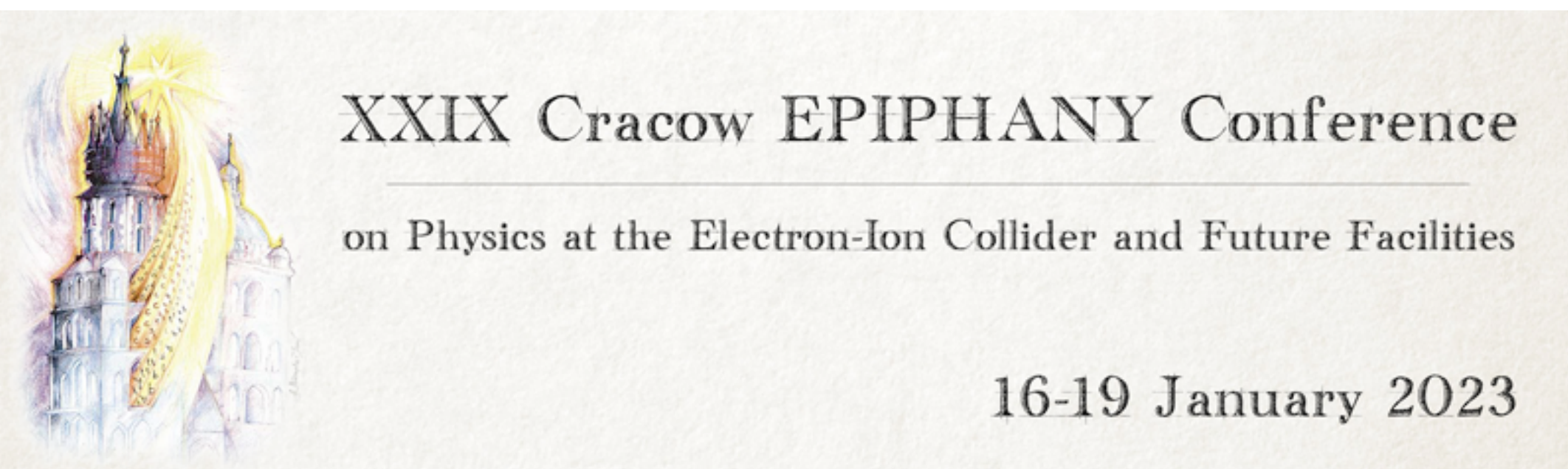


Polarised physics at the LHC

Pasquale Di Nezza



The LHC beams cannot be polarised



The only possibility to have polarised collisions is through a polarised fixed-target

Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before

The LHC beams cannot be polarised

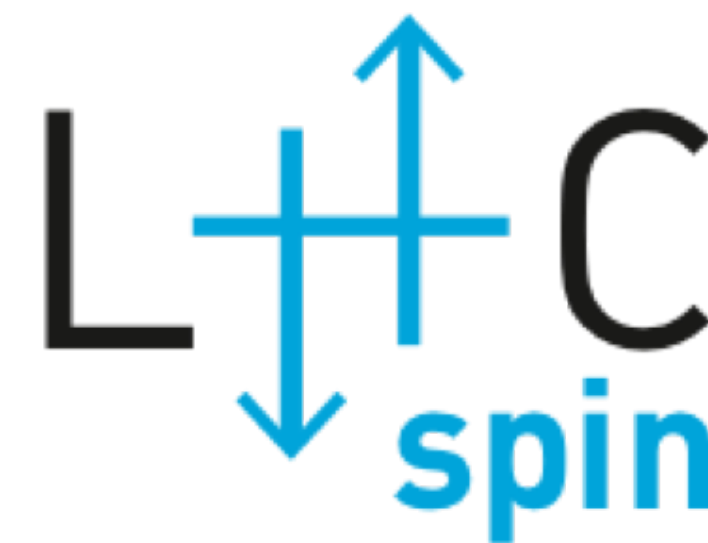


The only possibility to have polarised collisions is through a polarised fixed-target

Collisions provided by a TeV-scale beam (LHC) on fixed target will exploit a unique kinematic region poorly probed. Advanced detectors make available probes never accessed before

SMOG2

unpolarised gas target



polarised (+unpolarised) gas target

The LHCb detector

- LHCb is a general-purpose forward spectrometer, fully instrumented in $2 < \eta < 5$, and optimised for c and b hadron detection
- Excellent momentum resolution with VELO + tracking stations:

$$\sigma_p/p = 0.5 - 1.0 \% \quad (p \in [2, 200] \text{ GeV})$$

- Particle identification with RICH+CALO+MUON

$$\epsilon_\mu \sim 98 \% \quad \text{with} \quad \epsilon_{\pi \rightarrow \mu} \lesssim 1 \%$$

- Low momentum muon trigger:

$$p_{T_\mu} > 1.75 \text{ GeV (2018)}$$

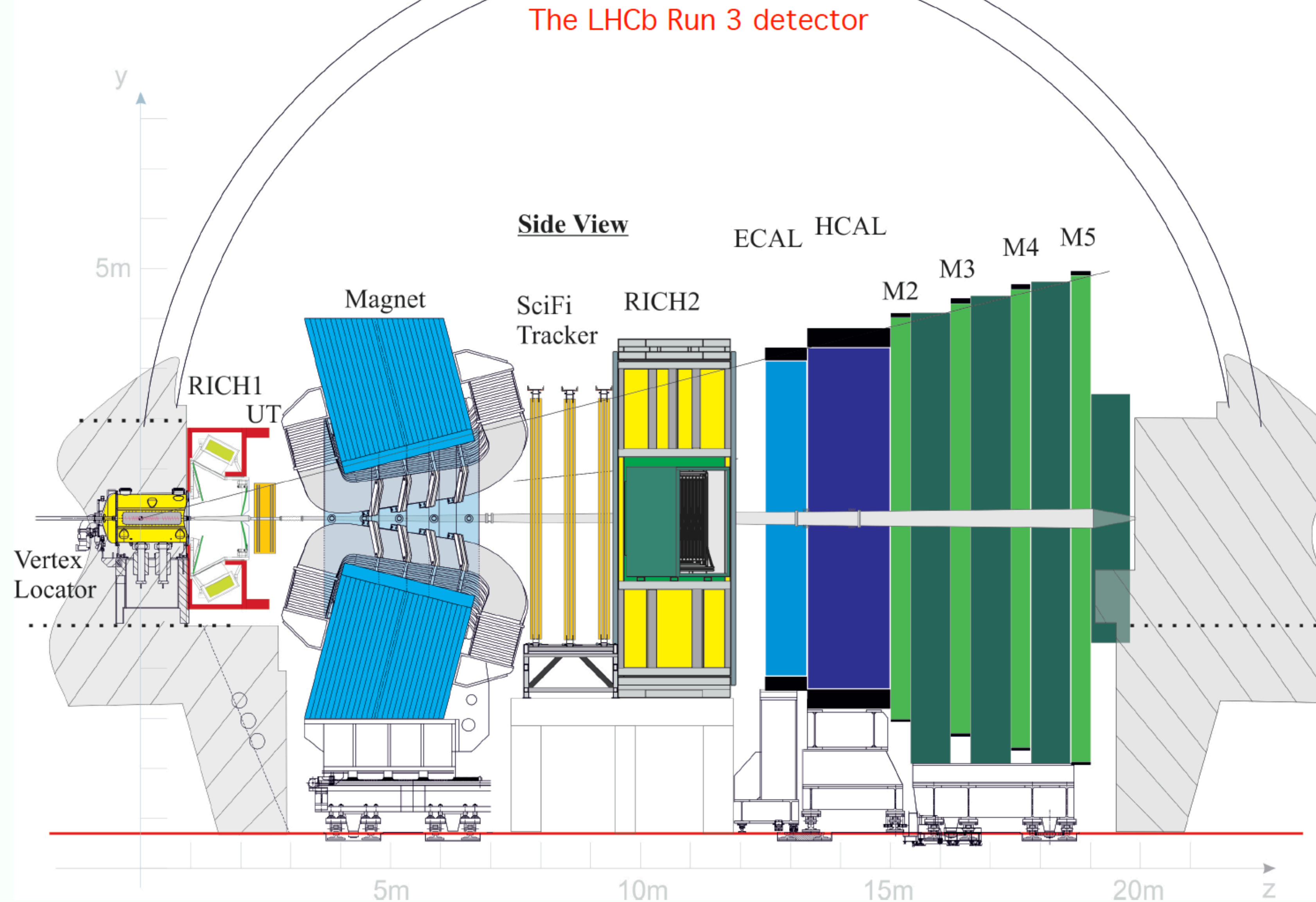
will be reduced thanks to the new fully-software trigger

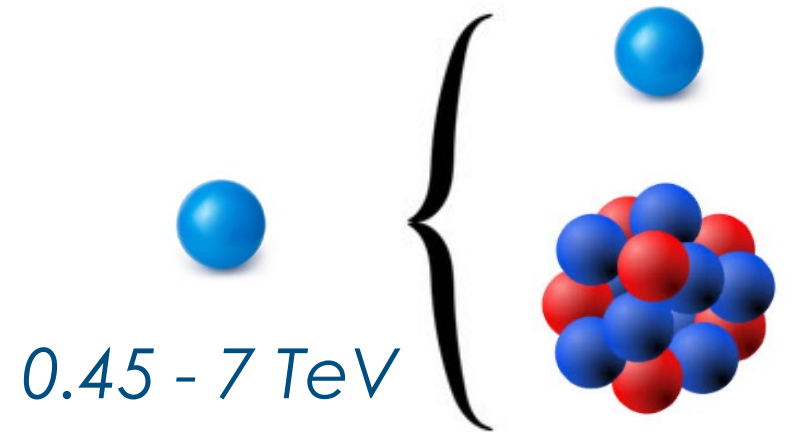
- Major detector upgrades performed during LS2 for the Run 3 (5x luminosity)

[[JINST 3 \(2008\) S08005](#)]

[[IJMP A 30, 1530022 \(2015\)](#)]

[[Comput Softw Big Sci 6, 1 \(2022\)](#)]

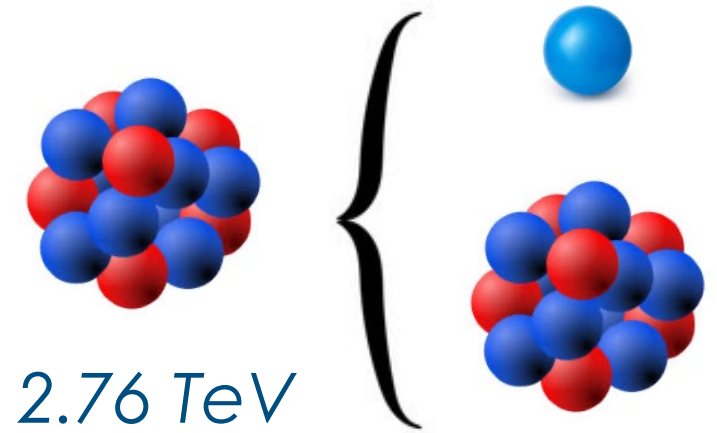




pp or pA collisions: 0.45 - 7 TeV beam on fix target

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 41 - 115 \text{ GeV}$$

$$y_{CMS} = 0 \rightarrow y_{lab} = 4.8$$



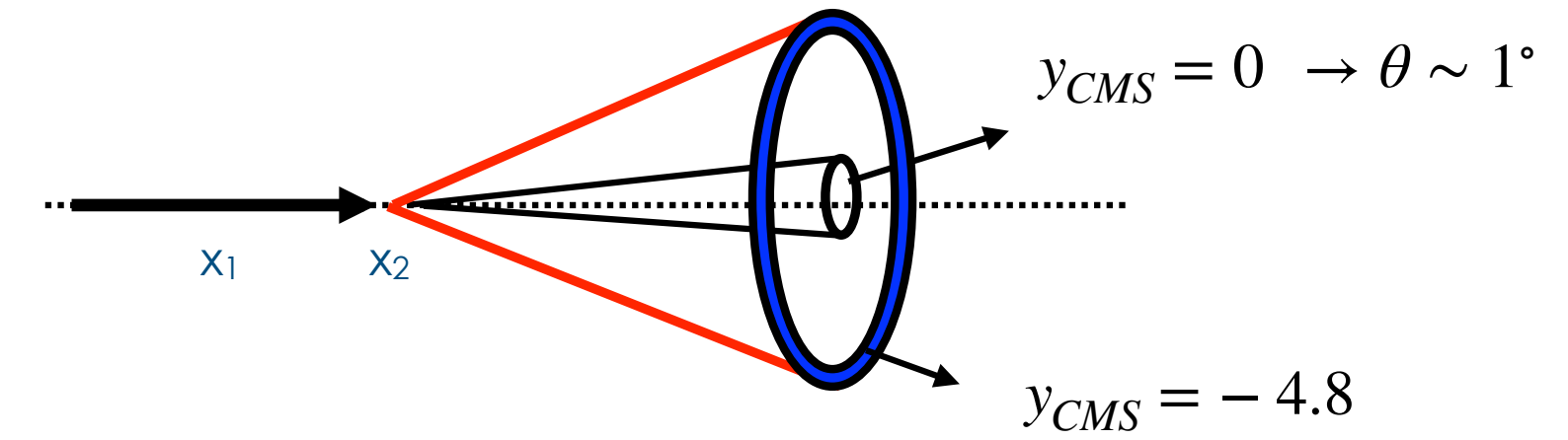
AA collisions: 2.76 TeV beam on fix target

$$\sqrt{s_{NN}} \simeq 72 \text{ GeV}$$

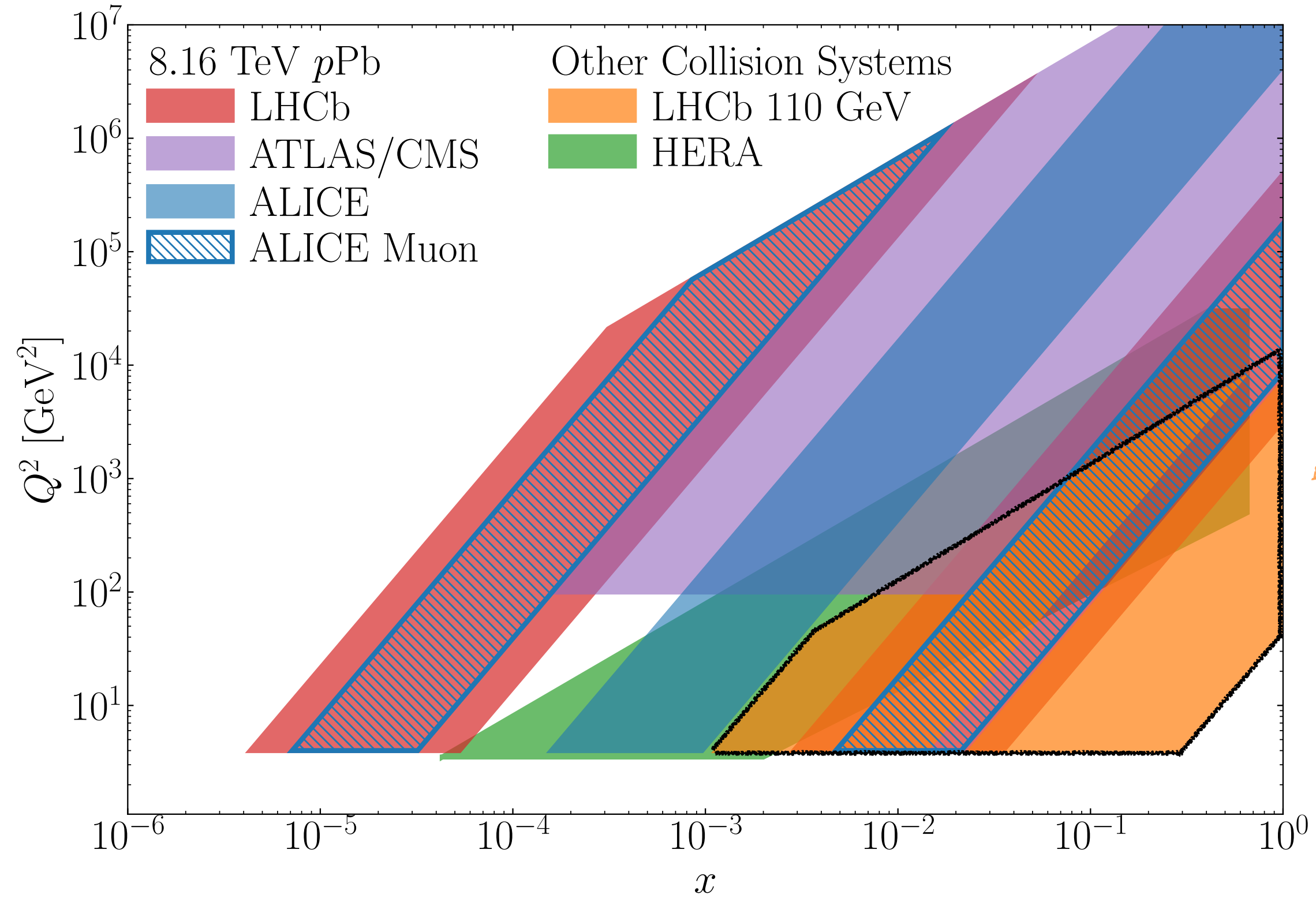
$$y_{CMS} = 0 \rightarrow y_{lab} = 4.3$$

1: beam; 2: target

Large CM boost, large x_2 values ($x_F < 0$) and small x_1



$$\gamma = \frac{\sqrt{s_{NN}}}{2m_p} \simeq 60$$



Broad and poorly explored kinematic range

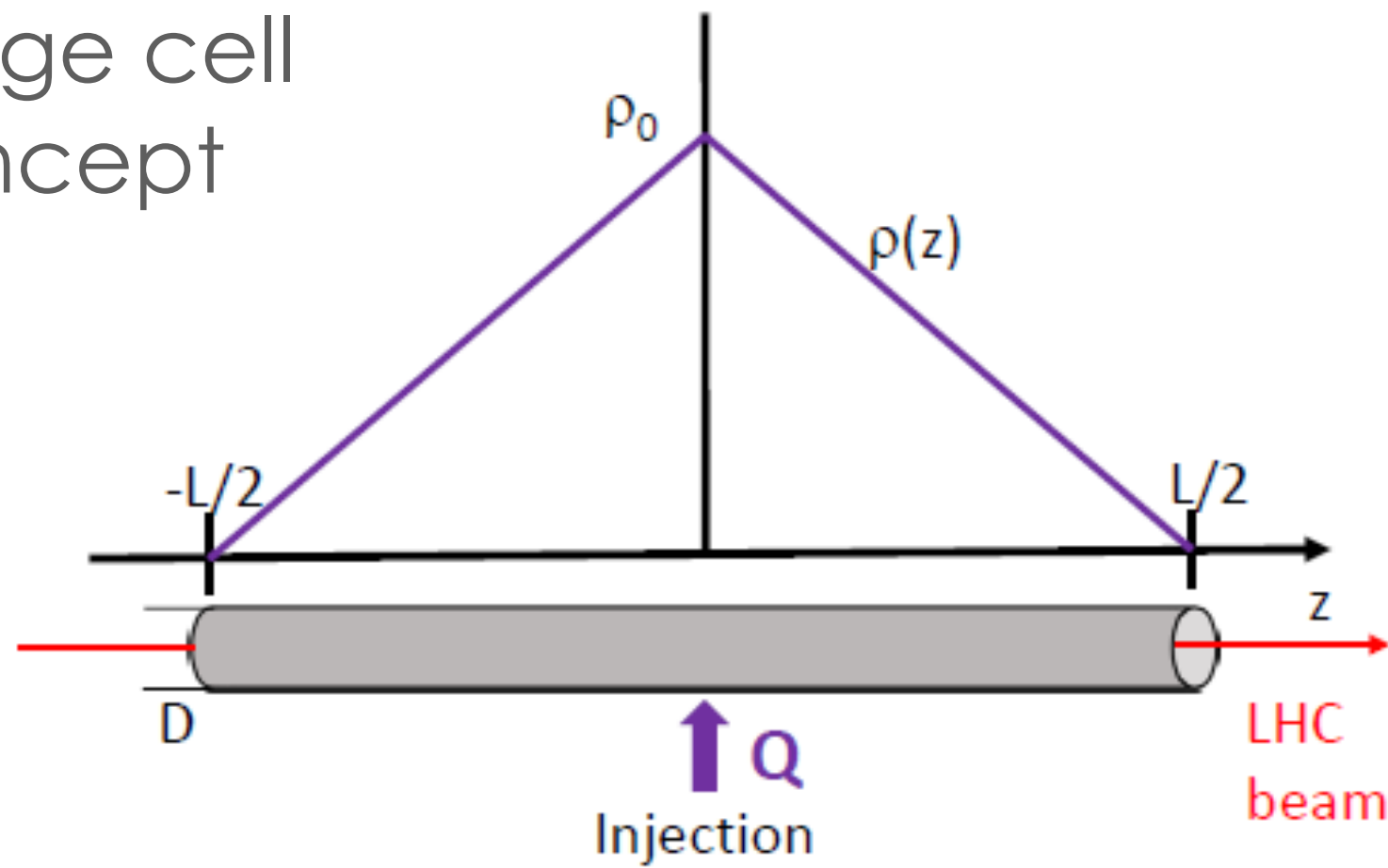
SMDQ2

an unpolarised target at

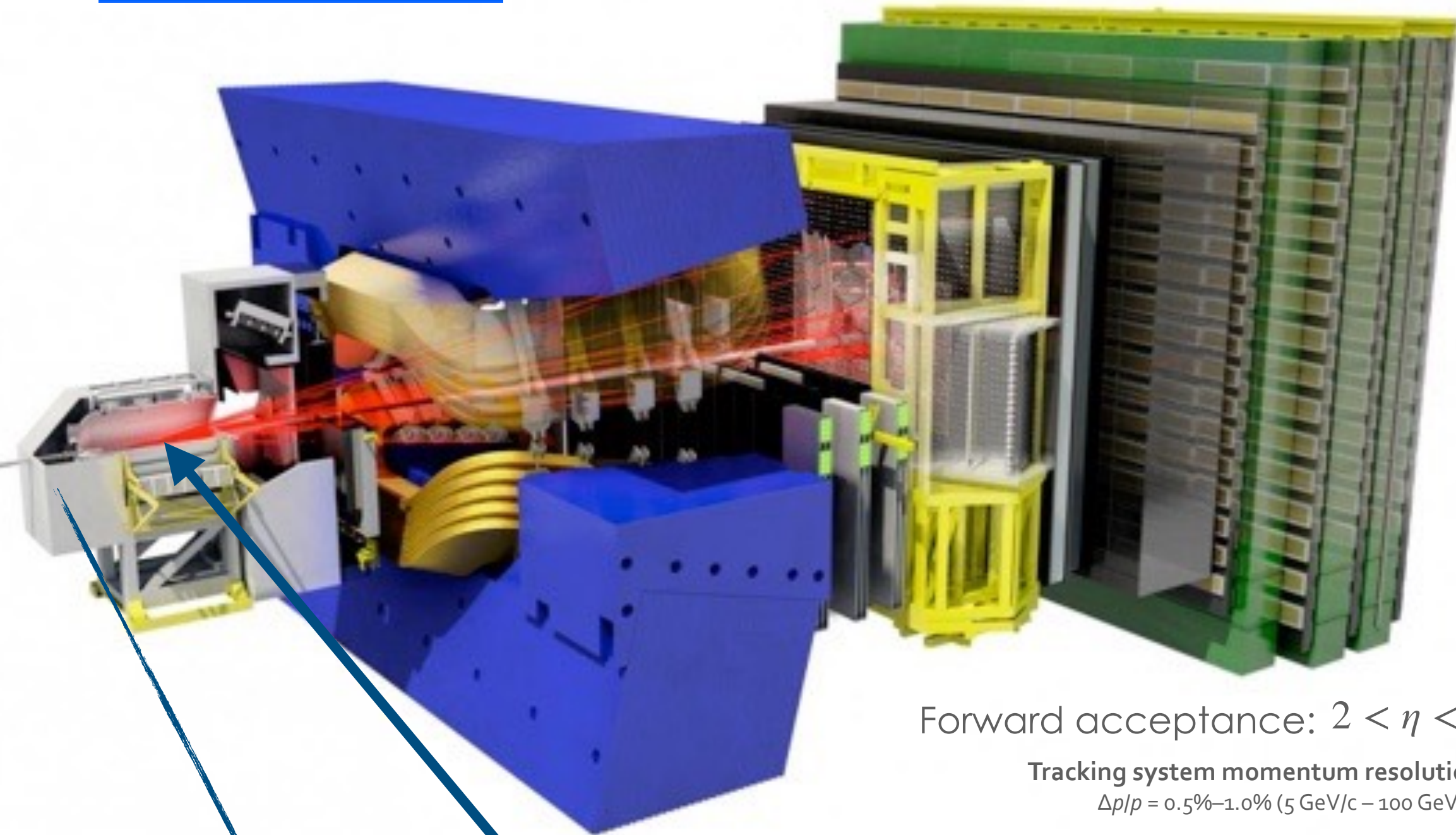


JINST 3 (2008) S08005
IJMPA 30 (2015) 1530022

Storage cell concept



LHC beam

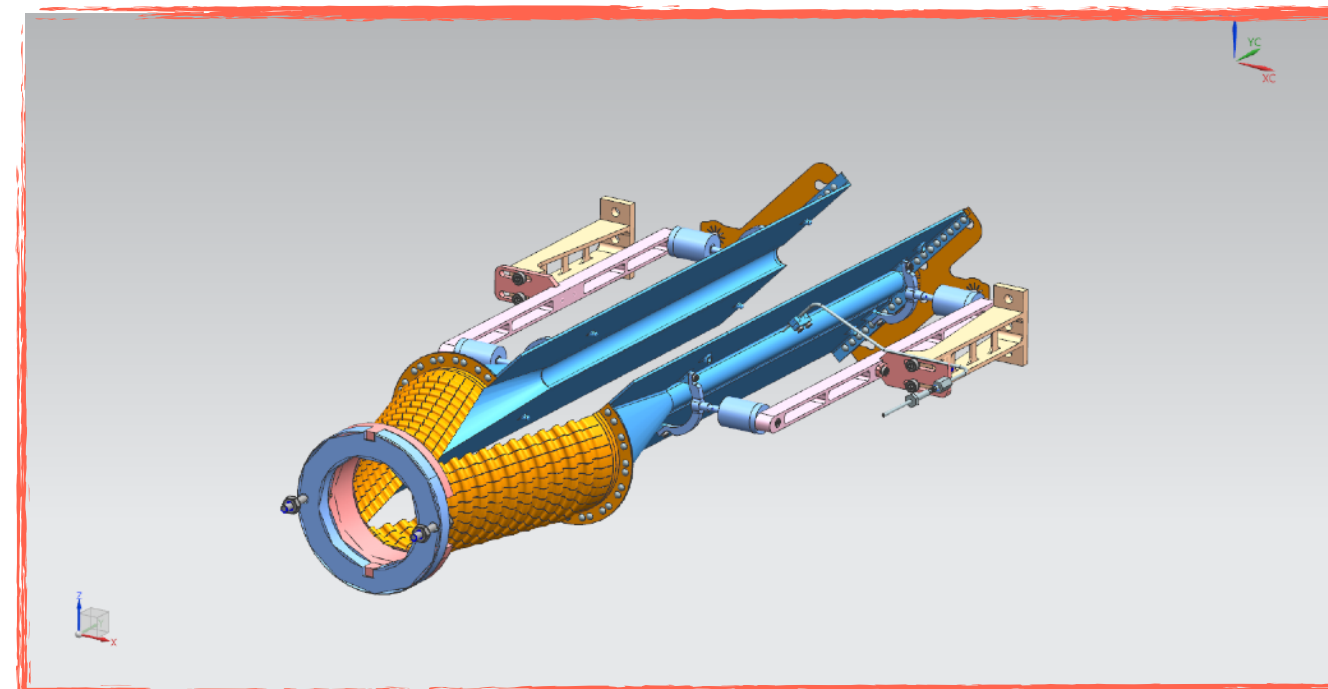
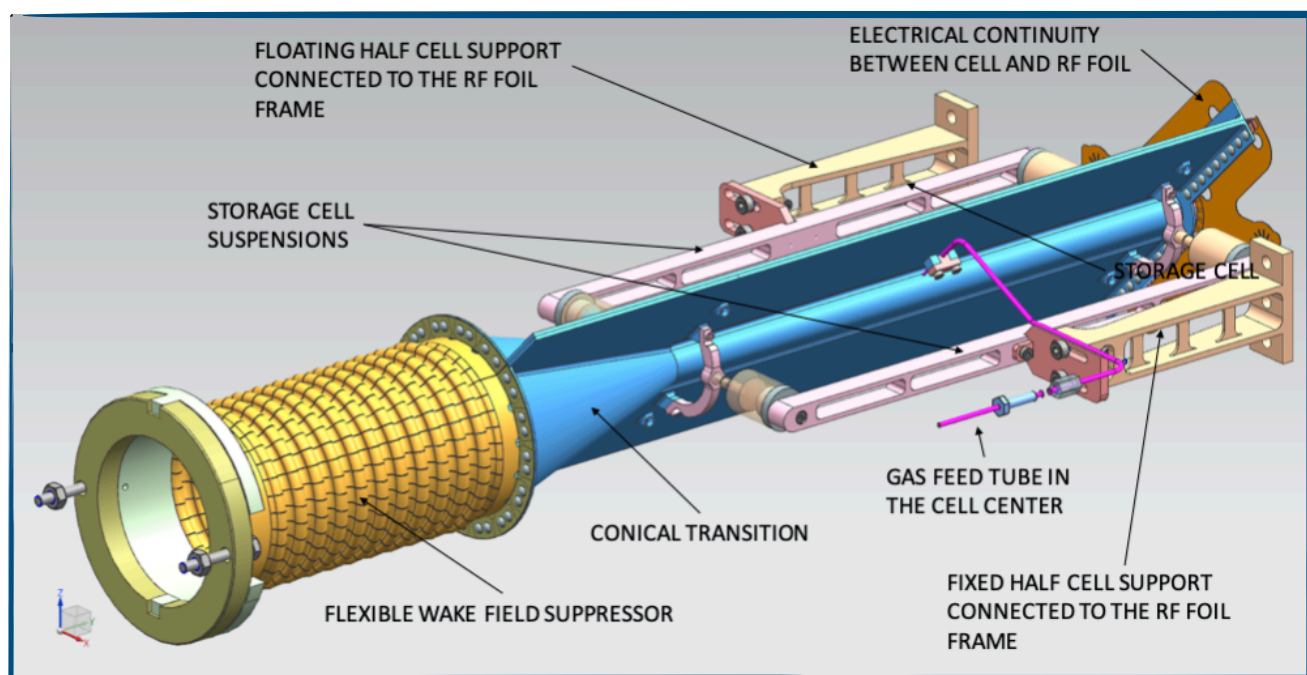


Forward acceptance: $2 < \eta < 5$

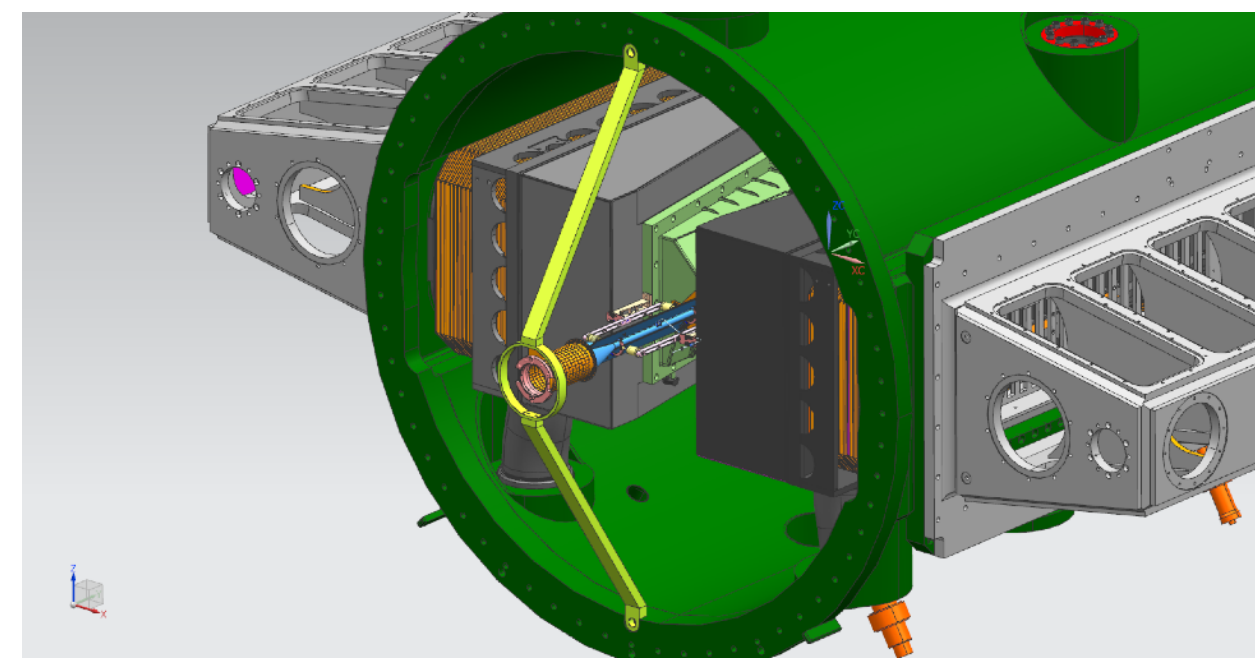
Tracking system momentum resolution
 $\Delta p/p = 0.5\% - 1.0\%$ (5 GeV/c - 100 GeV/c)

beam-beam collisions

Openable cell

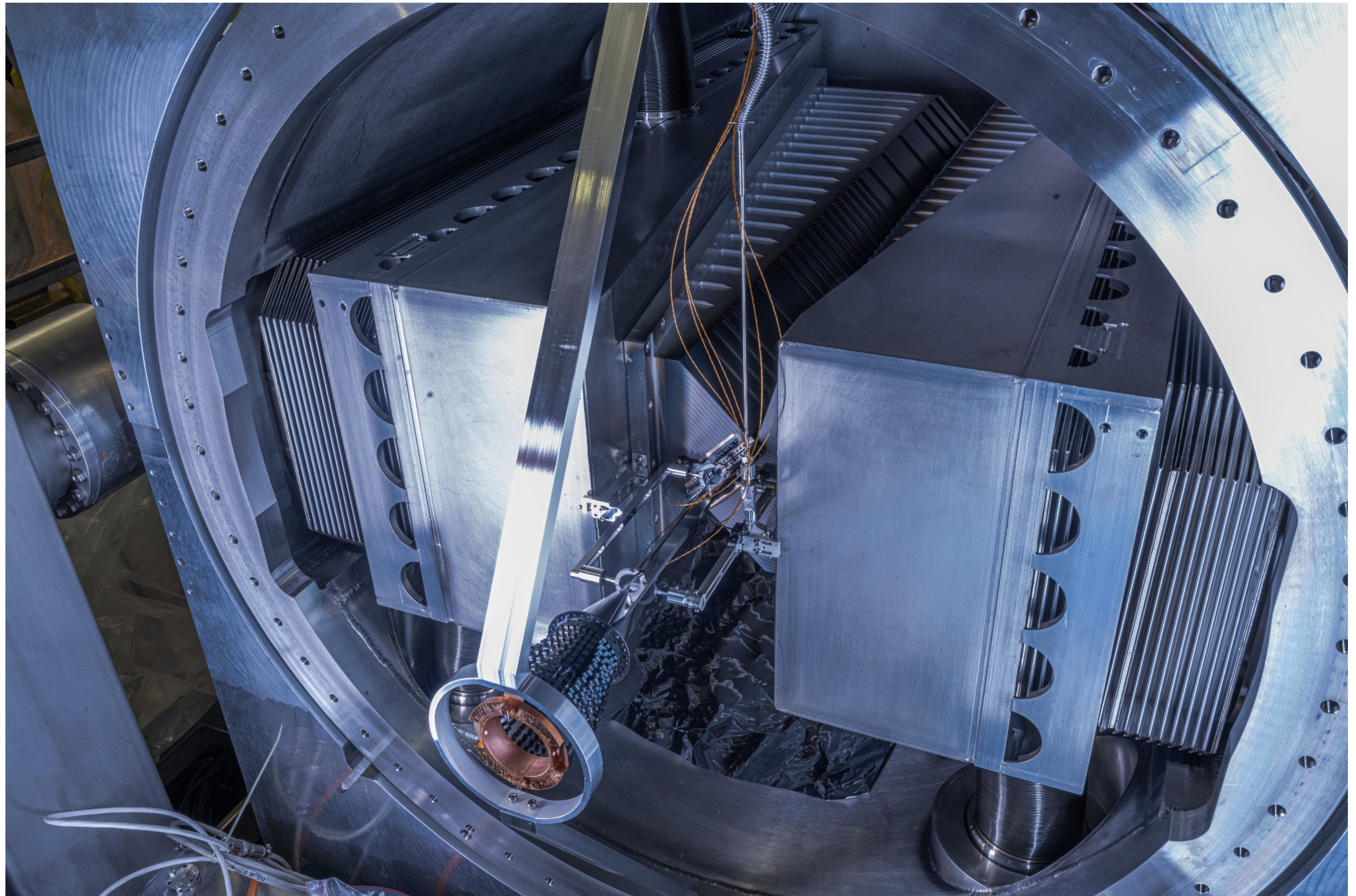


5 mm radius x 200 mm length



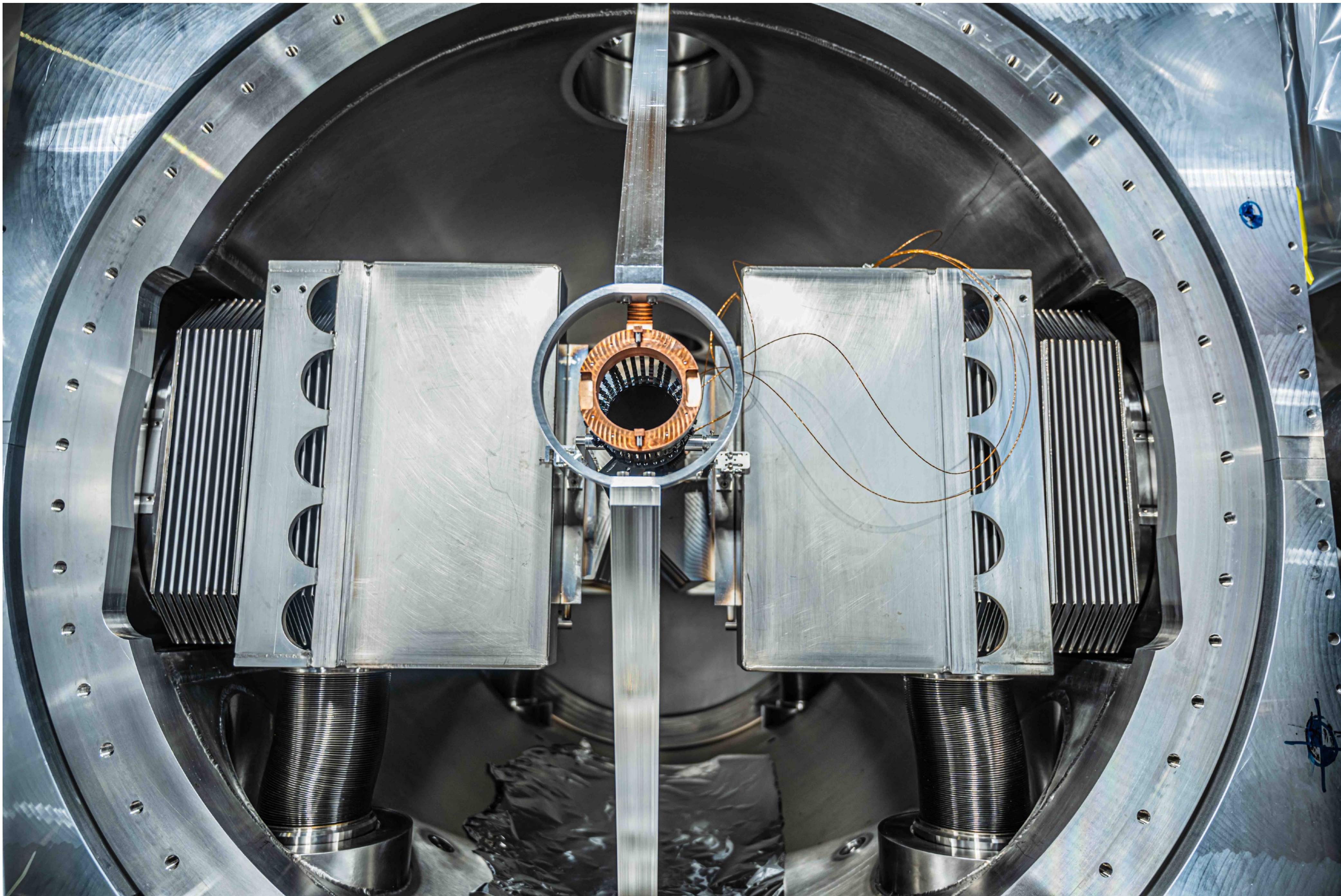
UNpolarised target
(beam-gas)

SMDQ2



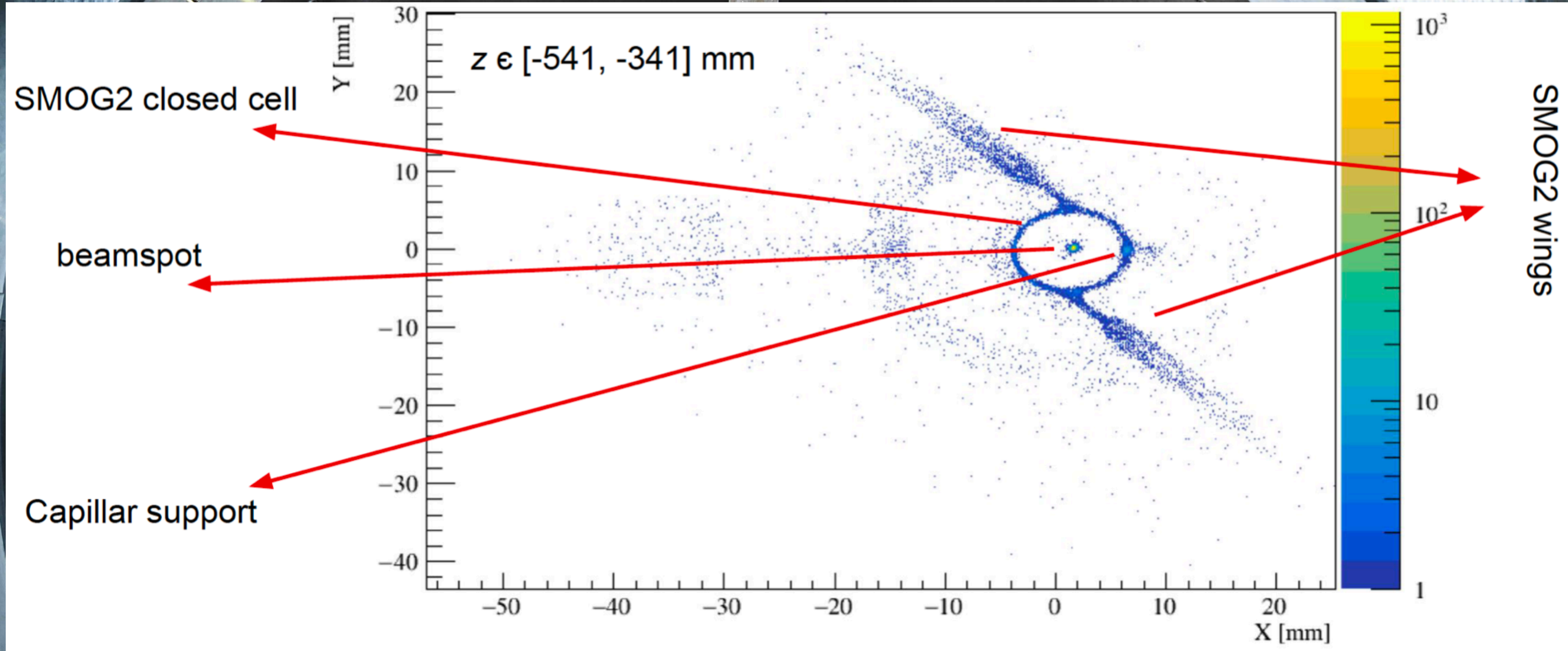
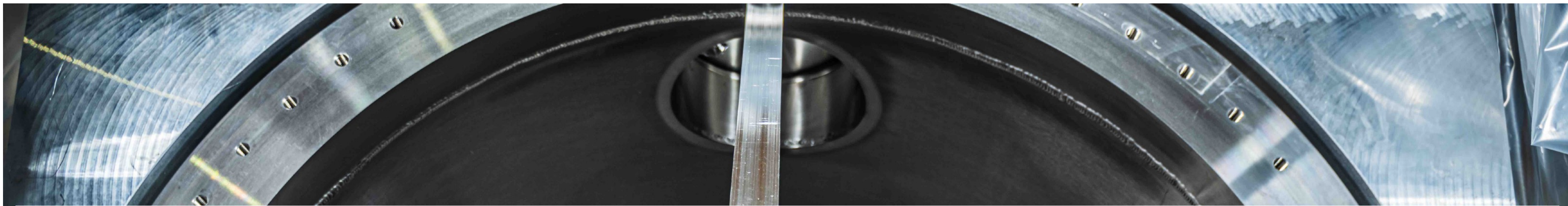
It is the only system present in the LHC primary vacuum

SMDQ2

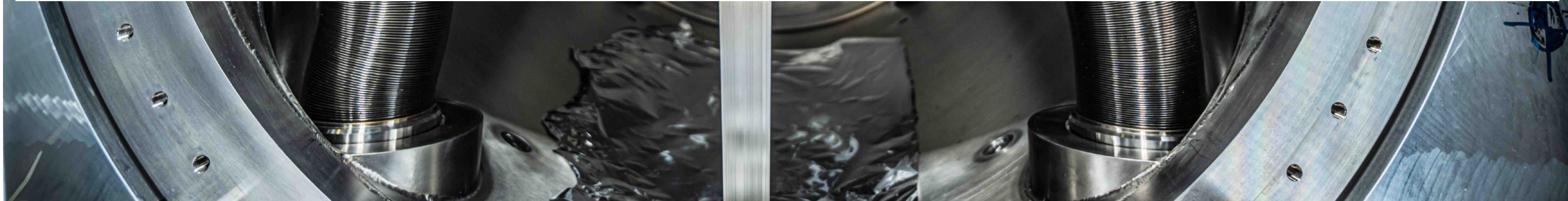


It is the only system present in the LHC primary vacuum

SMOG2

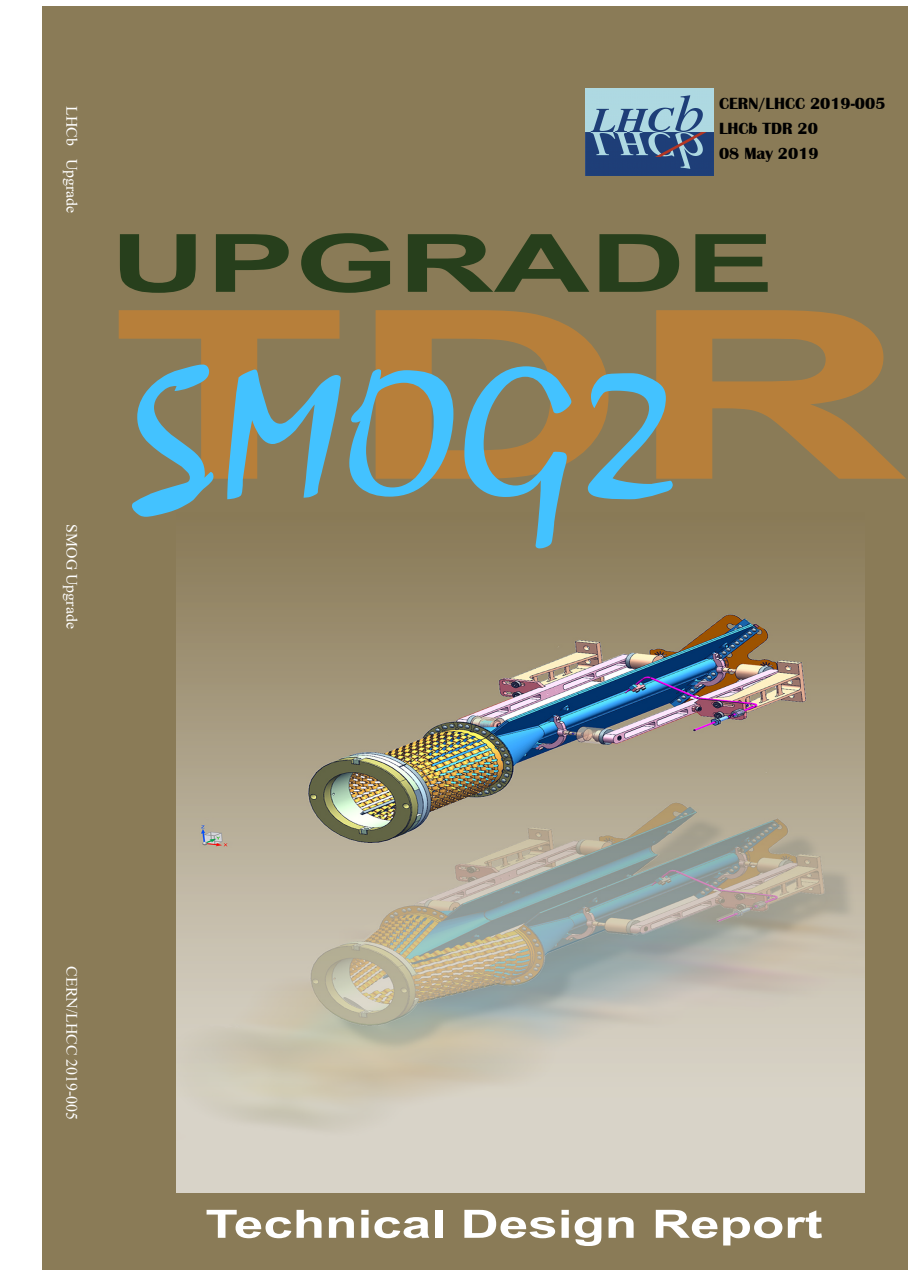


It is the only system present in the LHC primary vacuum

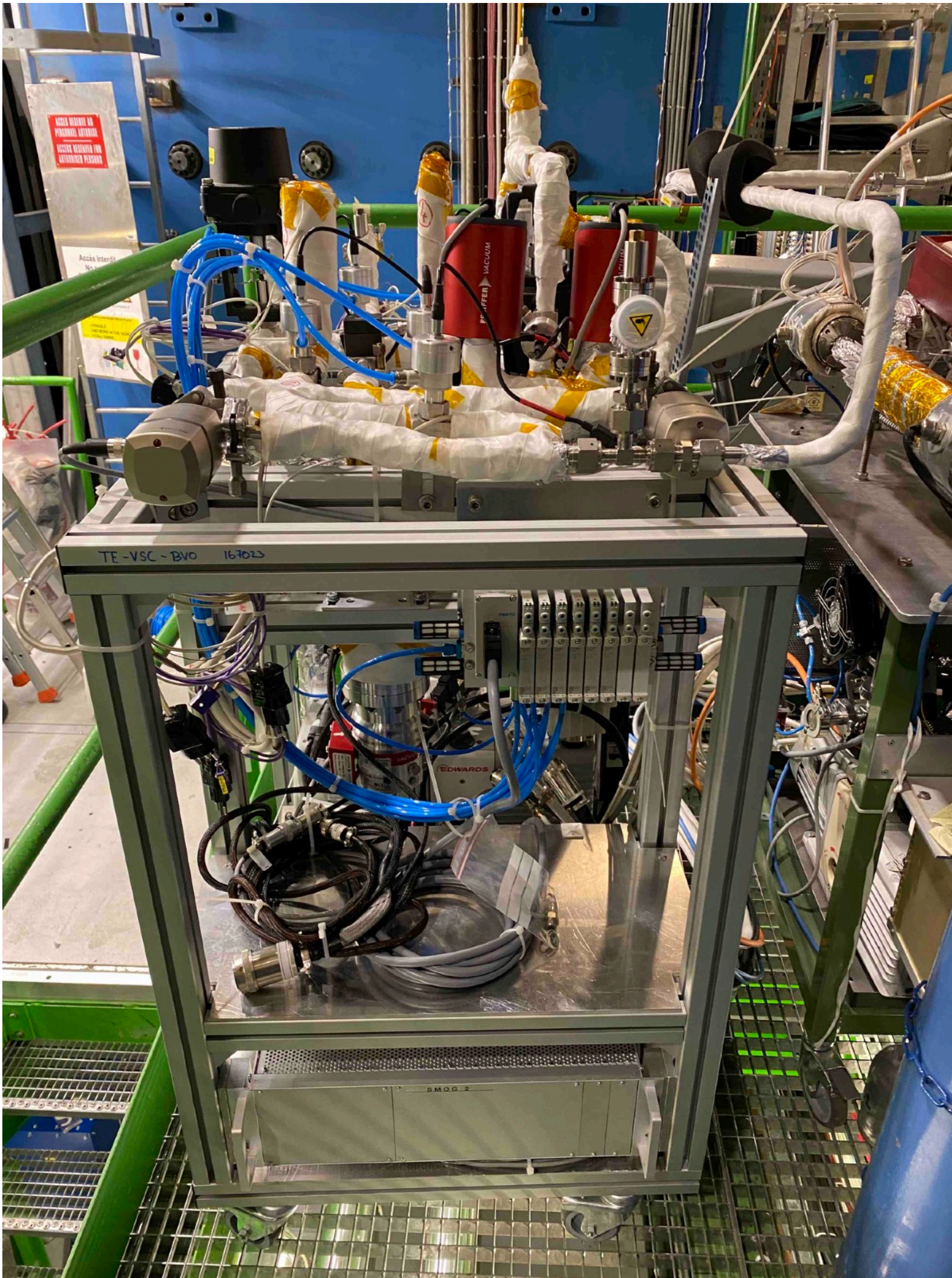


SMOG2

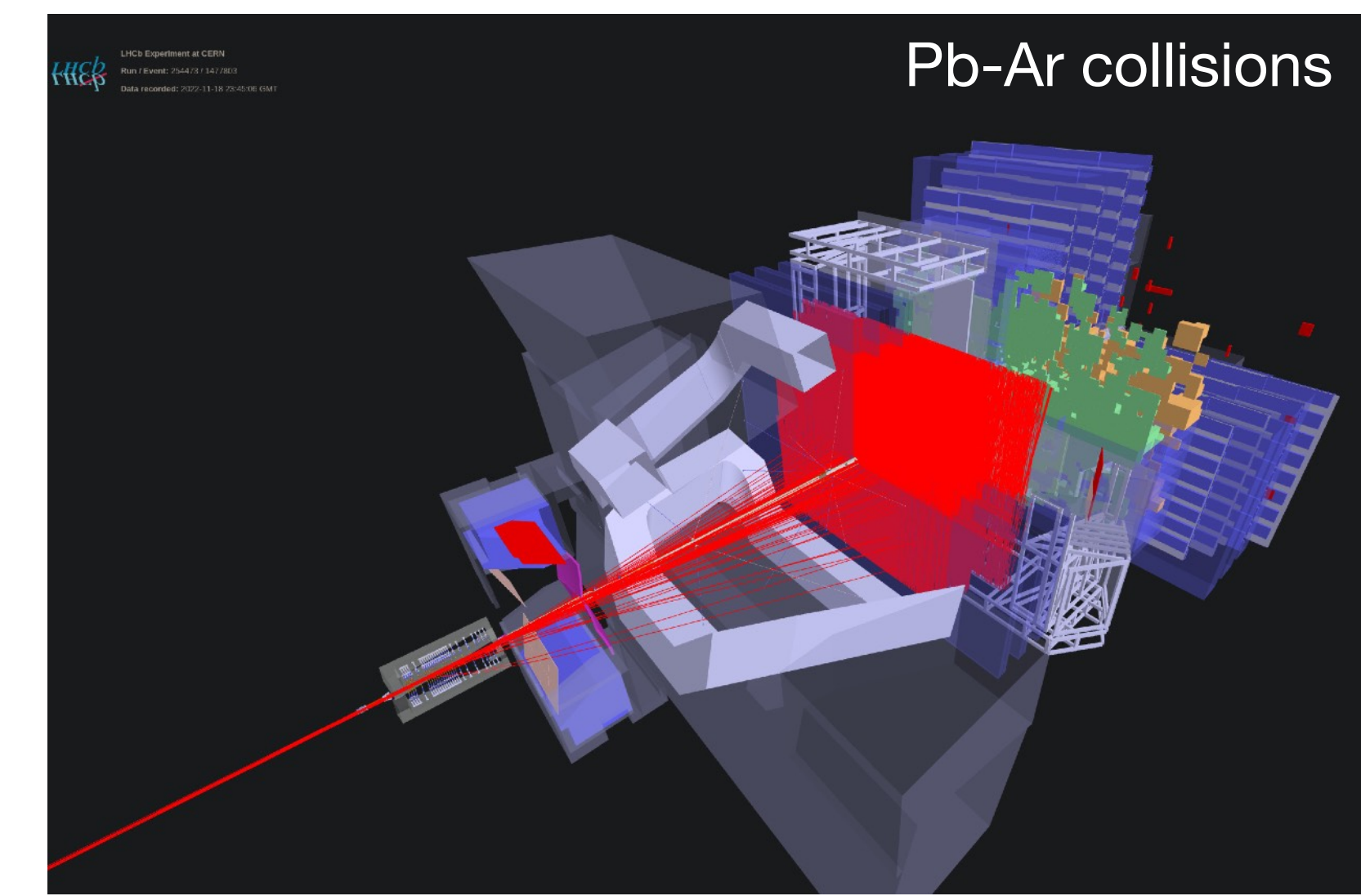
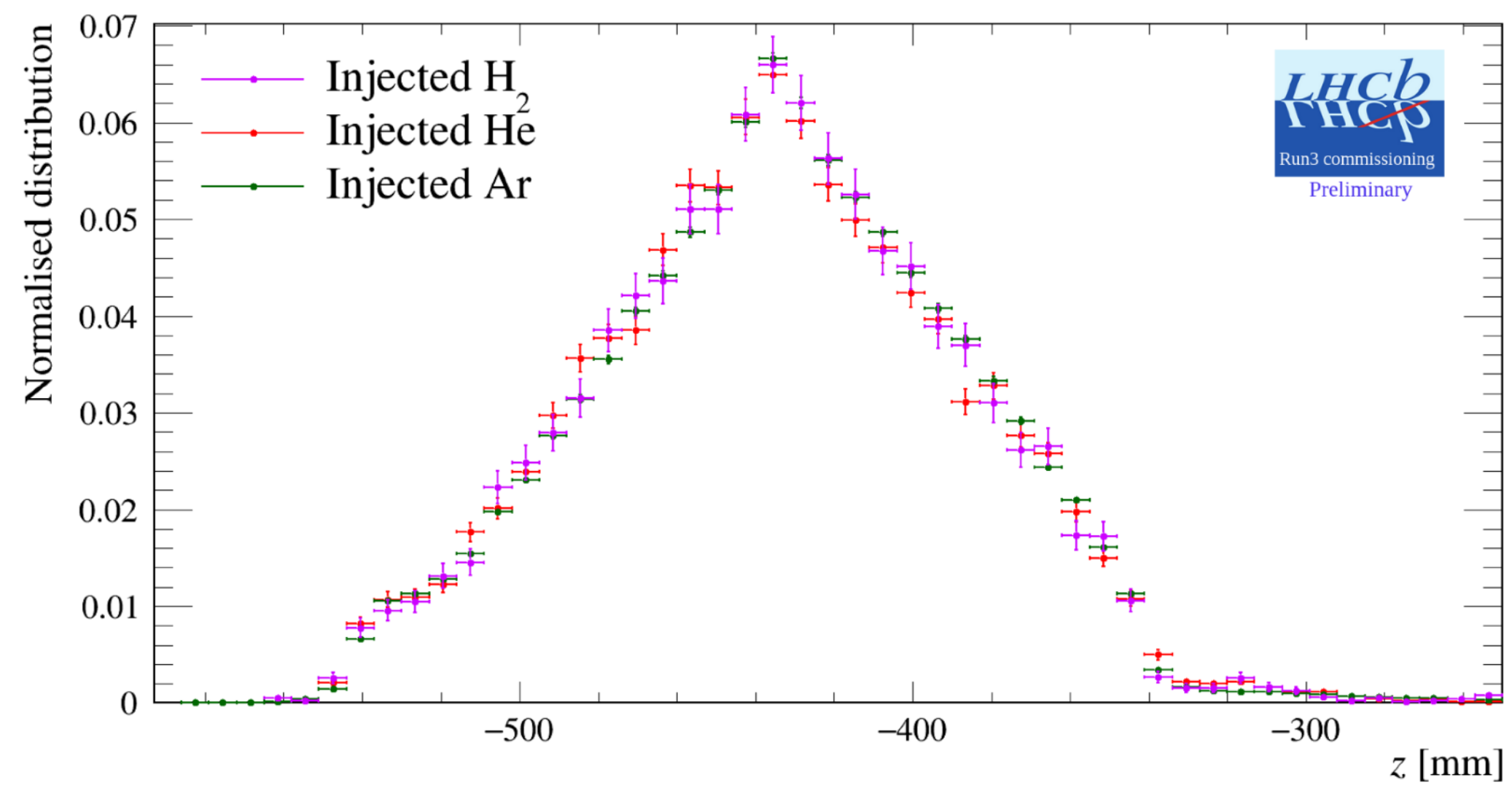
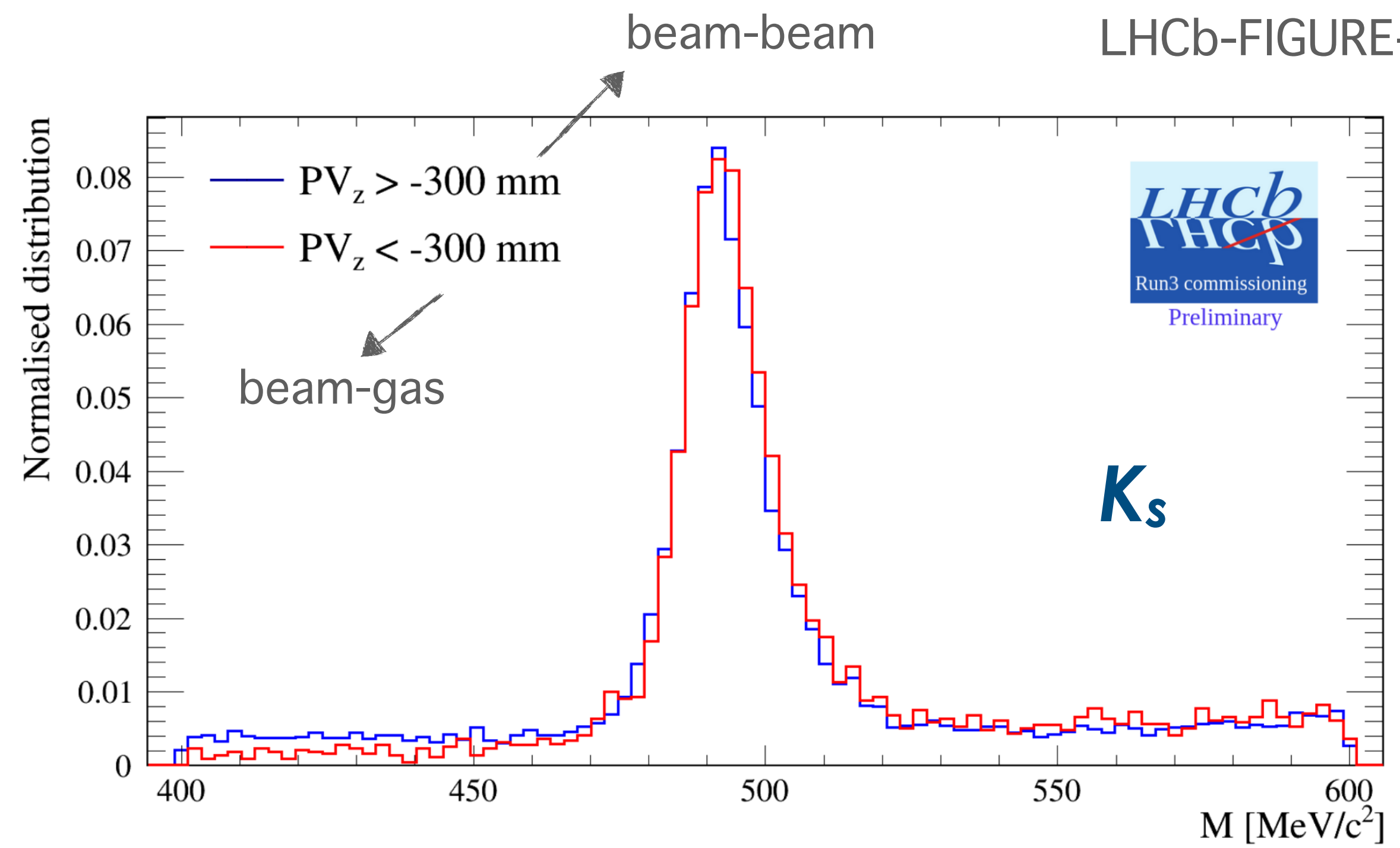
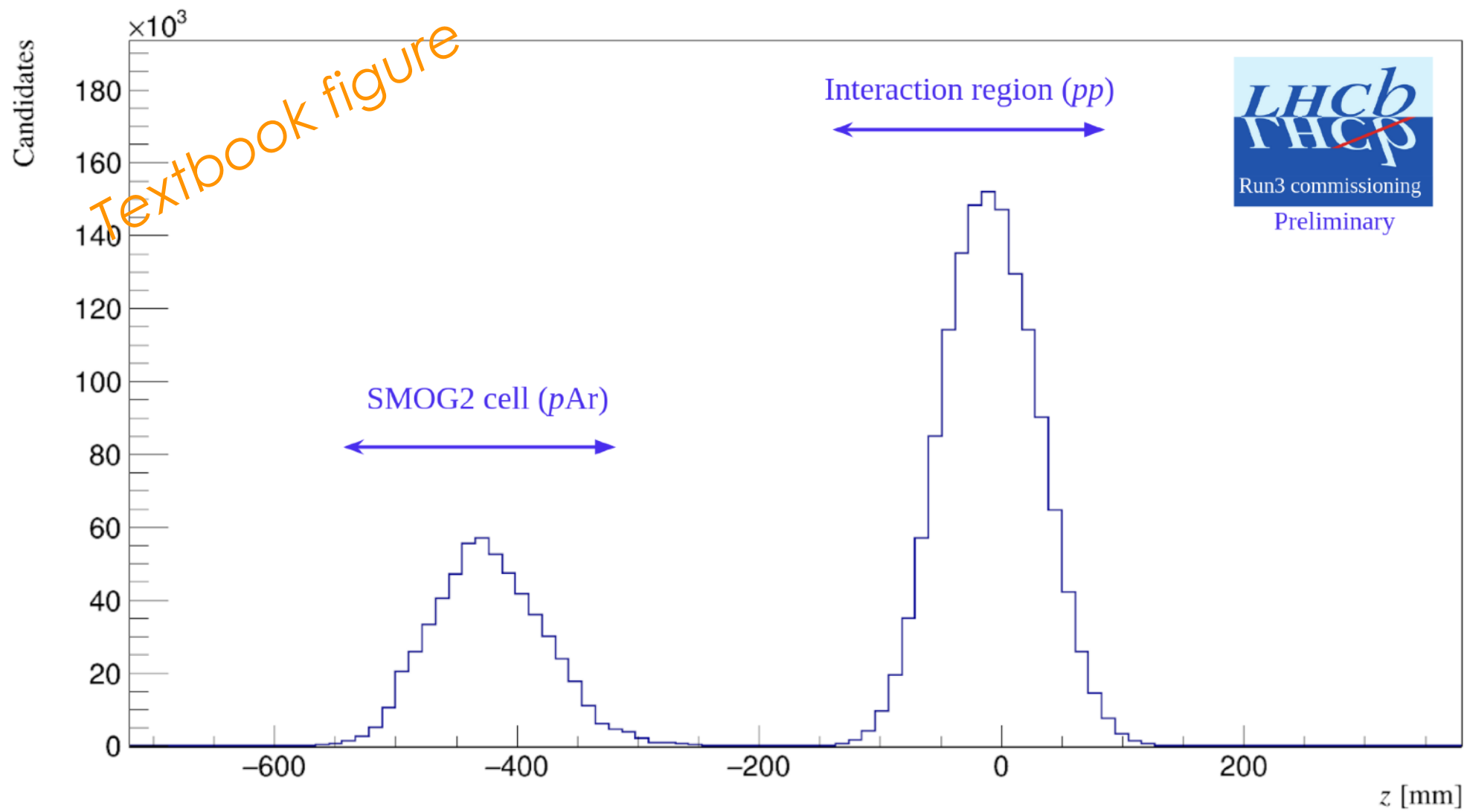
- The system is completely installed (storage cell + GFS + triggers + reconstruction)
- Negligible impact on the beam lifetime ($\tau_{beam-gas}^{p-H_2} \sim 2000$ days, $\tau_{beam-gas}^{Pb-Ar} \sim 500$ h)
- Injectable gases (3+1 reservoirs): H₂, D₂, N₂, O₂, He, Ne, Ar, Kr, Xe
- Flux known with 1% precision, measured relative contamination 10⁻⁴



GFS table installation



SMOG2 works!

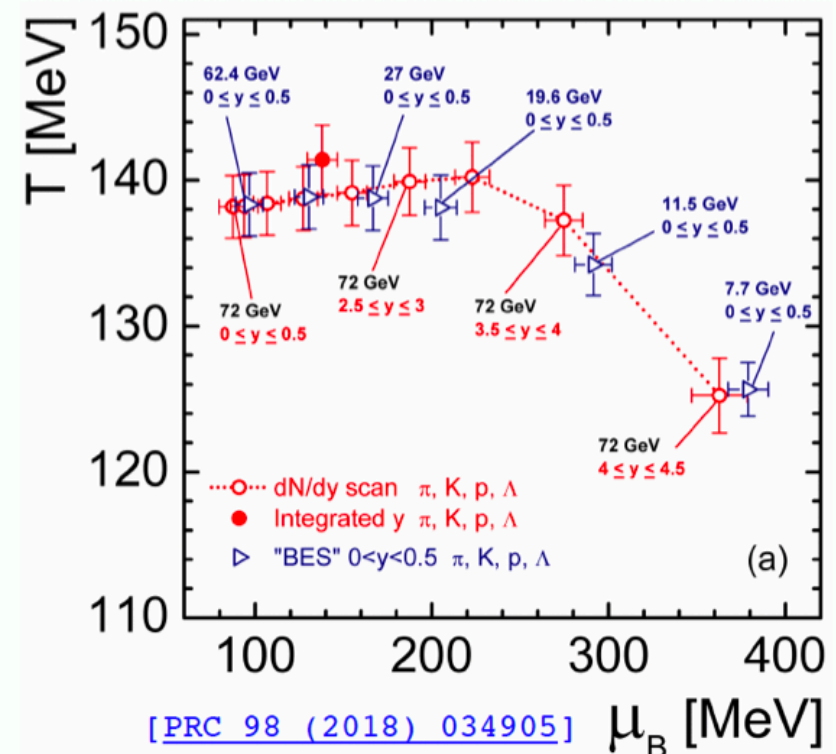
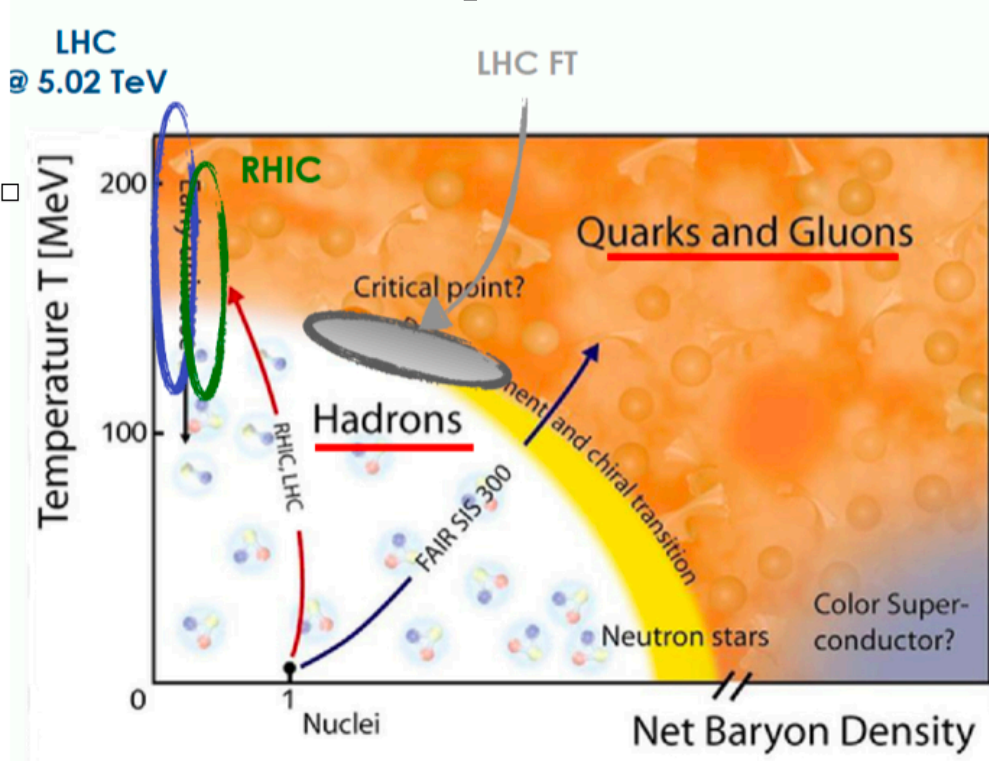


LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!

SMDQ2 ... few highlights

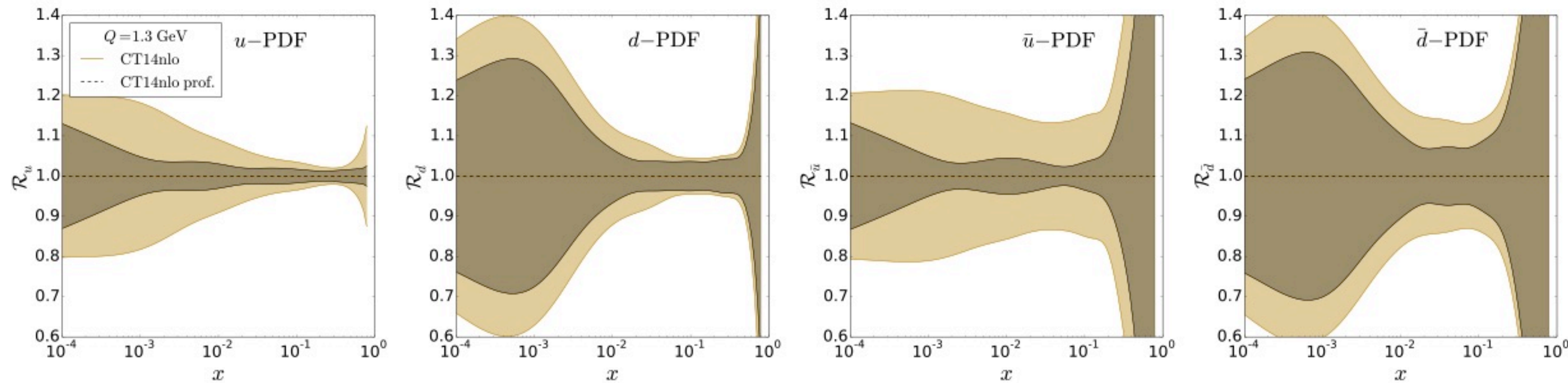
<http://cds.cern.ch/record/2649878/files/>

Heavy-Ion and QCD phase space



[PRC 98 (2018) 034905]

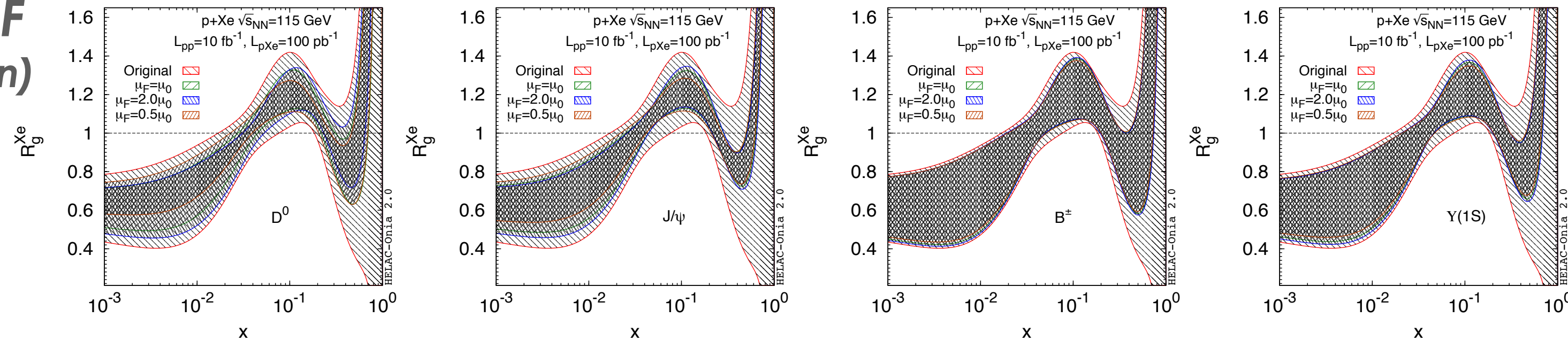
PDF



estimation with 10 fb⁻¹

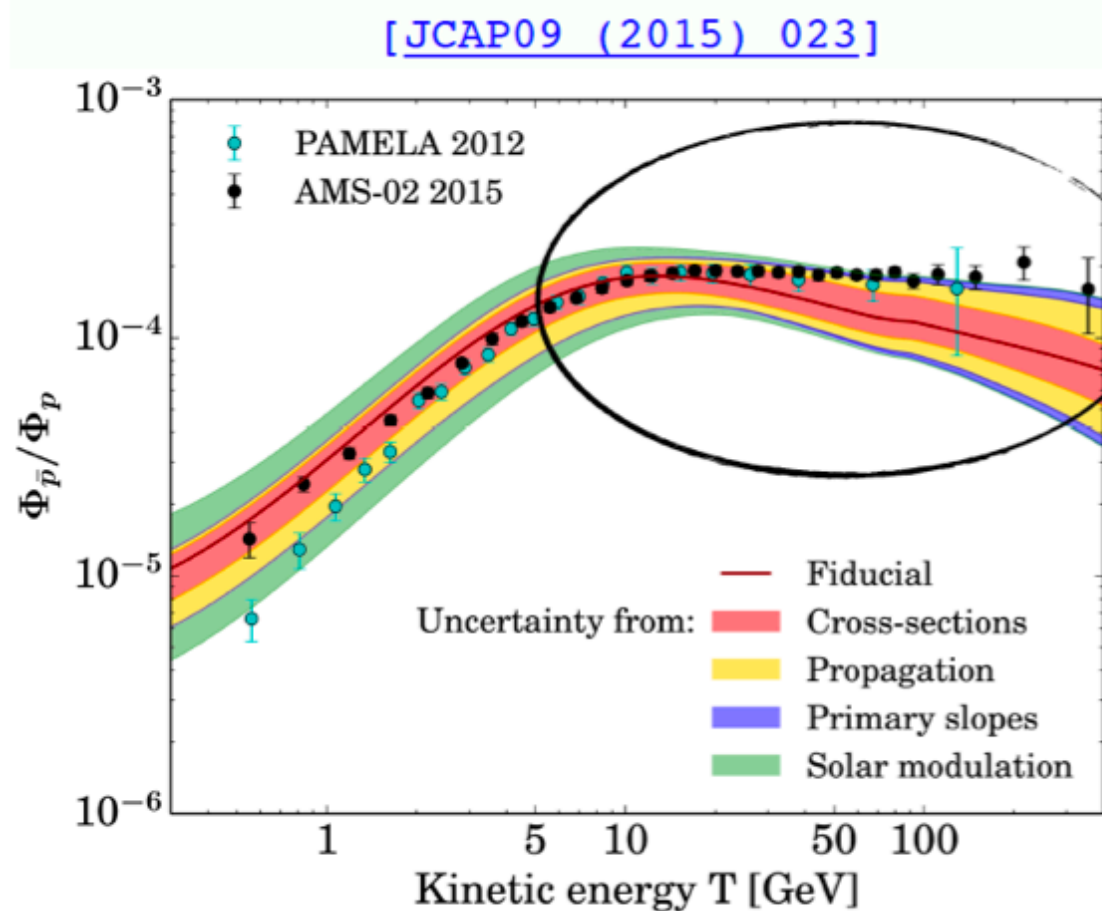
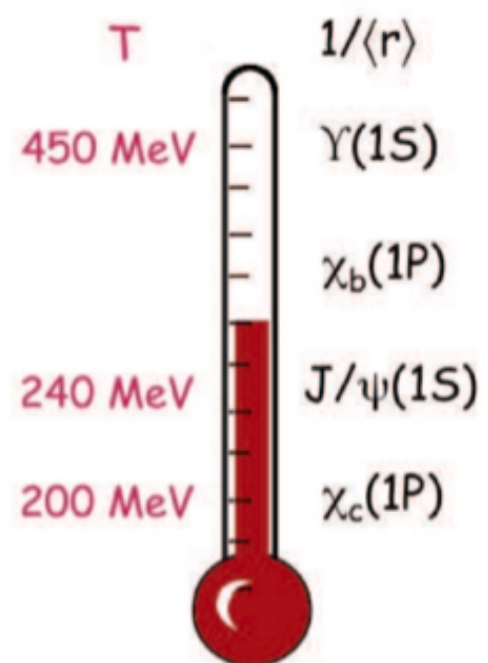
arXiv:1807.00603

nPDF (gluon)

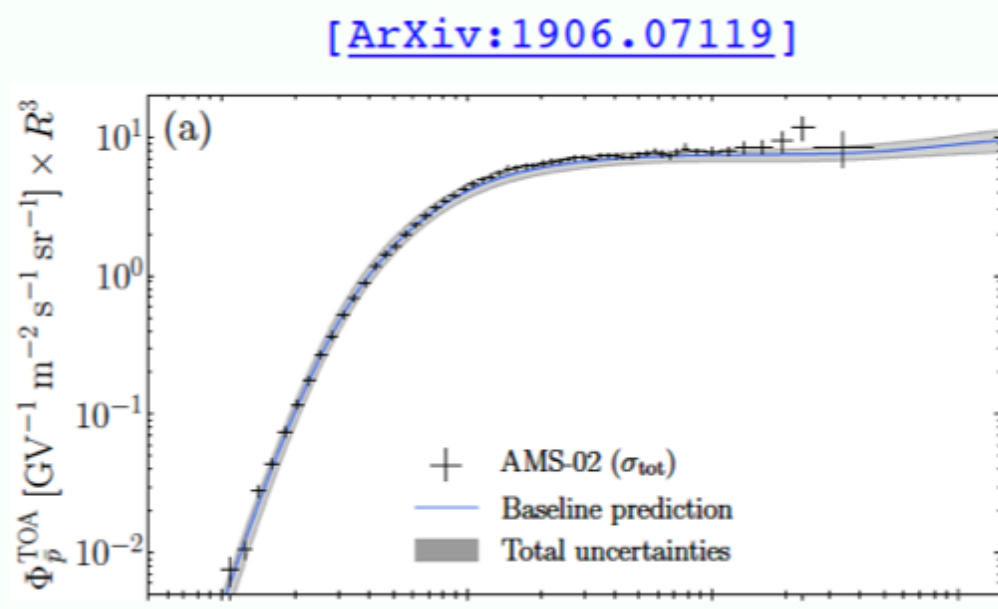


Astroparticle (DM and CR)

c \bar{c} bound states



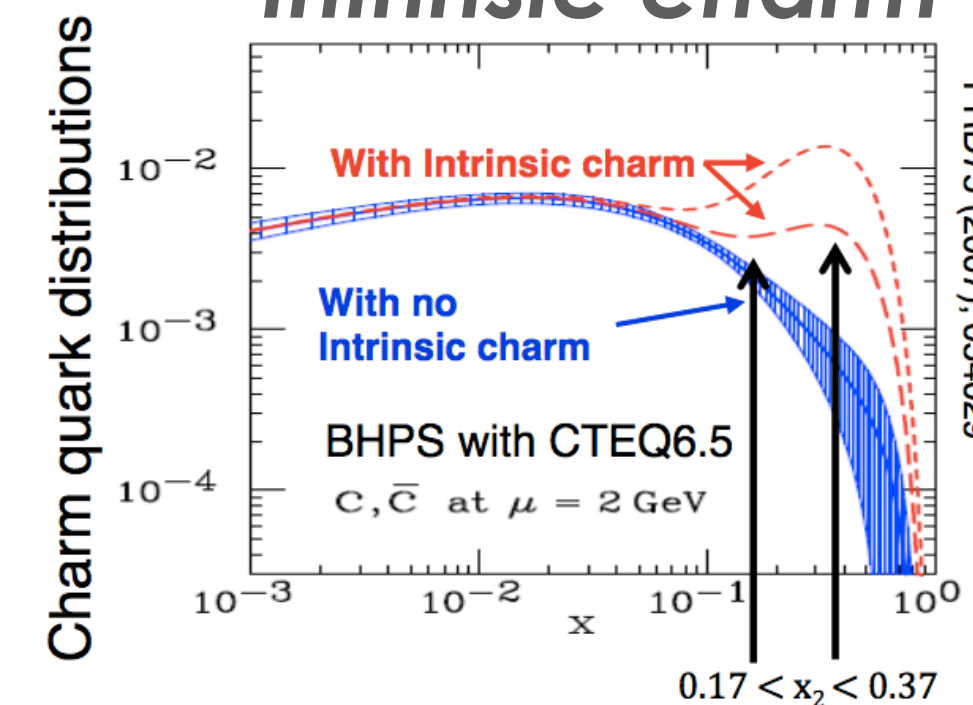
[JCAP09 (2015) 023]



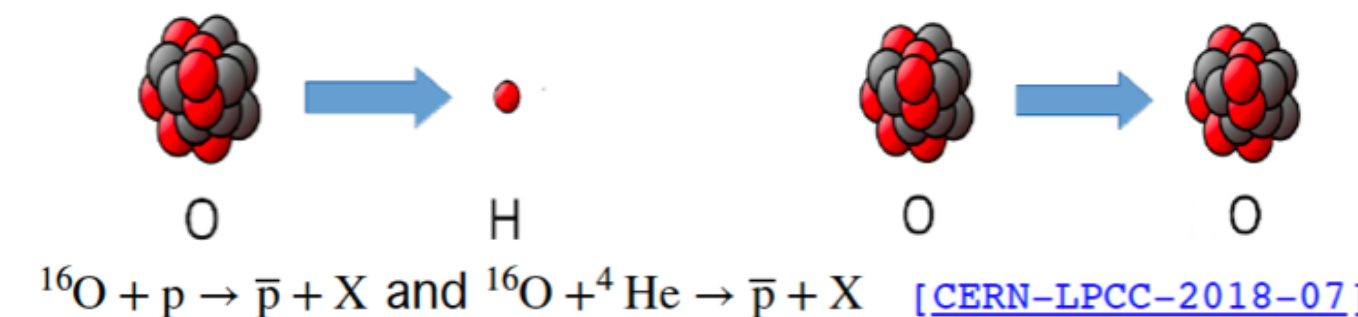
[ArXiv:1906.07119]

• Main uncertainty still due to cross sections!

Intrinsic charm

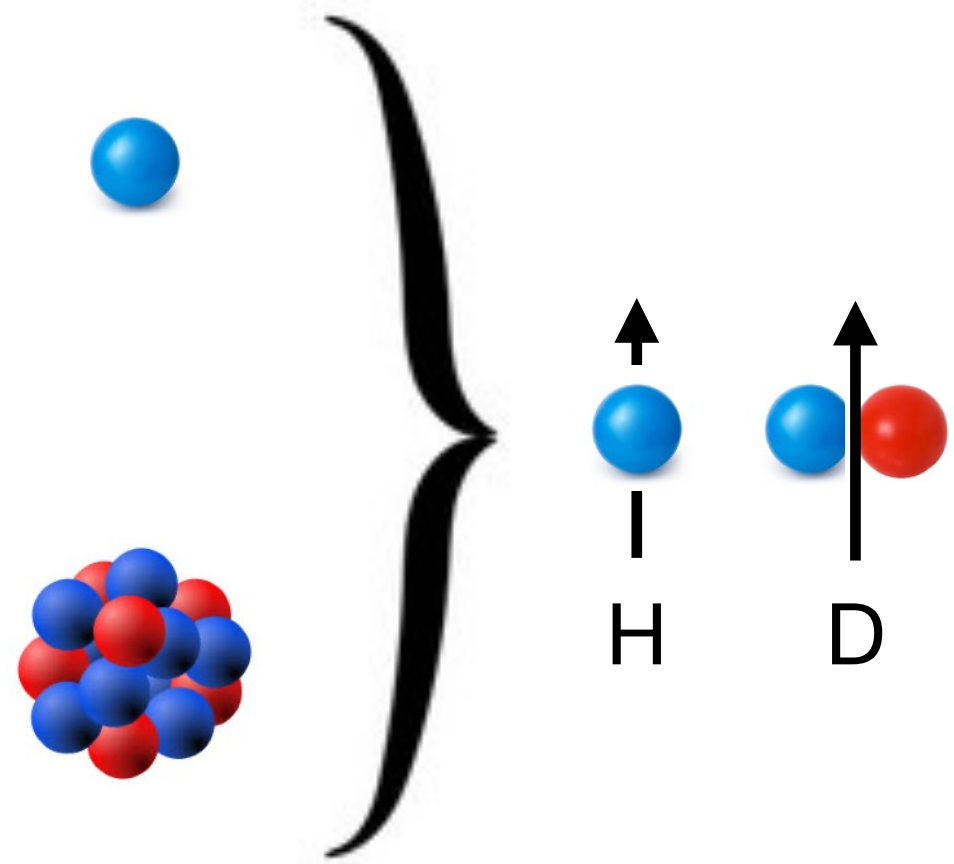


Special Runs



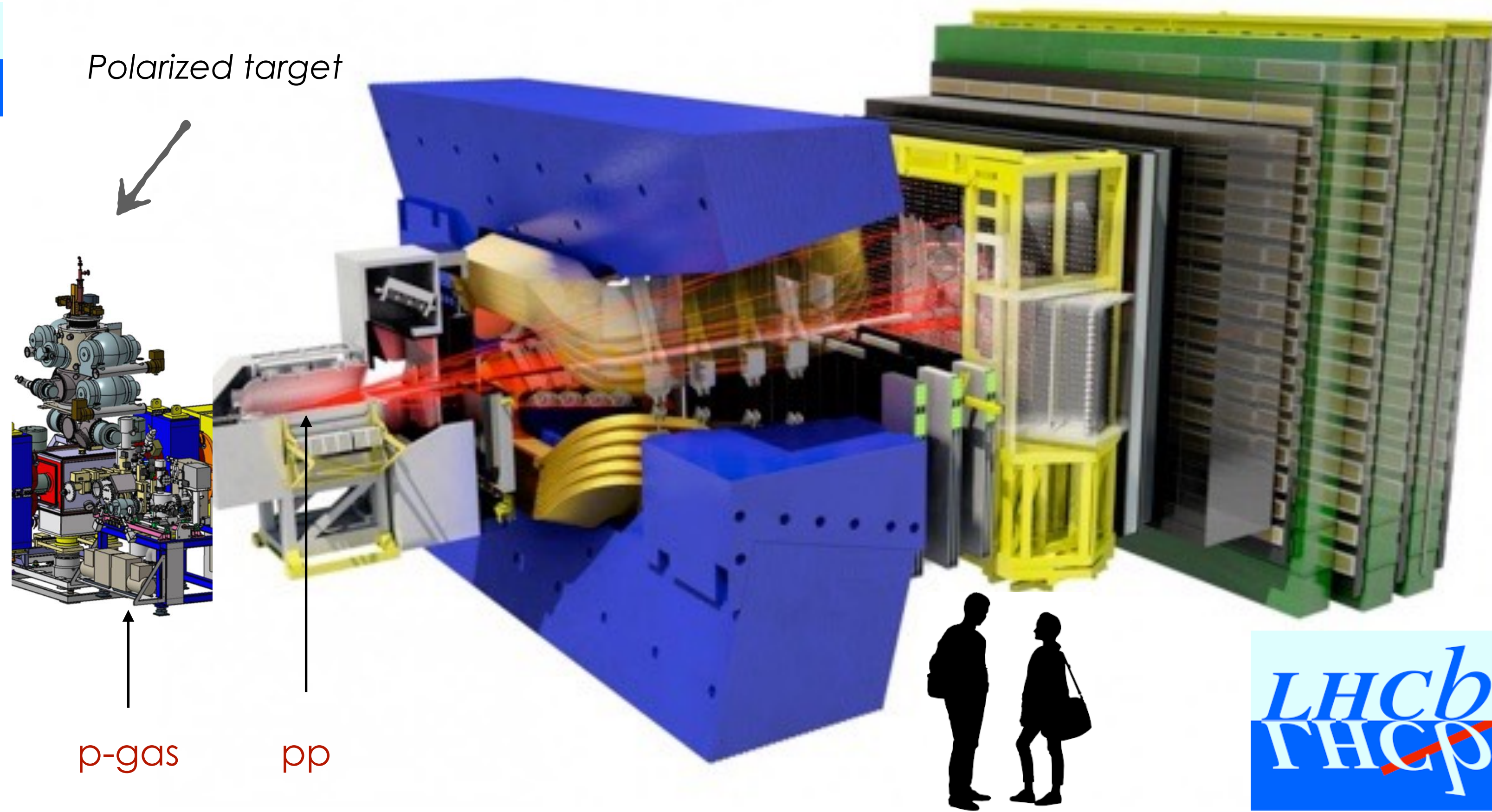
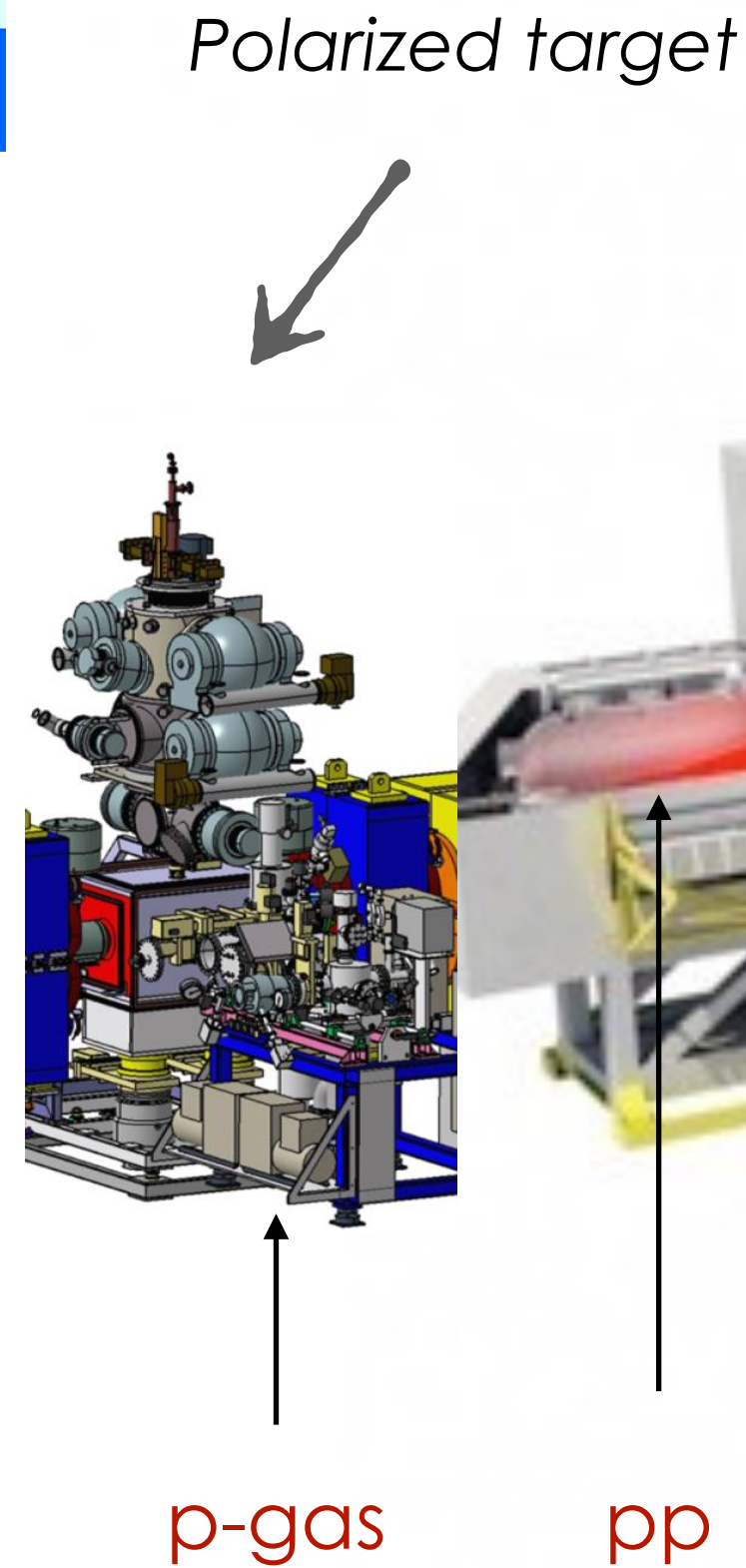
[CERN-LPCC-2018-07]

$L \uparrow \downarrow C$ spin a polarised target at 



Successful technology based on
HERA and COSY experiments

Challenge: develop a new
generation of polarized targets



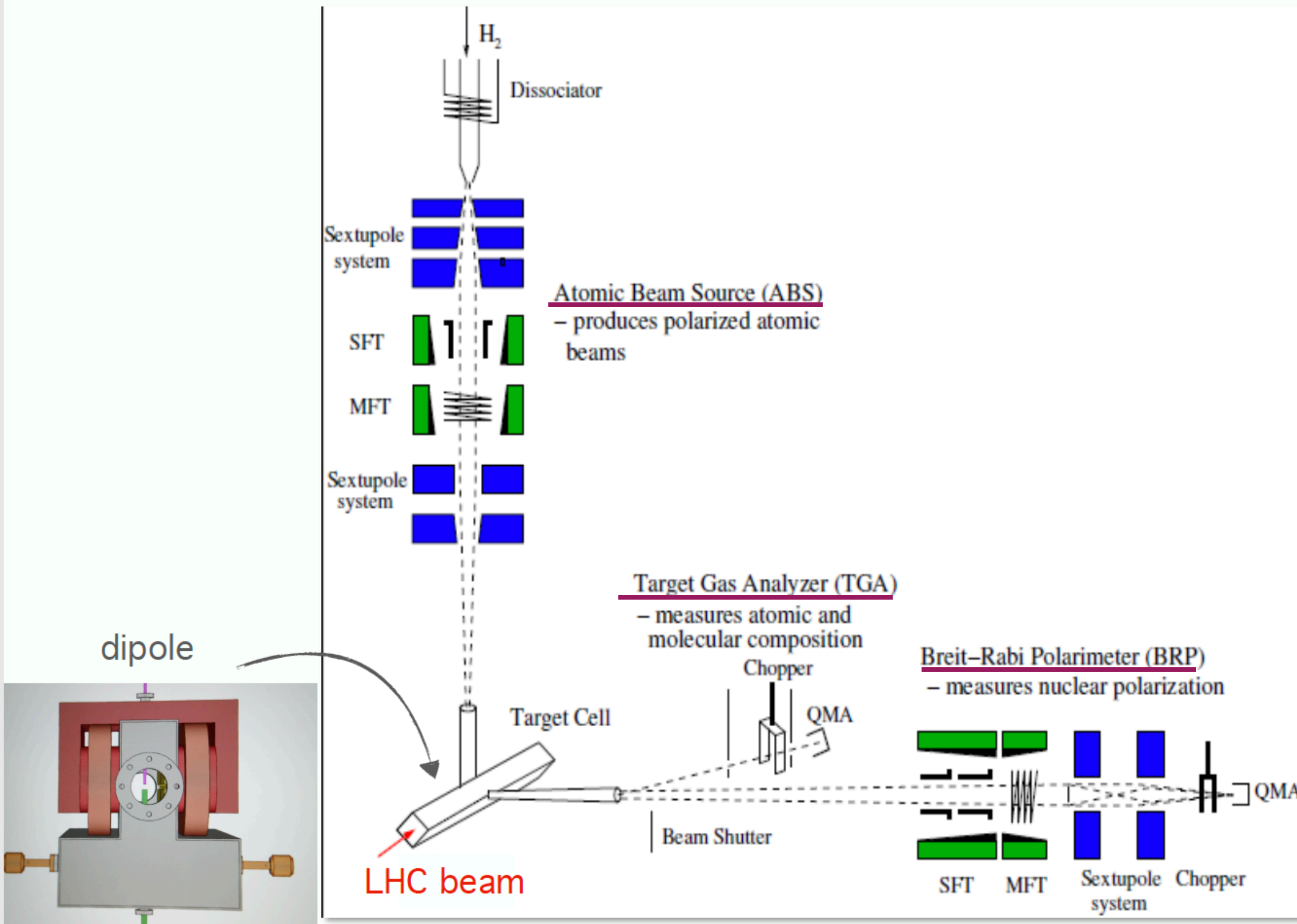
SMOG2 is not only a unique project itself,
but also a great playground for $L \uparrow \downarrow C$ spin

LHCspin experimental setup

Target density (H) = $7 \times 10^{13} \text{ cm}^{-2}$
 LHC beam (Run4) = $6.8 \times 10^{18} \text{ p s}^{-1}$

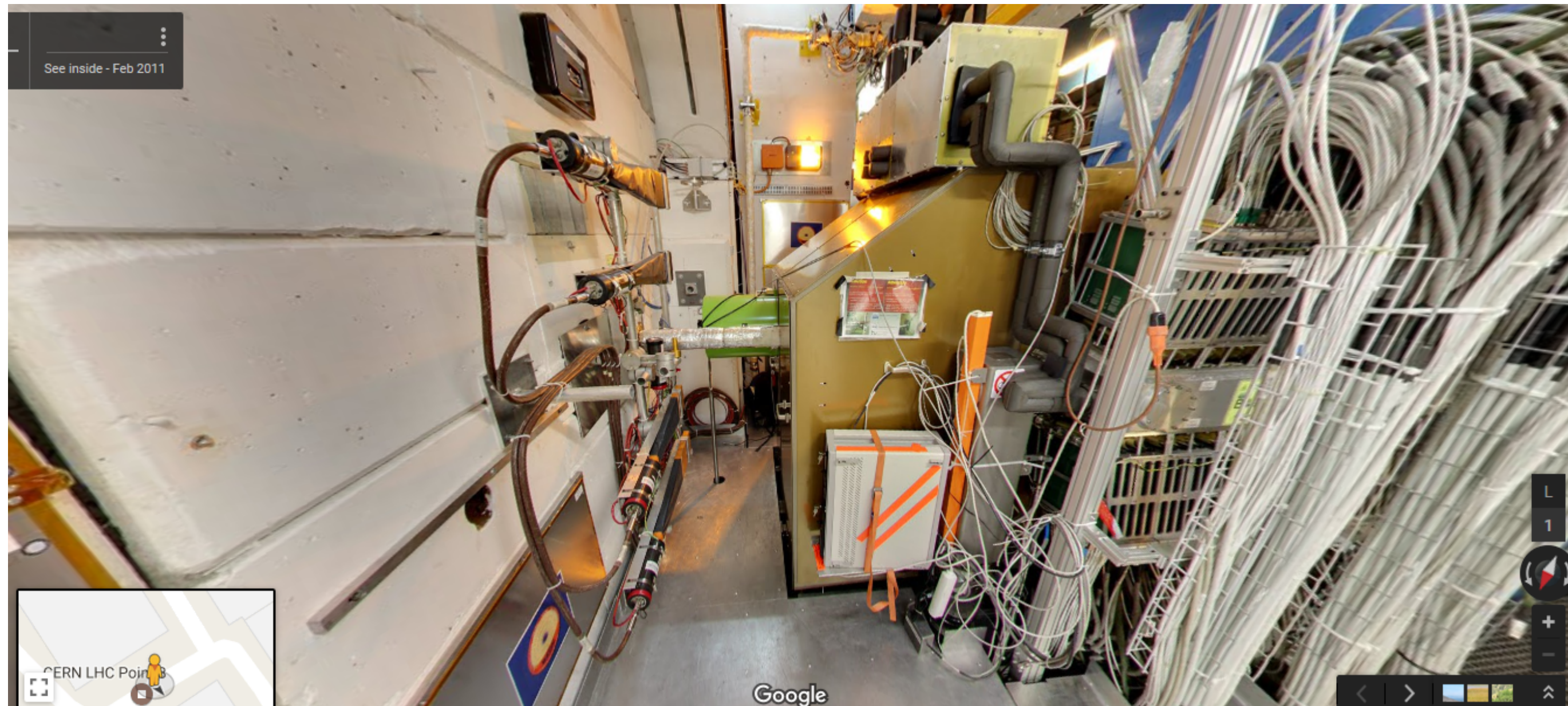
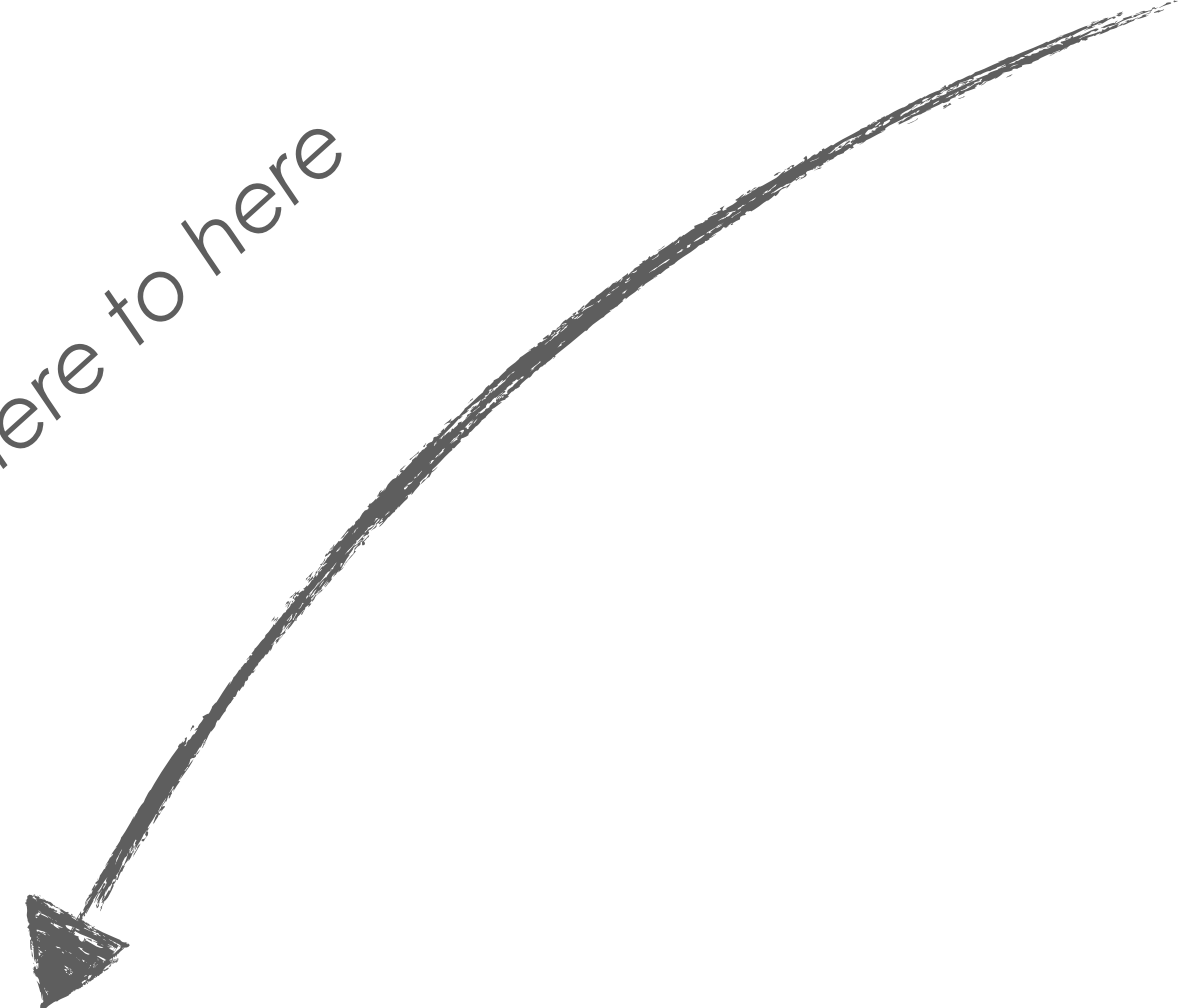
$$L_{pH} = 8 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-2}$$

- Start from the well established HERMES setup @ DESY...
- ... to create the next generation of fixed target polarisation techniques!



HERMES PGT

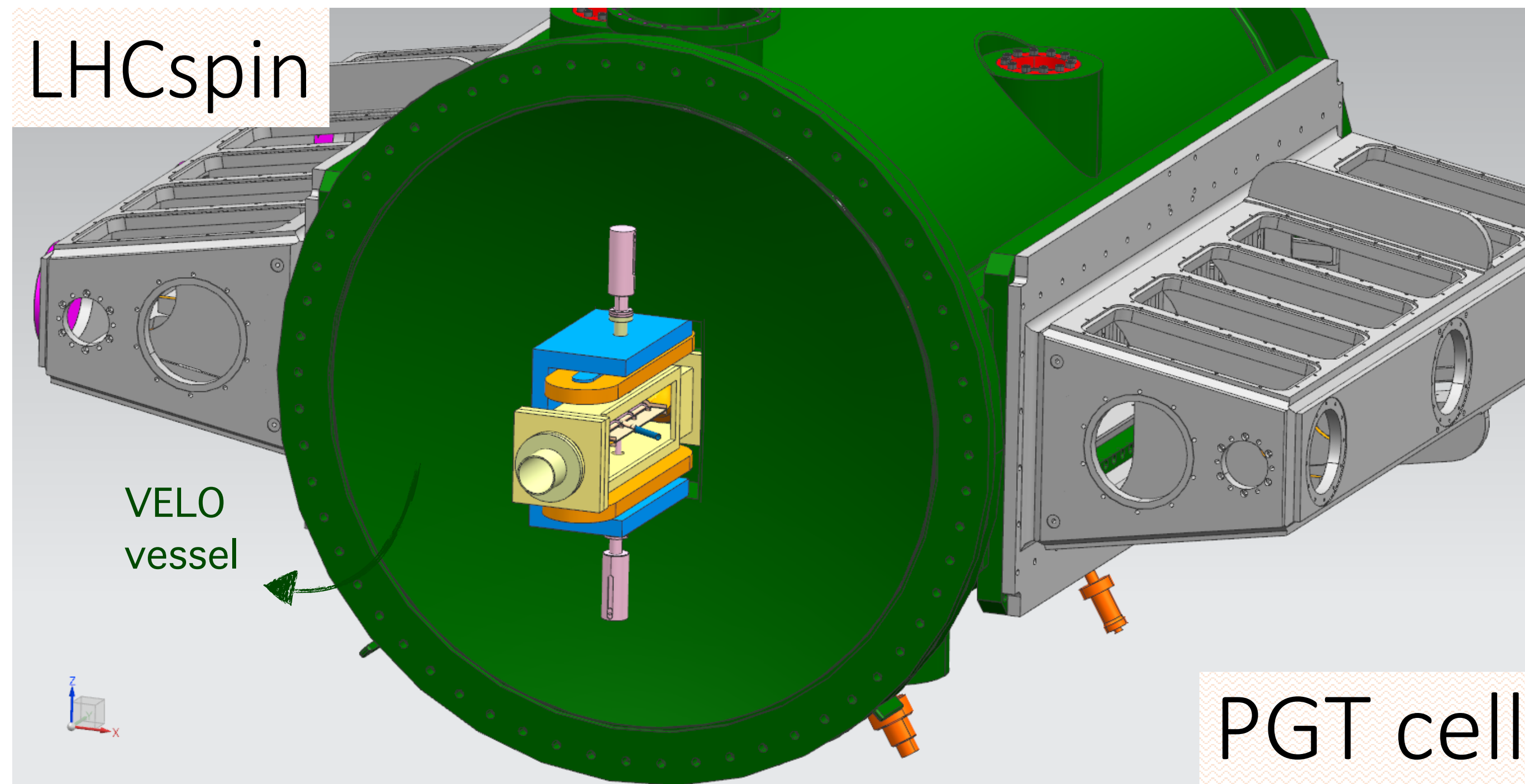
From there to here



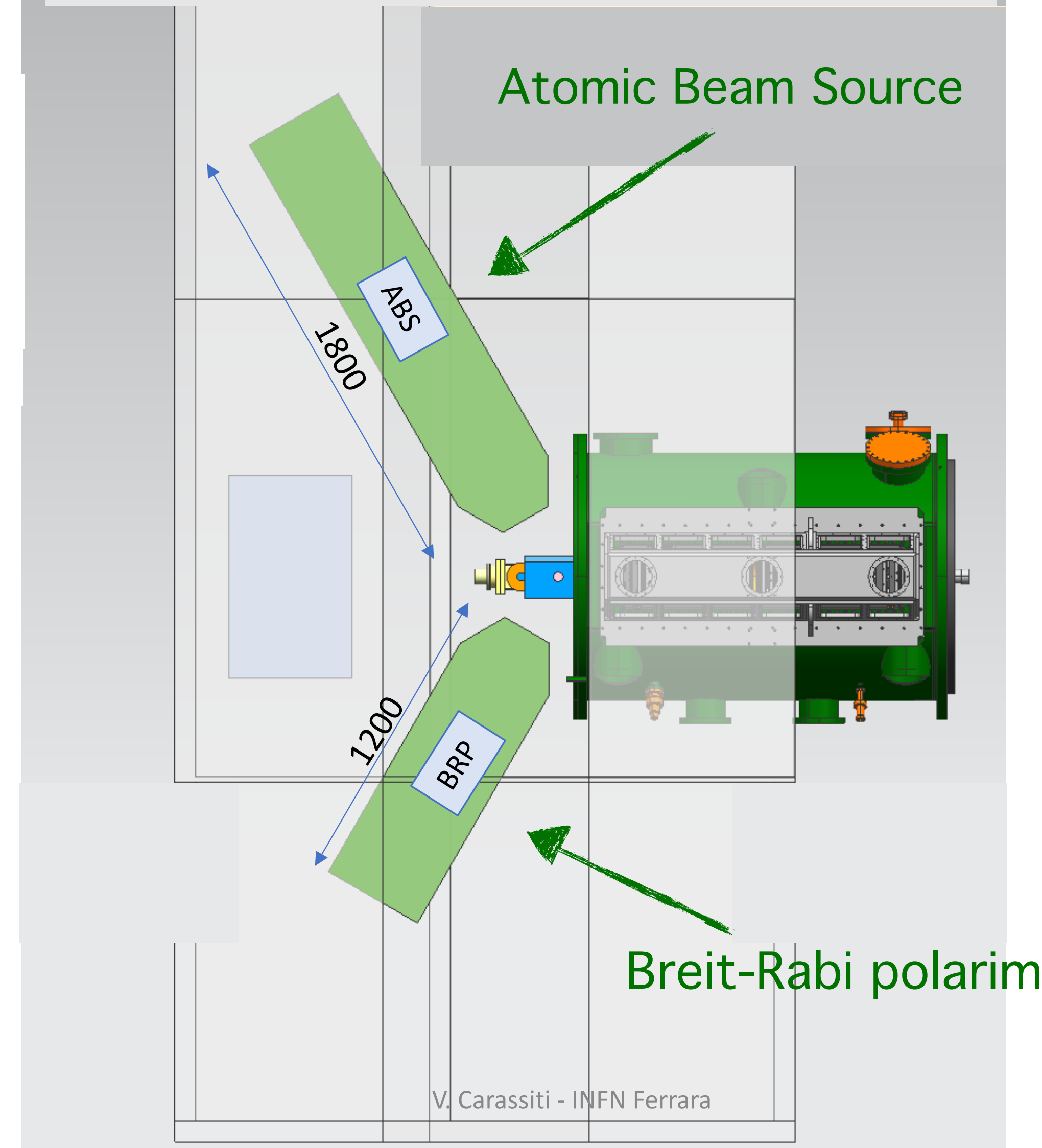
Space available in front of LHCb

PGT implementation into LHCb

- Cylindrical target cell with SMOG2 dimensions: $L = 20$ cm and $D = 1$ cm
- Full LHCb simulations show broader kinematic acceptance & higher efficiency in the same position of the SMOG2 cell



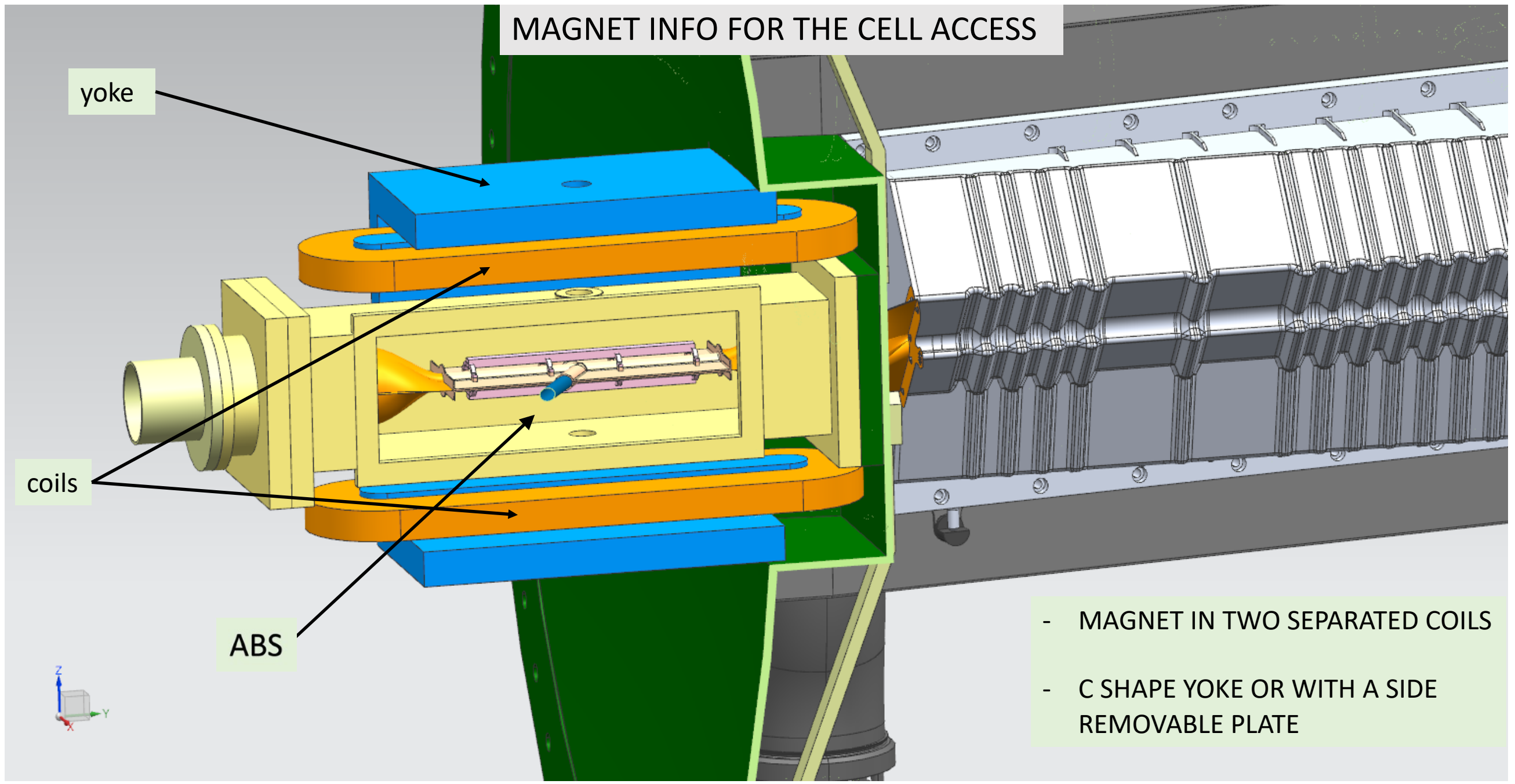
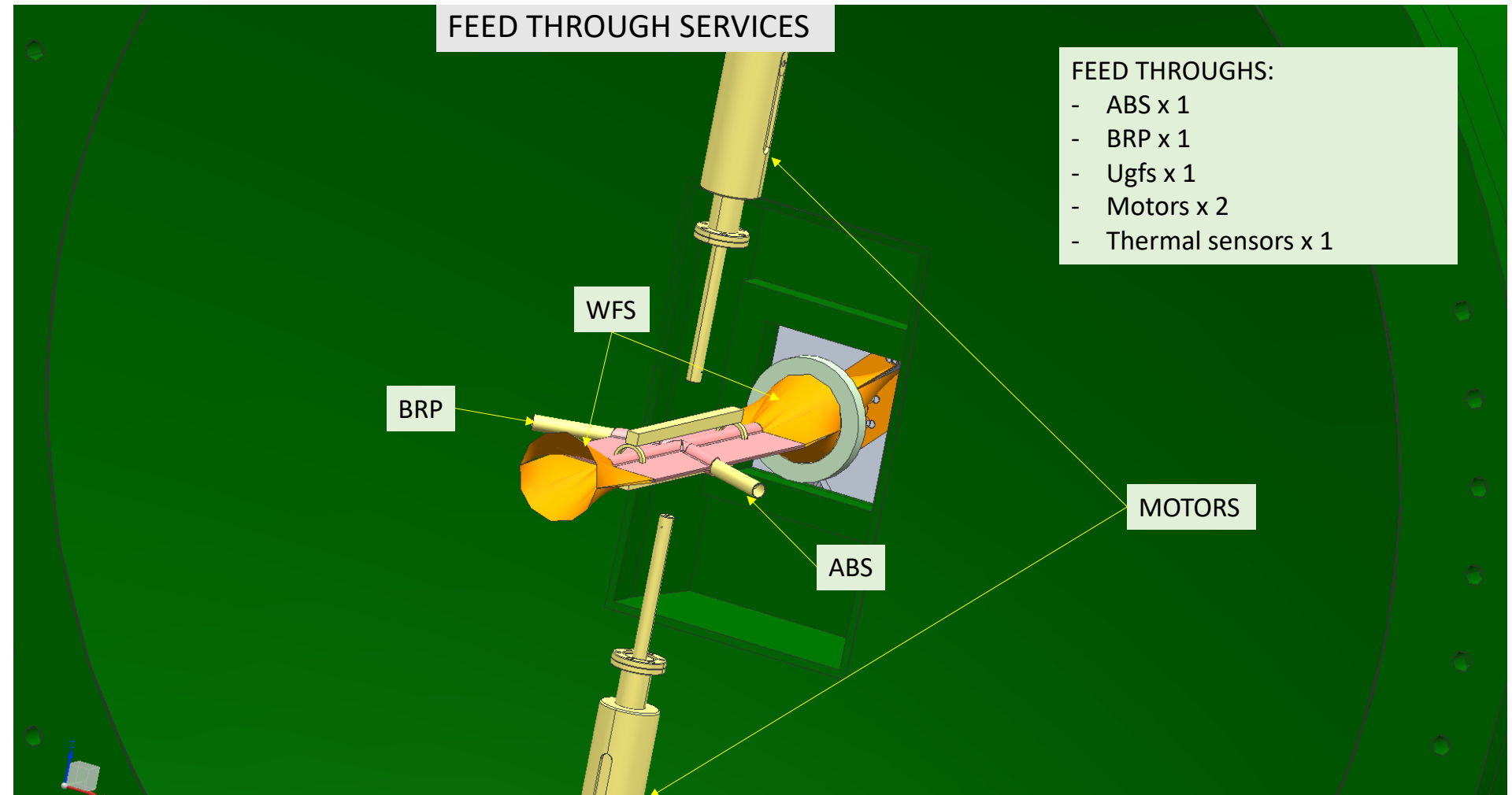
ABS & BRP IN VERTICAL LAYOUT – SIDE VIEW



PGT implementation into LHCb

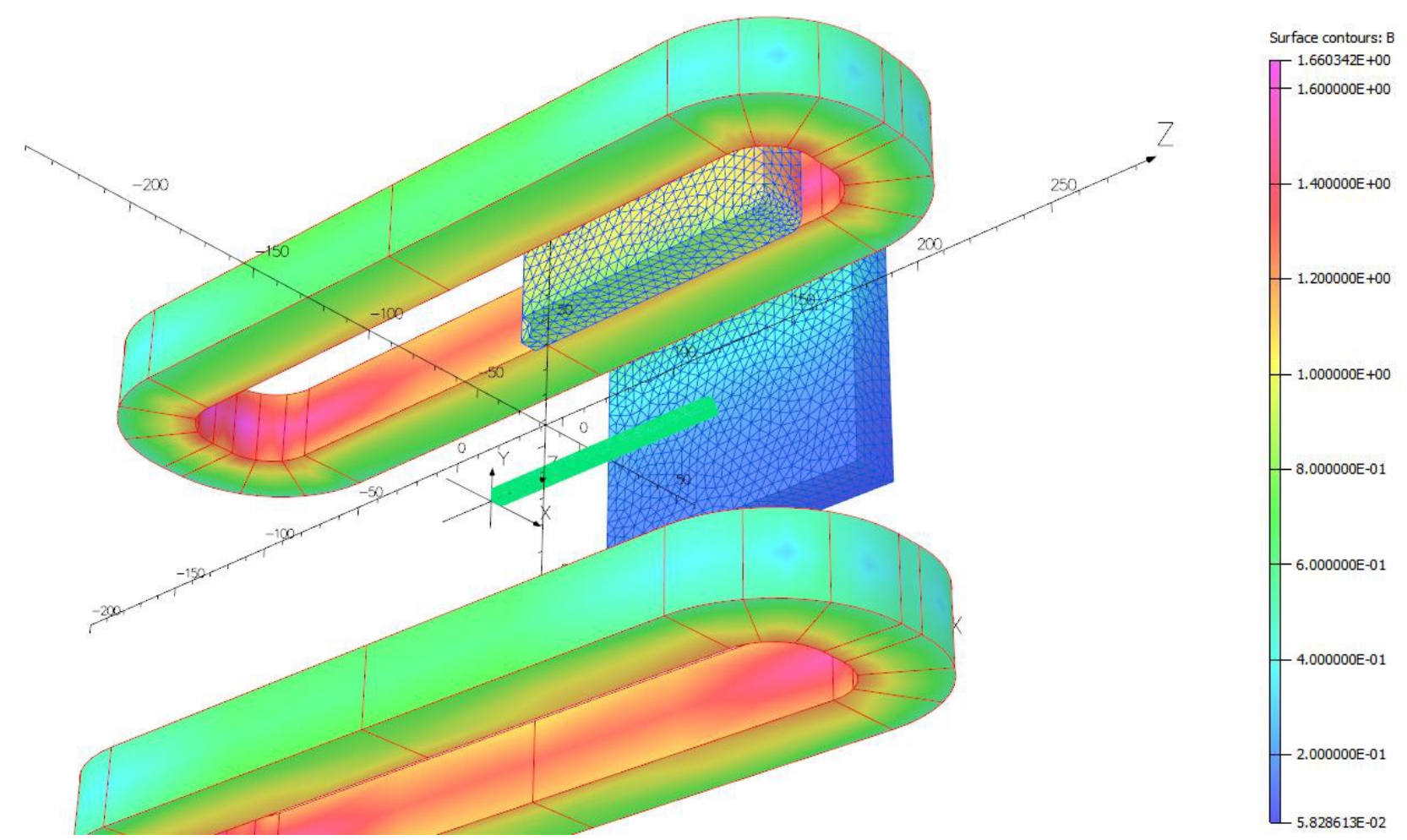
Transverse polarisation

- Inject polarised gas via ABS and unpolarised gas via UGFS



- Compact dipole magnet → static transverse field
- Superconductive coils + iron yoke configuration fits the space constraints
- $B = 300 \text{ mT}$ with polarity inversion, $\Delta B/B \approx 10\%$, suitable to avoid beam-induced depolarisation [[PoS \(SPIN2018\)](#)]

Possibility to switch to a solenoid and provide longitudinal polarisation (e.g. in LHC Run 5)



Role of the storage cell coating

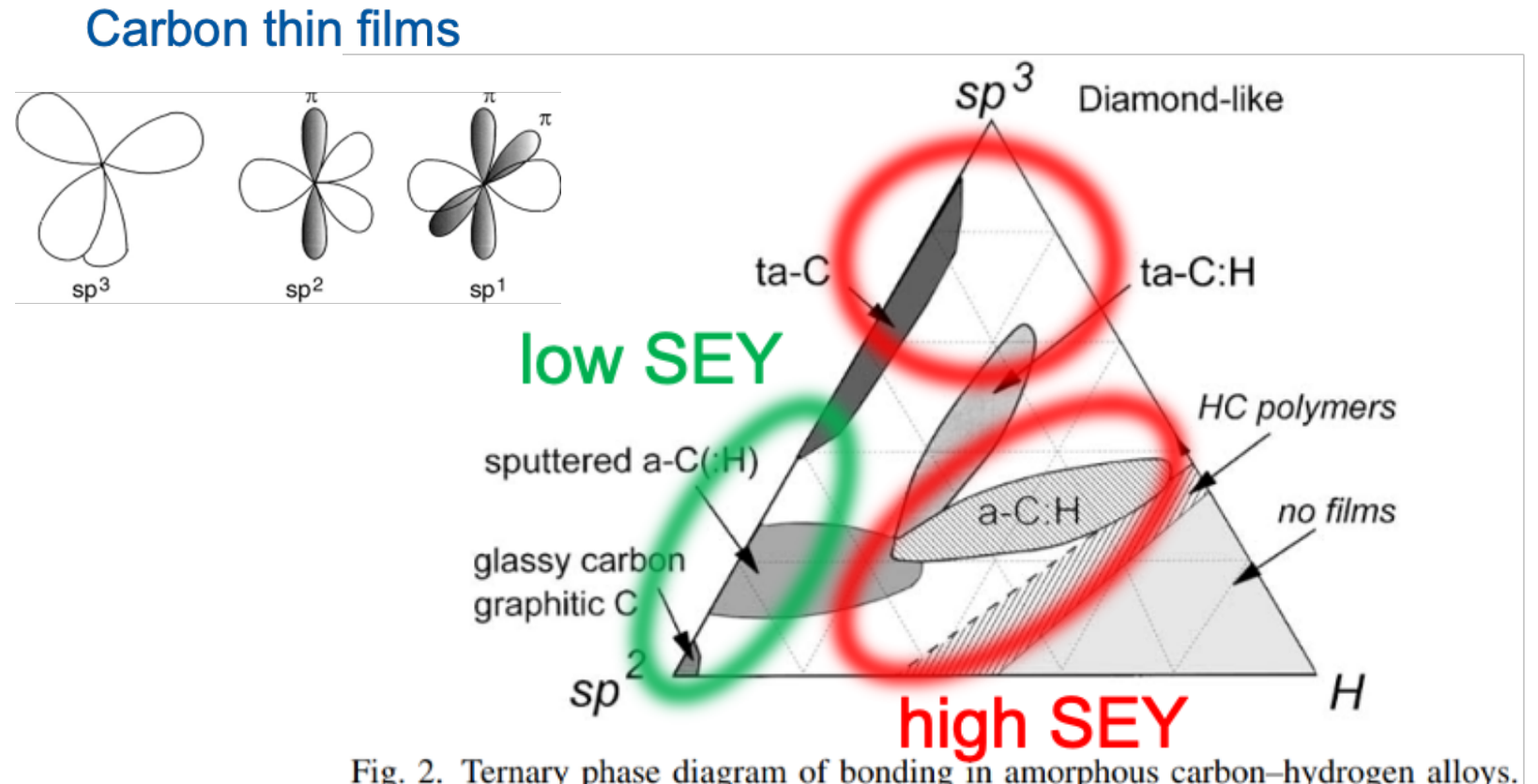
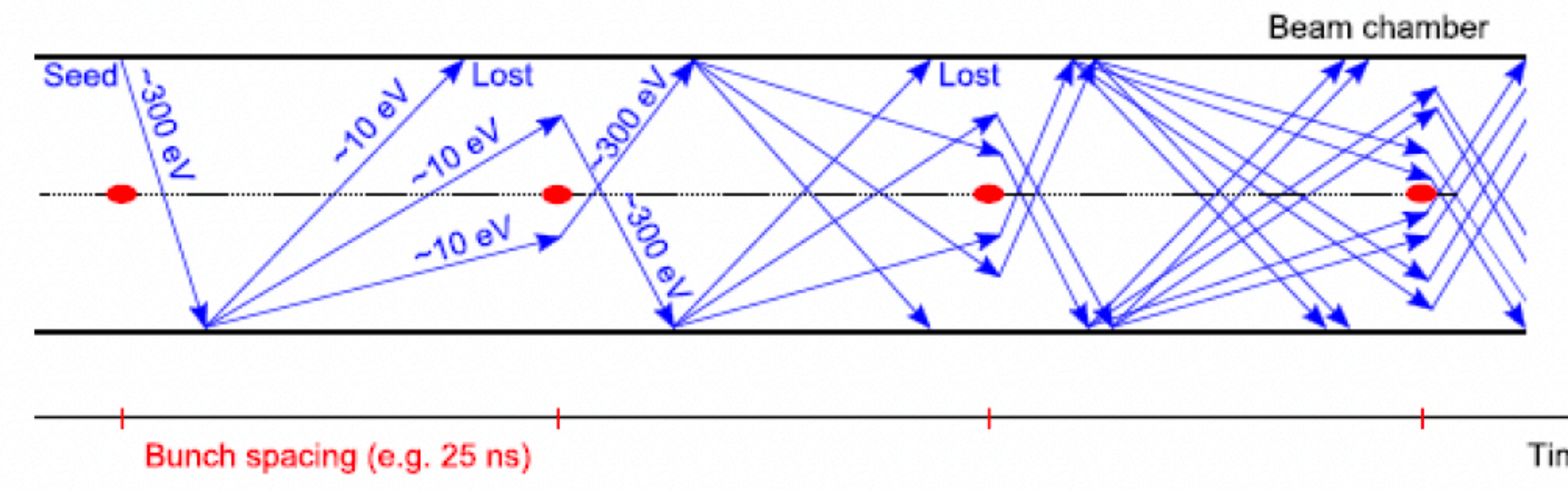


Fig. 2. Ternary phase diagram of bonding in amorphous carbon-hydrogen alloys.
J. Robertson / Materials Science and Engineering R 37 (2002) 129-281



The material of the cell walls must have a low Secondary Electron Yield (e-cloud)

As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?

Role of the storage cell coating

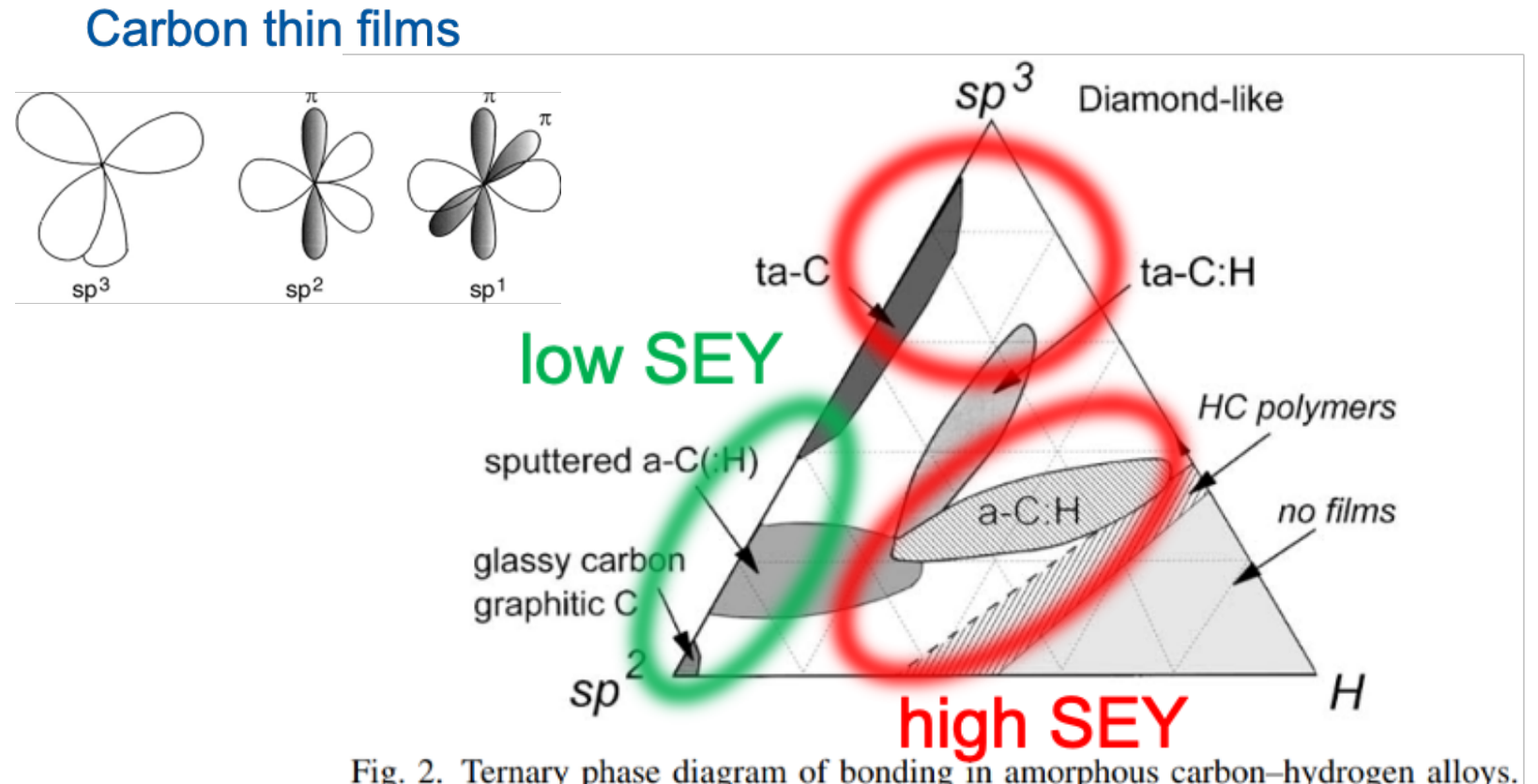
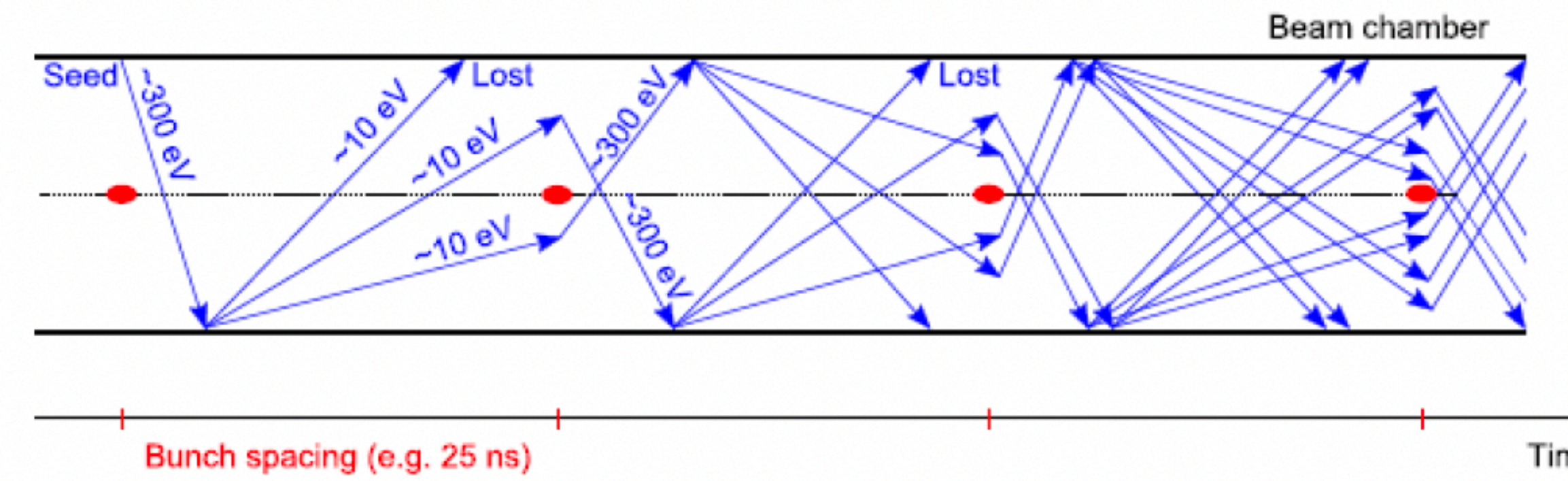


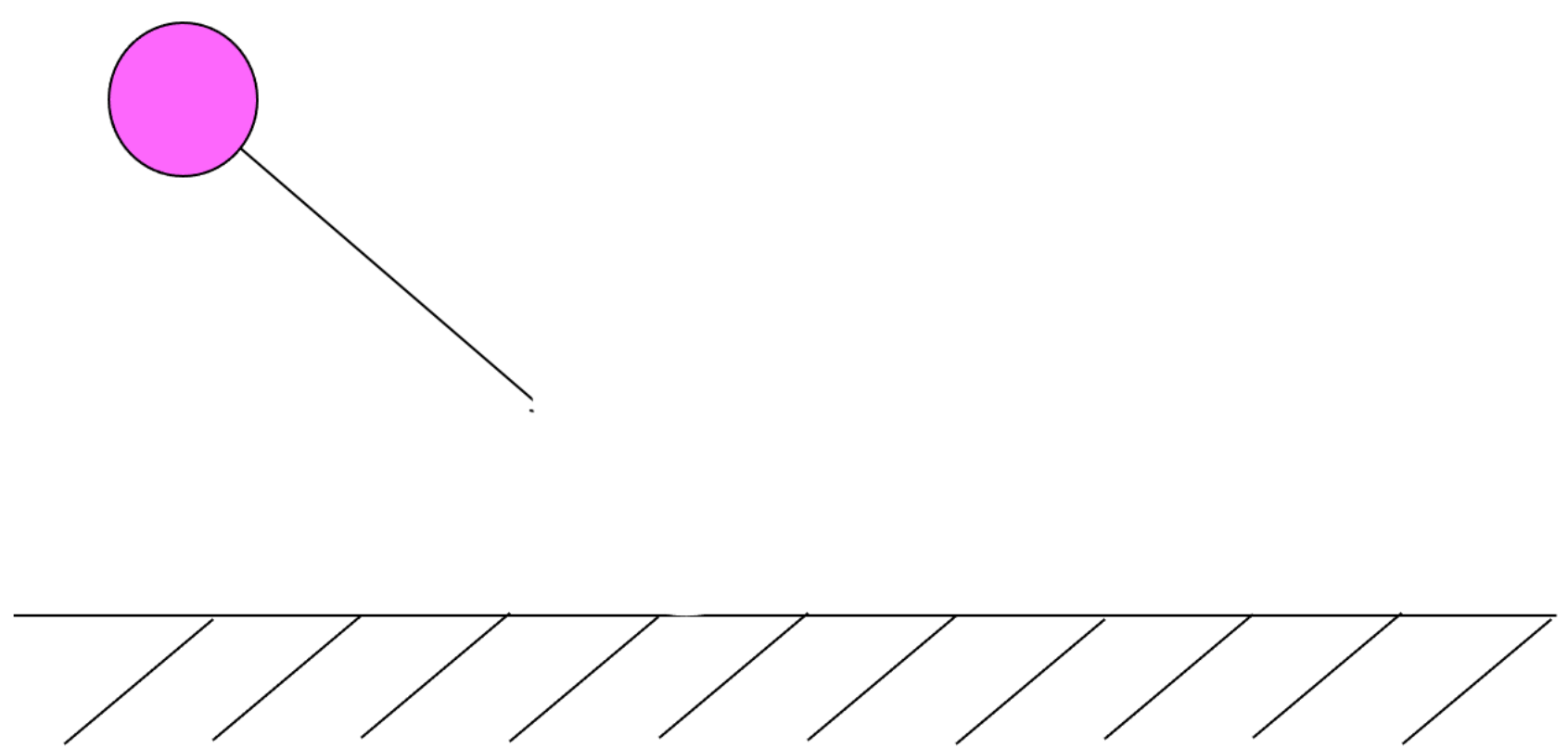
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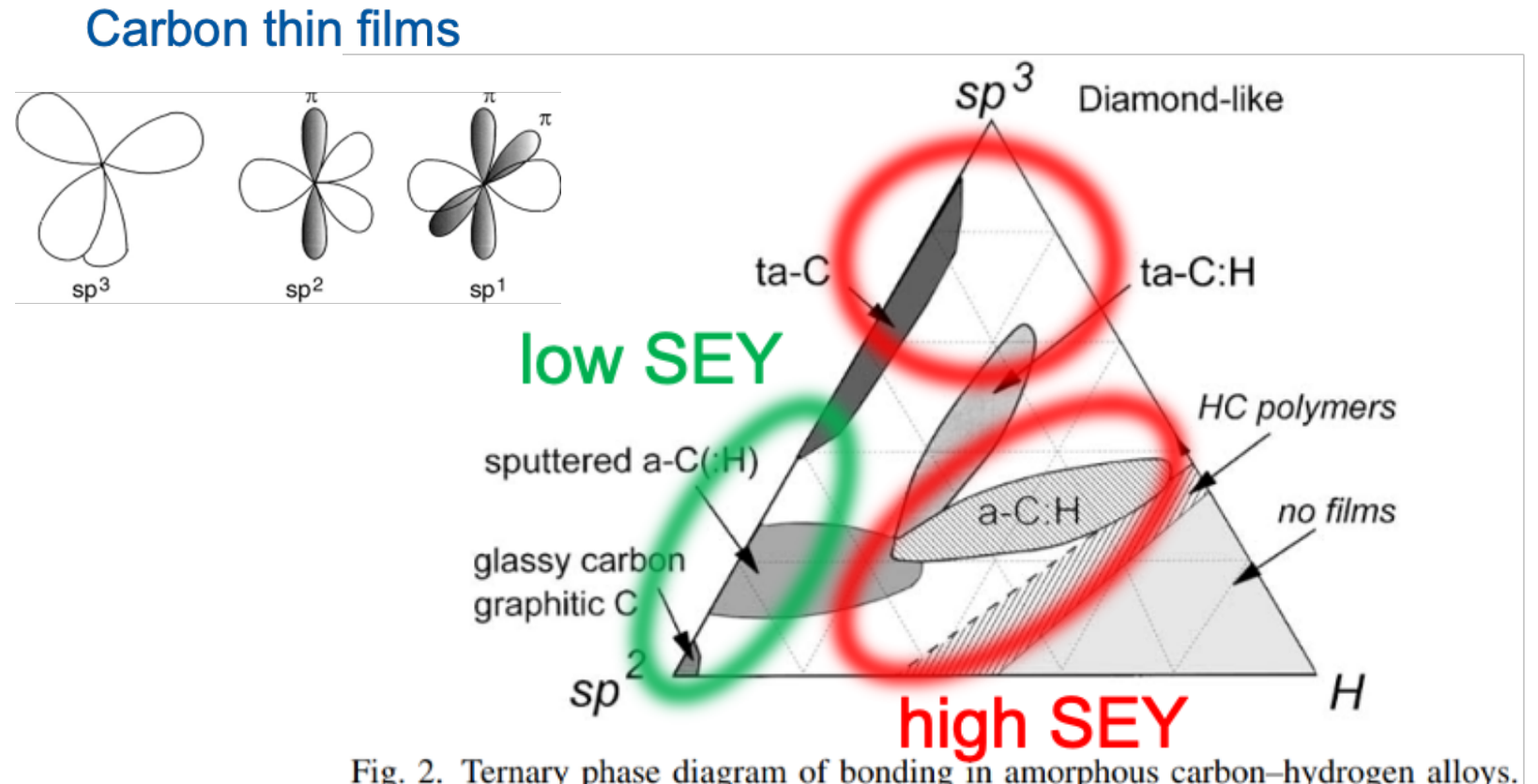
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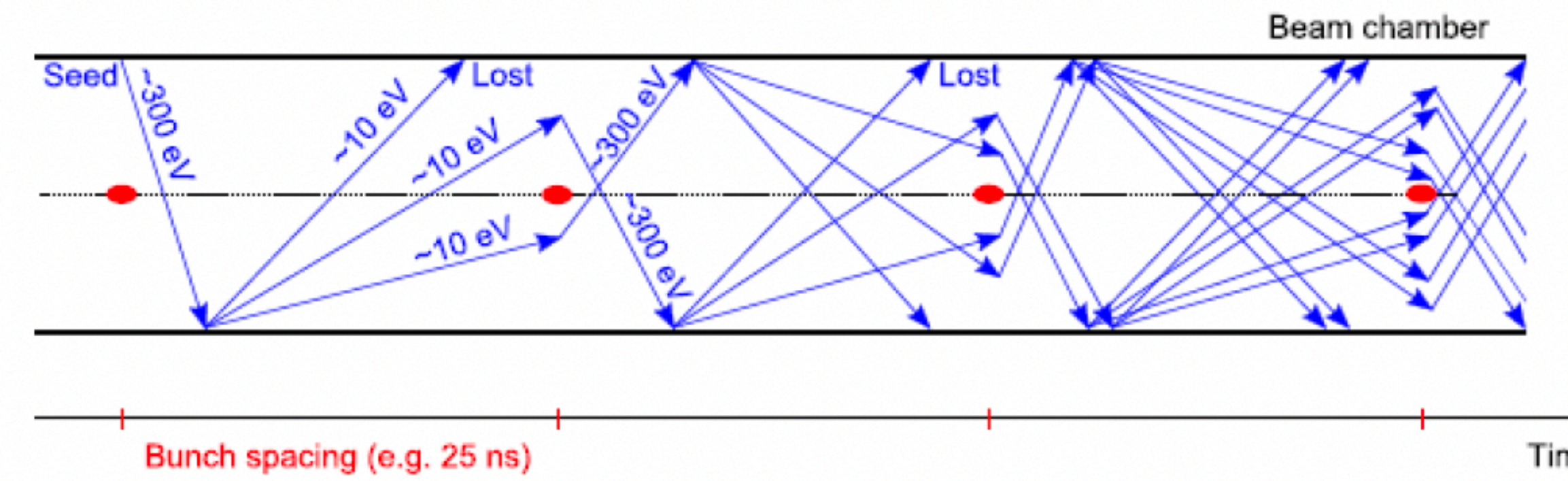
Eley-Rideal Mechanism



Role of the storage cell coating



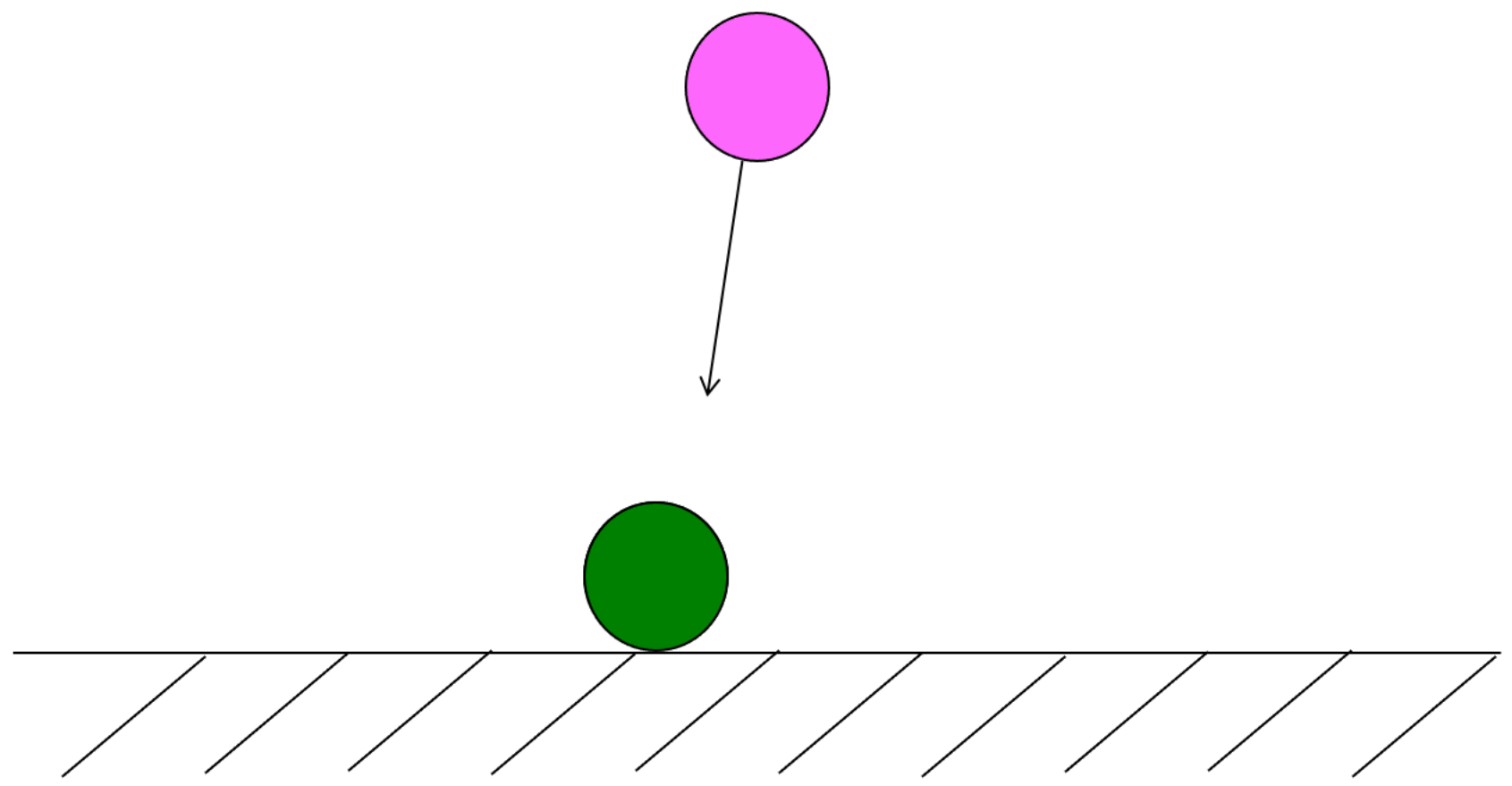
J. Robertson / Materials Science and Engineering R 37 (2002) 129-281



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Eley-Rideal Mechanism



Role of the storage cell coating

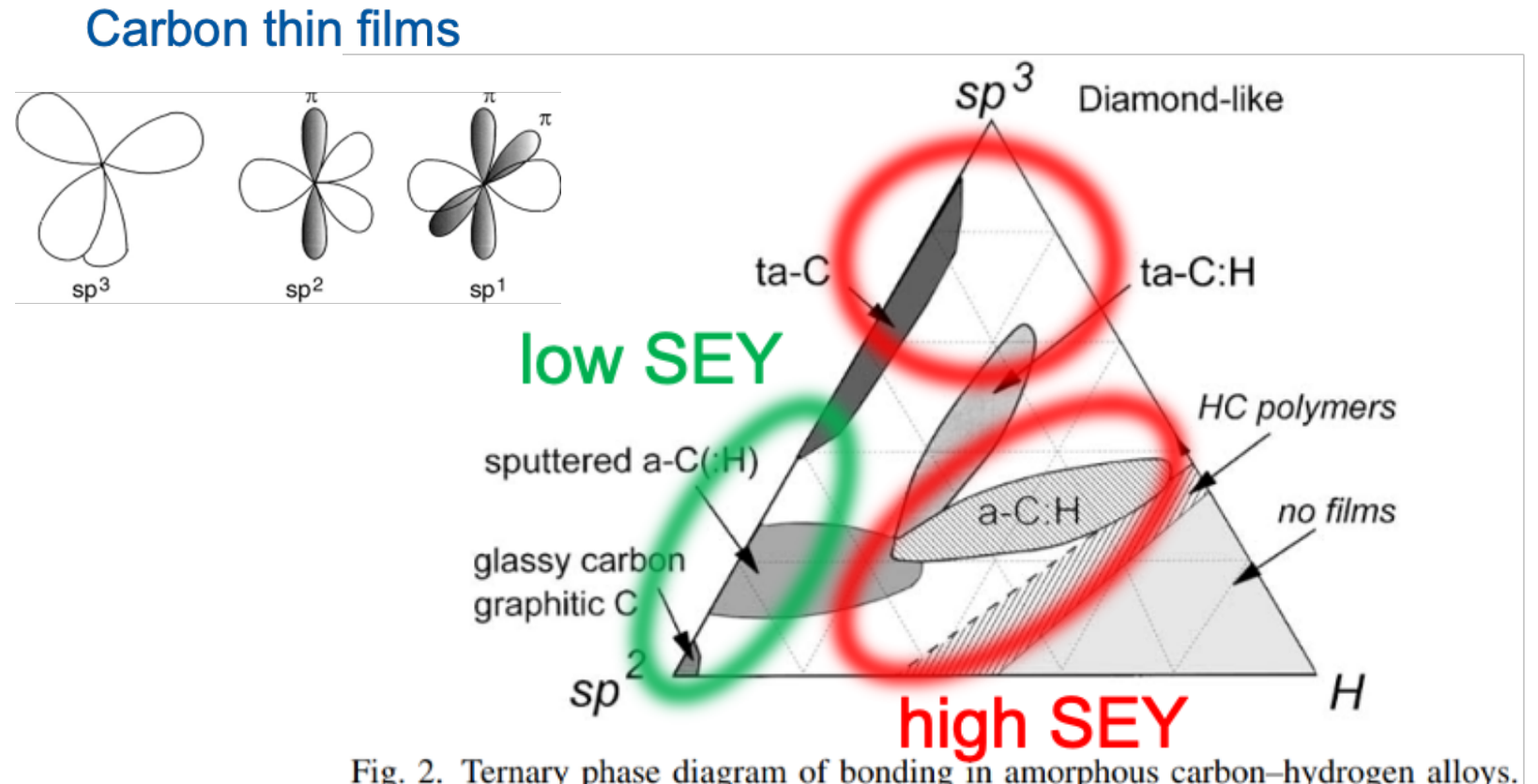
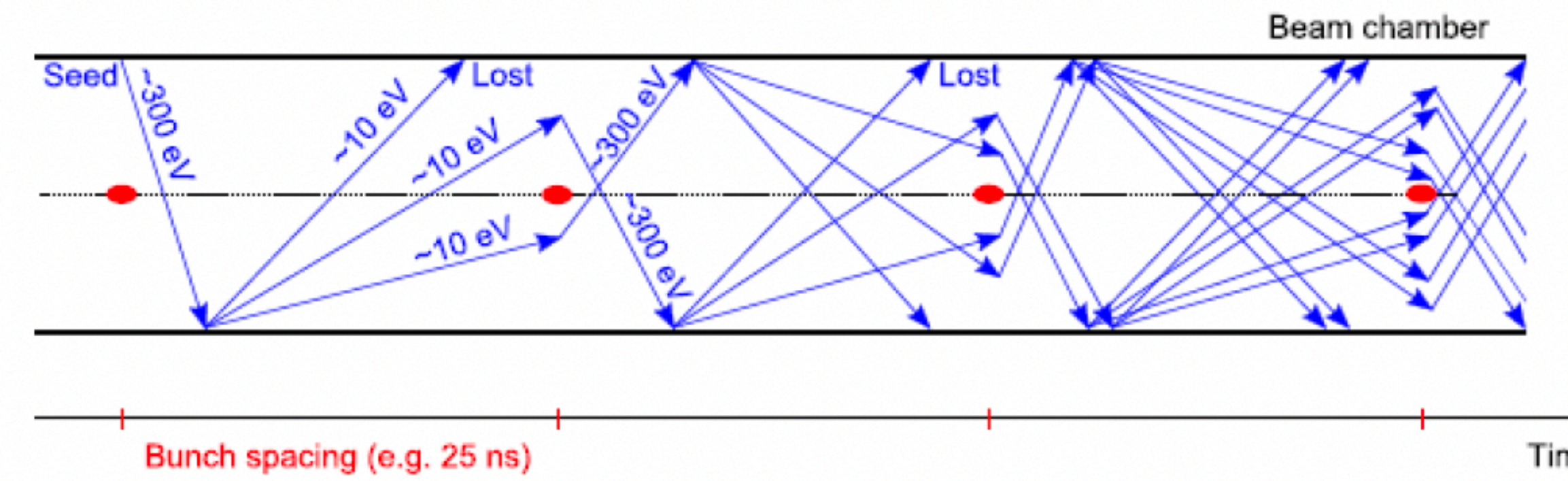


Fig. 2. Ternary phase diagram of bonding in amorphous carbon-hydrogen alloys.
J. Robertson / Materials Science and Engineering R 37 (2002) 129-281

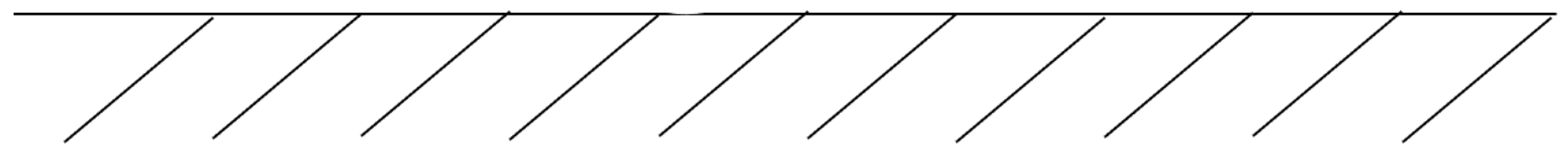
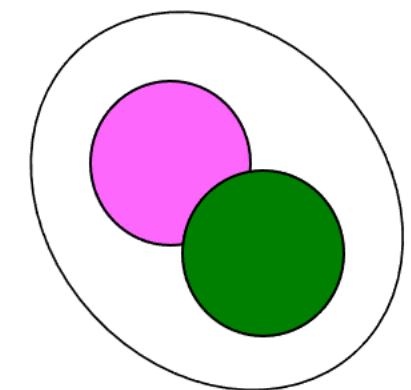


The material of the cell walls must have a low Secondary Electron Yield (e-cloud)

As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?

Eley-Rideal Mechanism

$$P_m = 0.5 P_a$$



Role of the storage cell coating

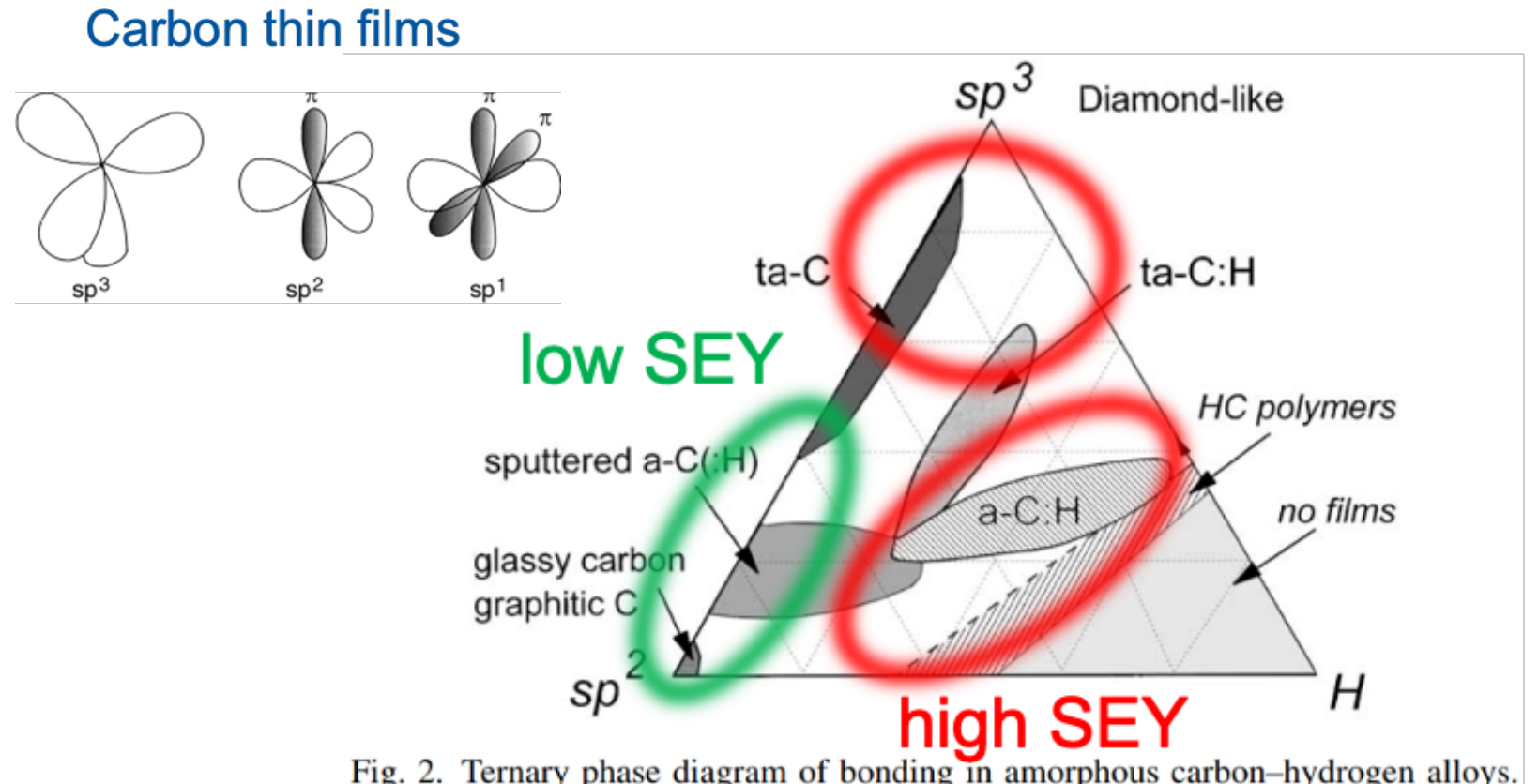
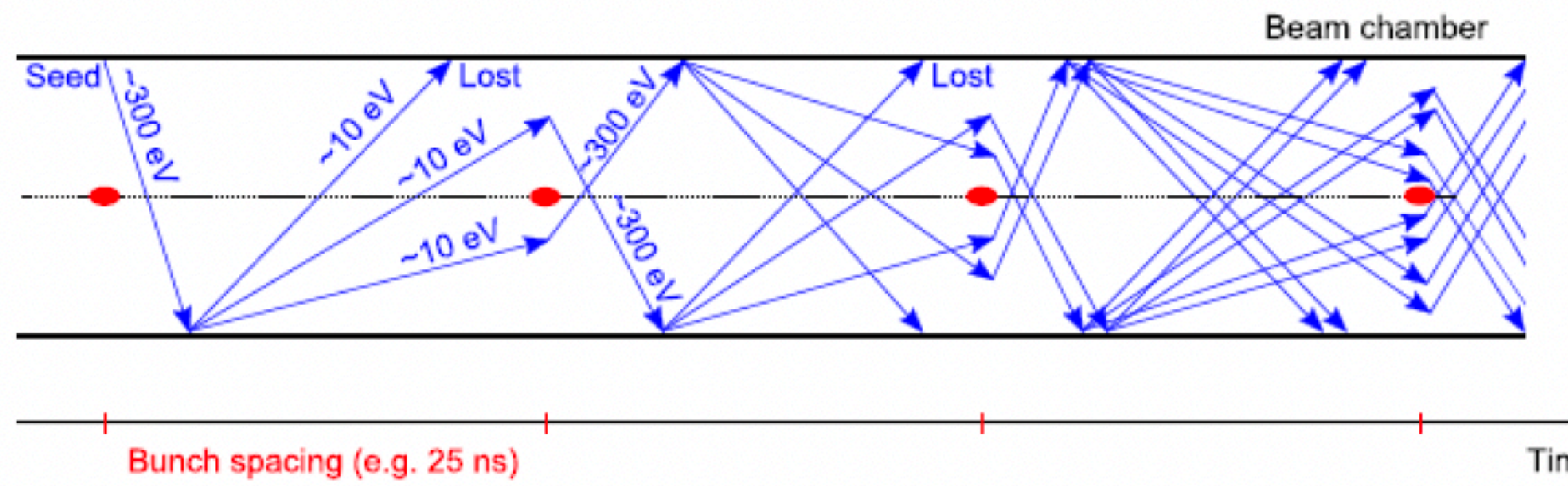
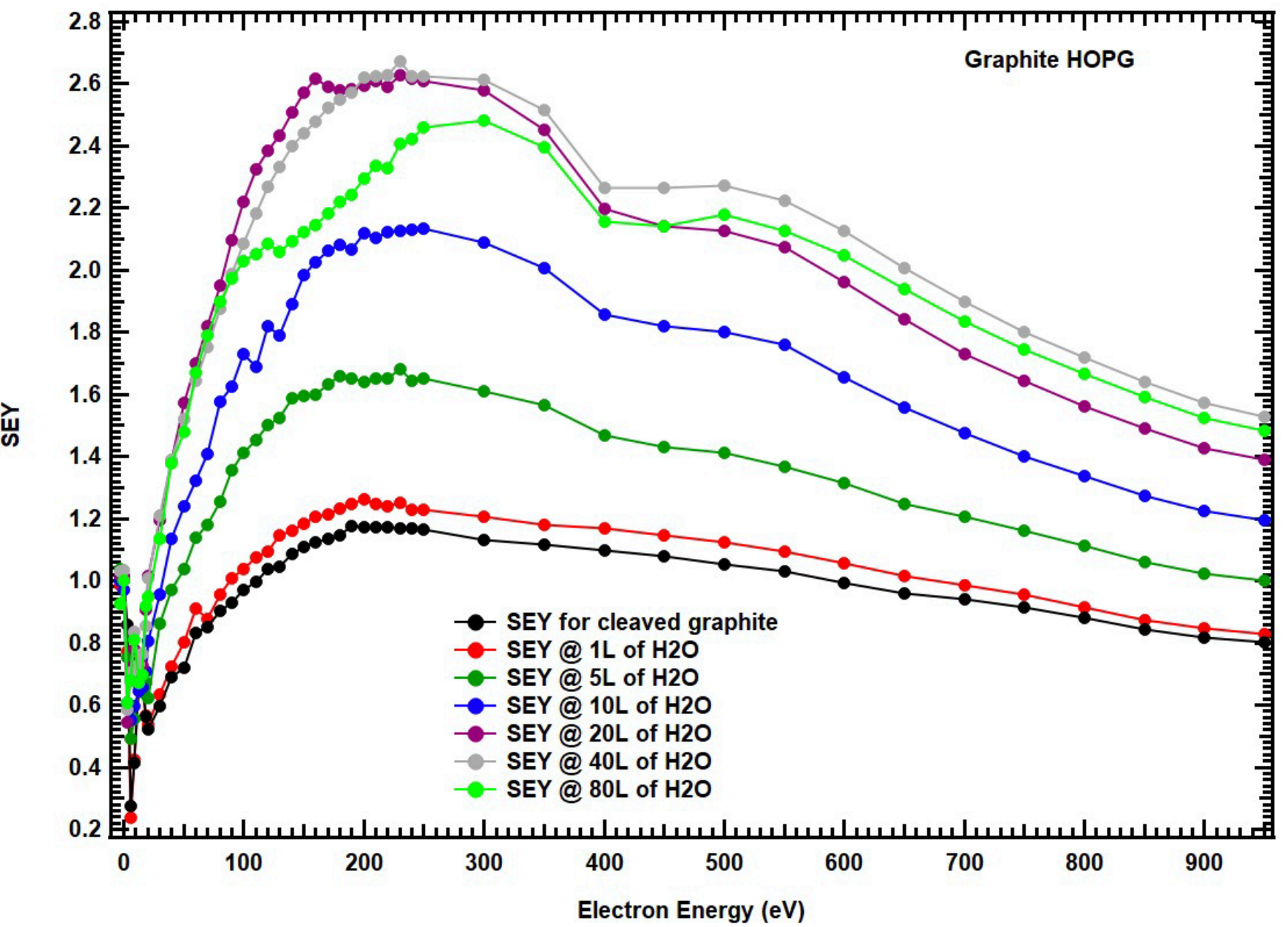


Fig. 2. Ternary phase diagram of bonding in amorphous carbon-hydrogen alloys.
J. Robertson / Materials Science and Engineering R 37 (2002) 129-281

The material of the cell walls must have a low Secondary Electron Yield (e-cloud)
 As for SMOG2, Amorphous Carbon is ok. Has it a low H recombination as well?

Studies ongoing in order to understand if carbon films with low secondary Electron Yield cope with the required “recombination” rate of polarized H atoms injected in the storage cell

... or follow the HERMES experience to have an ice coating (low SEY, low H recombination)

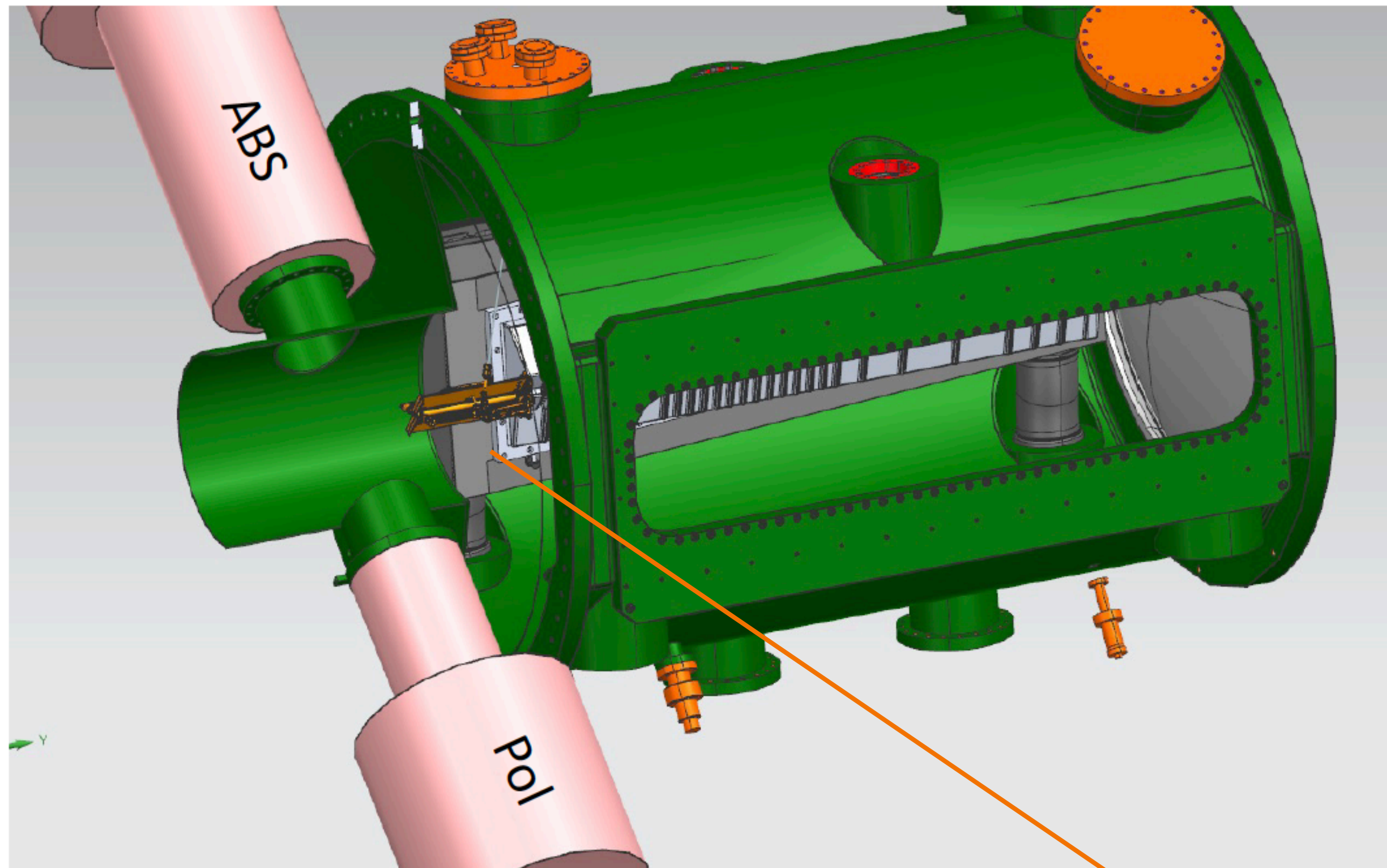


Backup solution is also being investigated: a jet target that provides lower density ($\sim 10^{12}$ atoms/cm²) but higher polarisation degree (up to 90%) and lower systematics

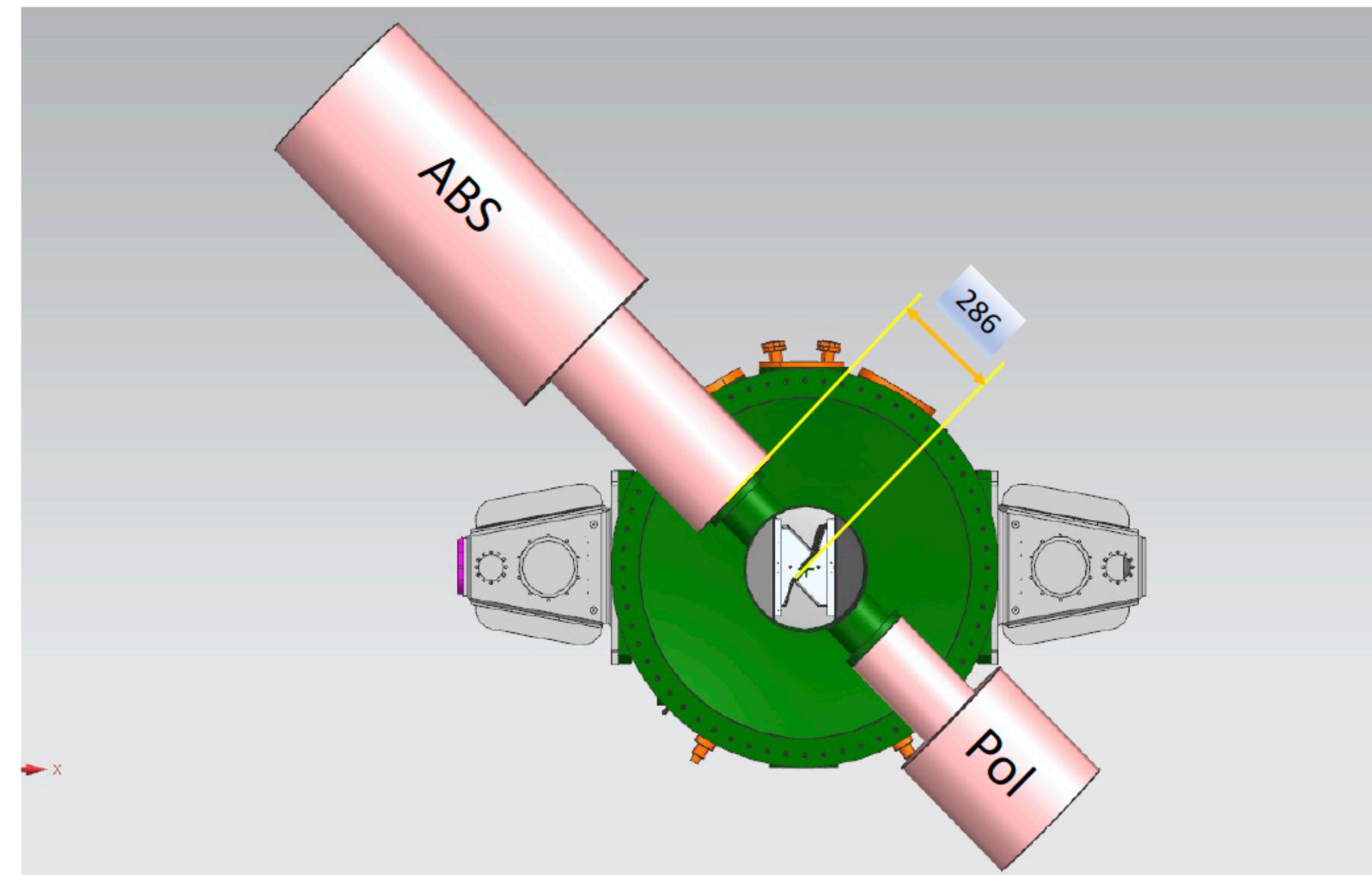
The jet target option

Alternative solution with **jet target** also under evaluation:

- lower density ($\sim 10^{12}$ atoms/cm²)
- higher polarization (up to 90%)
- lower systematics in P measurement (virtually close to 0)



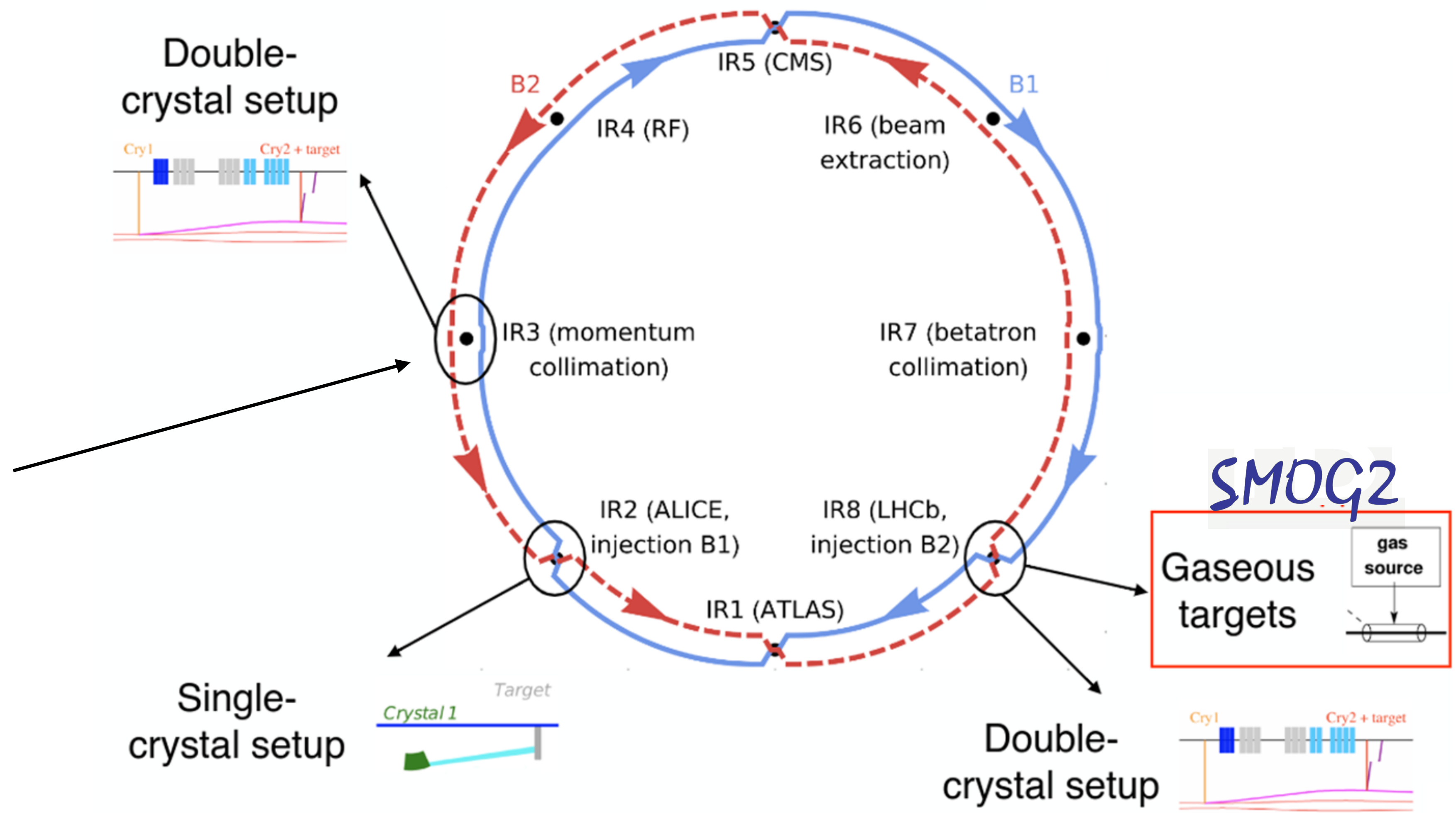
SMOG2



The LHC Interaction Region 3

- IR3 is a great opportunity to perform R&D on beam:
 - mutual polarised target - beam interaction (BID, impedance, aperture, ...)
 - hardware systems for LHCspin
 - ...

L C
 ↑ ↓ spin
 R&D



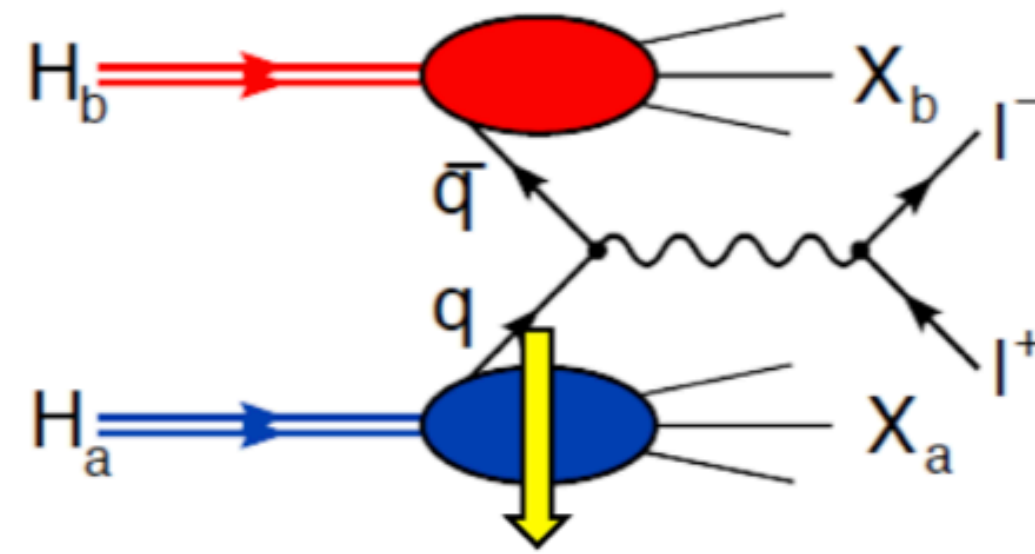
The physics goals of ... just a quick overview

- Multi-dimensional nucleon structure in a poorly explored kinematic domain
- Measure experimental observables sensitive to both **quarks and gluons TMDs**
- **Make use of new probes (charmed and beauty mesons)**
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs
- Measure exclusive processes to access GPDs

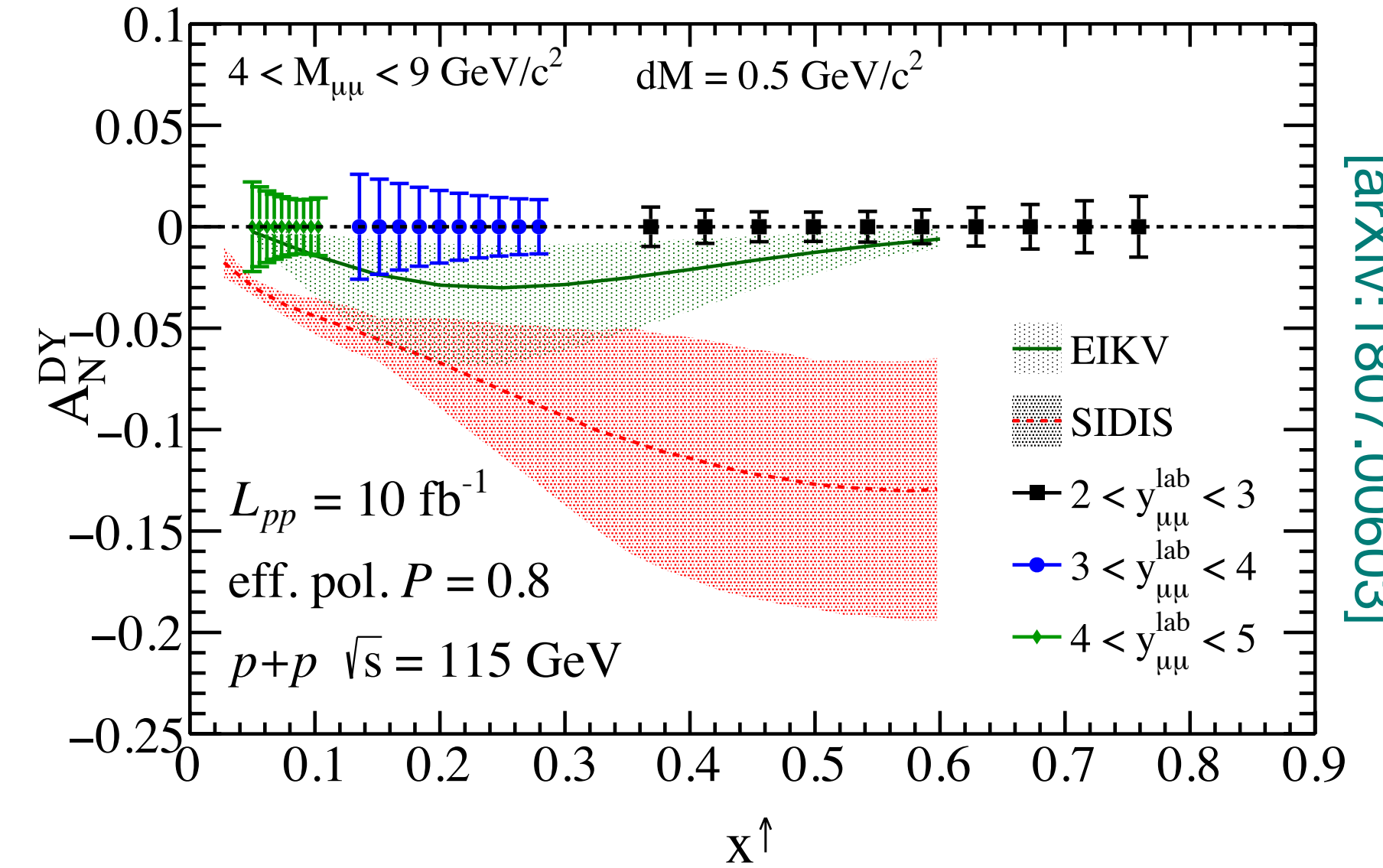
Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Transv. polarized Drell-Yan



Golden Channel



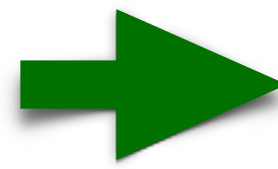
[arXiv:1807.00603]

- Sensitive to quark TMDs through TSSAs

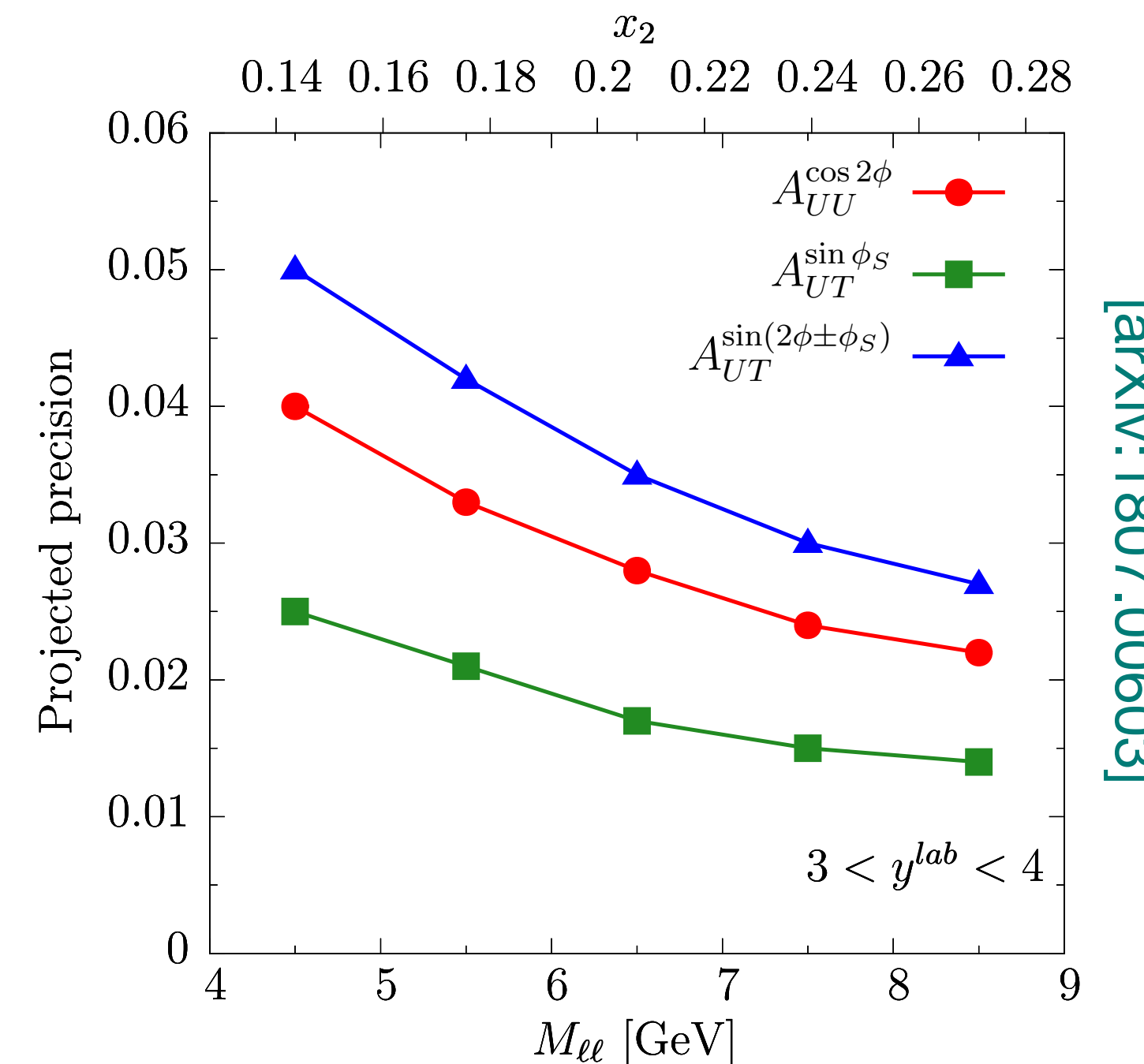
$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^\uparrow - \sigma_{DY}^\downarrow}{\sigma_{DY}^\uparrow + \sigma_{DY}^\downarrow} \Rightarrow A_{UT}^{\sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{\sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$

(ϕ : azimuthal orientation of lepton pair in dilepton CM)

LHCb has excellent μ -ID & reconstruction for $\mu^+\mu^-$



dominant: $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
 suppressed: $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$



[arXiv:1807.00603]

- Extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS $f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{SIDIS}$
- Test flavour sensitivity using both H and D targets

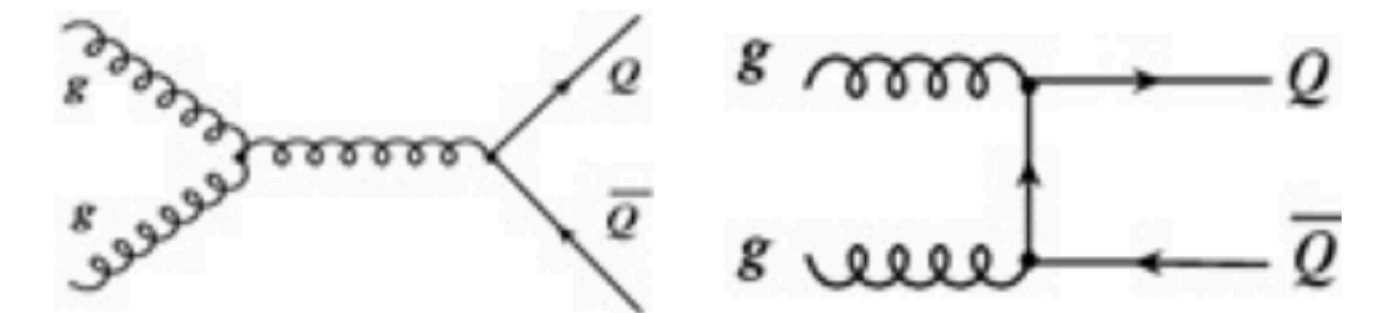
Gluon TMDs

Theory framework well consolidated, but experimental access still extremely limited

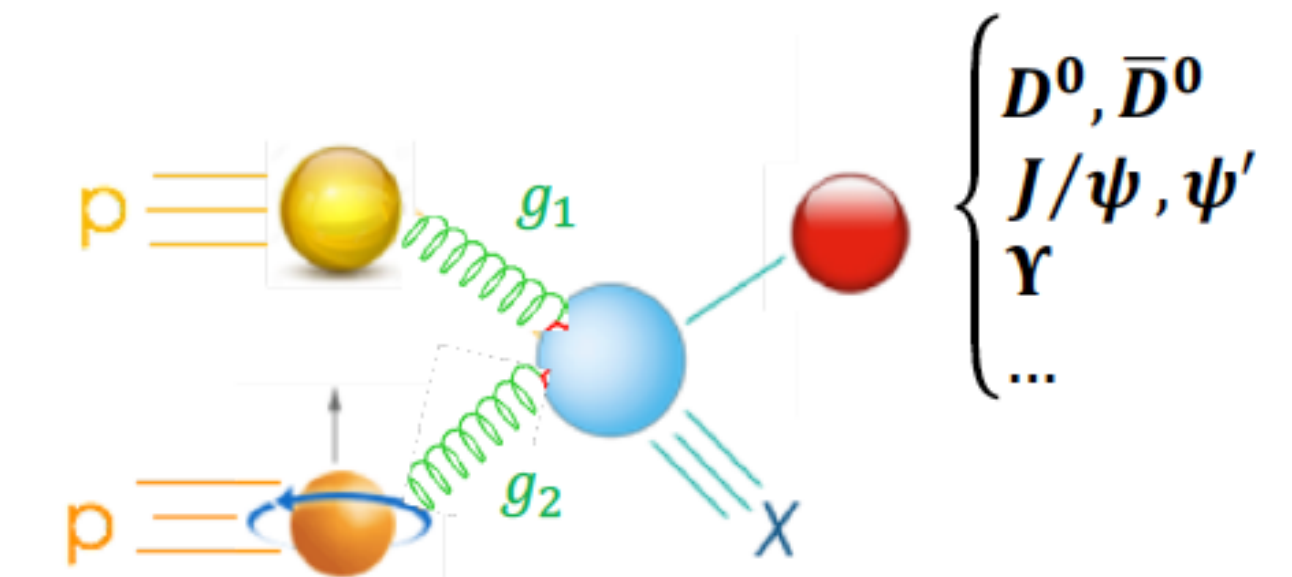
		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables.

At LHC heavy quarks are produced by the dominant gg fusion process



Inclusive quarkonia production in (un)polarized pp interaction turns out to be an ideal observable to access gTMDs



TMD factorisation requires $q_T(Q) \ll M_Q$:

- Can look at associate quarkonia production, where only relative q_T needs to be small (e.g. $pp^{(\uparrow)} \rightarrow J/\Psi + J/\Psi + X$)
- Due to the large masses, easier in case of bottomonium where factorisation can hold at large q_T

Gluon TMDs

Theory framework well consolidated, but experimental access still extremely limited

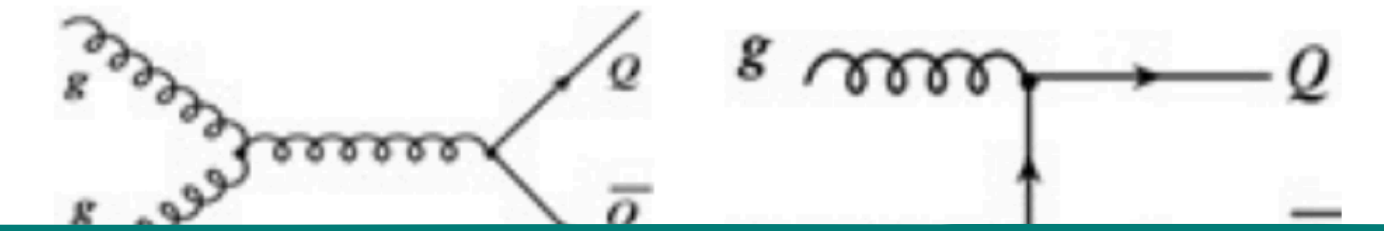
gluon pol.

	U	Circularly	Linearly
U	f_1^g		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

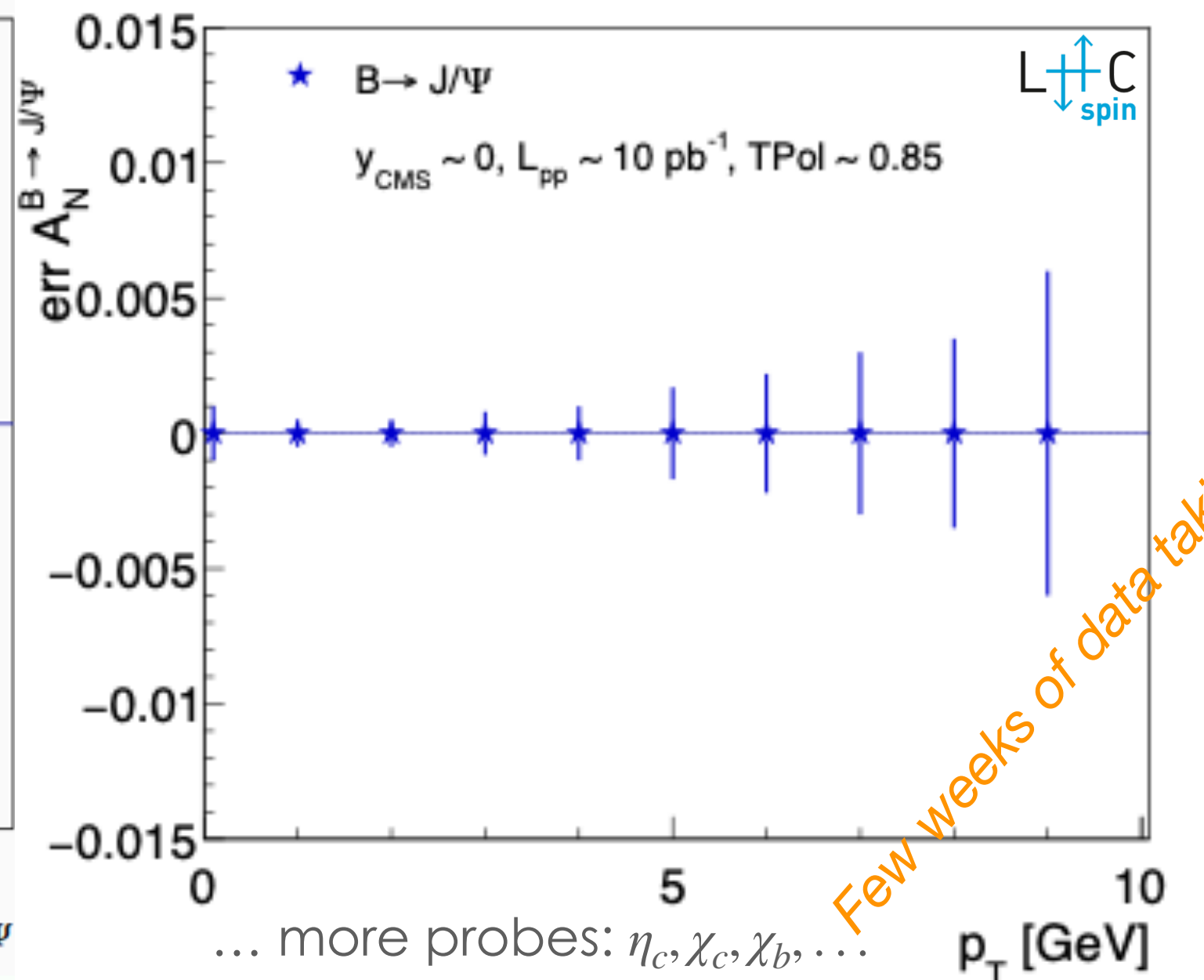
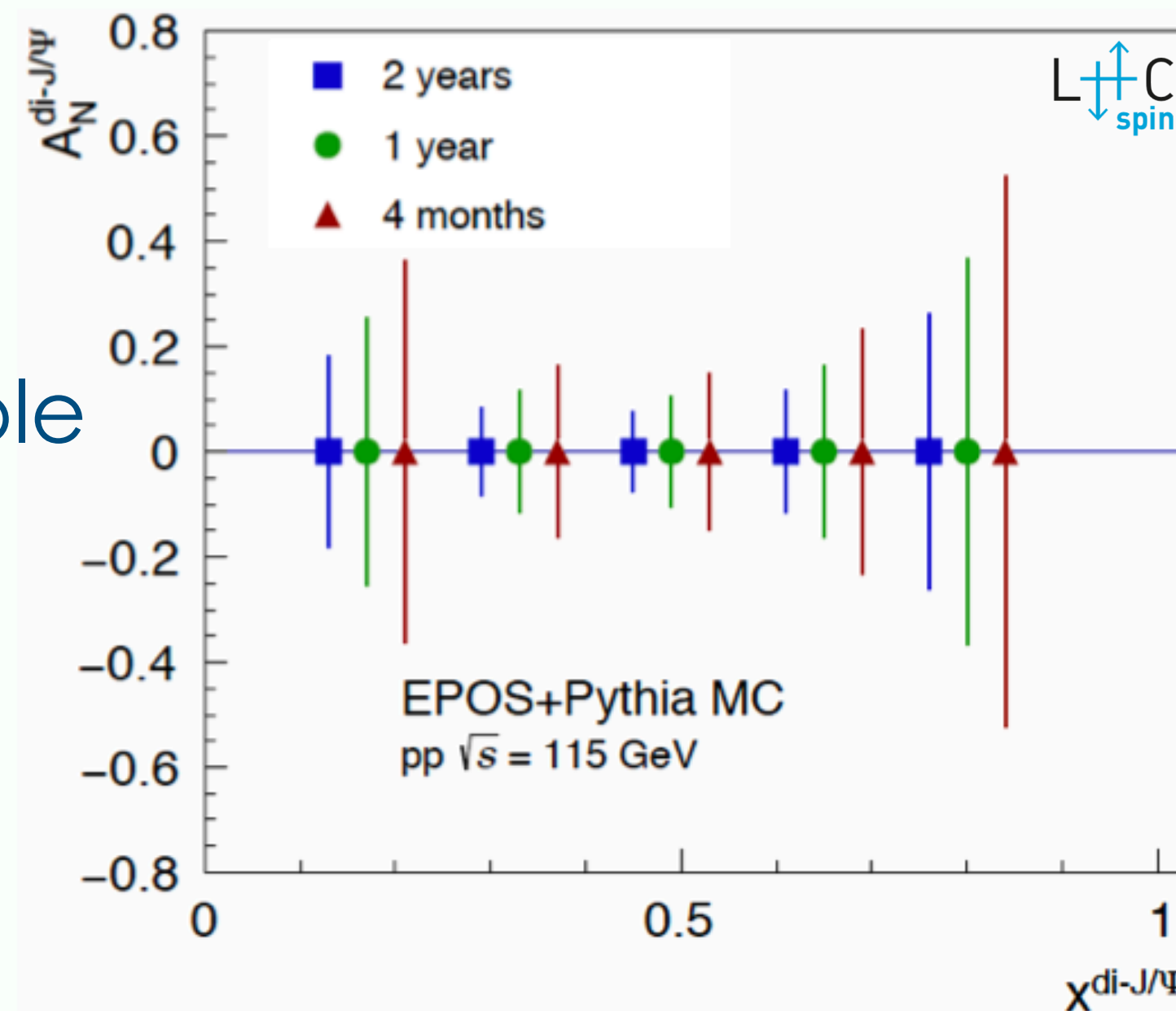
nucleon pol.

The most efficient way to access the gluon dynamics inside the proton at LHC is to measure heavy-quark observables.

At LHC heavy quarks are produced by the dominant $q\bar{q}$ fusion

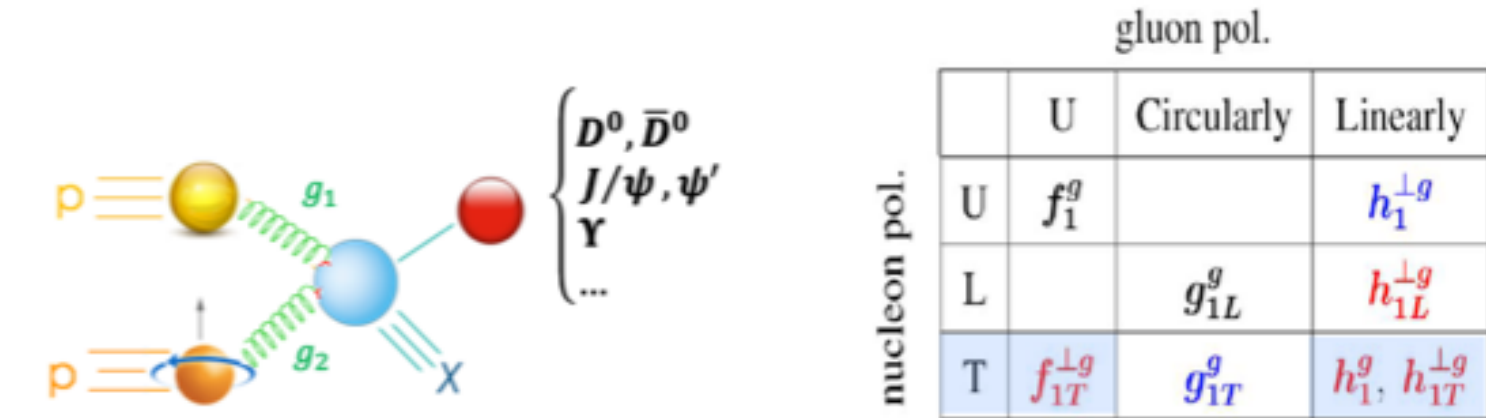


Gluon-induced asymmetries (unconstrained $h_1^{\perp g} + f_1^g$) accessible by, e.g., $di - J/\Psi$ or Υ production



factorisation can hold at large q_T

Probing the Sivers function

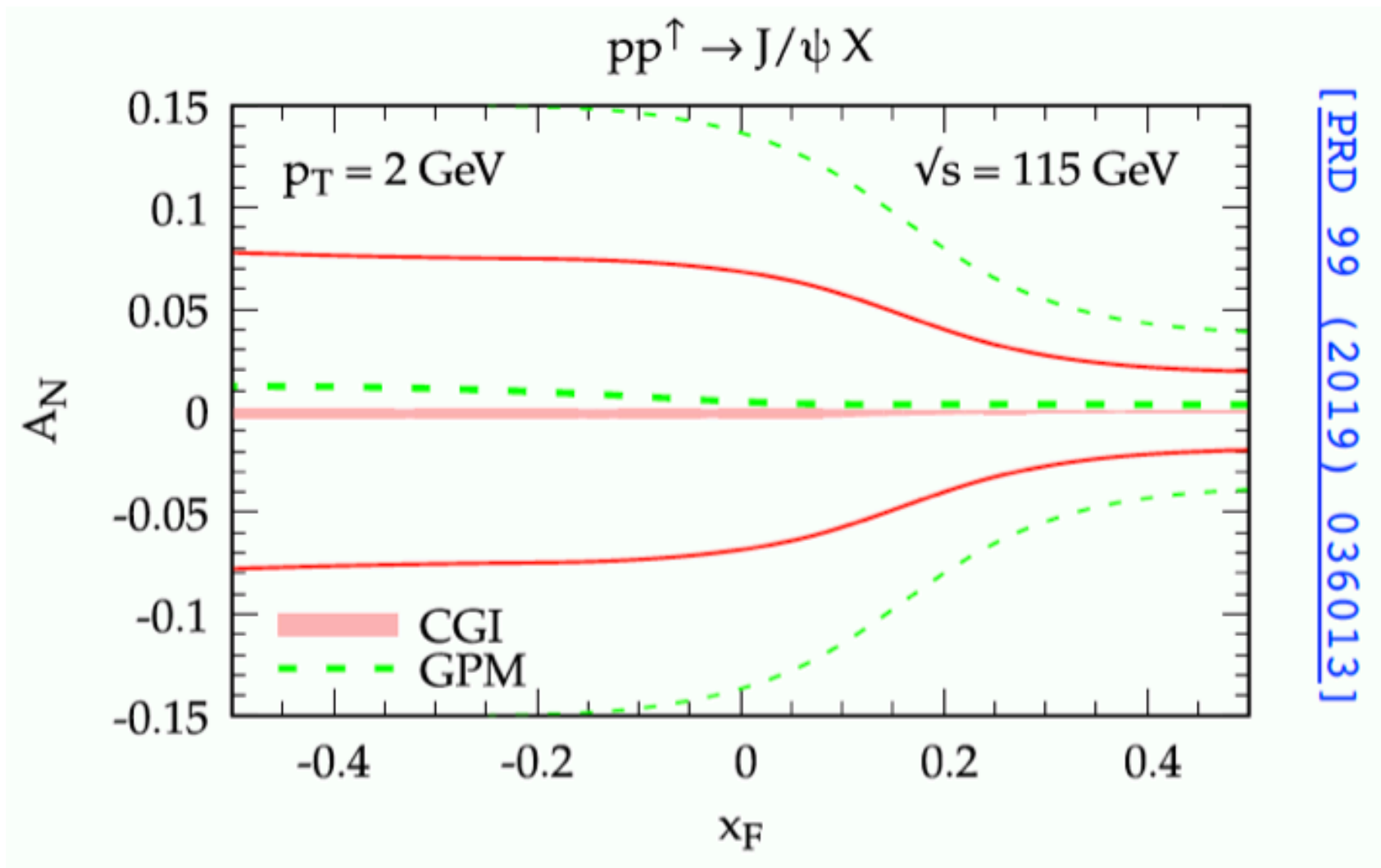


Can be accessed through the Fourier decomposition of the TSSAs for inclusive meson production

$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto \left[\underline{f_{1T}^{\perp g}}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \rightarrow QQg} \right] \sin \phi_S + \dots$$

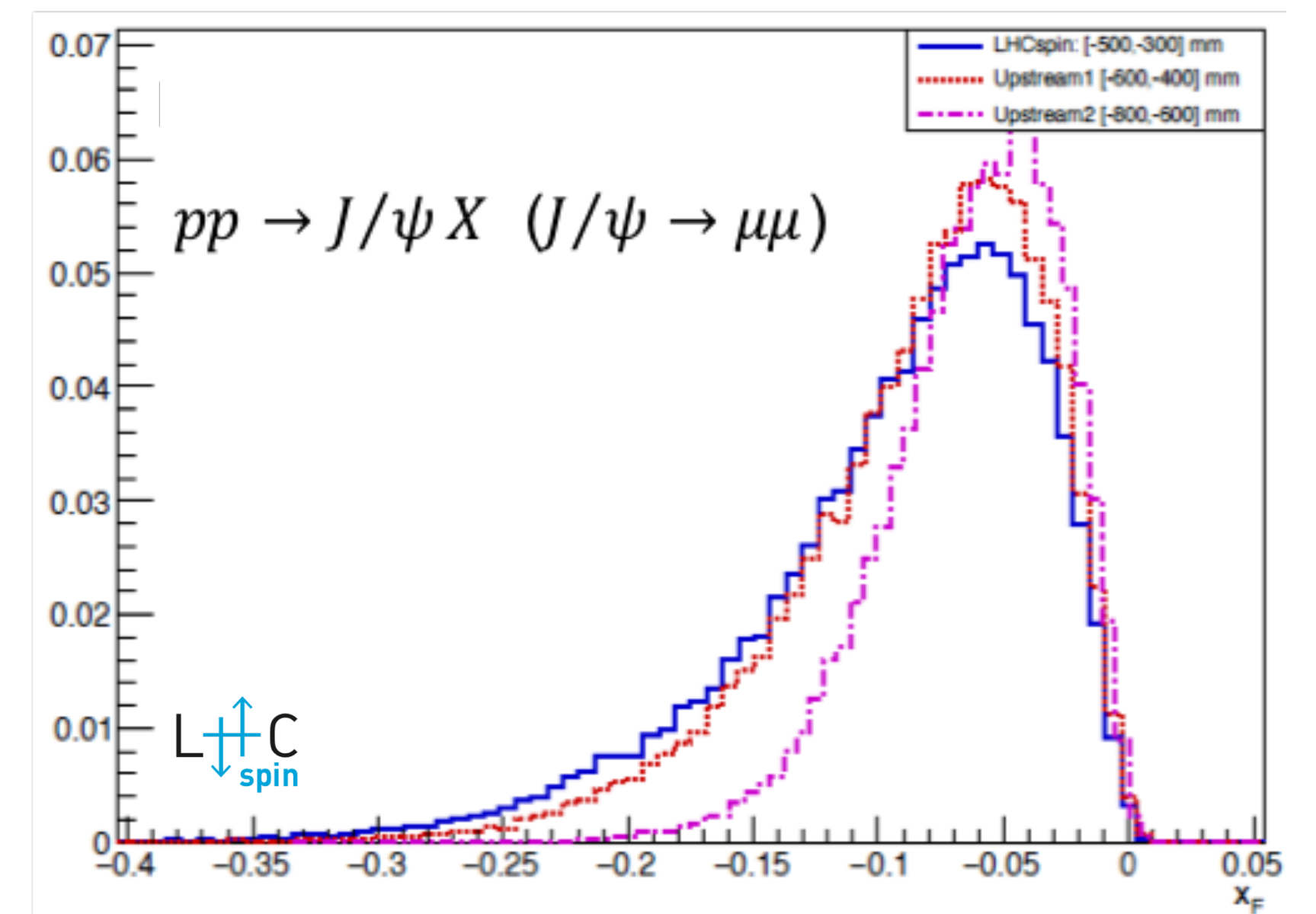
Sensitive to color exchange among IS and FS, and gluon OAM

Shed light on spin-orbit correlation of unpolarized gluons inside a transversely polarized proton



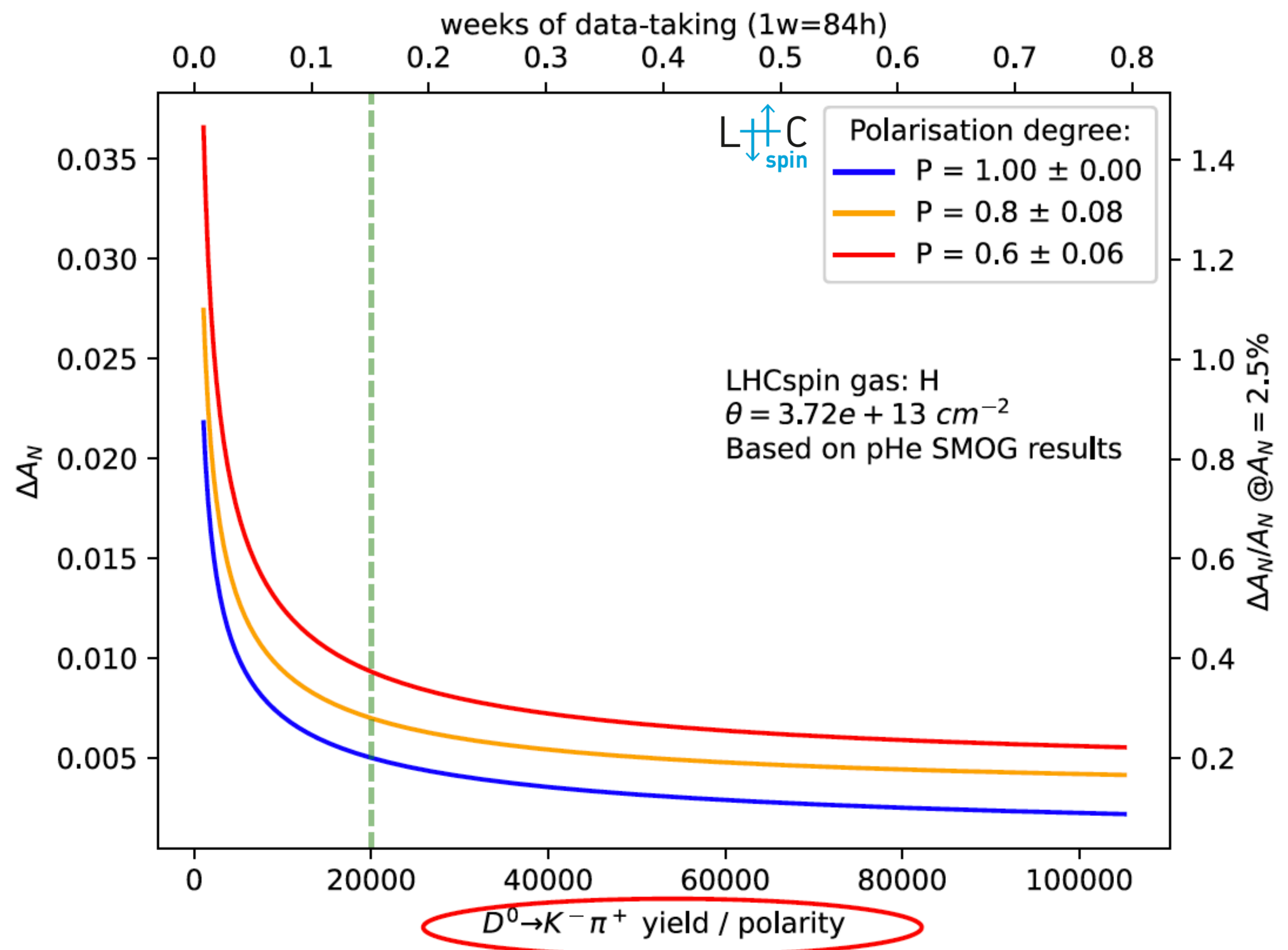
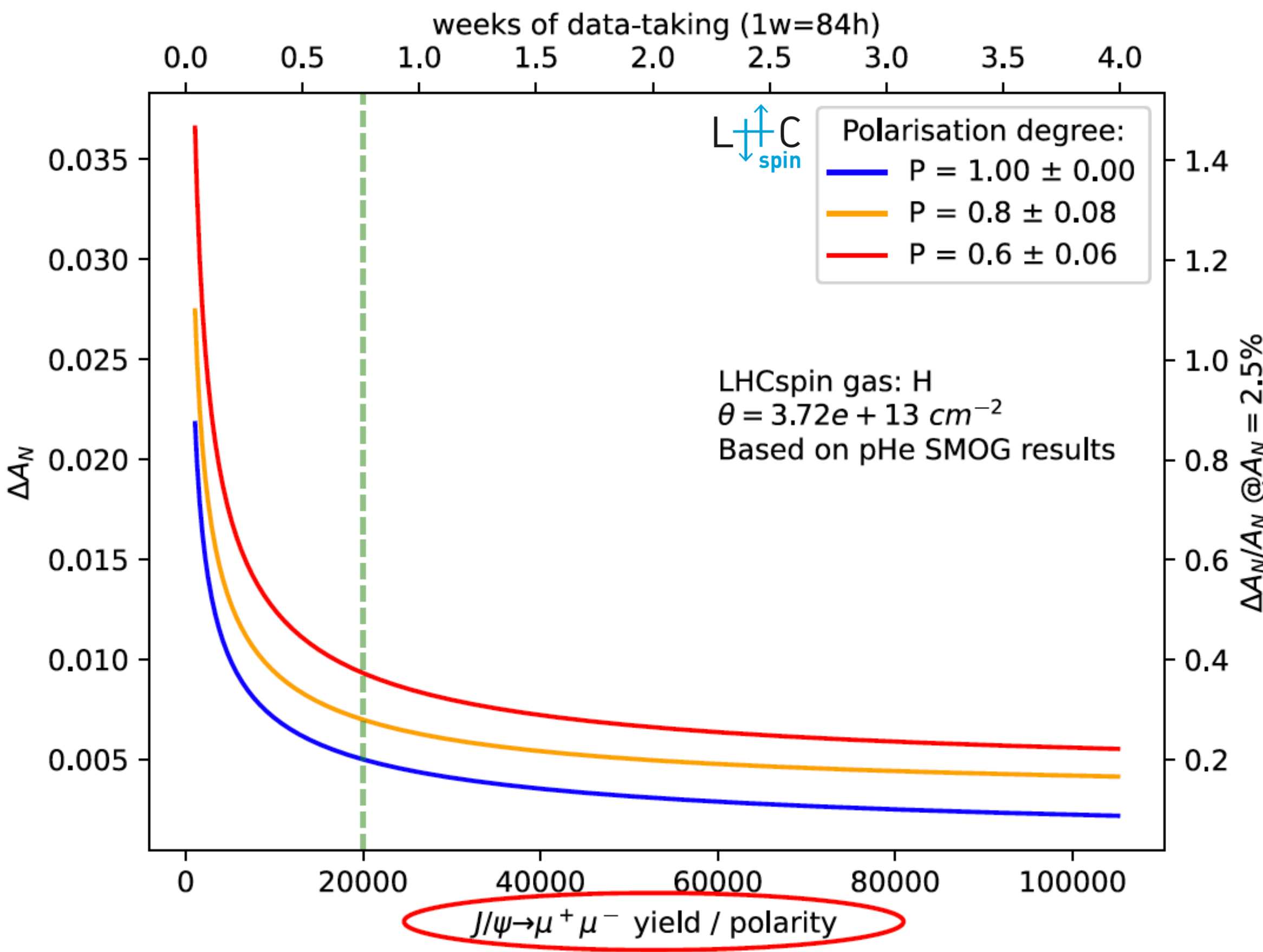
[PRD 99 (2019) 036013]

Predictions for J/Ψ production based on GPM & CGI-GPM
Expected amplitudes could reach 5-10% in the $x_F < 0$ region



LHCspin event rates

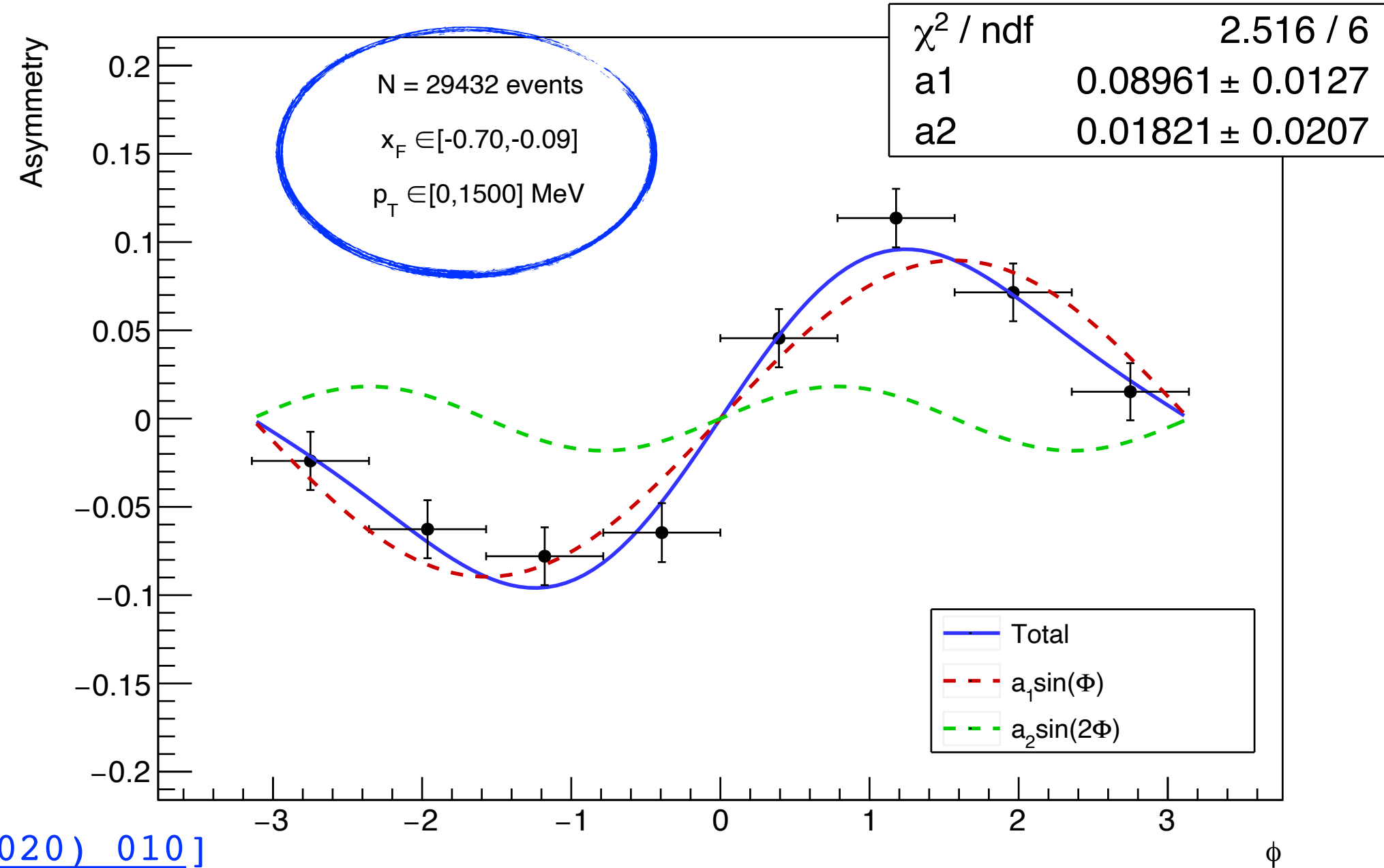
Precise spin asymmetry on $J/\Psi \rightarrow \mu^+ \mu^-$ and $D^0 \rightarrow K^- \pi^+$ for pH^\uparrow collisions in just few weeks with Run3 luminosity!
 Statistics further enhanced by a factor 3-5 in LHCb upgrade II



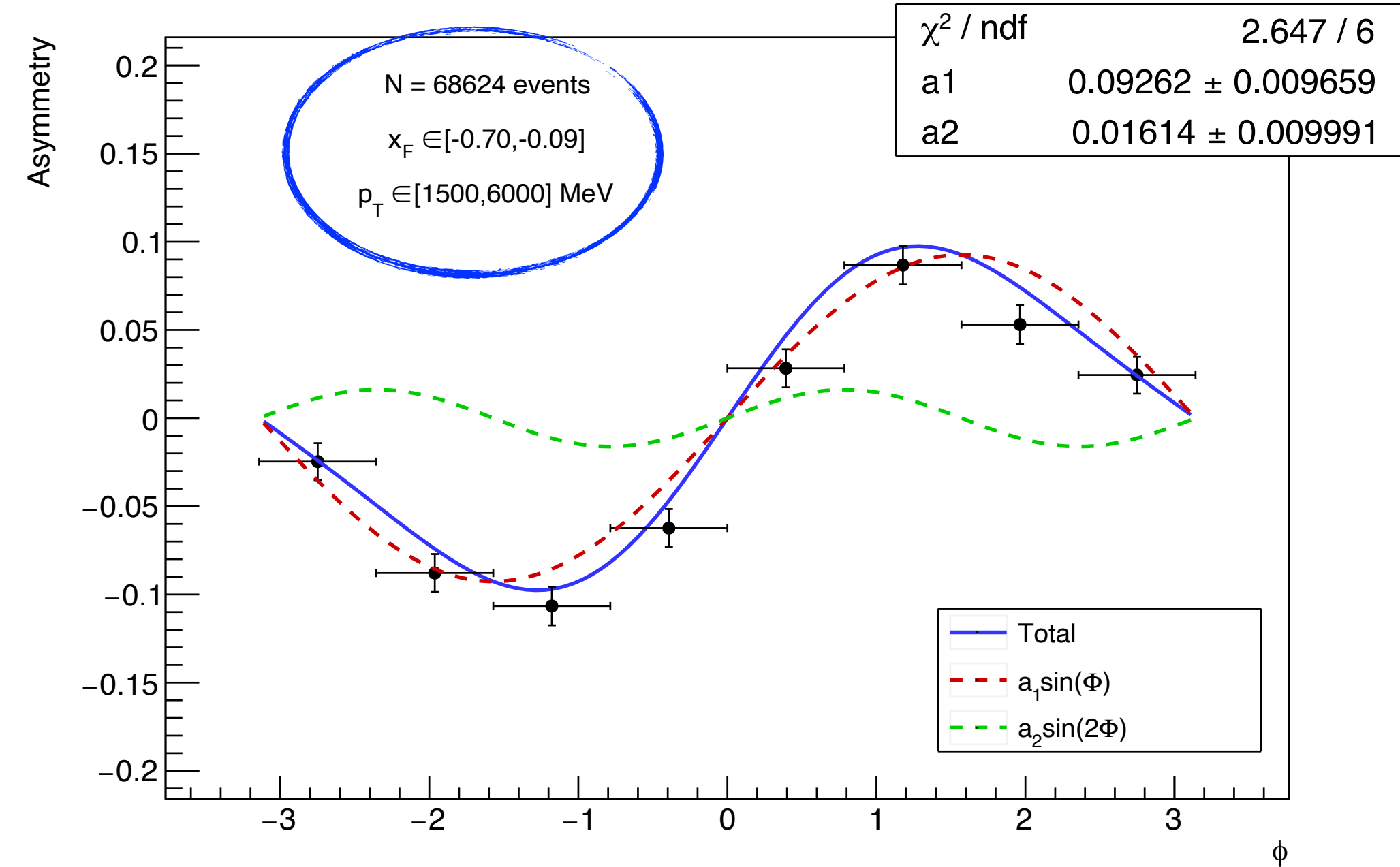
reconstructed particles

A TSSA analysis at LHCspin with $J/\Psi \rightarrow \mu^+\mu^-$ events (toy model)

TSSA on $J/\Psi \rightarrow \mu^+\mu^-$



TSSA on $J/\Psi \rightarrow \mu^+\mu^-$

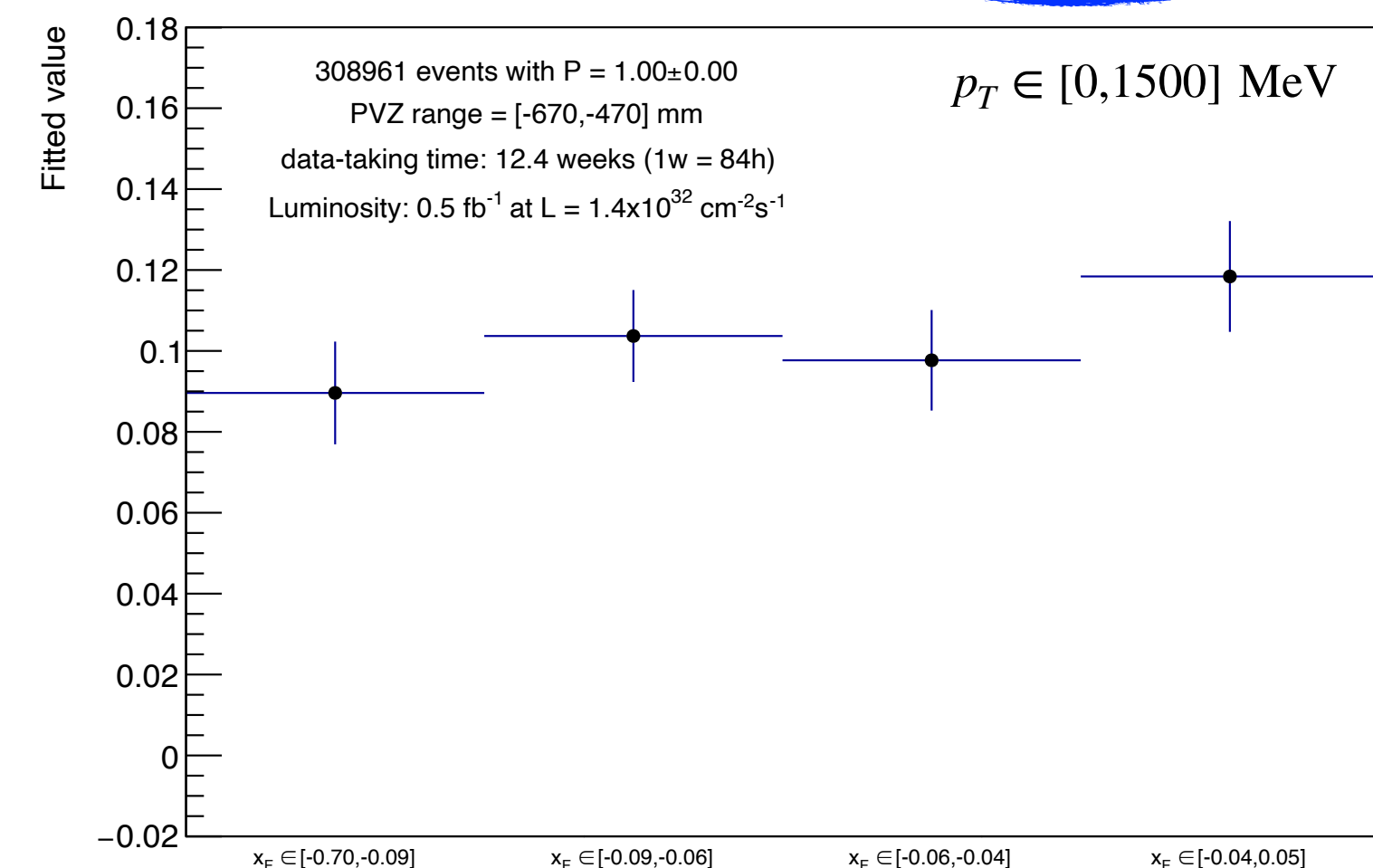


[JHEP 12 (2020) 010]

$$\rho = \frac{1}{2} \left[1 + \left(a_1 + a_2 \frac{x - \bar{x}}{x_{max}} + a_3 \frac{p_T - \bar{p}_T}{p_{T max}} \right) \sin \phi + \left(b_1 + b_2 \frac{x - \bar{x}}{x_{max}} + b_3 \frac{p_T - \bar{p}_T}{p_{T max}} \right) \sin 2\phi \right]$$

- Full LHCb simulations of $J/\Psi \rightarrow \mu^+\mu^-$ in pH collisions \rightarrow emulate the target polarisation by assigning a $\uparrow \downarrow$ tag according to a given model. In this example: 10% asymmetry on $\sin \phi$, 2% on $\sin 2\phi$ + mild x_F, p_T dependence
- Fit the polarised data with the sum of two Fourier amplitudes (a_1, a_2) in $4 x_F \times 2 p_T \times 8 \phi$ bins
- Within this statistics, corresponding to ~ 3 months of data-taking,

$J/\Psi \rightarrow \mu^+\mu^-$: fit results for parameter a_1



$A_N \sim 0.1 \pm 0.01$

Knowledge of the polarisation degree

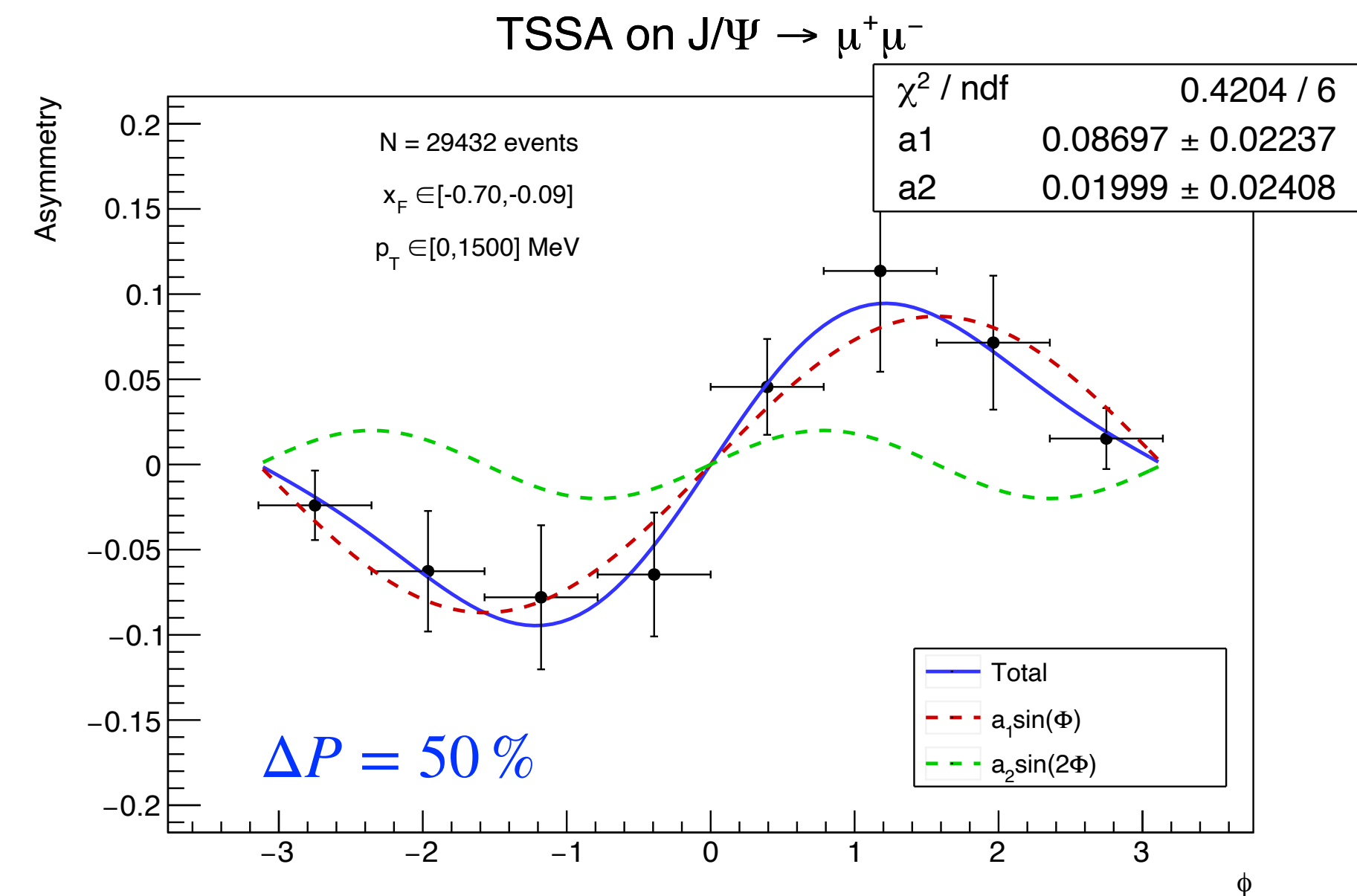
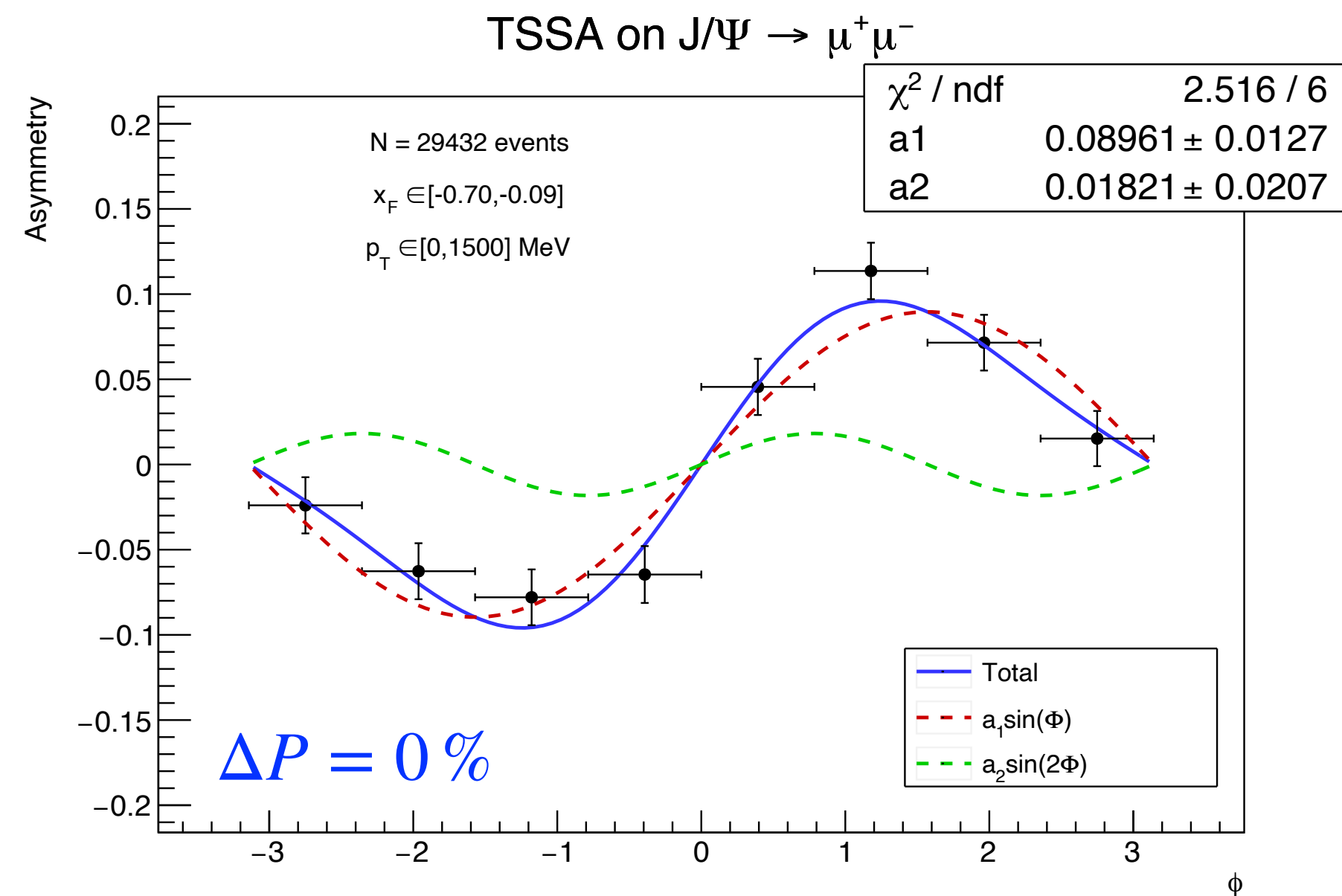
- To estimate the systematic error due to the measurement of the polarisation degree, the analysis is repeated with different ΔP
- Very relevant for the R&D (e.g. cell vs jet target). With the shown analysis* :
- 5% error (realistic value) \rightarrow negligible effect
- 20% error \rightarrow 30-40% of the stat. error
- 50% error \rightarrow syst. dominated

$\Delta P = 5\%$

p_T (MeV)	x_F	a_1
[0,1500]	[-0.70,-0.09]	0.089 ± 0.013
[0,1500]	[-0.09,-0.06]	0.104 ± 0.012
[0,1500]	[-0.06,-0.04]	0.098 ± 0.013
[0,1500]	[-0.04,0.05]	0.117 ± 0.014
[1500,6000]	[-0.70,-0.09]	0.092 ± 0.010
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.011
[1500,6000]	[-0.06,-0.04]	0.105 ± 0.012
[1500,6000]	[-0.04,0.05]	0.105 ± 0.012

$\Delta P = 20\%$

p_T (MeV)	x_F	a_1
[0,1500]	[-0.70,-0.09]	0.087 ± 0.014
[0,1500]	[-0.09,-0.06]	0.103 ± 0.016
[0,1500]	[-0.06,-0.04]	0.097 ± 0.016
[0,1500]	[-0.04,0.05]	0.114 ± 0.017
[1500,6000]	[-0.70,-0.09]	0.090 ± 0.013
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.015
[1500,6000]	[-0.06,-0.04]	0.104 ± 0.015
[1500,6000]	[-0.04,0.05]	0.102 ± 0.015

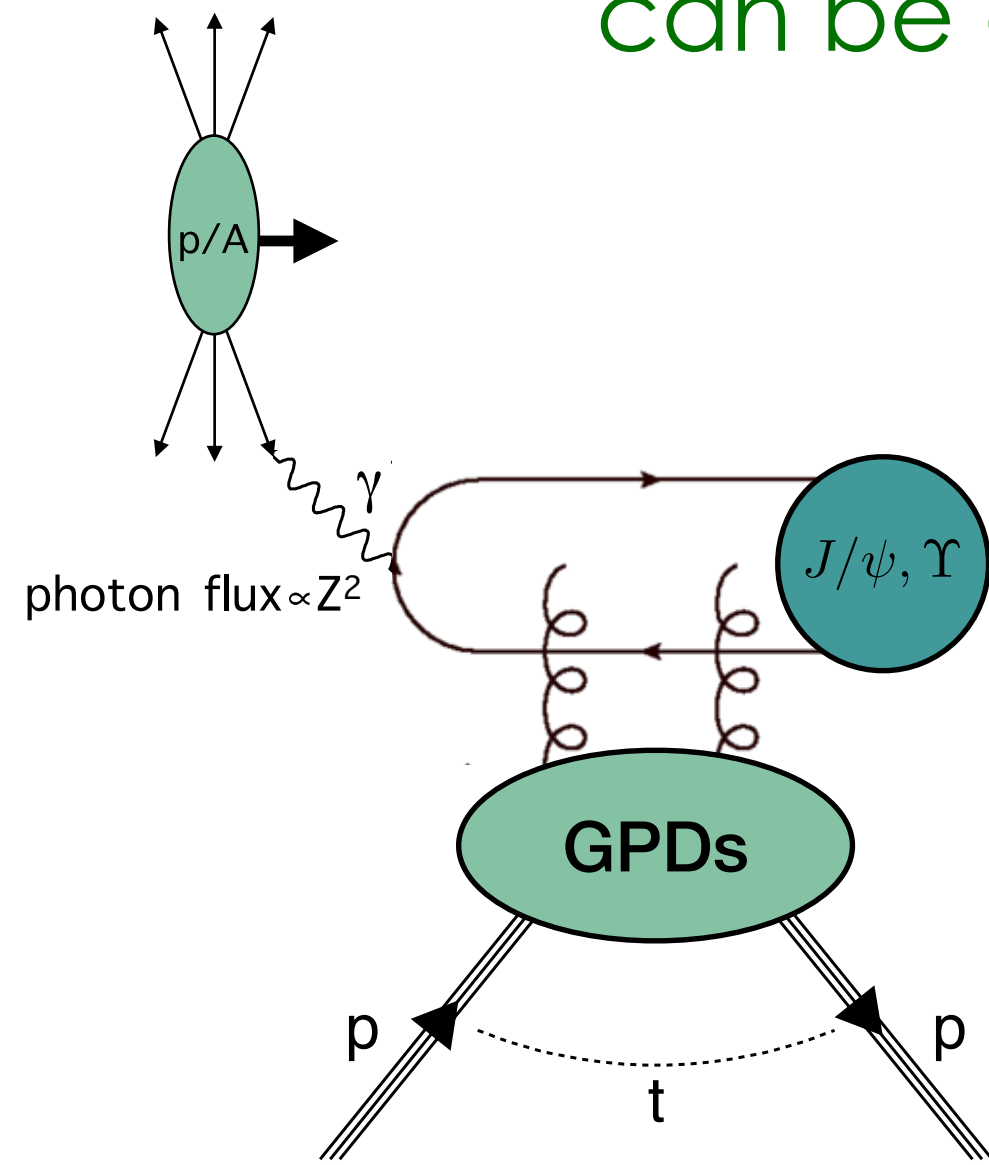


* i.e. \sim 3 months of data-taking with this example model, channel and kinematic binning

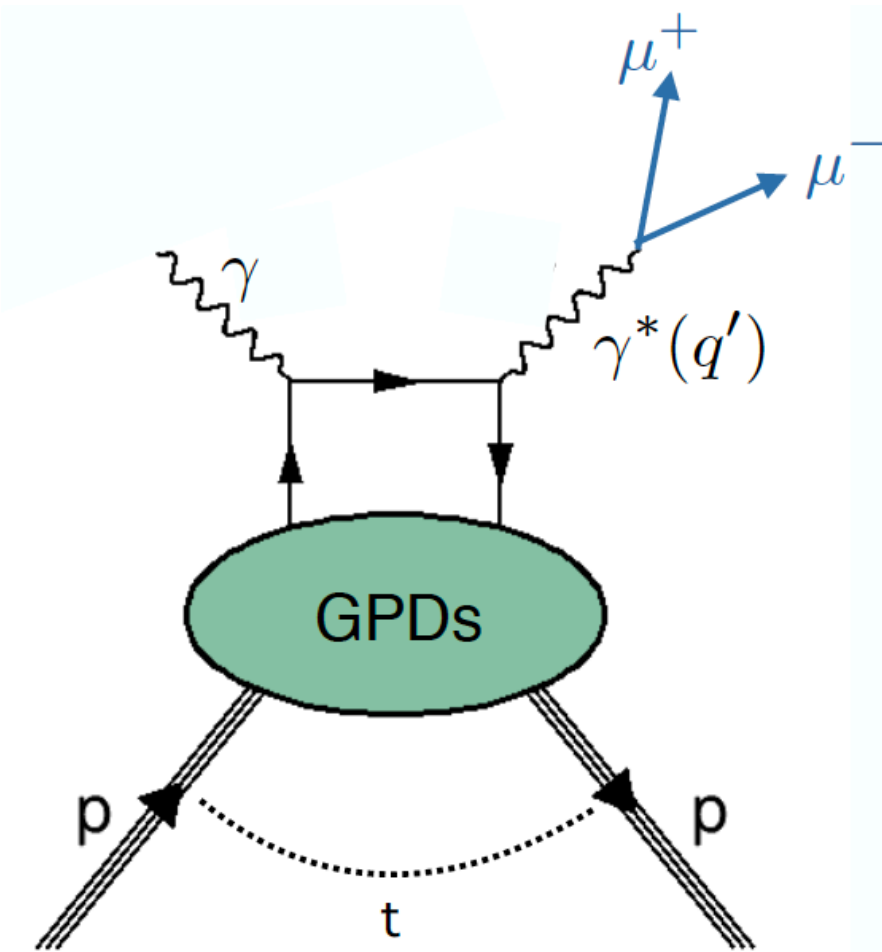
UPC and gGPDs

Accessible already with SMOG2 for the unpol part

can be accessed at LHC in Ultra-Peripheral collisions (UPC)



Exclusive meson production
hard scale = quark mass



Timelike Compton scattering (TCS)
(access via angular modulation)
hard scale = large q^2 (in practice few GeV^2)

Recall:
- barely explored high- x_B region
- moderate Q^2

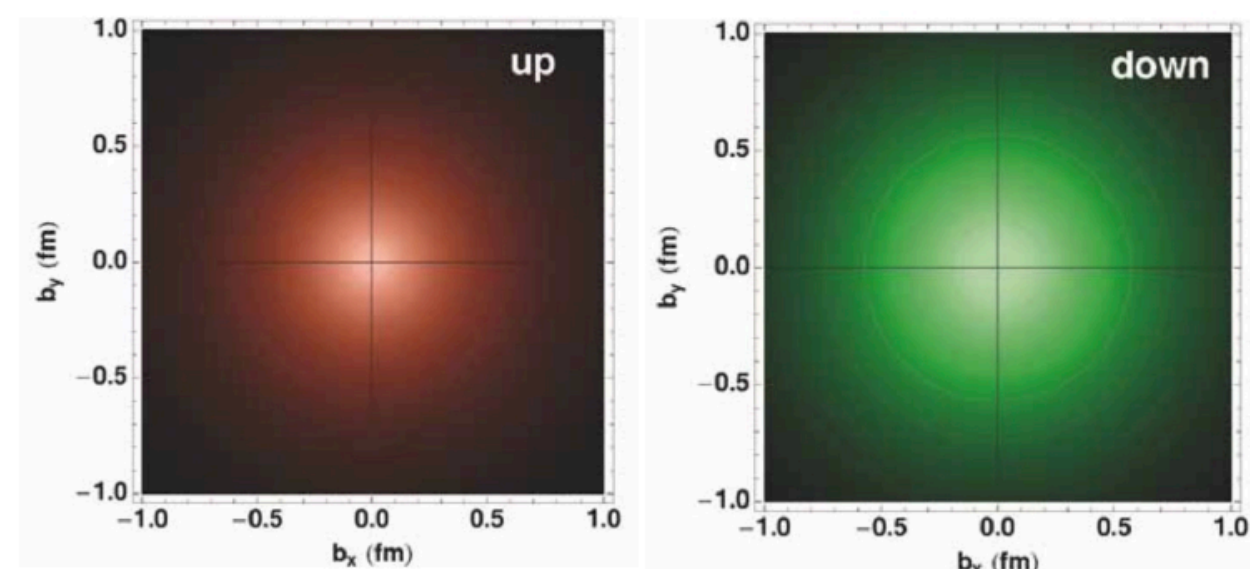
- Impact parameter larger than sum of radii
 - Process dominated by EM interactions
 - Gluon distributions probed by pomeron exchange
 - Exclusive quarkonia prod. sensitive to gluon GPDs
- [PRD 85 (2012), 051502]

LHCspin could allow to access the GPD E^g (a key ingredient of the Ji sum rule)

$$J^g = \frac{1}{2} \int_0^1 dx \left(H^g(x, \xi, 0) + E^g(x, \xi, 0) \right)$$

GPD	U	L	T
U	H		\mathcal{E}_T
L		\tilde{H}	\tilde{E}_T
T	E	\tilde{E}	H_T, \tilde{H}_T

3D maps of parton densities in coordinate space

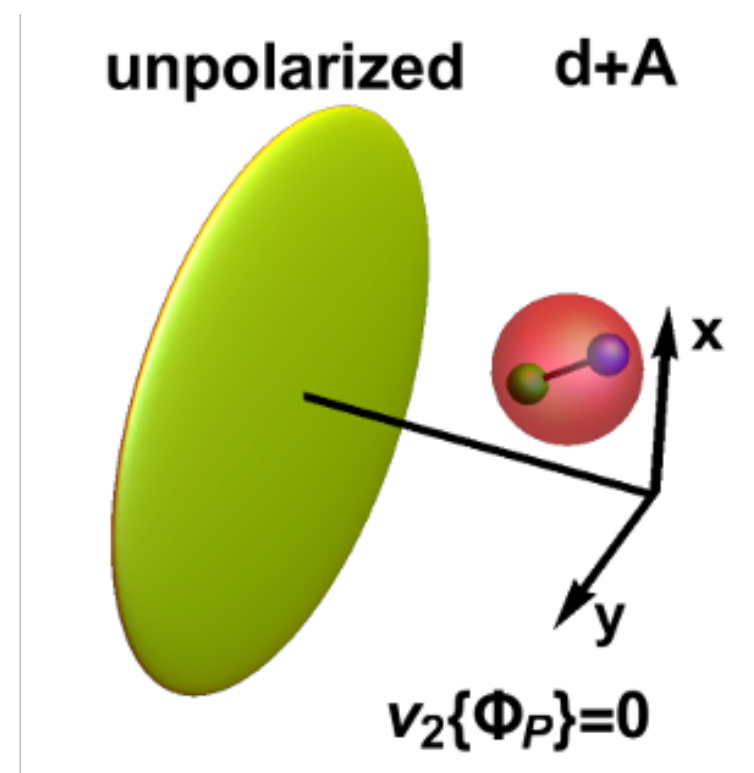
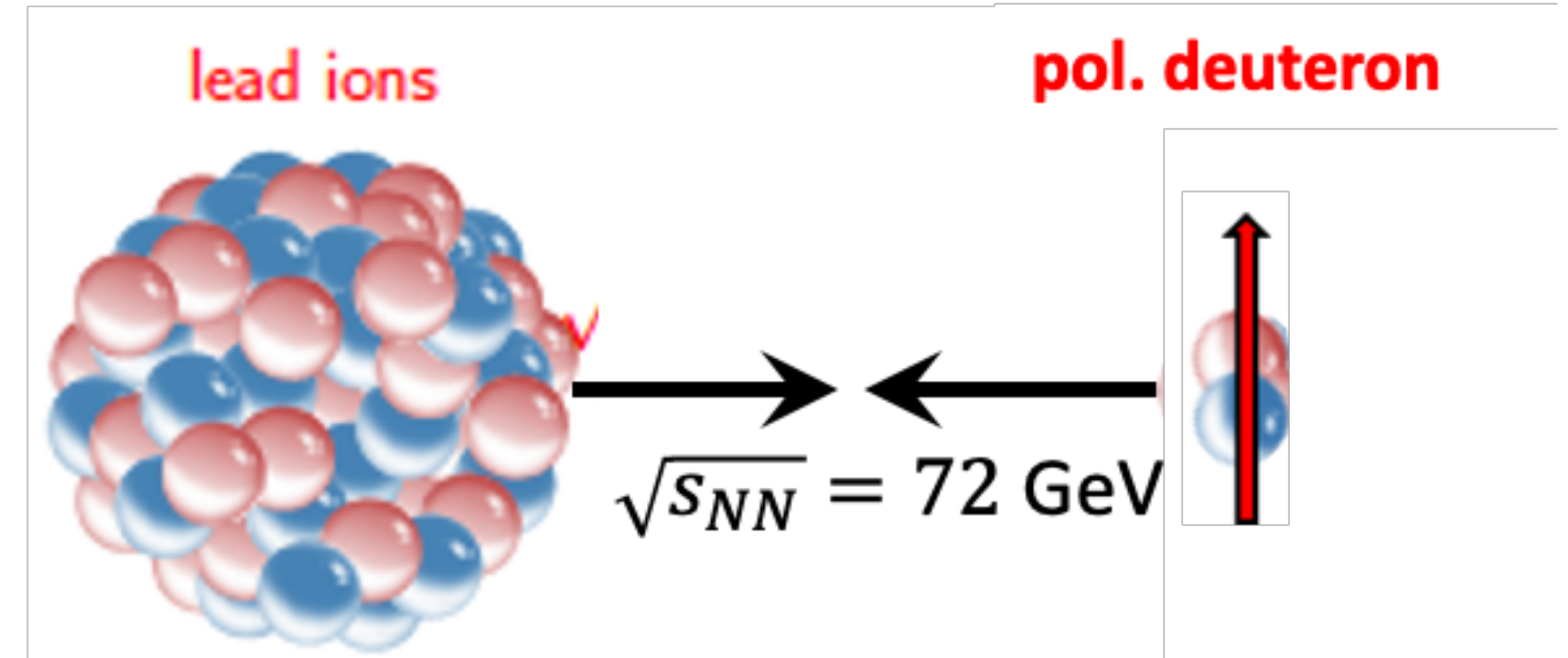


J/ψ , total uncertainty on cross section, assuming 4% uncertainty on luminosity

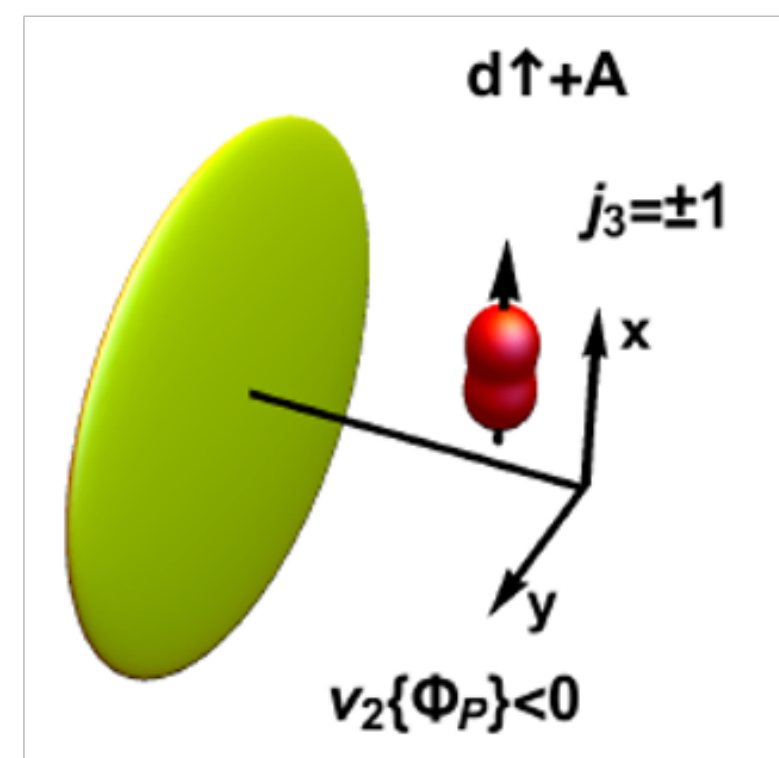
pp	pD	pAr	pKr	pXe
10 %	-	5 %	5 %	5 %
Pbp		PbAr		
-		5 %		

Spin physics in heavy-ion collisions

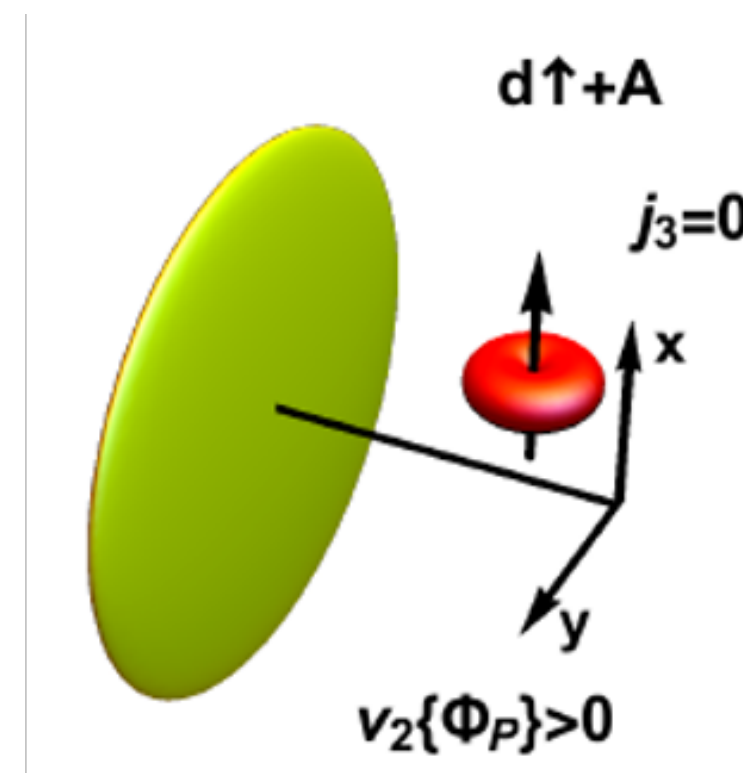
- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).



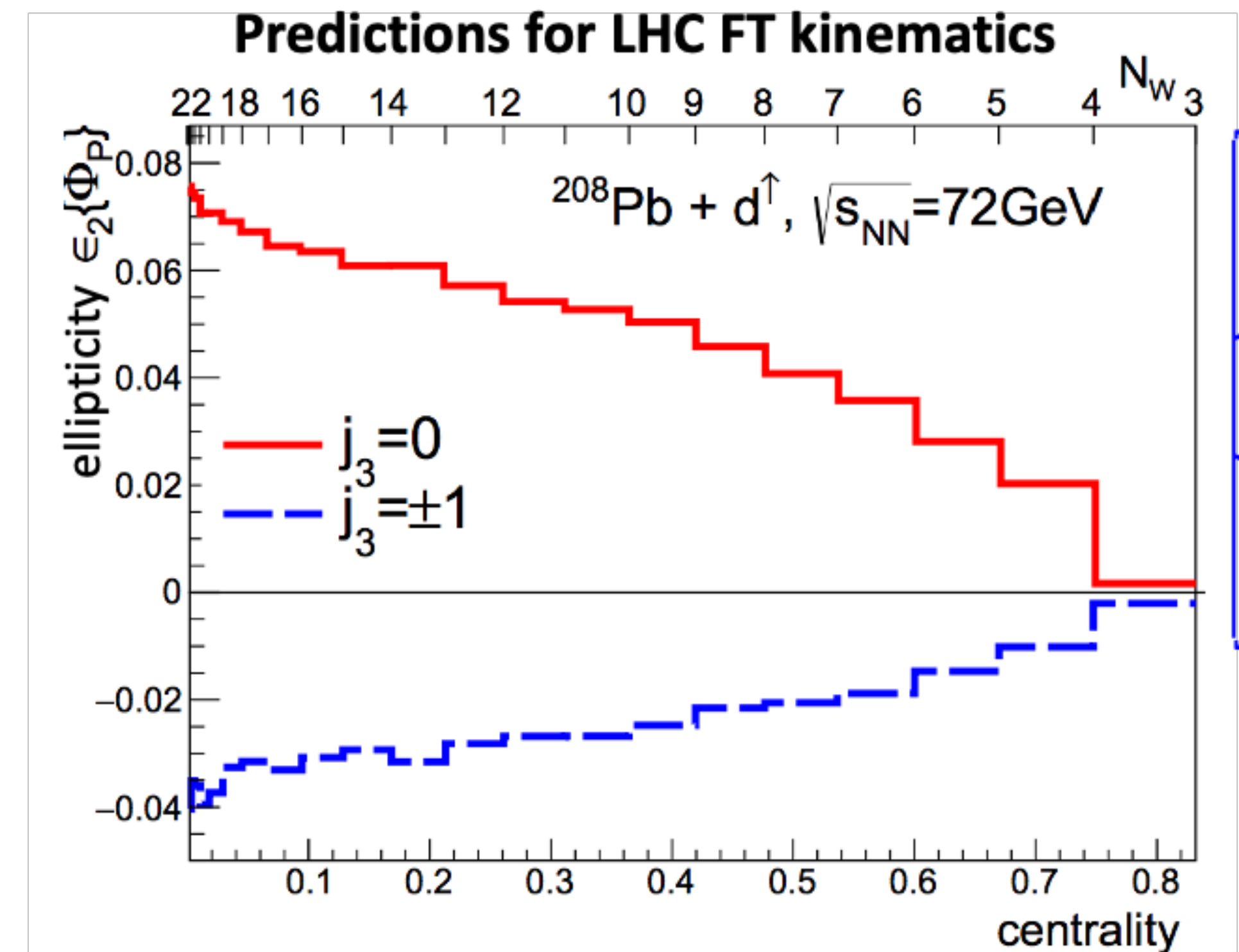
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.



$j_3 = \pm 1 \rightarrow$ prolate fireball stretched along the pol. axis, corresponds to $v_2 < 0$



$j_3 = 0 \rightarrow$ oblate fireball corresponds to $v_2 > 0$



[PRC 101 (2020) 024901]

Wojciech Broniowski, Piotr Bozek

International framework and feedback

Several experiments dedicated to spin physics, but with many limitations:

very low energy, no rare probes, no ion beam, ...

➔ LHCspin is unique in this respect

LHCspin is complementary to EIC

[D. Boer: [arXiv:1611.06089](https://arxiv.org/abs/1611.06089)]

unpolarized gluon TMD

	DIS	DY	SIDIS	$pA \rightarrow \gamma \text{ jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$f_1^g^{[+,+]}$ (WW)	×	×	×	×	✓	✓	✓
$f_1^g^{[+,-]}$ (DP)	✓	✓	✓	✓	×	×	×

linearly polarized gluon TMD

	$pp \rightarrow \gamma \gamma X$	$pA \rightarrow \gamma^* \text{ jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$h_1^{\perp g [+,+]}$ (WW)	✓	×	✓	✓	✓
$h_1^{\perp g [+, -]}$ (DP)	×	✓	×	×	×

TMDs (Sivers)

[D. Boer: [arXiv:1611.06089](https://arxiv.org/abs/1611.06089), D. Boer et al. HEPJ 08 2016 001]

	DY	SIDIS	$p^\dagger A \rightarrow h X$	$p^\dagger A \rightarrow \gamma^{(*)} \text{ jet } X$	$p^\dagger p \rightarrow \gamma \gamma X$ $p^\dagger p \rightarrow J/\psi \gamma X$ $p^\dagger p \rightarrow J/\psi J/\psi X$	$ep^\dagger \rightarrow e' Q \bar{Q} X$ $ep^\dagger \rightarrow e' j_1 j_2 X$
$f_{1T}^{\perp g [+,+]}$ (WW)	×	×	×	×	✓	✓
$f_{1T}^{\perp g [+, -]}$ (DP)	✓	✓	✓	✓	×	×

$f_{1T}^{\perp g [+,+]}$ (Weizsacker-Williams type or “f-type”) → antisymmetric colour structures

$f_{1T}^{\perp g [+, -]}$ (Dipole s type or “d-type”) → symmetric colour structures

☐ Can be measured at the Electron Ion-Collider (EIC)

☐ Can be measured at LHCspin

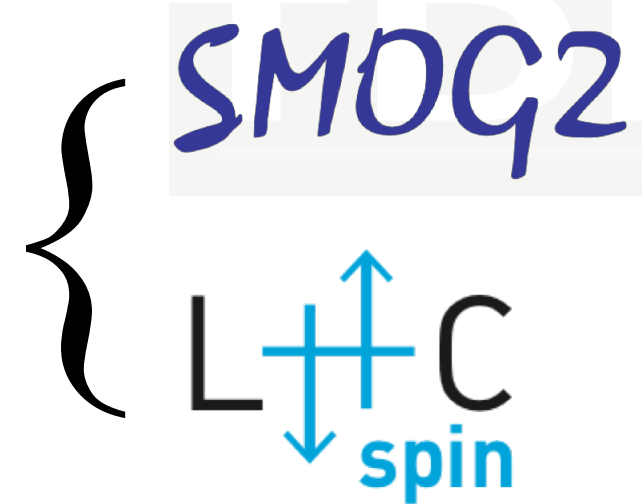
“Ambitious and long term LHC-Fixed Target research program. The efforts of the existing LHC experiments to implement such a programme, including specific R&D actions on the collider, **deserve support**” (European Strategy for Particle Physics)

“This would be **unique and highly complementary** to existing and future measurements in lepton-proton collisions, because the asymmetries in question have a process dependence between pp and lp that is predicted by theory” (CERN Physics Beyond Collider)

Recognised relevance

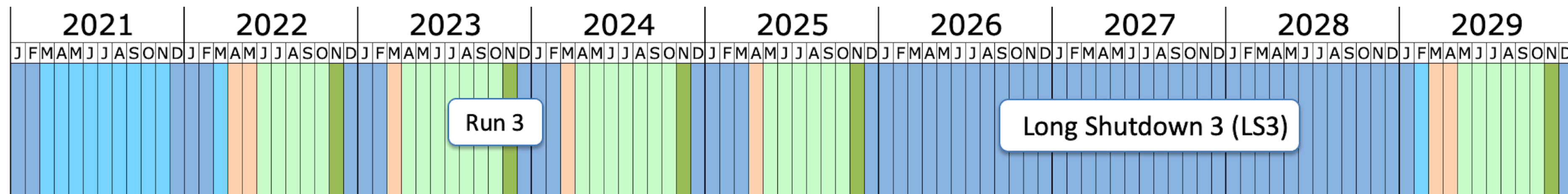
The polarised physics is very alive and will benefit of complementary probes, from LHC to EIC

Fixed target physics at LHC is an exiting reality



SMOQ2 already operative and taking unpolarised data

is an innovative and unique project conceived to bring polarized physics at the LHC. It is extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs



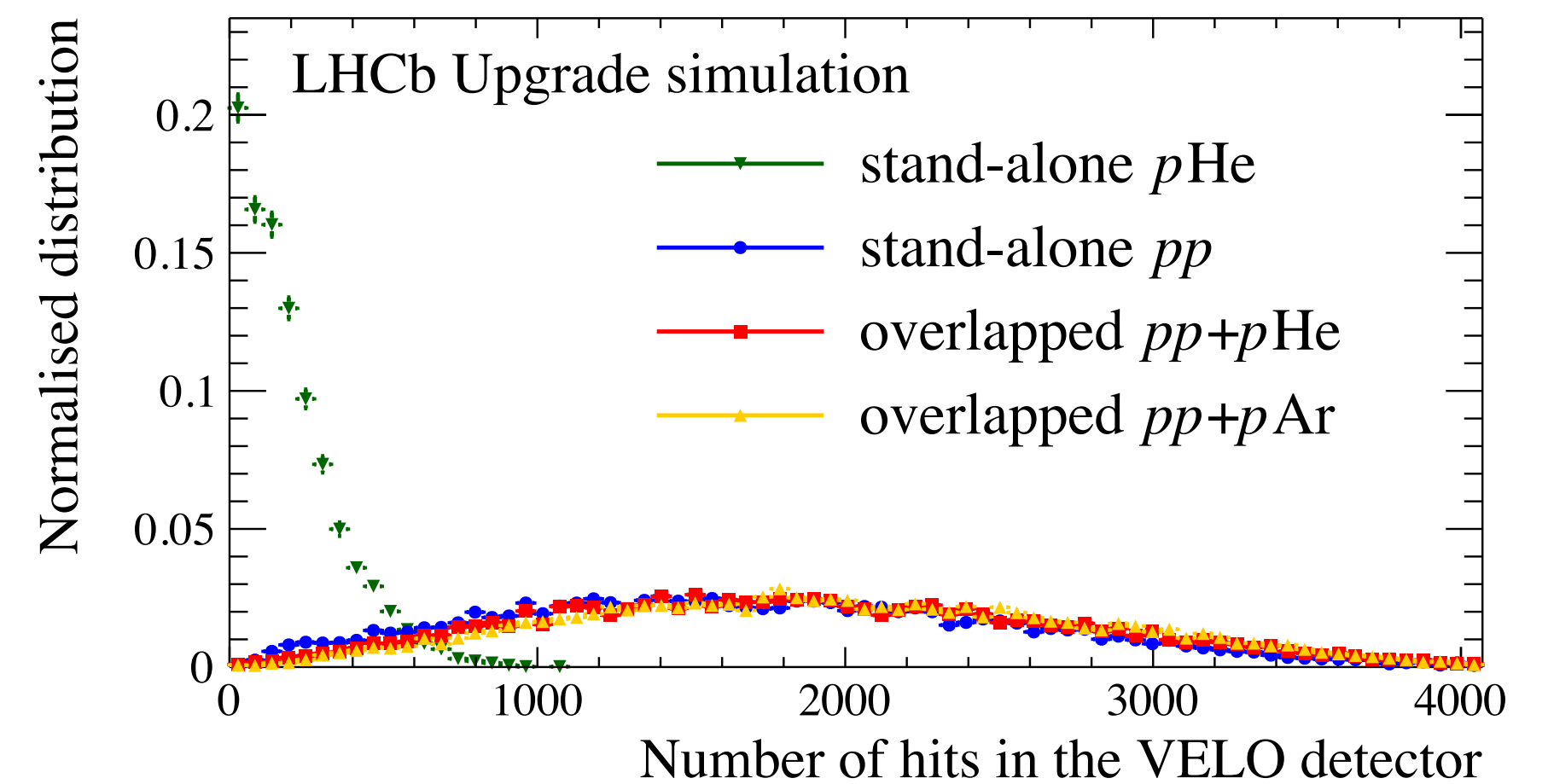
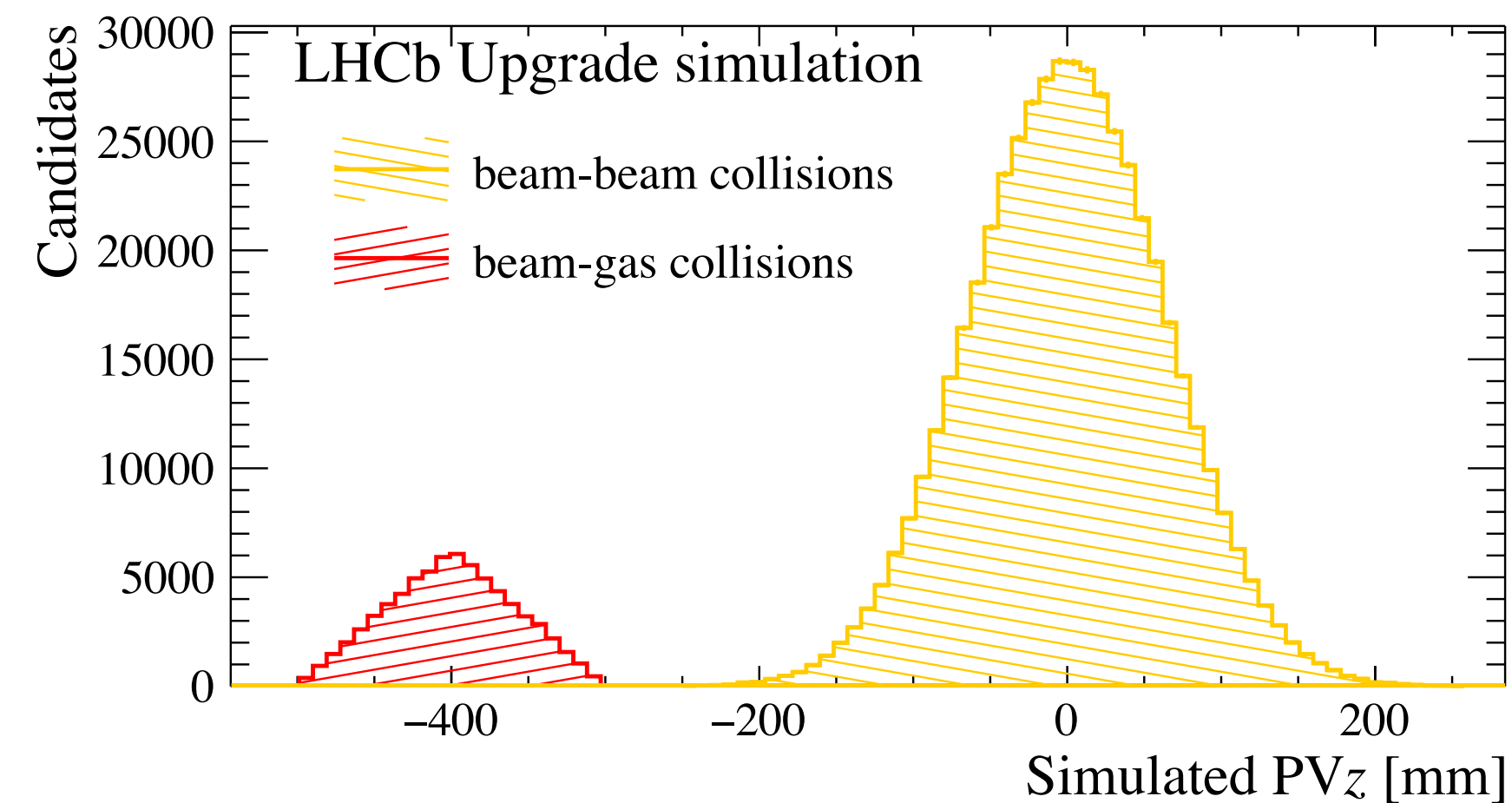
SMOQ2



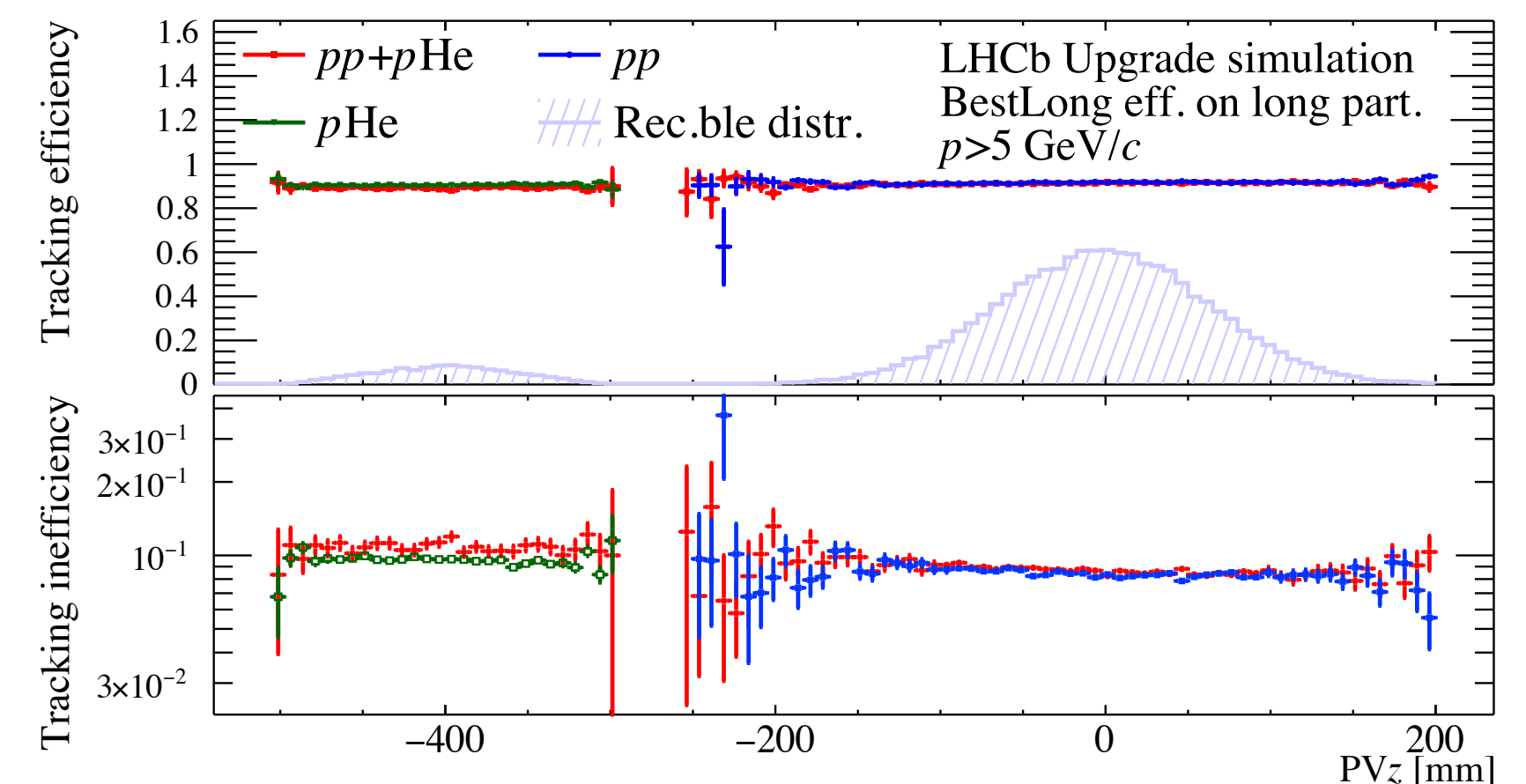
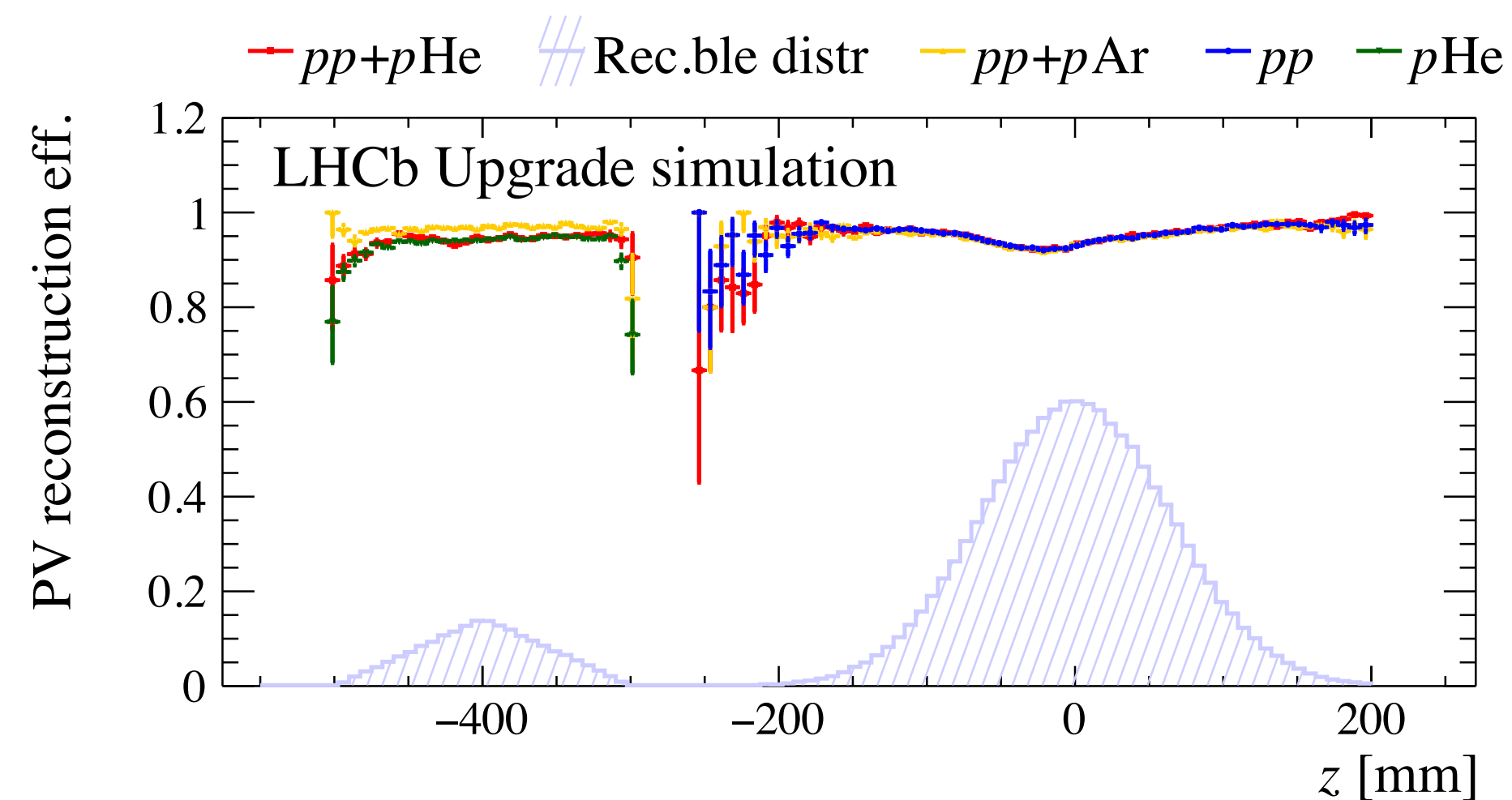
SMOG2/LHCspin performances

[LHCb-FIGURE-2022-002]

- beam-beam and beam-gas interaction regions are well detached
- Negligible increase of multiplicity: 1 – 3 % throughput decrease when adding beam-gas to the LHCb event reconstruction sequence



- Full reconstruction efficiency (PV & tracks) retained in the beam-gas region



LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!