Deeply-virtual and photoproduction of mesons at higher-order and higher-twist

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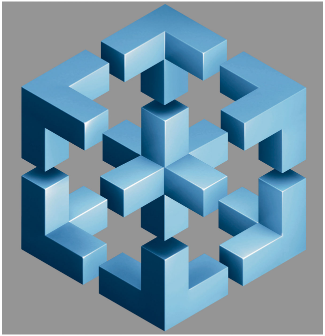
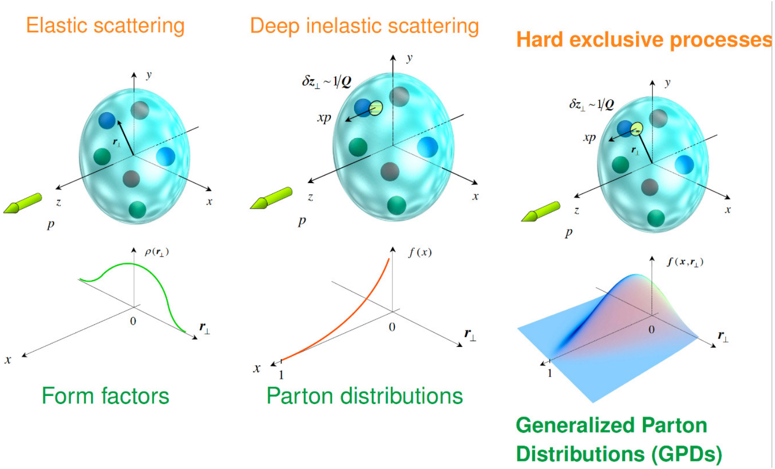
STRONG 2020

Outline

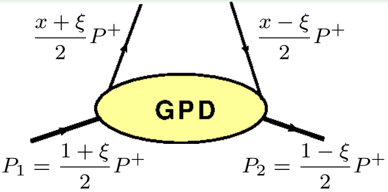


(Escher 3D, Al Borge)



(see J. Wagner talk)

Generalized Parton Distributions

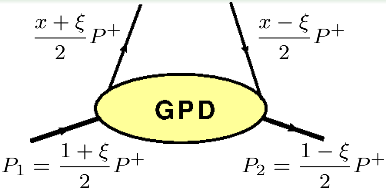


$$P = P_1 + P_2$$

$$\Delta = P_2 - P_1$$

x	parton's "average" longitudinal momentum fraction
$\xi = -\frac{\Delta^+}{P^+}$	longitudinal momentum transfer (skewness)
$\Delta^2 = t$	momentum transfer

Generalized Parton Distributions



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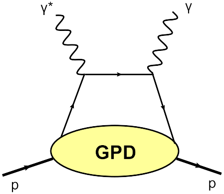
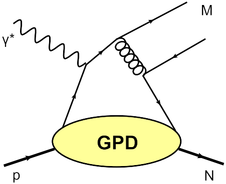
GPDs: $\boxed{F^a(x,\xi,t;\mu),\ a \in \{q,g\}}$ $\mu \ldots$ factorization scale

- vector (H^a, E^a) and axial-vector GPDs ($\widetilde{H}^a$, $\widetilde{E}^a$)
- transversity GPDs (H_T^a, E_T^a, $\widetilde{H}_T^a$, $\widetilde{E}_T^a$)

Handbag factorization



$\gamma^{(*)} N \to \gamma N$



$\gamma^{(*)} N \to M N'$

DEEPLY VIRTUAL	WIDE ANGLE
$Q^2 >>,\ -t <<$	$-t,\ -u,\ s >>$
DVCS (Compton scattering)	WACS
[Collins, Freund '99]	[Radyushkin '98]
	[Diehl, Feldman, Kroll, Jakob '98]
DVMP (meson production)	WAMP
[Collins, Frankfurt, Strikman '97]	[Huang, Kroll '00]

Deeply virtual:
- factorization $\mathcal{H} \otimes GPD$
- GPDs at small $(-t)$

Wide angle:
- arguments for factorization $\mathcal{H}(1/x \otimes GPD(\xi = 0))$
- GPDs at large $(-t)$

DVMP status

- DV (V_L) P:
  - tw-2 predictions ($\gamma_L^* N \to V_L N$) can describe the data
  - tw-3 calculations $\gamma_T^* N \to V_{L,T} N$ [Anikin, Teryaev '02], [Goloskokov, Kroll '13]
- DV (PS) P:
  - tw-2 predictions $\gamma_L^* N \to \pi N$ bellow the data [HERMES '09] [JLab '12,'16, '20] [COMPAS '19] $\Rightarrow$ importance of $\gamma_T^* N \to \pi N$

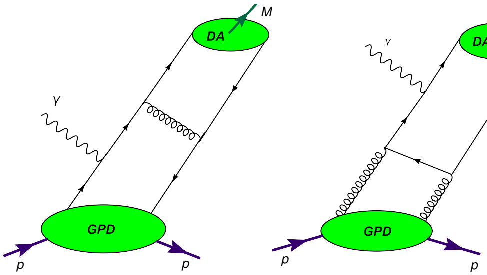
  $\Rightarrow$ tw-3 calculations with transversity (chiral-odd) GPDs (H_T^q...) [Goloskokov, Kroll '10] (2-body, i.e., WW approximation), [Ahmad, Goldstein Liuti '09, Goldstein, Hernandez, Liuti '13]
- WA (PS) P:
  - tw-2 results [Huang, Kroll '00] bellow the data [SLAC '76], [JLab '05, '18] for photoproduction ($Q^2 = 0$)
  - tw-3 2-body π photoproduction vanishes [Huang, Jakob, Kroll, P-K '03]

  $\Rightarrow$ tw-3 (2- and 3-body) prediction to π_0 photoproduction [Kroll, P-K '18] fitted to CLAS data [CLAS '18]; photoproduction of η, η' mesons [Kroll, P-K. '22] [preliminary GlueX '20]

  $\Rightarrow$ tw-3 prediction for $\pi^\pm, \pi^0$ photo- and electroproduction ($Q^2 < -t$) [Kroll, P-K. '21]; extension to DVMP is straightforward

DVMP

$$\gamma^* q \to (q\bar{q})q,\ \gamma^* g \to (q\bar{q})g$$



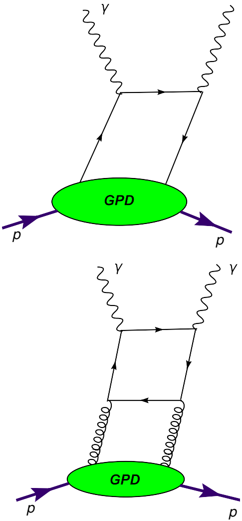
NLO DV PS^{+} prod.: [Belitsky and Müller '01]

NLO DV V$_L$ prod.: [Ivanov et al '04,]

NLO DV V$_L$ (corr.), PS, (S, PV$_L$) prod.: [Duplančić, Müller, P-K. '17]

(D)DVCS

$$\gamma^* q \to \gamma^{(*)}q,\ \gamma^* g \to \gamma^{(*)}g$$



GPDs from DVCS and DVMP

DVCS: Compton form factors (known up to NNLO)

$$ {}^{a}\mathcal{C}(\xi, t, \mathcal{Q}^2) = \int \mathrm{d}x \; C^a(x, \xi, \mathcal{Q}^2/\mu^2) \; F^a(x, \xi, t, \mu^2) $$

$a=q,g$ or NS,S(Σ,g)

DVMP: transition form factors (known up to NLO)

$$ {}^{a}\mathcal{T}(\xi, t, \mathcal{Q}^2) = \int \mathrm{d}x \int \mathrm{d}y \; T^a(x, \xi, y, \mathcal{Q}^2/\mu^2) \; F^a(x, \xi, t, \mu^2) \; \phi(y, \mu^2) $$

- *"Curse of the dimensionality"*
  When the dimensionality increases, the volume of the space increases so fast that the available data become sparse.

- complete deconvolution is impossible
- different modelling venues (momentum fractions space, conformal momentum space) and softwares (PARTONS [... Moutarde, Sznajder et al. '18], GeParD [Kumerički 2006-, '22])

DVMP

Transition form factors

$$ {}^{a}\mathcal{T}(\xi,t,Q^2) = \int \mathrm{d}x \int \mathrm{d}u\, T^a(x,\xi,u,\mu_\varphi,\mu_F)\, F^a(x,\xi,t,\mu_\varphi)\, \phi(u,\mu_F) $$

$a = q, g$ or NS,S(Σ, g)

hard-scattering amplitude (known up to NLO)

$$ T^a(x,\xi,u,\mu_\varphi,\mu_F) = \frac{\alpha_s(\mu_R)}{4\pi} T^{a(1)}(x,\xi,u) + \frac{\alpha_s^2(\mu_R)}{(4\pi)^2} T^{a(2)}(x,\xi,u,\mu_R,\mu_\varphi,\mu_F) + \cdots $$

distribution amplitude (DA) evolution, similary GPD (F^a) evolution (known up to NNLO)

$$ \phi(x;\mu_F,\mu_0) = \phi^{(0)}(u,\mu_F,\mu_0) + \frac{\alpha_s(\mu_F)}{4\pi}\phi^{(1)}(u,\mu_F,\mu_0) + \frac{\alpha_s^2(\mu_F)}{(4\pi)^2}\phi^{(2)}(u,\mu_F,\mu_0) + \cdots $$

$\rightarrow$ evolution simpler to implement in conformal momentum representation [Müller '98]

From x space to conformal momentum space

$$ {}^{a}\mathcal{T}(\xi, t, Q^2) = \int \mathrm{d}x \int \mathrm{d}u \, T^a(x, \xi, y, \mu^2)) \, F^a(x, \xi, t, \mu^2) \, \phi(u, \mu^2) $$

F...GPDs,a=q,g or NS,S(Σ,g)

conformal moments (analogous to Mellin moments in DIS $x^n \to C_n^{3/2}(x), C_n^{5/2}(x)$)

[Müller, Lautenschläger, P-K., Schäfer 2014] [Duplančić, Müller, P-K. 2017]

$$ {}^{a}\mathcal{T}(\xi, t, \mathcal{Q}^2) = \frac{1}{2i} \int_{c-i\infty}^{c+i\infty} dj \left[i \pm \left\{ \begin{array}{c} \tan \\ \cot \end{array} \right\} \left(\frac{\pi j}{2} \right) \right] \xi^{-j-1} $$

$$ \times \left[\mathrm{T}_{jk}(\mathcal{Q}^2/\mu^2) \overset{k}{\otimes} \phi_{\mathrm{M},k}(\mu^2) \right] F_j^{\mathrm{a}}(\xi, t, \mu^2) $$

all channels calculated to NLO :

$\mathcal{H}_M^{q(+)}, \mathcal{E}_M^{q(+)}, \mathcal{H}_M^{\mathrm{g}}, \mathcal{E}_M^{\mathrm{g}}$ $\widetilde{\mathcal{H}}_M^{q(-)}, \widetilde{\mathcal{E}}_M^{q(-)}$	1_L^{--}=V$_L$ 0^{-+}=PS	$\mathcal{H}_M^{q(-)}, \mathcal{E}_M^{q(-)}$ $\widetilde{\mathcal{H}}_M^{q(+)}, \widetilde{\mathcal{E}}_M^{q(+)}, \widetilde{\mathcal{H}}_M^{\mathrm{g}}, \widetilde{\mathcal{E}}_M^{\mathrm{g}}$	0^{++}=S 1_L^{+-}=PV$_L$

(x-space, conformal mom. space, imaginary parts for disp. relations)

NLO predictions

[Müller, Lautenschlager, P-K., Schäfer '14], [Duplančić, Müller, P-K., '17]

- large NLO corrections and model dependence
- results sensitive to the choice of DA
- LO evolution important
- NLO calculations should include NLO evolution
- evolution effects can be called moderate, except for H/E at small ξ
- NLO global DIS+DVCS+DVMP fits needed

NLO for DV V_L production

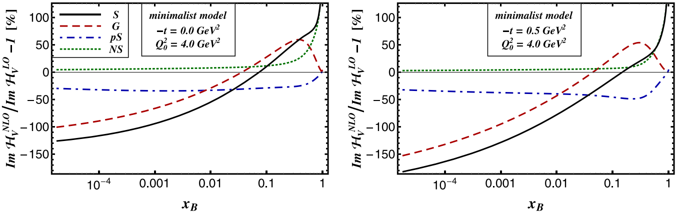


Fig. 6. Relative NLO corrections to the imaginary part of the flavor singlet TFF $\mathcal{F}_{\mathrm{V}}^{\mathrm{S}}$ (solid) broken down to the gluon (dashed), pure singlet quark (dash-dotted) and 'non-singlet' quark (dotted) at $t = 0\ \mathrm{GeV}^2$ (left panel) and $t = -0.5\ \mathrm{GeV}^2$ (right panel) at the initial scale $\mathcal{Q}_0^2 = 4\ \mathrm{GeV}^2$.

[Müller, Lautenschlager, P-K., Schäfer '14]

- big $\ln(1/\xi)$ terms for $\xi <<$, i.e, $j = 0$ pole,
  in gluon evolution and gluon coefficient function

NLO for DV PS/PV production

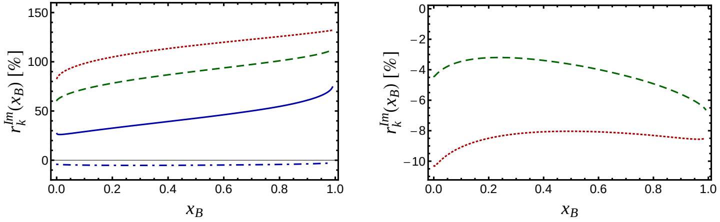


Figure 2: Relative NLO corrections (36) to the imaginary part of the TFF (21) versus x_B for the $k = 0$ (solid), $k = 2$ (dashed), $k = 4$ (dotted) partial waves arising from the quark-quark channel (left panel) and quark-gluon channel (right panel). The pure singlet quark contribution for $k = 0$ is shown as dash-dotted line in the left panel.

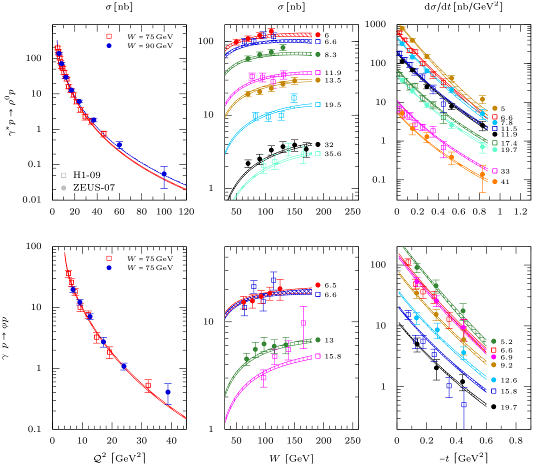
[Duplančić, Müller, P-K., '17]

- NLO corrections higher for higher DA conformal moments ⇒ important for non-asymptotic DAs
- the role of gluons (PV production) smaller since LO vanishes

Global NLO fits (DIS+DVCS+DVV_LP)

small-x global fits to HERA collider data

- LO: [Meskauskas, Müller '11] $(\chi^2/n_{\mathrm{d.o.f}} \approx 2)$
- NLO: [Lautenschlager, Müller, Schäfer '13] (normalization of experimental DVMP datasets treated as fitting parameters)



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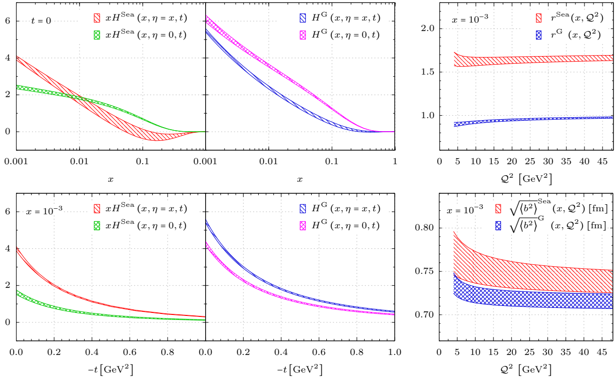
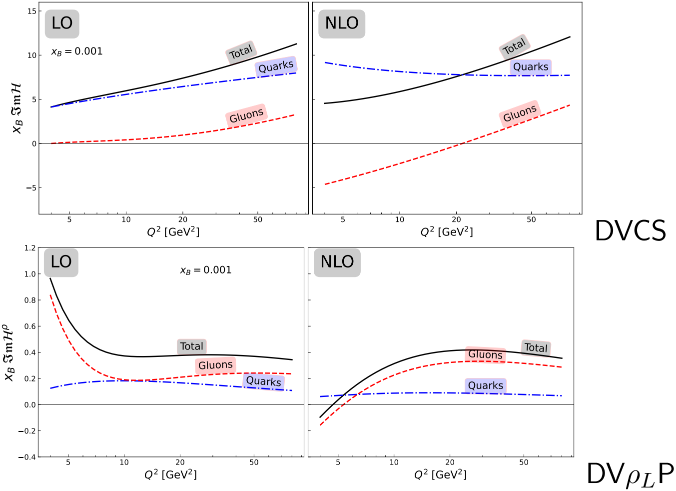


Global NLO fits (DIS+DVCS+DVV_LP)

- hard scattering amplitude corrected [Duplančić, Müller, P-K. '17]
- new NLO fit using GEPARD software: $\chi^2/n_{\rm d.o.f} = 254.3/231$



[preliminary K. Kumerički at Transversity 2022]

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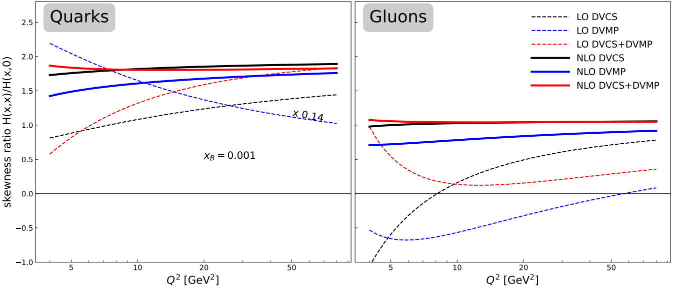


[preliminary K. Kumerički at Transversity 2022]

Helicity amplitudes $\mathcal{M}$ for WAMP

$$\mathcal{M}^P_{0+,\mu+} = \frac{e_0}{2}\sum_\lambda \Big[\mathcal{H}^P_{0\lambda,\mu\lambda}\left(R^P_V(t) + 2\lambda\, R^P_A(t)\right) \quad \to \text{twist-2}$$

$$-2\lambda\,\frac{\sqrt{-t}}{2m}\,\mathcal{H}^P_{0-\lambda,\mu\lambda}\,\bar{S}^P_T(t)\Big] \quad \to \text{twist-3}$$

$$\mathcal{M}^P_{0-,\mu+} = \frac{e_0}{2}\sum_\lambda \Big[\frac{\sqrt{-t}}{2m}\,\mathcal{H}^P_{0\lambda,\mu\lambda}\,R^P_T(t) \quad \to \text{twist-2}$$

$$-2\lambda\,\frac{t}{2m^2}\,\mathcal{H}^P_{0-\lambda,\mu\lambda}\,S^P_S(t)\Big] + e_0\mathcal{H}^P_{0-,\mu+}\,S^P_T(t) \quad \to \text{twist-3}$$

μ photon helicity, $\lambda\ldots$ quark helicities, $P \in \{\pi^\pm, \pi^0, \eta_8, \eta_1, \eta, \eta'\}$,

$$R^a_V(t) = \int \frac{dx}{x}\,H^a(x,\xi=0,t)$$ ... form factors

$$a \in \{u,d\} \Rightarrow R^{\pi^\pm}_V = R^u_V - R^d_V,\; R^{\pi^0}_V = \tfrac{1}{\sqrt{2}}\left(e_u R^u_V - e_d R^d_V\right)$$

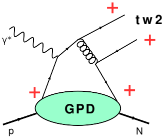
$$R^{\eta_8}_V \approx \tfrac{1}{\sqrt{2}} R^{\eta_1}_V \approx \tfrac{1}{\sqrt{6}}\left(e_u R^u_V + e_d R^d_V\right)$$

$(H, \tilde{H}, E) \to (R_V, R_A, R_T)$

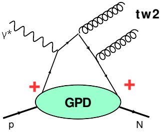
$(H_T, \tilde{H}_T, \bar{E}_T) \to (S_T, S_S, \bar{S}_T)$ transversity GPDs

Subprocess amplitudes

$\mathcal{H}^{P}_{0\lambda,\mu\lambda}$ ... non-flip subprocess amplitudes (twist-2)

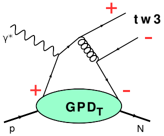


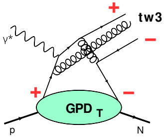
and



for η_1

$\mathcal{H}^{P}_{0-\lambda,\mu\lambda}$ ... flip subprocess amplitudes (twist-3)





($\mu_\pi = 2$ GeV)

distribution amplitudes (DAs):

twist-2 ($q\bar{q}$) : ϕ_P

2-body ($q\bar{q}$) twist-3 ϕ_{Pp}, $\phi_{P\sigma}$ 3-body ($q\bar{q}g$) twist-3 ϕ_{3P}

$\rightarrow$ connected by equations of motion (EOMs)

Subprocess amplitudes: twist-3

General structure:

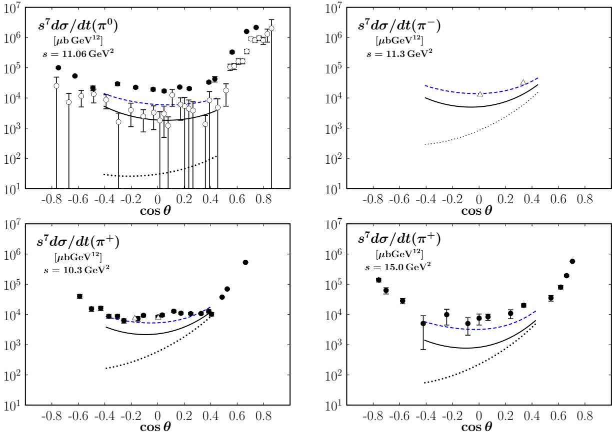
$$\begin{aligned}\mathcal{H}^{P,tw3} &= \mathcal{H}^{P,tw3,q\bar{q}} + \mathcal{H}^{P,tw3,q\bar{q}g} \\ &= \left(\mathcal{H}^{P,\phi_{Pp}} + \underbrace{\mathcal{H}^{P,\phi_{P2}^{EOM}}\right) + \left(\mathcal{H}^{P,q\bar{q}g,C_F}}_{} + \mathcal{H}^{P,q\bar{q}g,C_G}\right) \\ &= \mathcal{H}^{P,\phi_{Pp}} + \mathcal{H}^{P,\phi_{3P},C_F} + \mathcal{H}^{P,\phi_{3P},C_G}\end{aligned}$$

- 2- and 3-body contributions necessary for gauge invariance
- photoproduction ($Q \to 0$): $\mathcal{H}^{P,\phi_{Pp}} = 0$ [Kroll, P-K '18]
- DVMP ($\hat{t} \to 0$):
  - end-point singularities in $\mathcal{H}^{P,\phi_{Pp}}$ $\int_0^1 \frac{d\tau}{\bar{\tau}} \phi_{Pp}(\tau)$

    $\Rightarrow$ modified hard-scattering picture (with $k_\perp$)
  - complete twist-3 contribution [Kroll, P-K '21]
  - work in progress in modified and collinear picture

Photoproduction (π)

[Kroll, P-K '21]



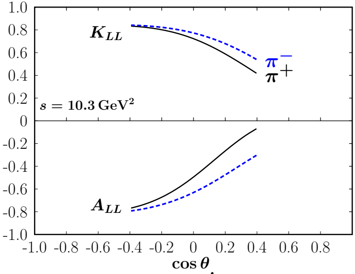
solid curves: complete twist-3
dotted curves: twist-2

exp data:
full circles [SLAC '76]
open circles [CLAS '17]
triangles [JLab, Hall A '05]

- twist-2 prediction well below the data [Huang, Kroll '00]

Spin effects - photoproduction

[Kroll, P-K '21]: $\pi^{\pm}$



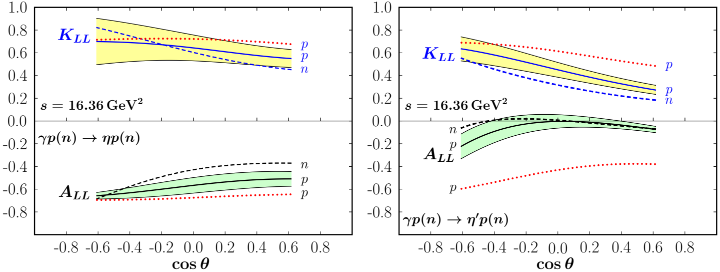
$A_{LL}(K_{LL})\ldots$ correlation of the helicities of the photon and incoming (outgoing) nucleon

$$A_{LL}^{P,tw2} = K_{LL}^{P,tw2}$$

$$A_{LL}^{P,tw3} = -K_{LL}^{P,tw3}$$

$\rightarrow$ characteristic signature for dominance of twist-3 (like $\sigma_T \gg \sigma_L$ in DVMP)

[Kroll, P-K '22]: η, η'



$\rightarrow$ in contrast to π and η, for η' dominance of twist-2 and sensitivity to gluons

Summary

- WA (PS) P:
  - meson's twist-3 contributions dominate for πs and η
  - different combinations of form factors $\Rightarrow$possibility of extraction $\Rightarrow$ large $-t$ behaviour of transversity GPDs (F_T^q)
- DV (PS) M
  - twist-3 dominates
  - complete (2- and 3-body) analysis underway
  - twist-2 NLO contributions available and should be tested
- DV (V$_L$) P
  - twist-2 contributions can describe the data
  - NLO tw2 contributions available for implementation; included in GeParD $\Rightarrow$ global DIS+DVCS+DVMP fits performed
- Experimental goals
  - clear L/T separation (eg., for DVπP JLab, Hall C)

pQCD prediction

$$\mathcal{M}(Q^2) = \mathcal{M}^{(0)}(Q^2) + \frac{\alpha_s(\mu_R)}{4\pi}\mathcal{M}^{(1)}(Q^2) + \frac{\alpha_s^2(\mu_R)}{(4\pi)^2}\mathcal{M}^{(2)}(Q^2, \mu_R) + \cdots$$

$Q^2 \ldots$ characteristic large scale of the process

$\mu_R \ldots$ renormalization scale

- finite order prediction!, renormalization scale and scheme dependence $\Rightarrow$ theoretical uncertainty
- higher-order corrections $(\mathcal{M}^{(2)}(Q^2, \mu_R), \ldots)$ are important: stabilizing effect reducing the dependence of the predictions on the scales and schemes

Discussing the finite order perturbative results

- assessment of the theoretical uncertainty
- optimal choice of the renormalization scale and scheme

$\downarrow$

analysis of the size of the higher-order corrections
and of the expansion parameter $\alpha_s(\mu_R)$

Popular renormalization scale settings:

- characteristic scale of the process: $\mu_R = Q^2$
- FAC (fastest apparent convergence) [Grunberg 1980]: $\mathcal{M}^{(2)}(Q^2, \mu_R) = 0$
- PMS (principle of minimum sensitivity) [Stevenson 1981]:
  $$\frac{d\mathcal{M}_{\text{finite order}}(Q^2, \mu_R)}{d\mu_R} = 0$$
- BLM scheme [Brodsky, Lepage, Mackenzie 1983]: $\mathcal{M}^{(2,\beta_0)}(Q^2, \mu_R) = 0$
  with $\mathcal{M}^{(2)}(Q^2, \mu_R) = \beta_0 \mathcal{M}^{(2,\beta_0)}(Q^2, \mu_R) + \mathcal{M}^{(2,\text{rest})}(Q^2)$
  $\rightarrow$ vacuum polarization effects from the β function resummed into $\alpha_s(\mu_R)$

- elementary hard-scattering amplitudes for twist-2 collinear approximation (t=0):
  - DVCS ($\gamma^* q \to \gamma^{(*)} q$)
    $\Leftrightarrow$ meson transition form factor ($\gamma^* \gamma^{(*)} \to (q\bar{q})$)
  - DVMP ($\gamma^* q \to (q\bar{q}) q$)
    $\Leftrightarrow$ meson electromagnetic form factor, i.e., meson-to-meson ff ($\gamma^* (q\bar{q}) \to (q\bar{q})$)
- bookkeeping of momentum fractions

$$\frac{\xi + x}{2\xi} = u \qquad \left(\frac{\xi - x}{2\xi} = 1 - u\right)$$

but u real so care with $i\epsilon$ in propagators, or a posteriori analytical continuation of energy, i.e., ξ and not u:

$$u \to \frac{\xi - i\epsilon + x}{2(\xi - i\epsilon)} = \frac{\xi + x}{2\xi} + i\epsilon \text{sign} x$$

M		J^{PC}	DA	GPDs
S, V$_L$	$(q_i\bar{q}_j)$	0^{++} , 1^{--} $(\lambda=0)$	$\phi_{\mathrm{asym}}, \phi_{\mathrm{sym}}$	(H, E)
S	(gg)	0^{++}	ϕ_{sym}	(H_g, E_g)
PS, PV$_L$	$(q_i\bar{q}_j)$	0^{-+}, 1^{+-} $(\lambda=0)$	$\phi_{\mathrm{sym}}, \phi_{\mathrm{asym}}$	$(\tilde{H}, \tilde{E})$
PS	(gg)	0^{-+}	ϕ_{asym}	$(\tilde{H}_g, \tilde{E}_g)$
V$_T$	$(q_i\bar{q}_j)$	1^{--} $(\lambda=\pm1)$	ϕ_{sym}	(H_T, E_T)
PV$_T$	$(q_i\bar{q}_j)$	1^{+-} $(\lambda=\pm1)$	ϕ_{asym}	$(\tilde{H}_T, \tilde{E}_T)$
T	(gg)	2^{++}	ϕ_{asym}	(H_{Tg}, E_{Tg})

$(q_i\bar{q}_j)$: $P=(-1)^{l+1}$, $C=(-1)^{l+s}$ $(i=j)$

(gg): $P=(-1)^l$, $C=1$

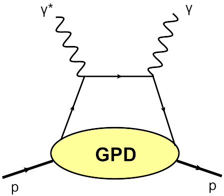
calculated at LO [...,Baier, Grozin'82, '85]

- $\gamma^*\gamma^{(*)} \to PS(S,T) \Rightarrow$ DVCS
- $\gamma_L^* M^\pm \to M^\pm$,
  $\gamma_L^* S(V_L) \to V_L(S)$, $\gamma_L^* PV_L(PS) \to PS(PV_L) \Rightarrow$ DVMP
- $\gamma\gamma \to M^\pm\ M^\pm$,
  $\gamma\gamma \to PS(S)\ PS(S)$, $\gamma\gamma \to S(PS)\ PS(S)$,
  $\gamma\gamma \to V(PV)\ V(PV)$, $\gamma\gamma \to V(PV)\ PV(V)$,
  $\gamma\gamma \to T(PS)\ PS(T)$, $\gamma\gamma \to T(S)\ S(T) \Rightarrow \gamma p \to \gamma M N$

and few at higher order:

Selected exclusive processes

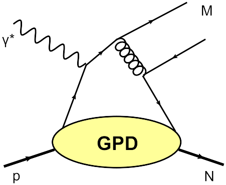
Deeply virtual Compton scattering (DVCS)



$$\gamma^* p \to \gamma p$$

factorization: [Collins, Freund '99]

Deeply virtual production of mesons (DVMP)

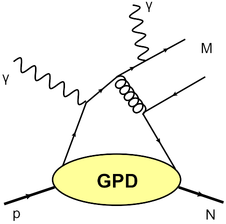


$$\gamma^* p \to M p$$

factorization:

[Collins, Frankfurt, Strikman '97]

Deeply virtual production of photon-meson pair



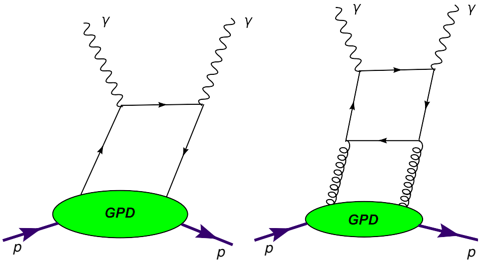
$$\gamma p \to \gamma M p$$

factorization: [Qiu, Yu '22]

(of crossed process)

(D)DVCS

$$\gamma^* q \to \gamma^{(*)} q,\ \gamma^* g \to \gamma^{(*)} g$$



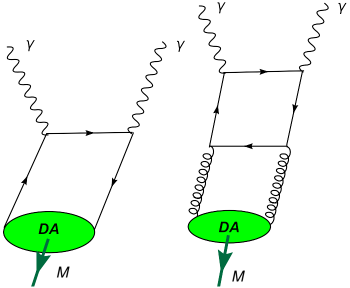
NLO: [Ji, Belitsky et al, Mankiewicz et al, '97]
[Pire, Szymanowski, Wagner '11]

β_0 proportional NNLO: [Belitsky, Schäfer '98]

NNLO from conf. sym: [Müller '05, Kumerički, Müller, P-K. '07]

Meson transition form factor

$$\gamma^* \gamma^{(*)} \to (q\bar{q}),\ \gamma^* \gamma^{(*)} \to (gg)$$



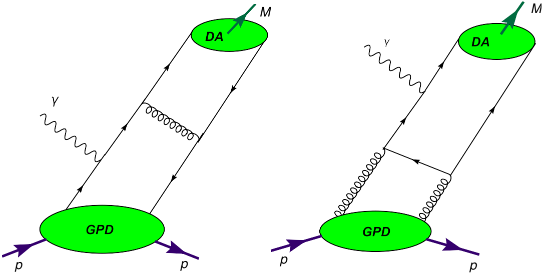
NLO: [..., Kroll, P-K '02] [Kroll, P-K '19]

β_0 proportional NNLO: [Melić, Nižić, Passek '01]

NNLO from conf. sym: [Melić, Müller, Passek '02]

DVMP

$$\gamma^* q \to (q\bar{q})q,\ \gamma^* g \to (q\bar{q})g$$



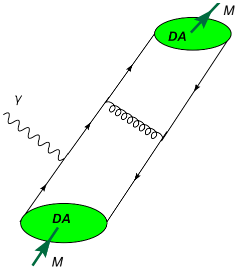
NLO DV PS$^+$ prod.: [Belitsky and Müller '01]

NLO DV V$_L$ prod.: [Ivanov et al '04,]

NLO DV V$_L$ (corr.), PS, (S, PV$_L$) prod.: [Duplančić, Müller, P-K. '17]

Meson em form factor

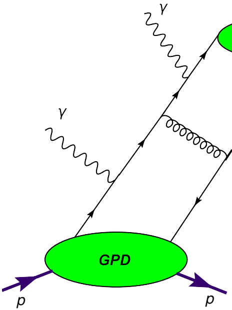
$$\gamma^*(q\bar{q}) \to (q\bar{q})$$



NLO: [..., Melić et al '99]

Photon-meson photoproduction

$$\gamma^* q \to \gamma q (q\bar{q})$$

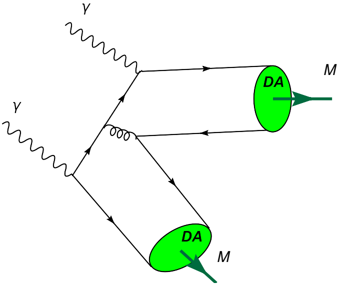


LO V mesons: [Boussarie, Pire, Szymanowsky, Wallon '16]

LO PS mesons: [Duplančić, P-K, Pire, Szymanowski, Wallon '18]

Meson pair production

$$\gamma^* \gamma \to (q\bar{q})(q\bar{q})$$



NLO: [Nižić '87, Duplančić, Nižić '06]

About "$\otimes$": DVCS

- factorization formula for singlet DVCS CFFs:

$$ {}^{S}\mathcal{H}(\xi, t, \mathcal{Q}^2) = \int \mathrm{d}x\, \boldsymbol{C}(x, \xi, \mathcal{Q}^2/\mu^2, \alpha_s(\mu))\, \boldsymbol{H}(x, \xi, t, \mu^2) $$

- ...in terms of conformal moments

(analogous to Mellin moments in DIS: $x^n \to C_n^{3/2}(x), C_n^{5/2}(x)$):

$$ = 2\sum_{j=0}^{\infty} \xi^{-j-1} \boldsymbol{C}_j(\mathcal{Q}^2/\mu^2, \alpha_s(\mu))\, \boldsymbol{H}_j(\xi = \eta, t, \mu^2) $$

$$ H_j^q(\eta, \ldots) = \frac{\Gamma(3/2)\Gamma(j+1)}{2^{j+1}\Gamma(j+3/2)} \int_{-1}^{1} \mathrm{d}x\, \eta^{j-1} C_j^{3/2}(x/\eta) H^q(x, \eta, \ldots) $$

...

H_j^a even polynomials in η with maximal power η^{j+1}

- series summed using Mellin-Barnes integral over complex j:

$$ = \frac{1}{2i} \int_{c-i\infty}^{c+i\infty} dj \left[i + \tan\left(\frac{\pi j}{2}\right)\right] \xi^{-j-1} \boldsymbol{C}_j(\mathcal{Q}^2/\mu^2, \alpha_s(\mu))\, \boldsymbol{H}_j(\xi, t, \mu^2) $$

[Müller 2006, Kumerički, Müller, P-K., Schäfer 2006, 2007]

Modelling conformal moments

$$\boldsymbol{H}_j(\eta,t) = \underbrace{\begin{pmatrix} N'_\Sigma F_\Sigma(t)\,\mathrm{B}\big(1+j-\alpha_\Sigma(0),8\big) \\ N'_\mathsf{G} F_\mathsf{G}(t)\,\mathrm{B}\big(1+j-\alpha_\mathsf{G}(0),6\big) \end{pmatrix}}_{\text{Leading partial wave}} + \begin{pmatrix} s_\Sigma \\ s_\mathsf{G} \end{pmatrix} \begin{pmatrix} \text{subleading partial waves, } \eta\text{-dependence!} \end{pmatrix}$$

- Leading wave – simplest case:
  (at NLO data can be fitted with leading wave only)
  - Regge-inspired ansatz

$$\alpha_a(t)=\alpha_a(0)+0.15t \qquad F_a(t) = \frac{j+1-\alpha(0)}{j+1-\alpha(t)}\left(1-\frac{t}{{M_0^a}^2}\right)^{-p_a}$$

  - for $t=0$ corresponds to x-space PDFs of the form

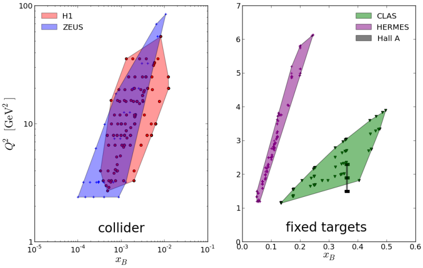
$$\Sigma(x) = N'_\Sigma\, x^{-\alpha_\Sigma(0)}\,(1-x)^7\;; \qquad G(x) = N'_\mathsf{G}\, x^{-\alpha_\mathsf{G}(0)}\,(1-x)^5$$

  - fit parameters: N_Σ, $\alpha_\Sigma(0)$, $\alpha_G(0)$ (DIS) and M_0^Σ (DVCS)

    ($M_0^G = \sqrt{0.7}$ GeV from J/Ψ prod.)

Experimental status

DVCS



[from Kumericki et al. 2015]

DVMP

- in the last decade: vector meson (ρ, J/Ψ, ϕ) production at H1 and ZEUS (HERA, DESY), COMPASS (CERN), pseudoscalar mesons (π, η) at CLAS (JLab) ...

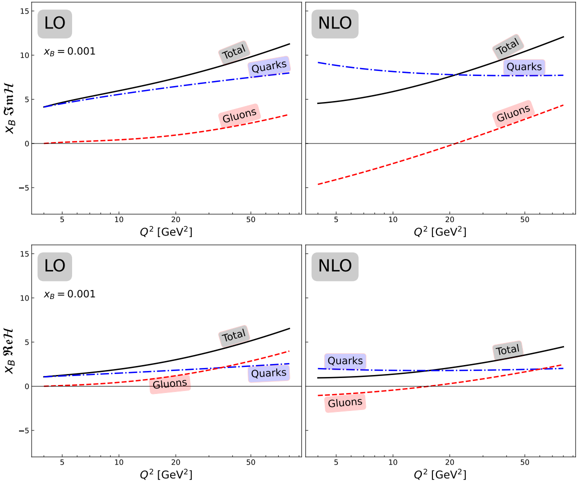
$\rightarrow$ new results from JLab@12 (2018)

COMPASS@LHC

EIC (Electron Ion Collider at Brokhaven, 2030)
LHeC proposed

Global NLO fits

- hard scattering amplitude corrected [Duplančić, Müller, P-K. '17]
- new NLO fit using GEPARD software: $\chi^2/n_{\rm d.o.f} = 254.3/231$



[K. Kumerički at Transversity 2022]

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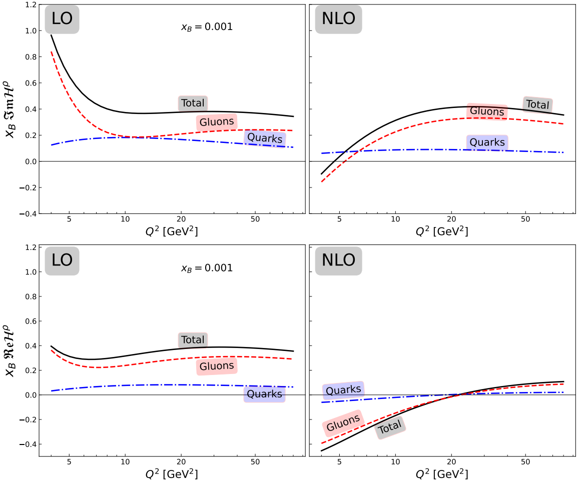


[K. Kumerički at Transversity 2022]

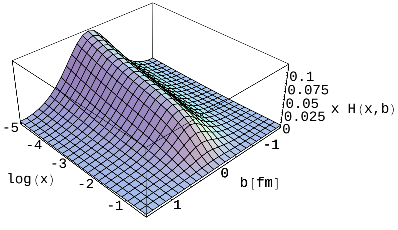
Parton probability density

- Fourier transform of GPD for $\eta = 0$ can be interpreted as probability density depending on x and transversal distance b
[Burkardt '00, '02]

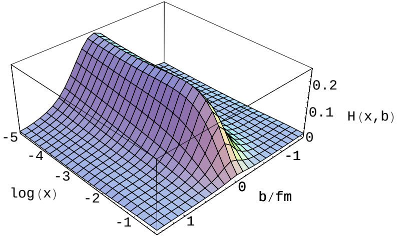
$$H(x, \vec{b}) = \int \frac{d^2 \vec{\Delta}}{(2\pi)^2} \, e^{-i\vec{b}\cdot\vec{\Delta}} H(x, \eta = 0, \Delta^2 = -\vec{\Delta}^2)$$

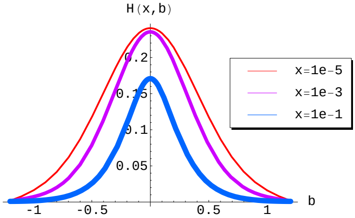
Three-dimensional image of a proton

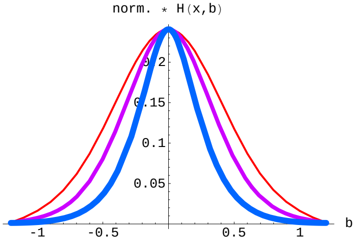
Quarks:



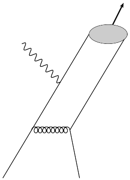
Gluons:







Subprocess amplitudes $\mathcal{H}$

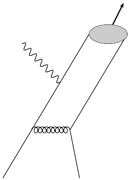


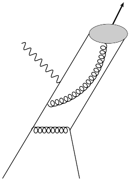
$q\bar{q} \to \pi$ projector [Beneke, Feldmann '00]

$$(\tau q' + k_\perp) + (\bar{\tau} q' - k_\perp) = q'$$

$$\begin{aligned}\mathcal{P}_2^\pi \quad \sim \quad & f_\pi \Big\{ \gamma_5 \not{q}' \phi_\pi(\tau, \mu_F) \\ & + \mu_\pi(\mu_F) \Big[\gamma_5 \, \phi_{\pi p}(\tau, \mu_F) \\ & - \frac{i}{6} \gamma_5 \, \sigma_{\mu\nu} \, \frac{q'^\mu n^\nu}{q' \cdot n} \, \phi'_{\pi\sigma}(\tau, \mu_F) \\ & + \frac{i}{6} \gamma_5 \, \sigma_{\mu\nu} \, q'^\mu \phi_{\pi\sigma}(\tau, \mu_F) \frac{\partial}{\partial k_{\perp\nu}} \Big] \Big\}_{k_\perp \to 0}\end{aligned}$$

Subprocess amplitudes $\mathcal{H}$





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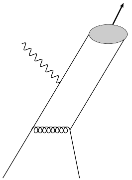
$q\bar{q}g \to \pi$ projector [Kroll, P-K '18]

$$\tau_a q' + \tau_b q' + \tau_g q' = q'$$

$$\mathcal{P}_3^\pi \sim f_{3\pi}(\mu_F) \, \frac{i}{g} \, \gamma_5 \, \sigma_{\mu\nu} q'^\mu g_\perp^{\nu\rho} \, \frac{\phi_{3\pi}(\tau_a, \tau_b, \tau_g, \mu_F)}{\tau_g}$$

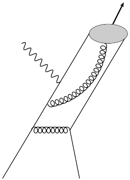
Subprocess amplitudes $\mathcal{H}$



$q\bar{q} \to \pi$ projector [Beneke, Feldmann '00]

$$(\tau q' + k_\perp) + (\bar{\tau} q' - k_\perp) = q'$$

$$\mathcal{P}_2^\pi \sim f_\pi \Big\{ \gamma_5 \, \not{q}' \phi_\pi(\tau, \mu_F) + \mu_\pi(\mu_F) \Big[\gamma_5 \, \phi_{\pi p}(\tau, \mu_F) - \frac{i}{6} \gamma_5 \, \sigma_{\mu\nu} \frac{q'^\mu n^\nu}{q' \cdot n} \phi'_{\pi\sigma}(\tau, \mu_F) + \frac{i}{6} \gamma_5 \, \sigma_{\mu\nu} \, q'^\mu \phi_{\pi\sigma}(\tau, \mu_F) \frac{\partial}{\partial k_{\perp\nu}} \Big] \Big\}_{k_\perp \to 0}$$



$q\bar{q}g \to \pi$ projector [Kroll, P-K '18]

$$\tau_a q' + \tau_b q' + \tau_g q' = q'$$

$$\mathcal{P}_3^\pi \sim f_{3\pi}(\mu_F) \, \frac{i}{g} \, \gamma_5 \, \sigma_{\mu\nu} q'^\mu g_\perp^{\nu\rho} \, \frac{\phi_{3\pi}(\tau_a, \tau_b, \tau_g, \mu_F)}{\tau_g}$$

$\mu_\pi = m_\pi^2/(m_u + m_d) \cong 2$ GeV, $f_{3\pi} \sim \mu_\pi$

distribution amplitudes (DAs):

twist-2 $(q\bar{q})$: ϕ_π

2-body $(q\bar{q})$ twist-3 $\phi_{\pi p}$, $\phi_{\pi\sigma}$ 3-body $(q\bar{q}g)$ twist-3 $\phi_{3\pi}$

$\to$ connected by equations of motion (EOMs)

DAs and EOMs

$$\begin{aligned} \tau\phi_{\pi p}(\tau) + \frac{\tau}{6}\,\phi'_{\pi\sigma}(\tau) - \frac{1}{3}\,\phi_{\pi\sigma}(\tau) &= \phi_{\pi 2}^{EOM}(\bar\tau) \\ \bar\tau\phi_{\pi p}(\tau) - \frac{\bar\tau}{6}\phi'_{\pi\sigma}(\tau) - \frac{1}{3}\phi_{\pi\sigma}(\tau) &= \phi_{\pi 2}^{EOM}(\tau) \end{aligned}$$

$$\phi_{\pi 2}^{EOM}(\tau) = 2\frac{f_{3\pi}}{f_\pi \mu_\pi}\int_0^{\bar\tau} \frac{d\tau_g}{\tau_g}\,\phi_{3\pi}(\tau, \bar\tau - \tau_g, \tau_g)$$

- EOMs and symmetry properties
  $\Rightarrow$ the subprocess amplitudes in terms of two twist-3 DAs and 2- and 3-body contributions combined
- combined EOMs $\rightarrow$ first order differential equation $\Rightarrow$ from known form of $\phi_{3\pi}$ [Braun, Filyanov '90] one determines $\phi_{\pi p}$ (and $\phi_{\pi\sigma}$)

Note: $q\bar q g$ projector and EOMs were derived using light-cone gauge for constituent gluon

Subprocess amplitudes: twist-2

Transverse photon polarization ($\mu = \pm 1$) T

$$\mathcal{H}^{\pi,tw2}_{0\lambda,\,\mu\lambda} \sim f_\pi\, C_F\, \alpha_s(\mu_R)\, \frac{\sqrt{-\hat{t}}}{\hat{s}+Q^2} \int_0^1 d\tau\, \phi_\pi(\tau) \Bigg[(2\lambda\mu+1)\left(\frac{(\hat{s}\tau+Q^2)(\hat{s}+Q^2)-\hat{u}Q^2\bar{\tau}}{\hat{s}\bar{\tau}(Q^2\bar{\tau}-\hat{t}\tau)}\, e_a + \frac{(\hat{s}\tau-Q^2)(\hat{s}+Q^2)-\hat{u}Q^2\bar{\tau}}{\hat{u}\tau(Q^2\tau-\hat{t}\bar{\tau})}\, e_b\right) + (2\lambda\mu-1)\left(\frac{\hat{u}\, e_a}{(Q^2\bar{\tau}-\hat{t}\tau)} + \frac{\hat{s}\bar{\tau}\, e_b}{\tau(Q^2\tau-\hat{t}\bar{\tau})}\right)\Bigg]$$

Longitudinal photon polarization L

$$\mathcal{H}^{\pi,tw2}_{0\lambda,\,0\lambda} \sim f_\pi\, C_F\, \alpha_s(\mu_R)\, \lambda\, \frac{Q\sqrt{-\hat{u}\hat{s}}}{\hat{s}+Q^2} \int_0^1 d\tau\, \phi_\pi(\tau) \left(\frac{\hat{u}\, e_a}{\hat{s}(Q^2\bar{\tau}-\hat{t}\tau)} - \frac{(\hat{t}+\tau\hat{u})\, e_b}{\tau\hat{u}(Q^2\tau-\hat{t}\bar{\tau})}\right)$$

→ photoproduction ($Q \to 0$): $\mathcal{H}^{\pi,tw2}_{L}\Big|_{Q\to 0} = 0$

$$\mathcal{H}^{\pi,tw2}_{T}\Big|_{Q\to 0} \sim f_\pi\, C_F\, \alpha_s(\mu_R)\, \frac{1}{\sqrt{-\hat{t}}} \int_0^1 \frac{d\tau}{\tau}\, \phi_\pi(\tau) \left((1+2\lambda\mu)\,\hat{s} - (1-2\lambda\mu)\,\hat{u}\right) \left(\frac{e_a}{\hat{s}} + \frac{e_b}{\hat{u}}\right)$$

→ DVMP ($\hat{t} \to 0$): $\mathcal{H}^{\pi,tw2}_{T}\Big|_{\hat{t}\to 0} = 0$

$\mathcal{H}^{\pi,tw2}_{L}\Big|_{\hat{t}\to 0}$: $\hat{s} = -\frac{\xi - x}{2\xi}\, Q^2\,,\ \hat{u} = -\frac{\xi + x}{2\xi}\, Q^2$ ⇒ well known LO result for DVMP

Subprocess amplitudes: twist-3

General structure:

$$\begin{aligned}\mathcal{H}^{P,tw3} &= \mathcal{H}^{P,tw3,q\bar{q}} + \mathcal{H}^{P,tw3,q\bar{q}g} \\ &= \left(\mathcal{H}^{P,\phi_{\pi p}} + \underbrace{\mathcal{H}^{P,\phi_{\pi 2}^{EOM}}\right) + \left(\mathcal{H}^{P,q\bar{q}g,C_F}} + \mathcal{H}^{P,q\bar{q}g,C_G}\right) \\ &= \mathcal{H}^{P,\phi_{\pi p}} \quad + \qquad \mathcal{H}^{P,\phi_{3\pi},C_F} \qquad + \mathcal{H}^{P,\phi_{3\pi},C_G}\end{aligned}$$

- 2-body twist-3 $\sim C_F$; 3-body C_F and C_G proportional parts
- C_G part is separately gauge invariant
- the sum of 2- and 3-body C_F parts is gauge invariant (QED and QCD)
- no end-point singularities for $\hat{t} \neq 0$!

Subprocess amplitudes: twist-3 at $Q <<$ or $\hat{t} <<$

General structure:

$$
\begin{aligned}
\mathcal{H}^{P,tw3} &= \mathcal{H}^{P,tw3,q\bar{q}} + \mathcal{H}^{P,tw3,q\bar{q}g} \\
&= \left(\mathcal{H}^{P,\phi_{\pi p}} + \underbrace{\mathcal{H}^{P,\phi_{\pi 2}^{EOM}}\right) + \left(\mathcal{H}^{P,q\bar{q}g,C_F}} + \mathcal{H}^{P,q\bar{q}g,C_G}\right) \\
&= \mathcal{H}^{P,\phi_{\pi p}} + \mathcal{H}^{P,\phi_{3\pi},C_F} + \mathcal{H}^{P,\phi_{3\pi},C_G}
\end{aligned}
$$

- $\mathcal{H}_L^{P,tw3} \sim Q\sqrt{-t} \to 0$ both for $Q \to 0$ and $\hat{t} \to 0$
- photoproduction $(Q \to 0)$:
  - $\mathcal{H}^{P,\phi_{\pi p}} = 0$ [Kroll, P-K '18]
- DVMP $(\hat{t} \to 0)$:
  - end-point singularities in $\mathcal{H}^{P,\phi_{\pi p}}$ [Goloskokov, Kroll '10]
  - $\mathcal{H}^{P,\phi_{\pi 2}^{EOM}} = 0$

Subprocess amplitudes: twist-3 at $Q \to 0$, $t \to 0$

photoproduction

$$\mathcal{H}^{P,tw3}_{0-\lambda,\,\mu\lambda}|_{Q^2\to 0} \sim (2\lambda-\mu)\, f_{3\pi}\, \alpha_S(\mu_R)\, \sqrt{-\hat{u}\hat{s}} \int_0^1 d\tau \int_0^{\bar\tau} \frac{d\tau_g}{\tau_g}\, \phi_{3\pi}(\tau,\bar\tau-\tau_g,\tau_g) \times \left[C_F \left(\frac{1}{\bar\tau^2} - \frac{1}{\bar\tau(\bar\tau-\tau_g)}\right)\left(\frac{e_a}{\hat{s}^2}+\frac{e_b}{\hat{u}^2}\right) + -C_G\, \frac{2}{\tau\tau_g}\, \frac{\hat{t}}{\hat{s}\hat{u}} \left(\frac{e_a}{\hat{s}}+\frac{e_b}{\hat{u}}\right)\right]$$

DVMP

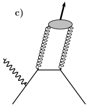
$$\mathcal{H}^{P,\phi_{\pi p}}_{0-\lambda,\mu\lambda}|_{\hat{t}\to 0} \sim (2\lambda+\mu)\; f_\pi \mu_\pi C_F \alpha_S(\mu_R) \sqrt{-\frac{\hat{u}}{\hat{s}}} \left[\frac{e_a}{\hat{s}} + \frac{\hat{s}}{\hat{u}}\frac{e_b}{\hat{u}}\right] \int_0^1 \frac{d\tau}{\bar\tau} \phi_{\pi p}(\tau)$$

$$\mathcal{H}^{P,C_F,\phi_{3\pi}}_{0-\lambda,\mu\lambda}|_{\hat{t}\to 0} \sim -(2\lambda+\mu)\, f_{3\pi}\, C_F \alpha_S(\mu_R) \sqrt{-\frac{\hat{u}}{\hat{s}}} \left(\frac{e_a}{\hat{s}} + \frac{\hat{s}}{\hat{u}}\frac{e_b}{\hat{u}}\right) \times \int_0^1 \frac{d\tau}{\bar\tau^2} \int_0^{\bar\tau} \frac{d\tau_g}{\tau_g(\bar\tau-\tau_g)}\, \phi_{3\pi}(\tau,\bar\tau-\tau_g,\tau_g)$$

$$\mathcal{H}^{P,qqg,C_G}_{0-\lambda,\mu\lambda}|_{\hat{t}\to 0} \sim (2\lambda+\mu)\; f_{3\pi}\, C_G \alpha_S(\mu_R) \frac{Q^2}{\sqrt{-\hat{s}\hat{u}}} \left(\frac{e_a}{\hat{s}} + \frac{e_b}{\hat{u}}\right) \times \int_0^1 \frac{d\tau}{\bar\tau} \int_0^{\bar\tau} \frac{d\tau_g}{\tau_g(\bar\tau-\tau_g)}\, \phi_{3\pi}(\tau,\bar\tau-\tau_g,\tau_g)$$

Subprocess amplitudes $\mathcal{H}^{\eta_8,\eta_1} \to \mathcal{H}^{\eta,\eta'}$

Novel features:

1. 

$|gg\rangle$ states contribute to twist-2

- $\mathcal{H}^{\pi,tw2} \Rightarrow \mathcal{H}^{\eta_8,tw2}, \mathcal{H}^{\eta_{1,q},tw2}$ $\qquad (\phi_\pi, f_\pi) \to (\phi_{\eta_8}, f_{\eta_8}), (\phi^q_{\eta_1}, f_{\eta_1})$

$\mathcal{H}^{\eta_1} = \mathcal{H}^{\eta_{1q},tw2} + \mathcal{H}^{\eta_{1g},tw2}$ $\qquad \phi^q_{\eta_1}$ and $\phi^g_{\eta_1}$ mix under evolution

- $\mathcal{H}^{\pi,tw3} \Rightarrow \mathcal{H}^{P,tw3}$ $\qquad (\phi_{3\pi}, f_\pi, f_{3\pi}) \to (\phi_{3P}, f_P, f_{3P})$

2. flavour-mixing:
   - simplest: flavour-mixing embedded in the decay constants

$$f^8_\eta = f_8 \cos\theta_8 \qquad f^1_\eta = -f_1 \sin\theta_1$$

$$f^8_{\eta'} = f_8 \sin\theta_8 \qquad f^1_{\eta'} = f_1 \cos\theta_1$$

[review Feldmann '00]

Pion distribution amplitudes

Twist-2 DA: $\phi_\pi(\tau, \mu_F) = 6\tau\bar{\tau}\left[1 + a_2(\mu_F)\, C_2^{3/2}(2\tau - 1)\right]$

Twist-3 DAs:

$$\begin{aligned}\phi_{3\pi}(\tau_a, \tau_b, \tau_g, \mu_F) &= 360\tau_a\tau_b\tau_g^2\Big[1 + \omega_{1,0}(\mu_F)\,\frac{1}{2}(7\tau_g - 3)\\ &+ \omega_{2,0}(\mu_F)\,(2 - 4\tau_a\tau_b - 8\tau_g + 8\tau_g^2)\\ &+ \omega_{1,1}(\mu_F)\,(3\tau_a\tau_b - 2\tau_g + 3\tau_g^2)\Big]\end{aligned}$$

[Braun, Filyanov '90]

using EOMs [Kroll, P-K '18]:

$$\begin{aligned}\phi_{\pi p}(\tau, \mu_F) &= 1 + \frac{1}{7}\frac{f_{3\pi}(\mu_F)}{f_\pi\mu_\pi(\mu_F)}\Big(7\,\omega_{1,0}(\mu_F) - 2\,\omega_{2,0}(\mu_F) - \omega_{1,1}(\mu_F)\Big)\\ &\times \Big(10\,C_2^{1/2}(2\tau - 1) - 3\,C_4^{1/2}(2\tau - 1)\Big)\,, \quad \phi_{\pi\sigma}(\tau) = \ldots\end{aligned}$$

Parameters:

- $a_2(\mu_0) = 0.1364 \pm 0.0213$ at $\mu_0 = 2$ GeV [Braun et al '15] (lattice)
- $\omega_{10}(\mu_0) = -2.55\,, \omega_{10}(\mu_0) = 0.0$ and $f_{3\pi}(\mu_0) = 0.004\ \mathrm{GeV}^2$. [Ball '99]
- $\omega_{20}(\mu_0) = 8.0$ [Kroll, P-K '18] fit to π^0 photoproduction data [CLAS '17]

Evolution of the decay constants and DA parameters taken into account.
Choice of scales: $\mu_R{}^2 = \mu_F{}^2 = \hat{t}\hat{u}/\hat{s}$

η, η' distribution amplitudes

Twist-2 DA:

$$\phi_8(\tau,\mu_F) = 6\tau\bar{\tau}\left[1 + a_2^8(\mu_F)\, C_2^{3/2}(2\tau-1)\right]$$

$$\begin{aligned}\phi_{1,q}(\tau,\mu_F) &= 6\tau\bar{\tau}\left[1 + a_2^1(\mu_F)\, C_2^{3/2}(2\tau-1)\right]\\ \phi_{1,g}(\tau,\mu_F) &= 30\tau^2\bar{\tau}^2\left[1 + a_2^g(\mu_F)\, C_1^{5/2}(2\tau-1)\right]\end{aligned}$$

Twist-3 DAs:

assumption

$$\phi_{38}(\tau_a,\tau_b,\tau_g,\mu_F) = \phi_{31}(\tau_a,\tau_b,\tau_g,\mu_F) \approx \phi_{3\pi}(\tau_a,\tau_b,\tau_g,\mu_F)$$

Parameters:

- $a_2^8(\mu_0) = -0.039$, $a_2^1(\mu_0) = -0.057$, $a_2^g(\mu_0) = 0.038$ [Kroll, KPK '13], and other choices tested
- $f_{38}(\mu_0) = 0.86 f_{3\pi}(\mu_0) \Leftarrow$ [Ball '99; Braun, Filyanov '90]
- $f_{31}(\mu_0) = 0.86 f_{3\pi}(\mu_0) \Leftarrow \eta$ exp: [GlueX preliminary '20]
- mixing parameters from [Feldmann, Kroll, Stech '98]

Form factors and GPDs

$R_i \ldots 1/x$ moment of $\xi = 0$ GPD (K_i)

- $R_V(\leftarrow H)$, $R_T(\leftarrow E)$ from nucleon form factor analysis [Diehl, Kroll '13]
- $R_A(\leftarrow \tilde{H})$ form factor analysis and WACS KLL asymmetry [Kroll '17]
- $S_T(\leftarrow H_T)$, $\bar{S}_T(\leftarrow \bar{E}_T)$ low $-t$ from DVMP analysis [Goloskokov, Kroll '11]
- $S_S(\leftarrow \tilde{H}_T) \cong \bar{S}_T/2$ ($\bar{E}_T = 2\tilde{H}_T + E_T$)

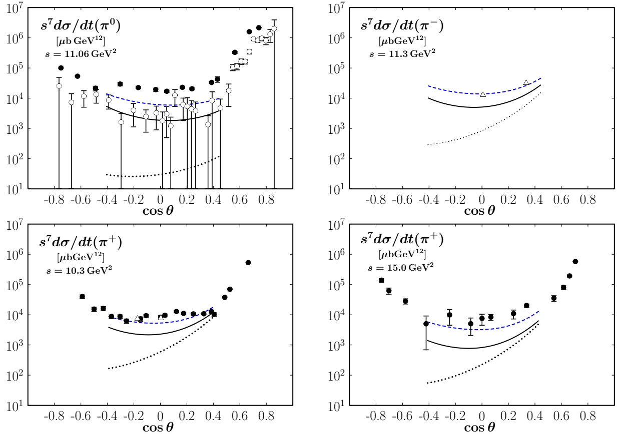
GPD parameterization [Diehl, Feldmann, Jakob, Kroll '04, Diehl, Kroll '13]

$$K_i^a = k_i^a(x) \exp\left[t f_i^a(x)\right], \; f_i^a(x) = \left(B_i^a - {\alpha'}_i^a \ln x\right)(1-x)^3 + A_i^a x(1-x)^2$$

- strong $x - t$ correlation
- power behaviour for large $(-t)$
- choice for transversity GPDs $A = 0.5$ GeV^{-2}

Photoproduction (π)

[Kroll, P-K '21]



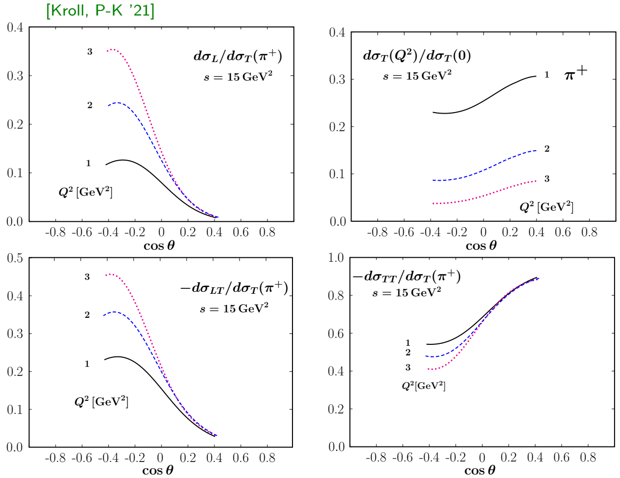
theoretical predictions with parameters from [Kroll, P-K '18] (fit of π^0 twist-3 prediction to [CLAS '17] data)

solid curves: complete twist-3
dotted curves: twist-2
dashed curves: $\omega_{20} = 10.3$
$\mu_R = \mu_F = 1$ GeV

exp data:
full circles [SLAC '76]
open circles [CLAS '17]
triangles [JLab, Hall A '05]

- twist-2 prediction well beyond the data [Huang, Kroll '00]
- scaling: s^{-7} (s^{-8}) twist-2 (twist-3) $\rightarrow$ effective s^{-9} $\rightarrow$ too strong

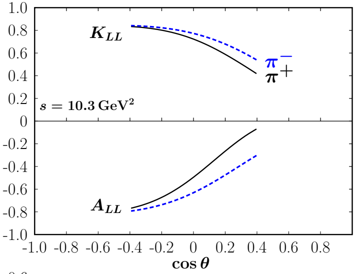
Electroproduction (π)

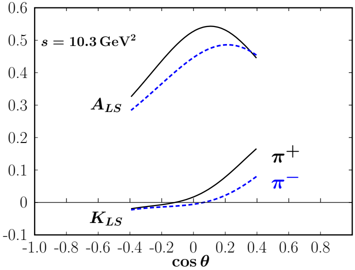


- both for σ_L and σ_{LT} no twist-2 and twist-3 interference $\Rightarrow$ information on S_T (H_T)
- information on S_S ($\tilde{H}_T$) from σ_{TT} (suppressed for DVMP)

Spin effects - photoproduction

[Kroll, P-K '21]





$A_{LL}(K_{LL})\ldots$ correlation of the helicities of the photon and incoming (outgoing) nucleon

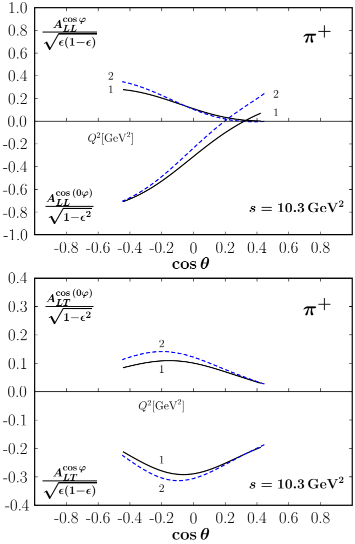
$$A_{LL}^{P,tw2} = K_{LL}^{P,tw2}$$
$$A_{LL}^{P,tw3} = -K_{LL}^{P,tw3}$$

$\rightarrow$ characteristic signature for dominance of twist-3 (like $\sigma_T \gg \sigma_L$ in DVMP)

$A_{LS}(K_{LS})\ldots$ correlation of the helicities of the photon and sideway polarization of the incoming (outgoing) nucleon

Spin effects - electroproduction

[Kroll, P-K '21]



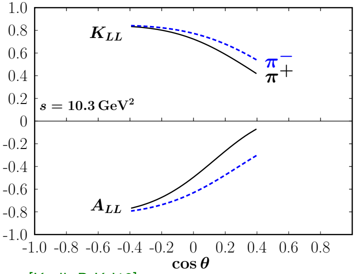
$A_{LL}(K_{LL})$ have two modulations for electroproduction

($\rightarrow$ measured for DVMP [CLAS '15])

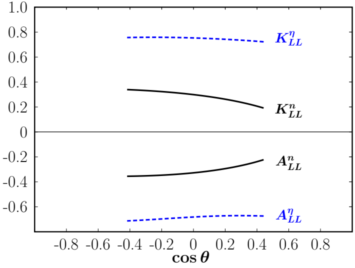
$A_{LT}(K_{LT})\ldots$ correlation between the lepton helicity and transversal target polarization

Spin effects - photoproduction

[Kroll, P-K '21]



[Kroll, P-K '18]



$A_{LL}(K_{LL})\ldots$ correlation of the helicities of the photon and incoming (outgoing) nucleon

$$A^{P,tw2}_{LL} = K^{P,tw2}_{LL}$$
$$A^{P,tw3}_{LL} = -K^{P,tw3}_{LL}$$

$\rightarrow$ characteristic signature for dominance of twist-3 (like $\sigma_T \gg \sigma_L$ in DVMP)

$A_{LL}(K_{LL})$ for π^0 photoproduction on neutron and η photoproduction