Diffractive and forward physics measurements by CMS and CMS-TOTEM

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What does CMS have to say on small-*x* and diffraction?

Selection of results in pp and pPb collisions by CMS from the last 2+ years (personal bias). *G. Krintiras will cover CMS diffraction and UPC results in PbPb; M. Pitt discussed photon-induced processes with proton tagging.*



CMS detector



Tracker & muon chambers acceptance up to $|\eta| < 2.5$; $p_T > 200$ MeV for tracks hadronic calorimeter coverage up to $|\eta| < 5.2$; noise threshold E \ge 5 GeV in fwd region

Jet reconstruction spans wide range in $|\eta| < 4.7$ and as low as $p_{\tau} > 20$ GeV

Exclusive vector meson production in pPb collisions

Quasi-real photon exchange from Pb ion fluctuates into a color dipole that probes the



At LO in pQCD,

$$\sigma(\gamma^* p \rightarrow V p) \sim [g(x,Q^2 = m_V^2)]^2$$

In principle, strong sensitivity to small-*x* evolution of gluon PDFs

(NB: this is no longer true at NLO, as discussed by V. Guzey in this conference)

Physics a la EIC at the LHC (energetic "photon" beam, smaller photon virtualities)

Exclusive vector meson production in 5.02 TeV pPb collisions



Exclusive ρ requires special treatment of backgrounds (+ interference with other states).

Different masses of vector mesons allows us to scan different color dipole sizes and different pQCD scales \rightarrow *Necessary if we want to establish universality properties.*

|t| distributions (corrected for smearing effects)

One can use the p_{τ}^2 of the vector meson as a proxy for |t|.

Provides info on impact parameter space (*b* and *t* are Fourier conjugate variables)



Cross section in y*p frame

Cross sections in pPb frame can be "unfolded" to photon-proton frame using the photon flux from the Pb ion as an input for a given rapidity bin.



HERA+LHC data consistent with linear evolution. To probe non-linear evolution effects, one needs to increase beam energy or increase number of nucleons (<u>see G. Krintira's talk on PbPb</u>).

Hard diffraction with intact protons detected in Roman pots of TOTEM



Intact proton is an unambiguous signature of diffraction

Proton detection gives direct access to:

- Four-momentum transfer at the proton vertex [t] $(0.03 < |t| < 1 \text{ GeV}^2)$
- Fractional momentum loss ξ (x_{D} in HERA notation), proxy for the energy carried away by the pomeron/reggeon exchange. $(0.0 < \xi < 0.1$ for Run-1 analysis)

CMS-TOTEM setup



Roman pots: Near-beam Si tracker detectors

CMS:

- General purpose detector at IP5 of the CERN LHC.
- Jets with R = 0.4 reconstructed within $|\eta^{jet}| < 4.7$.

TOTEM:

▶ Roman pots: Forward tracking detectors at ≈ 220m w.r.t. IP5 that measure the protons scattered at small angles w.r.t. the beam.

|t| distribution for single-diffractive jets

Exponential slope $b = 6.5 \pm 0.6 \text{ GeV}^{-2}$ consistent with other hard diffraction probes.

Bare POMWIG overshoots data (requires survival probability of 7.4%), stronger factorization breaking compared to CDF.

PYTHIA8 predictions systematically off by a factor of ~2 at low |t|.

PYTHIA8 with dynamical gap (DG) model

correctly describes the rate and shape of the distribution, *no additional correction factor.*



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Fractional momentum loss ξ (x_p in HERA notation) <u>CMS-TOTEM, EPJC 80, 1164 (2020)</u>

H1 2006 fit B

Significantly extending reach based on forward gaps only $\xi < 0.01$.

Pomeron and reggeon exchange (**POMWIG**) yield the same shapes as pomeron-only (**PYTHIA8**).



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Data corrected to particle-level

Proxy for parton momentum fraction can be estimated from jets kinematics:

$$x^{\pm} = rac{\sum_{ ext{jets}} \left(E^{ ext{jet}} \pm p_z^{ ext{jet}}
ight)}{\sqrt{s}}$$
 ,

POMWIG (with a survival probability of 7.4%) describe qualitatively well the shapes. **PYTHIA8** predictions off at high- and low-*x*.

PYTHIA8 with dynamical gap correctly describes the rate in data, no additional suppression factor is needed.

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Suppression of single-diffractive jets as a function of \sqrt{s}

Fraction of diffractive jets decreases with energy (**Tevatron** \rightarrow **LHC**), qualitatively expected from survival probability dependence on \sqrt{s} .



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Mueller-Tang jets (a.k.a., "jet-gap-jet")



t-channel color-singlet exchange between partons \rightarrow *rapidity gaps between final-state jets*

In the high-energy limit (large $\Delta \eta_{jj}$), it is expected to be mediated by **BFKL pomeron** exchange. A. Mueller and W-K. Tang, PLB 284 (1992) 123.

Experimentally, a signal with a controllable QCD background.

CMS event displays (low pileup data)



Color-exchange event candidate (Background-like) **Color-singlet exchange** event candidate (Signal-like)

tracks with pT > 0.2 GeV are plotted here.

Rapidity gap definition

Number of charged-particles with p_T> 200 MeV in -1 < η < 1 is measured, *rapidity gap corresponds to absence of N_{tracks}*.



Each jet has $|\eta_{iet}| > 1.4$, with $\eta_{iet1}^* \eta_{iet2} < 0$, with $p_T > 40$ GeV.



Residual color-octet background is subtracted with data-driven methods. ¹⁶

Color-singlet fraction f_{CSE} by CMS at 13 TeV



About 0.6% of dijets are produced by hard color-singlet exchange, contribution neglected in modern MC generators.

Pure BFKL predictions (or pure BFKL + MPI) get the trend with data wrong as a function of $\Delta \eta_{jj}$ (Royon, Marquet, Kepka, PRD 83:034036, 2011)

BFKL + soft-color interaction for gap survival probability correctly describes $\Delta \eta_{jj}$ trend (<u>Ekstedt, Enberg</u>, <u>Ingelman, Motyka, arXiv:1703.10919</u>)

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Number of particles in the gap for color-singlet signal events



ISR = on \rightarrow more particles between the jets.

 $ISR = off \rightarrow fewer particles between the jets (unclustered wide-angle hadrons).$

Unexpected sensitivity to ISR in central pseudorapidities.

CB, P. Gonzalez, M. Klasen, J. Salomon, C. Royon, JHEP 08 (2022) 250

PDFs \otimes color structure for color-singlet exchange \otimes BFKL kinematical dependence





At 13 TeV, we are more sensitive to ISR effects for gluon-gluon processes.

Suppression of jet-gap-jet fraction with \sqrt{s}

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Decrease from Tevatron to LHC energies, consistent with diffractive dijet trend



Partial restoration of factorization; intact proton enhances the probability that the central gap "survives" the collision.

Analogous to restoration of factorization observed by <u>CDF Collaboration for</u> <u>double-pomeron exchange/single-diffractive dijets</u>.

Gap between jets with intact proton

CMS-TOTEM PRD 104, 032009 (2021)



Forward-backward dijet configuration (BFKL limit) \rightarrow Phase-space for additional radiation.

Multiple-gluon emissions induce large angular decorrelations between the forward-backward jets.

Fourier coefficients as a function of $\Delta y = CMS, JHEP 08 (2016) 139$



Predictions based on **BFKL resummation at NLL** are consistent with the data.

MC predictions based on fixed order pQCD + parton shower envelop the data in the entire Δy range.

New observables required? Or could 7 TeV \rightarrow 13.6 TeV change help probe BFKL limit?

Ratios of Fourier coefficients for further discrimination



BFKL calculations consistent with data. Nevertheless, fixed order pQCD + PS still does a great job throughout the entire range.

Would 7 TeV \rightarrow 13.6 TeV help? New observables? Asymmetric jet p_T cuts can help suppress DGLAP-like radiation.

Forward single jet spectra in proton-Pb collisions



High-*x* parton from proton, low-*x* parton from Pb ion.

Saturation scale enhanced by $A^{\frac{1}{3}} \sim 6$ relative to proton, potentially making it experimentally accessible at the LHC.

Requires very forward calorimetry + jet calibration under control.

CASTOR calorimeter to extend η acceptance

Forward calorimeter with unique access to low-*x* and low Q² kinematics at the LHC (-6.6 < η < -5.2, jet p_T ~ 3 GeV).

Calorimeter installed in special runs (low PU pp, pPb, and PbPb)

Calorimeter introduced initially to search for Centauro events, not optimized for jet physics. *No η-segmentation.*





Inclusive jet cross section in pPb mode CMS, JHEP 05 (2019) 043



EPOS and **QGSJETII** incorporate saturation effects via pomeron self-interactions. **HIJING** implements nuclear shadowing parametrically.

KATIE predictions (TMD-based, $2 \rightarrow 1$ matrix elements linear evolution with BFKL eqn. or non-linear evolution with BK eqn.)

AAMQS MV (TMD-based, $2 \rightarrow 1$ matrix elements, **non-linear evolution with BK eqn.**)

Predictions based on saturation effects undershoot the data. Missing fragment remnants in forward region?



Large uncertainties at higher energies due to alignment uncertainty and unfolding model dependence.

Summary

• Numerous measurements of hard-scale forward and diffractive processes, corrected to stable-particle level.

• So far no "smoking gun" signature of small-x evolution in forward jet data or diffractive vector meson production in pp and pPb collisions in CMS.

• Trade-off between clean experimental signatures/observables and control over the phenomenology.

High-energy limit of QCD

In the limit $\ \hat{s} \gg - \hat{t} \gg \Lambda_{\rm QCD}^2$, the fixed order pQCD expansion breaks down.

It can be rearranged (symbolically) as,

$$\begin{split} \mathrm{d}\hat{\sigma} &\simeq \alpha_s^2 \sum_{n=0}^{\infty} \alpha_s^n \ln^n \left(\frac{\hat{s}}{|\hat{t}|}\right) + \alpha_s^3 \sum_{n=0}^{\infty} \alpha_s^n \ln^n \left(\frac{\hat{s}}{|\hat{t}|}\right) + \alpha_s^4 \sum_{n=0}^{\infty} \alpha_s^n \ln^n \left(\frac{\hat{s}}{|\hat{t}|}\right) + \dots \\ \text{such that } \alpha_s^n \ln^n \left(\hat{s}/|\hat{t}\right) &\lesssim 1 \end{split}$$

Large logarithms are resummed with Balitsky–Fadin–Kuraev–Lipatov (BFKL) resummation to all orders in $\alpha_{_{\! S}}$.

Same gluon radiation pattern emerges in the proton/nucleus *(small-x evolution*).

