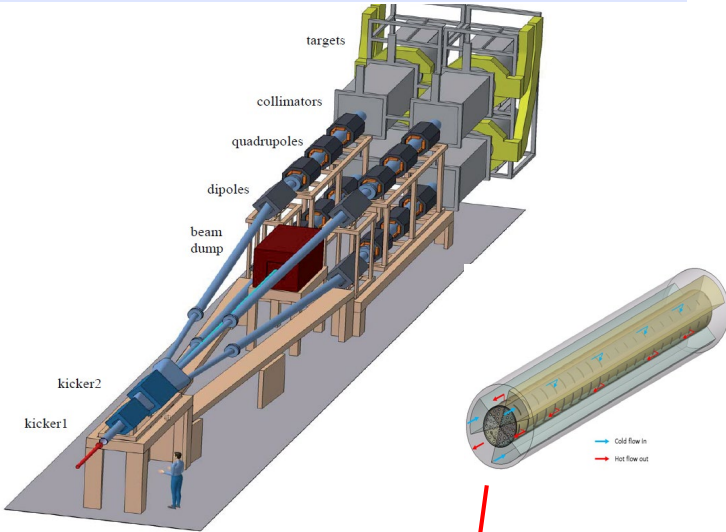
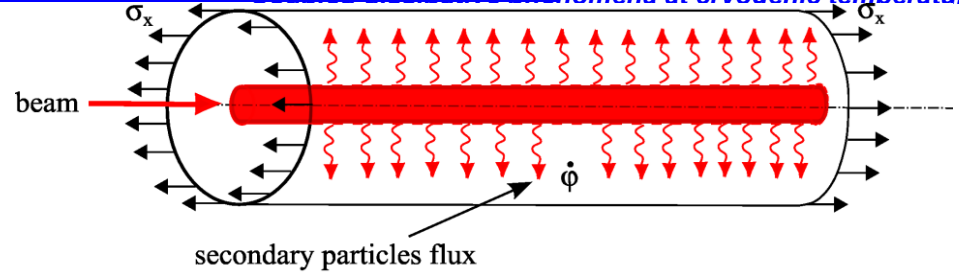
$$\underline{\underline{D}} = \underline{\underline{D}}_m + \underline{\underline{D}}_r = \underline{\underline{D}}_m + \frac{1}{3} D_r \underline{\underline{I}}$$



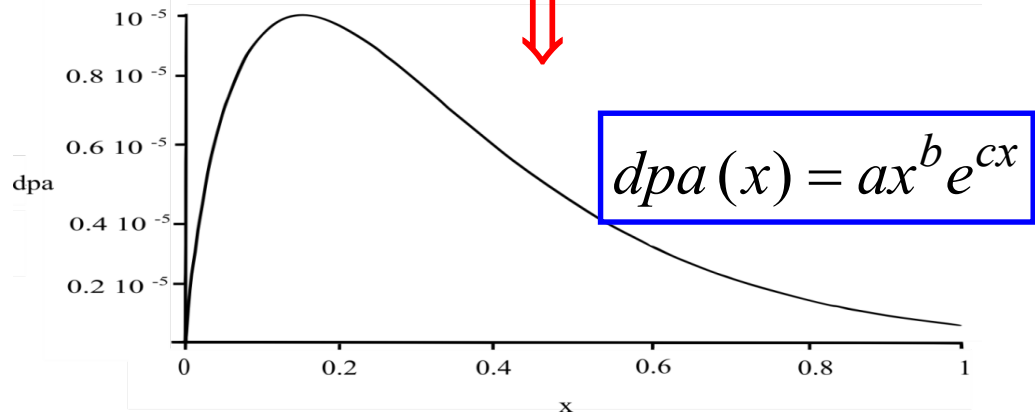
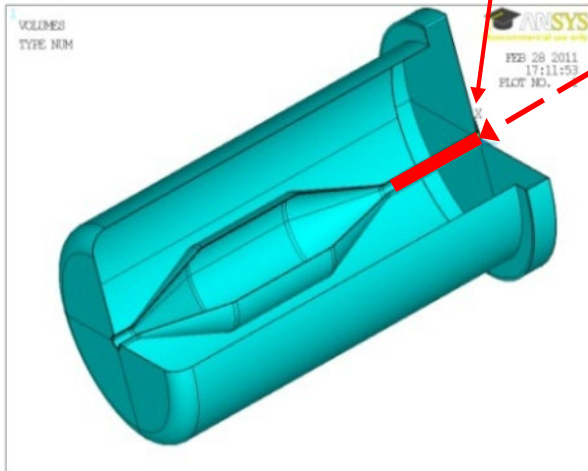
# Lifetime estimation for irradiated components

Secondary particles flux:  $\gamma$ ,  $n$ ,  $p^+$ ,  $\pi^\pm$  and  $e^\pm$

phenomena at cryogenic temperatures



Typical distribution of particle flux along the target axis





# Lifetime estimation for irradiated components

phenomena at cryogenic temperatures

## Kinetics of evolution of radiation induced damage (clusters of voids) under mechanical loads

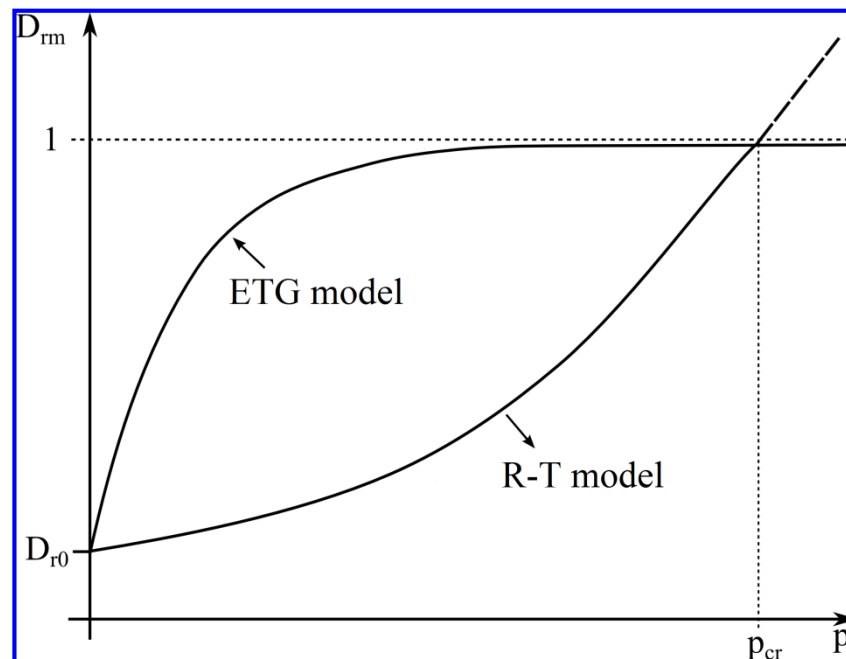
Rice&Tracey (R-T) model:

$$dr_c = r_c \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right) dp$$

Gurson (ETG) model:

$$d\xi = (1 - \xi) dp$$

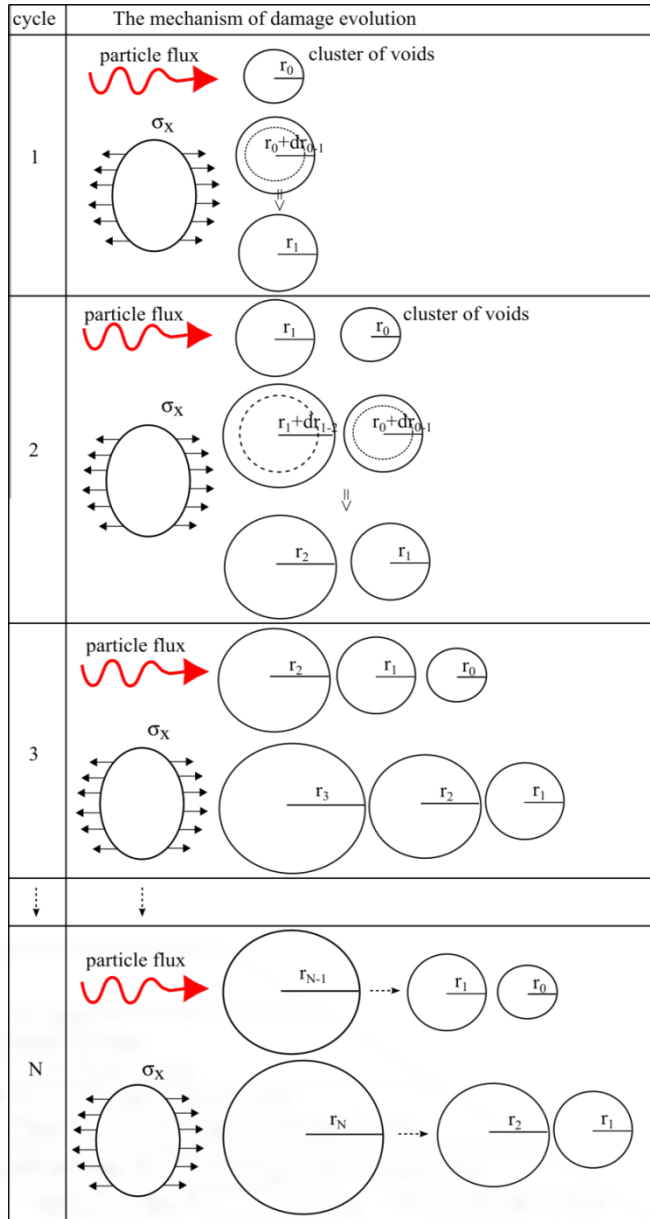
$$\dot{p} = \sqrt{\frac{2}{3} \dot{\epsilon}_{ij}^p \dot{\epsilon}_{ij}^p}$$





# Lifetime estimation for irradiated components

coupled dissipative phenomena at cryogenic temperatures



## Rice & Tracey law

$$\int_{D_i}^{D_{i+1}} dD = q_A 2\pi \int_{r_i}^{r_{i+1}} r dr$$



$$\Delta D_{i \rightarrow i+1} = q_A \pi (r_{i+1}^2 - r_i^2)$$

$$\int_{r_i}^{r_{i+1}} \frac{dr_c}{r_c} = \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right) \int_0^{\tilde{p}} dp$$



$$r_{i+1} = r_i e^{A\tilde{p}} \quad A := \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right)$$



$$\Delta D_{i \rightarrow i+1} = q_A \pi r_i^2 (e^{2A\tilde{p}} - 1)$$



# Lifetime estimation for irradiated components

Phenomena at cryogenic temperatures

$$D_{r0} = q_A \pi r_{c0}^2$$

$$D_{rm1} = D_{r0} + \Delta D_{rm(0 \rightarrow 1)} = D_{r0} + q_A \pi r_{c0}^2 (e^{2A\tilde{p}} - 1)$$

$$D_{rm2} = D_{rm1} + \Delta D_{rm(1 \rightarrow 2)} + D_{r0} + \Delta D_{rm(0 \rightarrow 1)} = 2D_{r0} + q_A \pi r_{c0}^2 e^{2A\tilde{p}} + q_A \pi r_{c0}^2 e^{4A\tilde{p}} - 2q_A \pi r_{c0}^2$$

⋮

$$D_{rmi+1} = D_{rmi} + D_{r0} + \Delta D_{rm(i \rightarrow i+1)} + \Delta D_{rm(i-1 \rightarrow i)} + \dots + \Delta D_{rm(0 \rightarrow 1)}$$

$$D_{rmN} = \underbrace{q_A \pi r_{c0}^2 e^{2A\tilde{p}} + q_A \pi r_{c0}^2 e^{4A\tilde{p}} + q_A \pi r_{c0}^2 e^{6A\tilde{p}} + \dots + q_A \pi r_{c0}^2 e^{2NA\tilde{p}}}_{\text{Geometric series}}$$

Geometric series

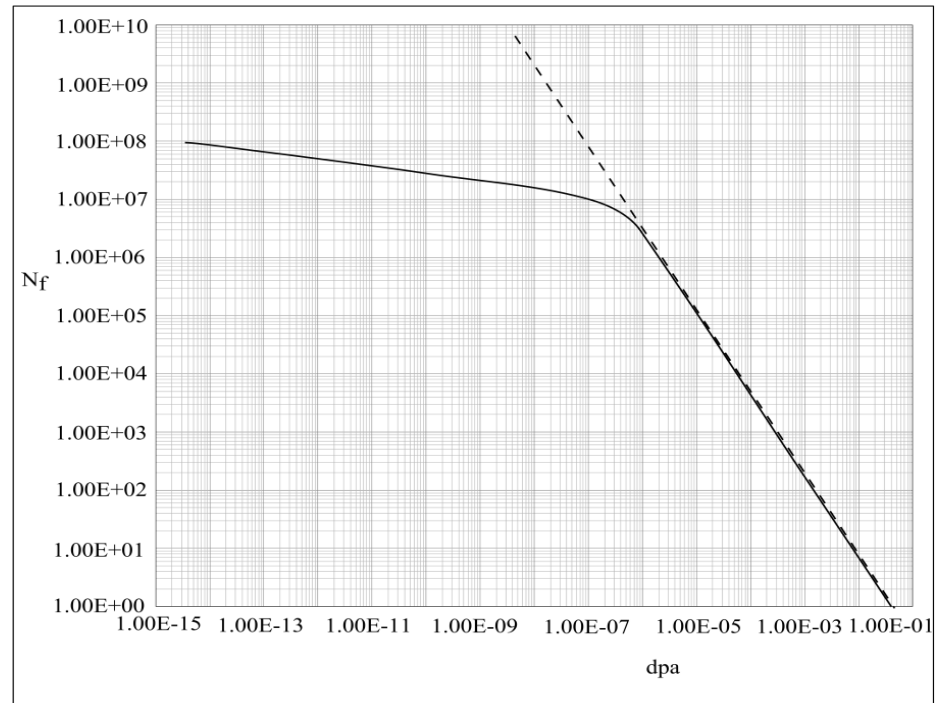
$$D_{rmN} = q_A \pi r_{c0}^2 \sum_{n=1}^N e^{2nA\tilde{p}}$$

$$S_N = a_1 \frac{1 - q^N}{1 - q} \quad S_N = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1 - e^{2ApN}}{1 - e^{2Ap}}$$

$$D_{rmN} = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1 - e^{2ApN}}{1 - e^{2Ap}}$$

$$D_{rmN_f} = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1 - e^{2A\tilde{p}N_f}}{1 - e^{2A\tilde{p}}} = D_{cr}$$

Number of cycles to failure  $N_f$  based on the critical damage criterion:  $D_{rm} = D_{cr}$

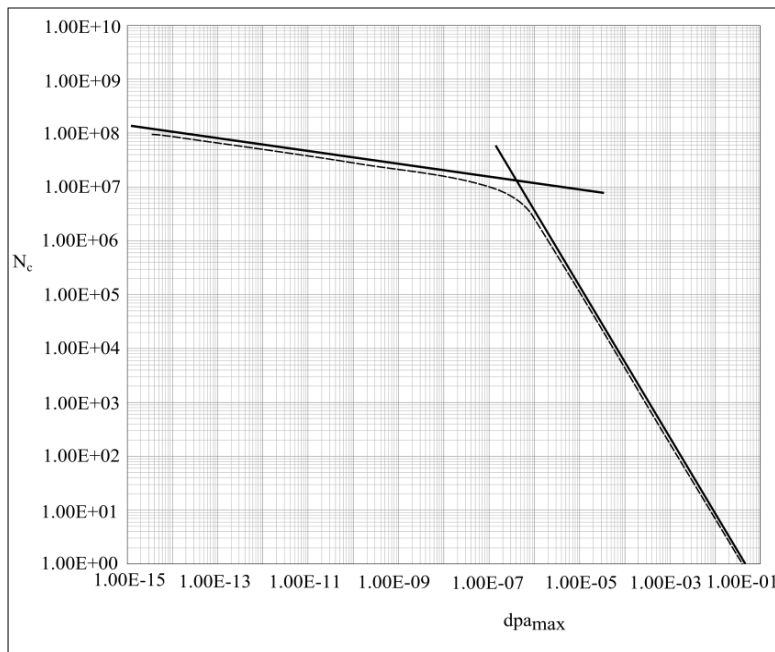




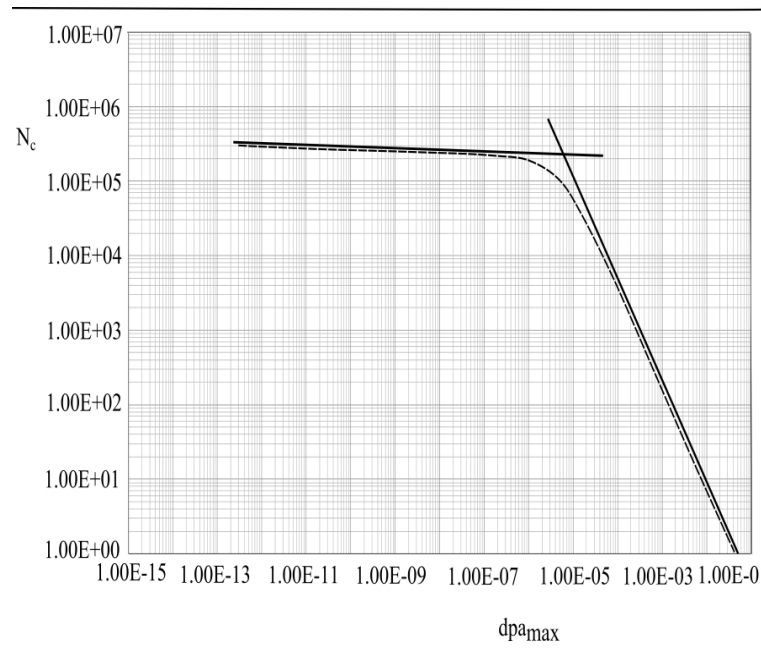
# Bilinear approximations for R-T and Gurson models

temperatures

## Rice & Tracey model



## Gurson model



$$\log(N_c) = a + b \log(dpa_{\max})$$

**Analytical formula - useful tool for estimation of number of cycles to failure**

$$N_c = 10^a dpa_{\max}^b$$

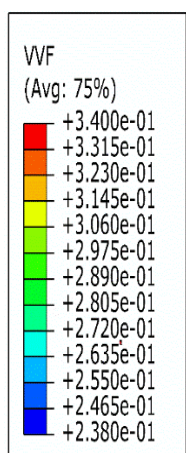
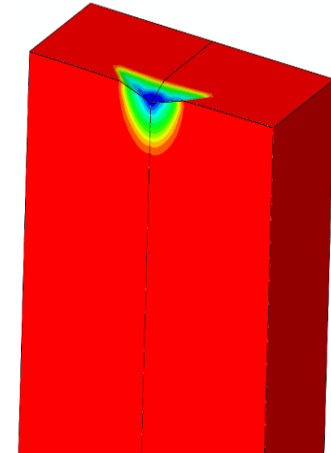
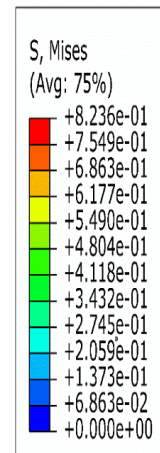
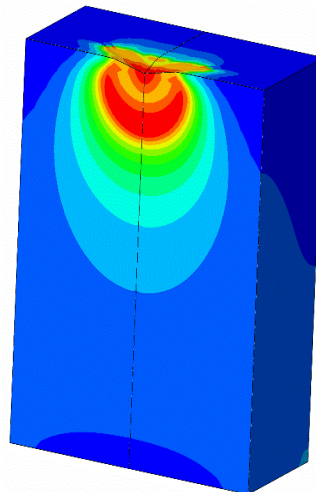
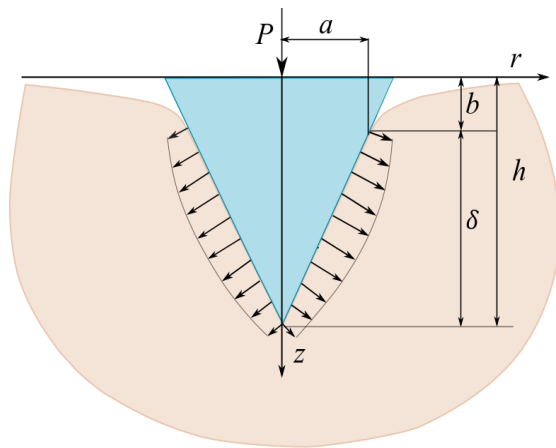
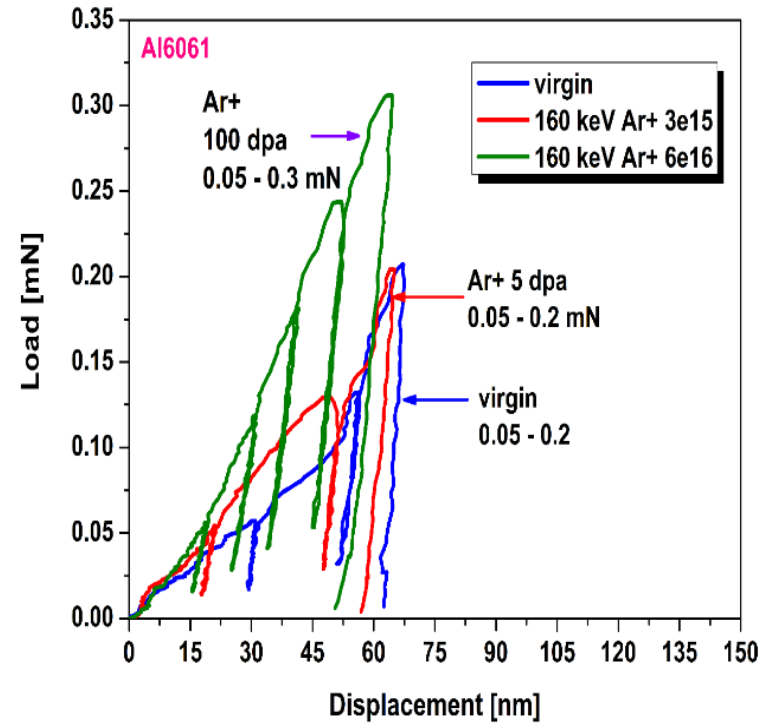
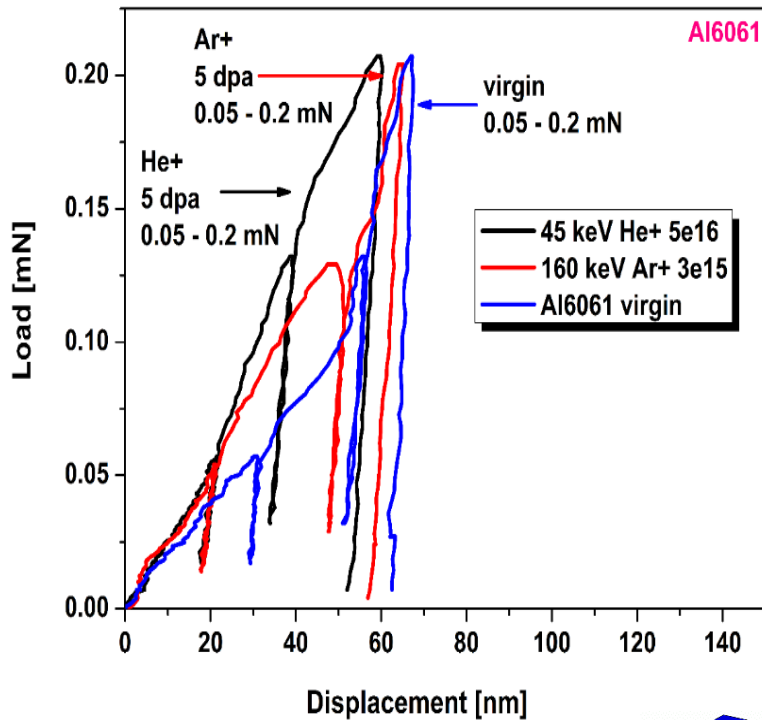
$$N_c = \begin{cases} 10^{-1.9} dpa_{\max}^{-1.4} & \text{for } dpa_{\max} \geq 10^{-6} \\ 10^{6.1} dpa_{\max}^{-0.13} & \text{for } dpa_{\max} < 10^{-6} \end{cases}$$

$$N_c = \begin{cases} 10^{-1.9} dpa_{\max}^{-1.4} & \text{for } dpa_{\max} \geq 10^{-5} \\ 10^{5.43} dpa_{\max}^{-0.016} & \text{for } dpa_{\max} < 10^{-5} \end{cases}$$



# Nanoindentation of irradiated Al6061

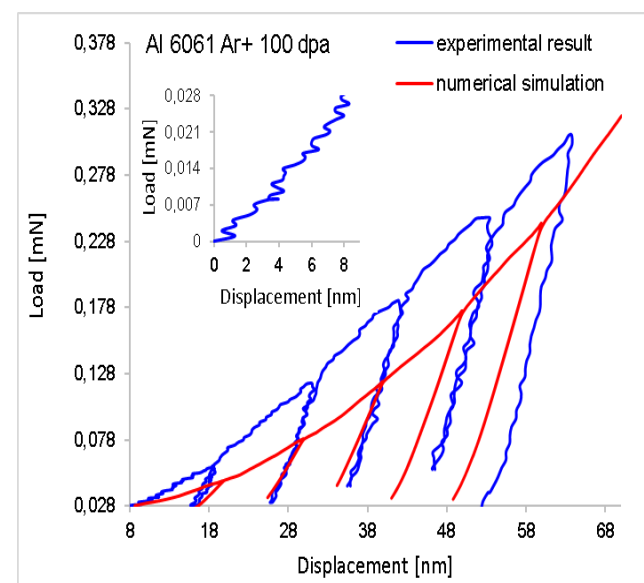
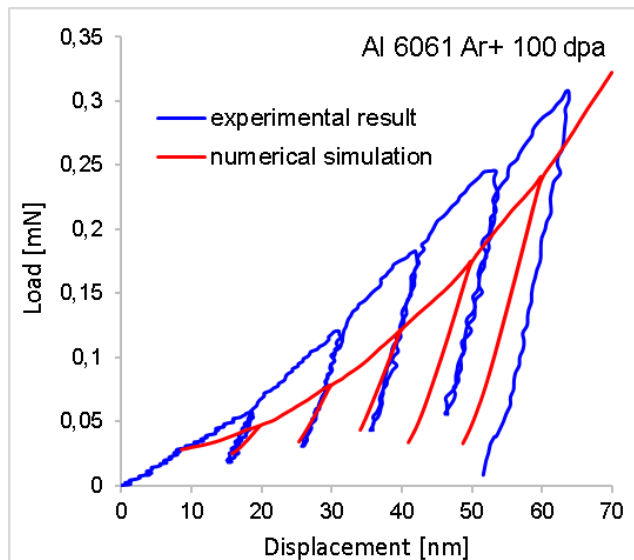
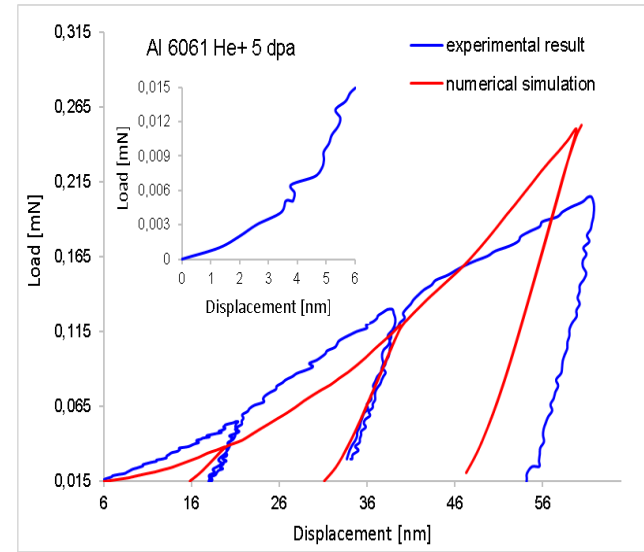
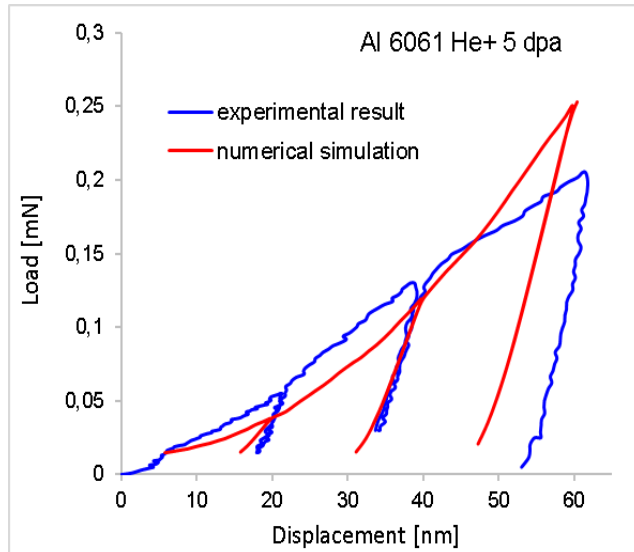
... phenomena at cryogenic temperatures





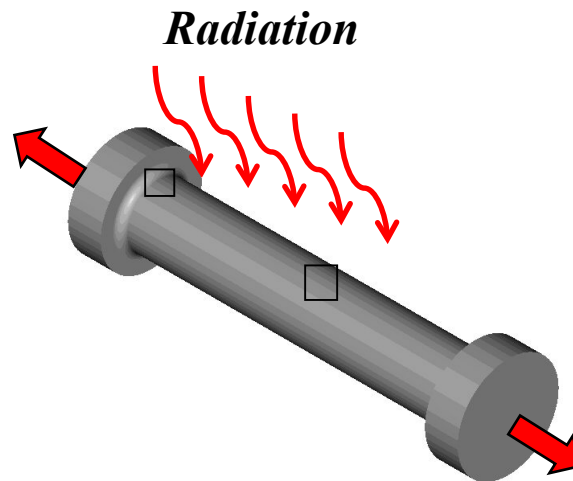


# Nanoindentation of irr. Al6061: experiment vs. Gurson model





# Radiation induced hardening



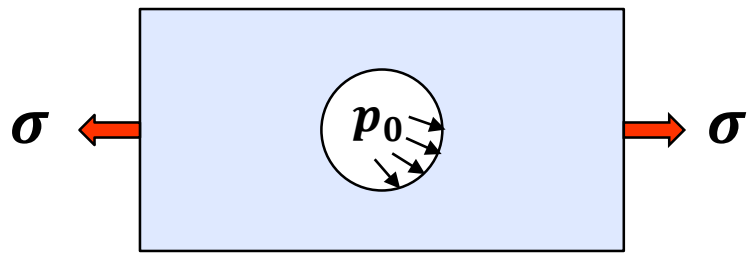


# Type Eshelby entities: the equivalent inclusion

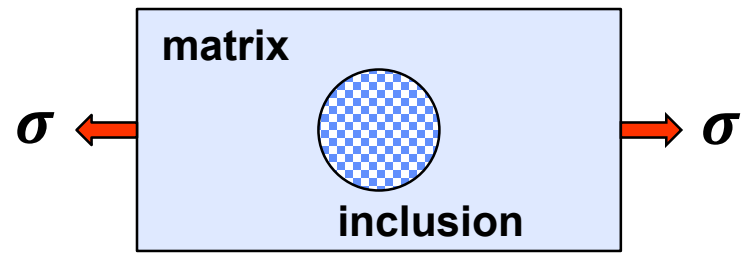
...ena at cryogenic temperatures

## Assumptions:

- small strains approach
- perfect gas inside the void at a constant temperature  $T$
- pressurized void is equivalent to inclusion subjected to hydrostatic stress



$$\Delta p = -3p_0 \Delta \varepsilon$$

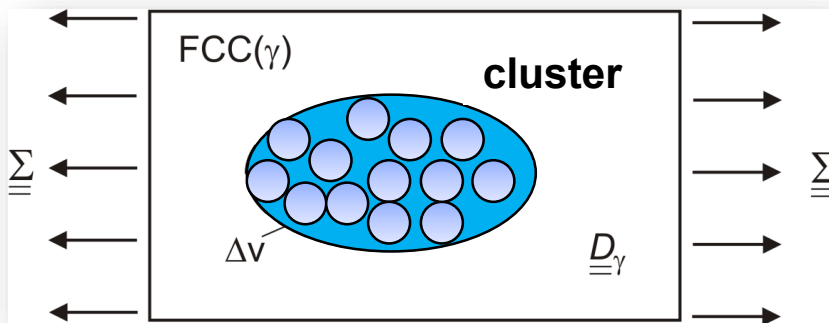


$$E_{pijkl} = 3k_p J_{ijkl} + 2\mu_p K_{ijkl}$$

$$\mu_p = 0 ; \nu_p = 0 ; k_p \neq 0$$

$$\Delta \sigma = E_p \Delta \varepsilon ; E_p = -3p_0$$

$$\Delta \sigma = -3p_0 \Delta \varepsilon$$





# Interaction of dislocations with voids

*Dissipative phenomena at cryogenic temperatures*

The Orowan mechanism:

$$\tau_p = \frac{\mu b^3}{d} \sqrt{\frac{6\xi_0}{\pi}} \left( 1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right)$$

$$\Delta\xi = 3C_1\xi_0\Delta p$$

$$\tau_p = \frac{\mu b^3}{d} \sqrt{\frac{6\xi_0}{\pi}} (1 + C_1\Delta p)$$

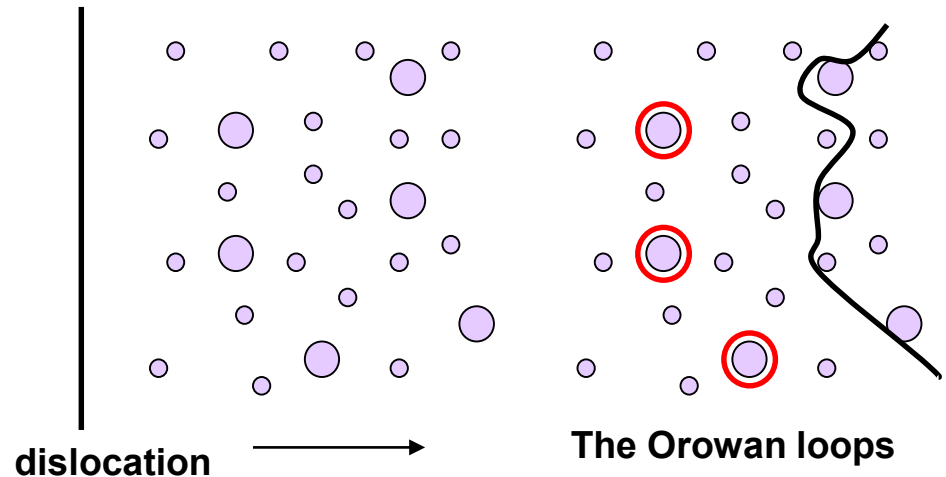
Using the Taylor factor:

$$\sigma_p = M\tau_p = MA_0\sqrt[3]{\xi_0} \left( 1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right)$$

Hardening modulus:

$$C = C_0 \left( 1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right)$$

$$C = C_0 (1 + h\Delta\xi)$$





## Uniaxial case – tension/compression

*dissipative phenomena at cryogenic temperatures*

$$d\sigma = d\sigma_i + d\sigma_{MT}$$

General

$$d\sigma_i = C_0(1 + h\Delta\xi)d\varepsilon^p \quad ; \quad d\sigma_{MT} = E_H d\varepsilon^p = C_{MT} d\varepsilon^p$$

$$\Delta\sigma_i = C_0(\xi_0) \left( \varepsilon^p + \frac{1}{2} C_1(\varepsilon^p)^2 \right)$$

Interaction

$$\Delta\sigma_{MT} = -\frac{5}{2} \mu \eta_0 \frac{\xi_0}{C_1} \left[ \frac{1}{4} (\chi^2 - 1) - \frac{4}{81} \xi_0 (\chi^3 - 1) + -\frac{2}{81} \xi_0^2 (\chi^4 - 1) \right]$$

$$\chi = 1 + 3C_1\varepsilon^p \quad \eta_0 = \frac{C_i}{E} = \frac{M \frac{\mu b}{d} \sqrt[3]{\frac{6}{\pi}} \sqrt[3]{\xi_0}}{E}$$

MT homogenization

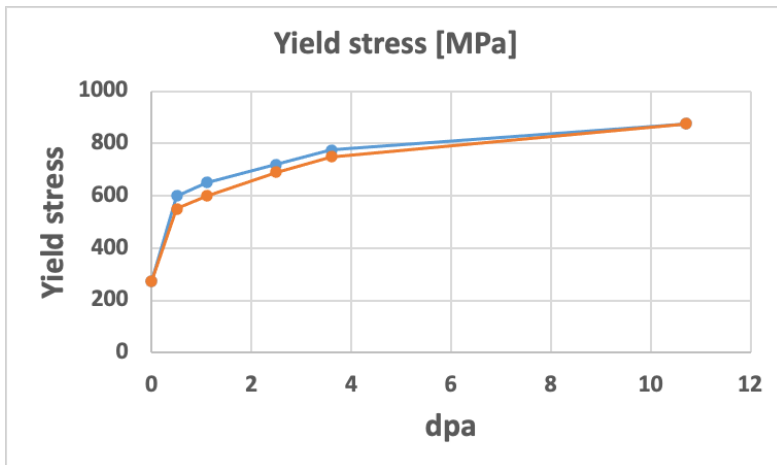
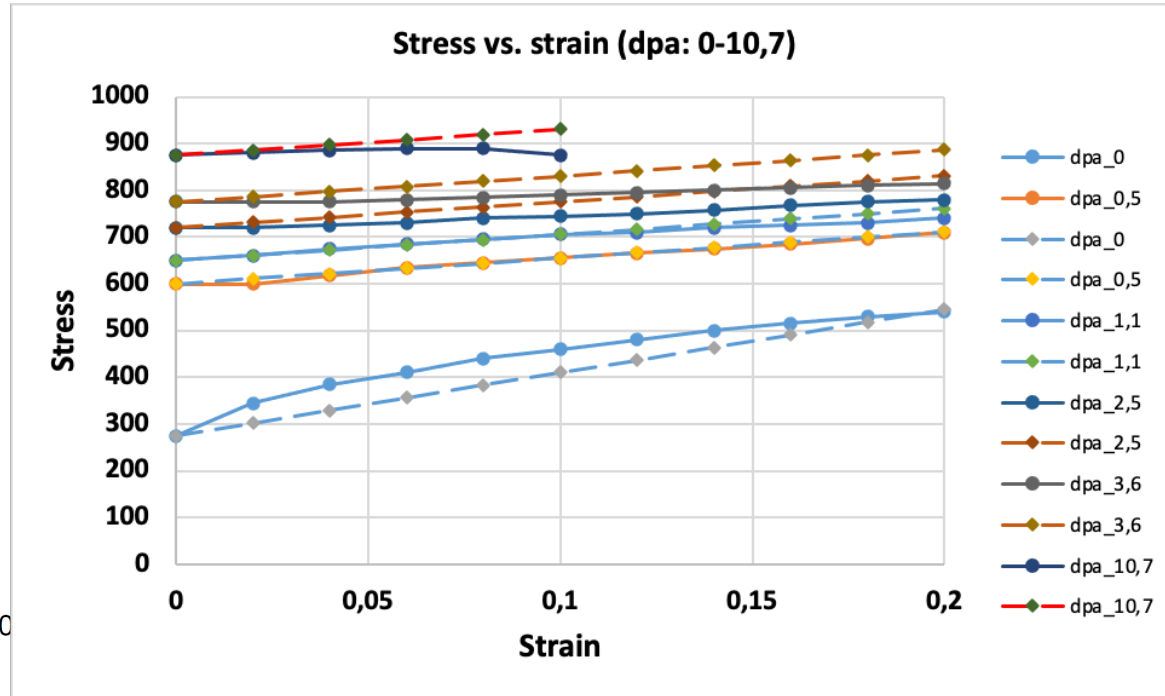
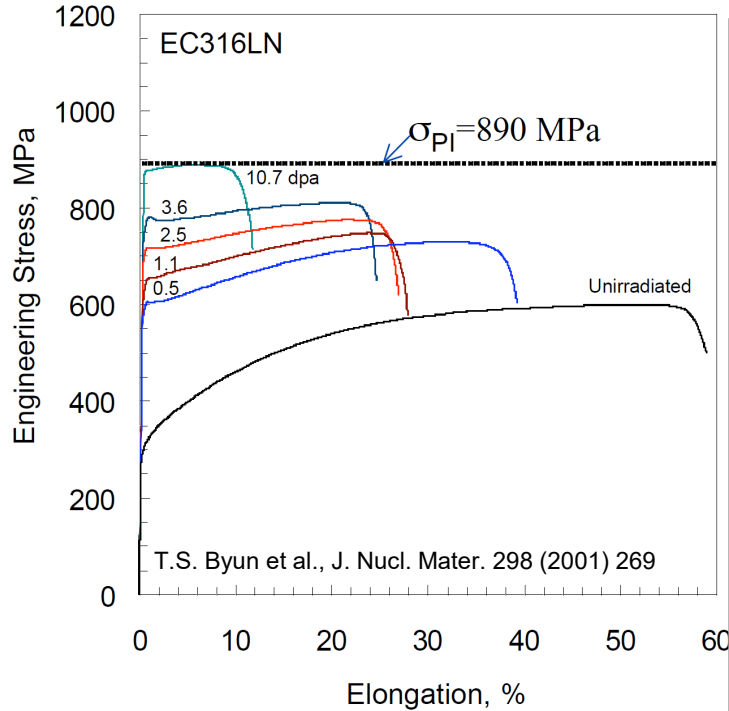
$$\sigma = \sigma_0 + C_0(\xi_0) \left( \varepsilon^p + \frac{1}{2} C_1(\varepsilon^p)^2 \right) - \frac{5}{2} \mu \frac{C_i(\xi_0)}{E} \frac{\xi_0}{C_1} \left[ \frac{1}{4} (\chi^2 - 1) - \frac{4}{81} \xi_0 (\chi^3 - 1) + -\frac{2}{81} \xi_0^2 (\chi^4 - 1) \right]$$





# Uniaxial case: 316LN stainless steel

*dissipative phenomena at cryogenic temperatures*

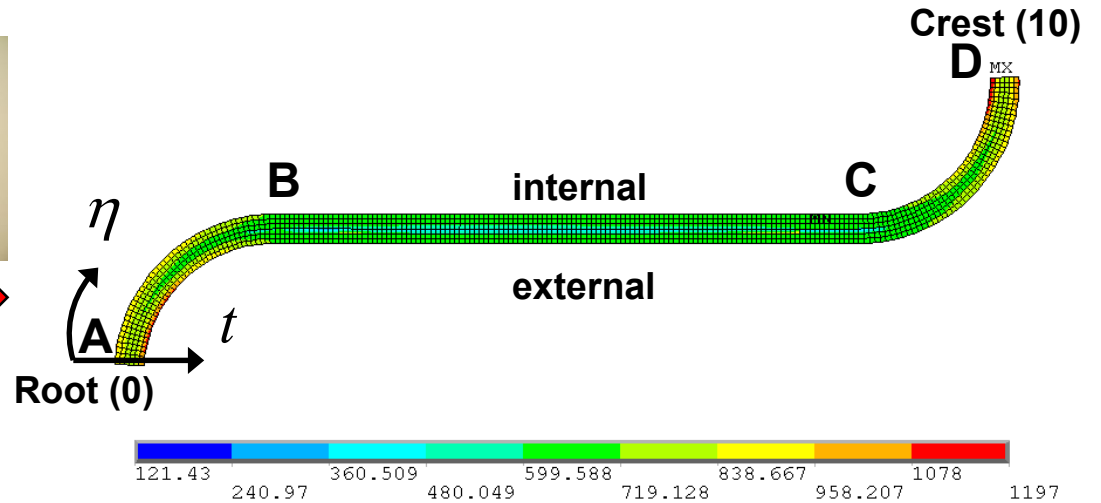
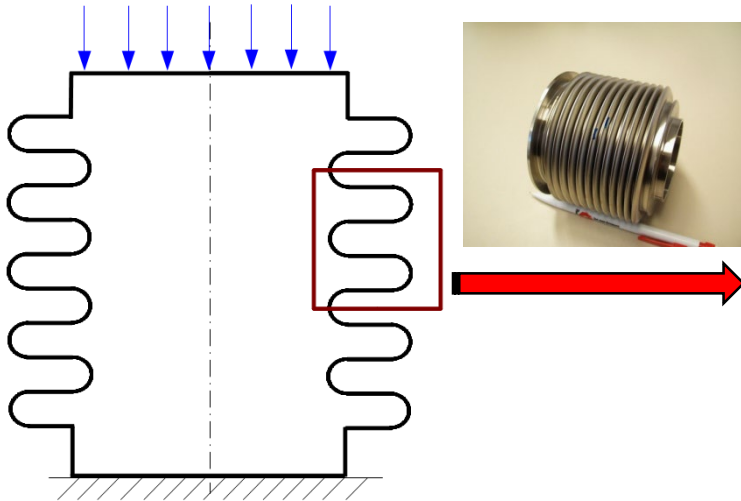


**Radiation induced hardening comprising:**

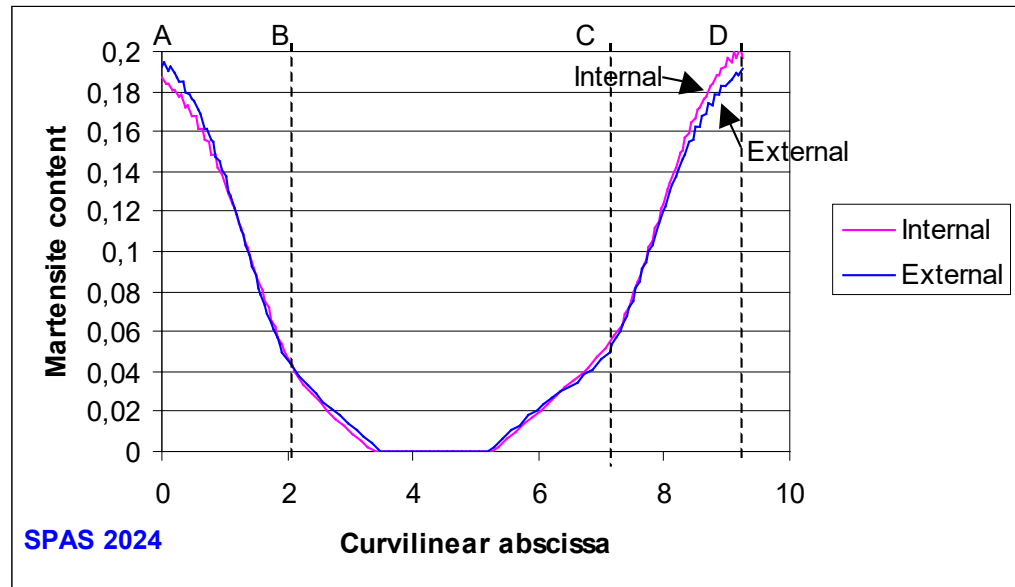
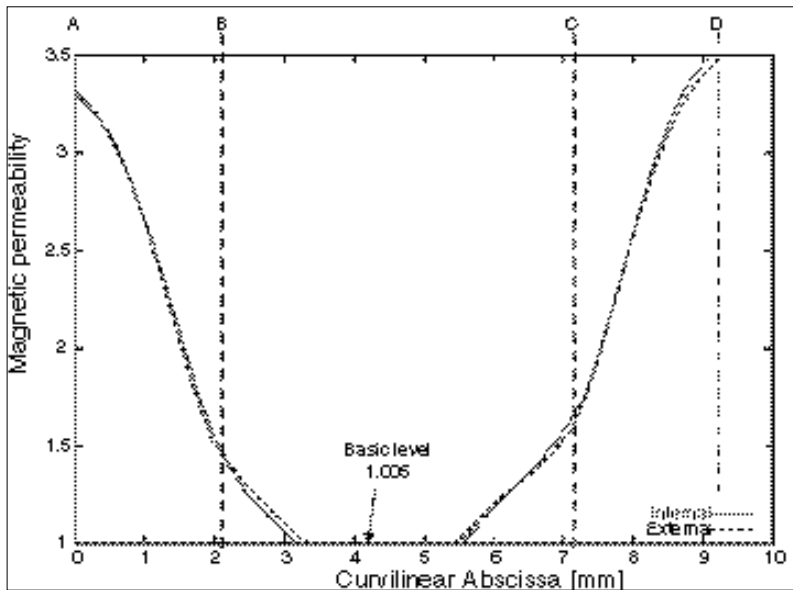
- massive interaction of dislocations with the pressurized voids,
- evolution of tangent stiffness expressed by the Mori-Tanaka homogenization.



# Problem 10: Compensation system of Large Hadron Collider (LHC)



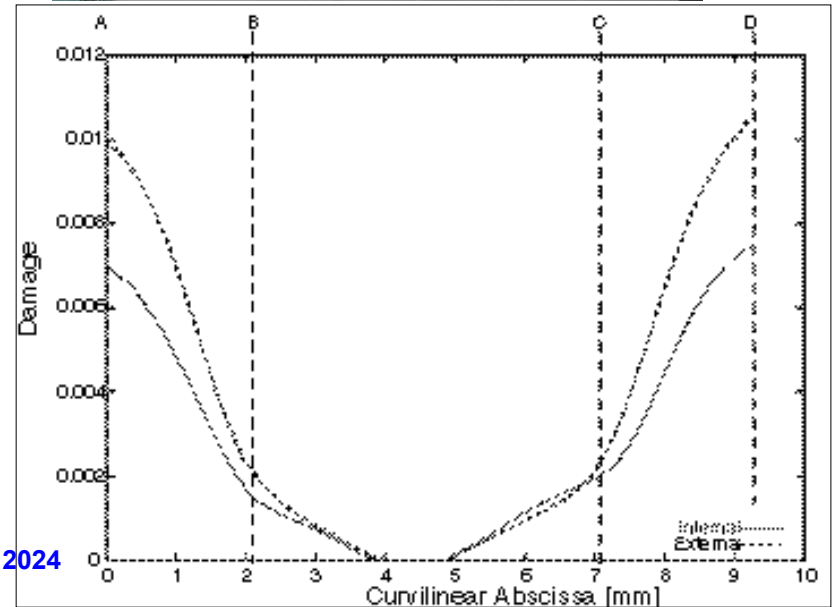
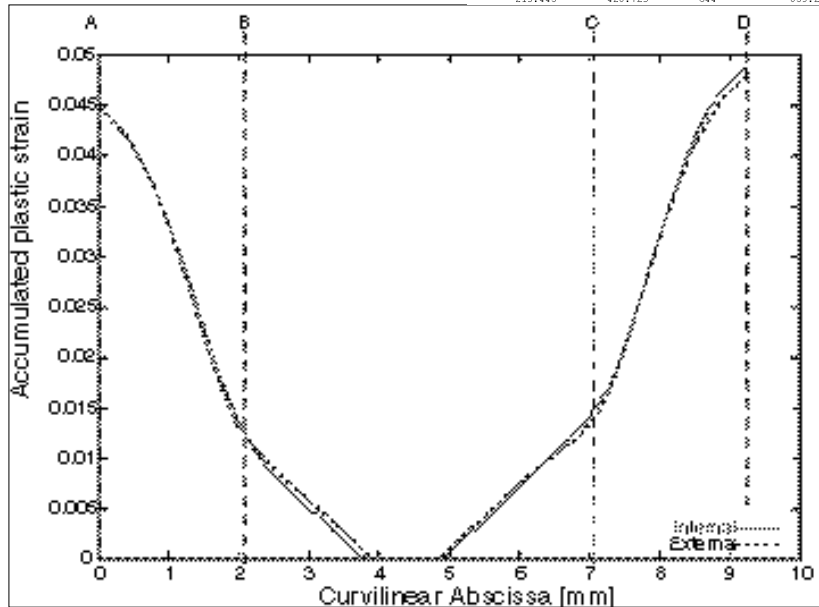
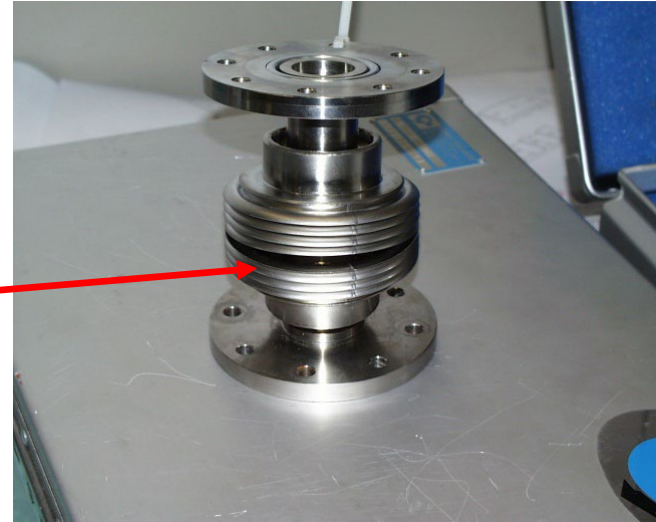
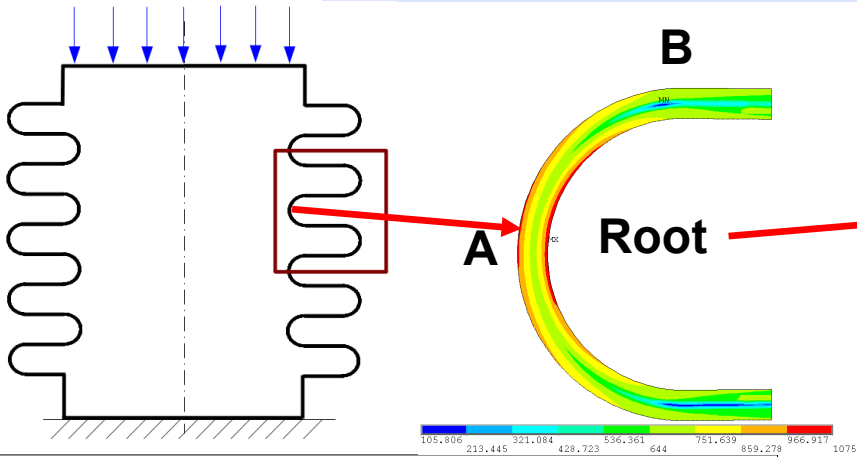
Profile of  $\alpha'$  phase at 77 K





# Problem 10: Compensation system of Large Hadron Collider (LHC)

## Evolution of microdamage







## **Conclusions:**

- 1. Radiation induced defects in the lattice constitute obstacles for the motion of dislocations.**
- 2. Microvoids filled with impurities (gas) induce two physical effects: hardening and swelling.**
- 3. Hardening is related to the interaction of dislocations with the defects, in particular with voids filled with impurities.**
- 4. Tangent stiffness corresponds to the proportion between the volume fraction of matrix, and the volume fraction of voids with impurities.**
- 5. Good correlation between the experiment and the numerical results was obtained.**



**Thank you for your  
attention!**