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Systematic analysis of nuclear reactions with a neutron rich projectile on multiple targets at intermediate energies



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Intermediate energy heavy-ion collisions

Nuclear Equation of State

Isospin transport

FAZIA detector

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Intermediate Energy Heavy-Ion Collisions

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<u>Intermediate Energy Heavy-Ion Collisions</u>

 \rightarrow The energy dissipation w.r.t. relative velocity (v_{rel}) b/w the initial partners in any nuclear reaction can be linked to a wavelength, $\lambda = -\frac{\lambda}{\hbar}$

$$\lambda = \frac{\lambda}{2\pi} = \frac{\hbar}{m_u v_{rel}}$$

 \rightarrow If order of λ exceeds the mean nucleon-nucleon distance (~ 2 fm), the reaction corresponds to low energy regime (Beam energy, $E_p < 20$ MeV/A).

 \rightarrow If order λ is shorter than the mean n-n distance, the reaction falls into the high energy regime (E_B> 100 MeV/A).

→When order of λ is comparable to the mean n-n distance, the reaction falls in the intermediate or Fermi energy range, (20 < E_R< 100 MeV/A).



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Intermediate Energy Heavy-Ion Collisions



Intermediate Energy Heavy-Ion Collisions



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 \rightarrow The nuclear Equation of State (EoS) describes relation between energy, temperature, nuclear density and neutron-proton (N/Z) ratios (isospin).

→ Staring with the Liquid Drop Model, the Bethe-Weizsacker formula for binding energy,

$$B.E. = a_V A - a_S A^{2/3} - a_C \frac{Z^2}{A^{1/3}} - a_A \frac{(A - 2Z)^2}{A} \pm \delta(A, Z)$$

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where, A=N+Z is the mass number and a_V , a_S , a_C and a_A are coefficients.

 \rightarrow Valid for a system at zero temperature and nuclear saturation density (~0.16 fm⁻³).

 \rightarrow For asymmetric nuclear matter, the binding energy can be written as a function of neutron and proton densities (ρ_n, ρ_n),

$$\frac{E}{A}(\rho, I) = \frac{E}{A}(\rho) + \frac{E_{sym}}{A}(\rho)I^2 + \dots$$
$$E_{sym}(\rho) = \frac{1}{2}\frac{\partial^2 E(\rho, I)}{\partial I^2}\Big|_{I=0}$$

where, $\rho = \rho_n + \rho_p$ is the nucleonic density, $I = (\rho_n - \rho_p)/\rho$ is the isospin

asymmetry and \mathbf{E}_{sym} is the symmetry energy.

Symmetry energy as a function of density as predicted by different models. The left panel shows the low density region while the right panel displays the high density range. $\rho_0 \rightarrow$ saturation density (~0.16 fm⁻³)



Importance of studying E_{sym}:

 \rightarrow Taylor series expansion of $\mathrm{E}_{\mathrm{sym}}$ around ρ_{0} ,

$$E_{sym}(\rho) = E_{sym}(\rho_0) - L\epsilon + 1/2K_{sym}\epsilon^2 + O[\epsilon^3]$$

where $\epsilon \equiv (\rho_0 - \rho)/3\rho_0$, L and K_{sym} are the slope and curvature at ρ_0 .

 \rightarrow The slope L provides a link to the astrophysical regime as,

$$L = 3\rho_0 \frac{dE_{sym}(\rho)}{d\rho} \bigg|_{\rho_0} = \frac{3P_0}{\rho_0}$$

where P_0 is the pressure from E_{sym} in pure neutron matter and also contributes to the pressure in a neutron star.

 \rightarrow There have been multiple terrestrial and astrophysical experiments to estimate L (59 ± 17 MeV) and E_{sym} (31.6 ± 0.9 MeV).



 \rightarrow The density dependence of the E_{sym} can be visualised by using, for example, the stochastic mean-field (SMF) transport approach. The E_{sym} can be written as a sum of its kinetic and potential parts as,

 $E_{sym} = E_{sym}(kinetic) + E_{sym}(potential)$

 \rightarrow There are three different density parametrizations of the symmetry potentials:

- strong density dependence ($\propto \rho^2$) Asy-superstiff
- linear dependence on density Asy-stiff
- weaker density dependence with saturation around ρ_0 and further decreasing Asy-soft
- $\rightarrow E_{svm}$ is larger for Asy-soft case below ρ_0 .

 \rightarrow Thus, isospin effects will be greater in Asy-soft case at sub-normal densities.

! This is also the region of concern for our research work !



Isospin Transport

<u>Isospin Transport</u>

Nucleon current relation:

 $\boldsymbol{j}_{n} - \boldsymbol{j}_{p} = (D_{n}^{\rho} - D_{p}^{\rho}) \nabla \rho - (D_{n}^{I} - D_{p}^{I}) \nabla I$

 $D^{\rho, I} \rightarrow$ transport coefficients $\rho = \rho_n + \rho_p \rightarrow$ nucleon density

 $I = (\rho_{\rm n} - \rho_{\rm p}) / \rho \rightarrow \text{isospin asymmetry}$

First term \rightarrow <u>Isospin drift</u>: Nucleon movement in density gradient ($\nabla \rho$) Second term \rightarrow <u>Isospin diffusion</u>: Nucleon movement in isospin gradient (∇l) D^{ρ} and D^{I} are related to the symmetry energy term (\mathbf{E}_{sym}) of the nuclear equation of state (EoS)

$$D_n^{\rho} - D_p^{\rho} = 4I \frac{\partial E_{sym}}{\partial \rho}$$
 $D_n^I - D_p^I = 4\rho E_{sym}$

Isospin drift and diffusion phenomena are collectively called "isospin transport"

The isospin (N/Z) equilibration of the system is governed by the isospin transport with the help of E_{sym}

The degree of isospin equilibration can be calculated from the isospin transport ratio:

$$R_i = \frac{2X_i^M - X_i^H - X_i^L}{X_i^H - X_i^L} \qquad (X \to \langle N \rangle / Z)$$

Asymmetric nuclear matter (N/Z≠1) with large N/Z can be used to study the isospin transport phenomena

Isospin transport and E_{sym} can be studied to put constraints on the nuclear EoS which is not well understood far from ground state

FAZIA Detector

- \rightarrow FAZIA (<u>F</u>orward-angle <u>A</u> & <u>Z</u> <u>I</u>dentification <u>A</u>rray) is a charged particle multi-detector with an excellent mass resolution
- \rightarrow Basic detection module of FAZIA is called a FAZIA Block
- → Block consists of 16 detection telescopes, each made of two Si layers (300/500 µm and 500 µm) and one CsI scintillator (10 cm)
- \rightarrow Si-Si-CsI telescopes have 2x2 cm² active area and a block covers 0.01 sr at 1 m from target
- \rightarrow Two telescopes connected to one front-end electronics (FEE) card each total 8 FEE cards.
- \rightarrow Block card for output to data acquisition system and input for power supply



 \rightarrow FAZIA collaboration improved the charged particle identification techniques and the detector is now capable to identify charged particles with mass (A) resolution up to charge (Z) ~ 20.

- \rightarrow There is no neutron (N) identification with the detector.
- \rightarrow N can be calculated using N = A-Z up to the A identification limit.
- \rightarrow N/Z can then be further studied !

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 Particle Identification

Systematic analysis of nuclear reactions with a neutron rich projectile on multiple targets at intermediate energies Particle Identification



There can be 3 scenarios for a charged particle entering the detector telescope:

- 1. The particle stops in Si₁
- 2. The particle deposits a fraction of energy in Si₁ and stops in Si₂
- 3. The particle deposits fractions of energy in Si₁ and Si₂ and stops in CsI



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- Simulation & analysis of heavy-ion collision data in 5-100 MeV/nucleon energy range.
- C++ framework based on ROOT.





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FAZIA-PRE experiment

Systematic analysis of nuclear reactions with a neutron rich projectile on multiple targets at intermediate energies



 \rightarrow Analyse the FAZIA-PRE data w.r.t. target mass and beam energy for basic reaction observables such as charge (Z), mass (A), multiplicity of charged particles (M_{tot}) and parallel velocity (v_{II}) and the isospin observable, $\langle N \rangle / Z$.

 \rightarrow Compare the observations of experimental data with HIPSE event generator to determine whether HIPSE can be used to study isospin physics.

FAZIA-PRE Data Analysis

FAZIA-PRE Data Analysis



Z vs v_{\parallel} correlation, beam velocities v_B (black dashed lines) and centre-of-mass velocities v_{CM} (red dashed lines) are marked The identification limit of FAZIA (Z~20) is sufficient to study full range of fragments here: projectile at Z=20 (⁴⁸₂₀Ca)

FAZIA-PRE Data Analysis



FAZIA-PRE Data Analysis



Basic Reaction Observables \rightarrow w.r.t. Target Mass

 \rightarrow For Z \approx 15 (A \approx 30):

Z & A relative yield increases with target mass

Increasing fragmentation with increasing target mass due to increased number of nucleons

 \rightarrow For Z \approx 15 (A \approx 30):

Increase in PLFs and heavier fragment detection with due to increasing elastic scattering events (increasing θ_{gr})

$$\rightarrow$$
 For $M_{tot} = 1$:

Relative yield should also increase with target mass for same reason

 \rightarrow For all data:

 ^{27}Al target data has lower relative yield for Z<6, A<15 and $M_{tot}\!>1$

Possible explanation from *M. Wang et al.*,*CPC (HEP & NP)*, 2012, **36**(12): 1603–201430 :

²⁷Al has lowest n and p separation energies
 Light fragments escape in all directions before reaching the detector placed at forward angles

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FAZIA-PRE Data Analysis



Basic Reaction Observables \rightarrow w.r.t. Beam Energy

 \rightarrow More projectile-like-fragments (Z~20) at 25 MeV/A due to dominance of elastic scattering

 \rightarrow Increased abundance of fragments (Z<20 / A<40) shows dominance of multi-fragmentation at higher beam energy (40 MeV/A)

 \rightarrow Multifragmentation increases the total charged particle multiplicity, thus 40 MeV/A systems have higher M_{tot}

<u>FAZIA-PRE Data Analysis</u>

Isospin effects



 \rightarrow N distribution using N = A - Z

 \rightarrow The trends are similar to the A and Z distributions

 \rightarrow The relative yield of neutrons increases with increasing target mass: more dissipative collisions with increasing target mass.

 \rightarrow Systems at 40 MeV/A have higher relative yield than systems at 25 MeV/A up to Z~16 and lesser for Z>16, as one approaches projectile Z (here, = 20): increased multi-fragmentation at higher beam energy.

 \rightarrow For Z \leq 6, relative yield of ²⁷Al target systems is the least at both beam energies due to its lowest proton and neutron separation energies.

 \rightarrow For Z=20, the relative yield is highest for all systems at N=28, showing the presence of projectile-like fragments (PLFs).

<u>FAZIA-PRE Data Analysis</u>

Isospin Effects → w.r.t. Target Mass



 $\rightarrow \langle N \rangle \! / Z$ plotted from N distributions as a function of Z

 \rightarrow The range of fragment $\langle N \rangle \! / Z$ stays between that of projectile and target

 \rightarrow The fragment $\langle N \rangle /Z$ observed to be decreasing with increasing target mass: with increasing target mass, more dissipative collisions lead to higher rate of isospin equilibration, decreasing the fragment $\langle N \rangle /Z$

 \to Checking the difference between $\langle N \rangle \!/ Z$ of different target systems confirms the observation

 \rightarrow Also confirmed from total N/Z of systems

System	N/Z total
$^{48}\mathrm{Ca}+^{12}\mathrm{C}$	1.31
$^{48}\mathrm{Ca}+^{27}\mathrm{Al}$	1.27
$^{48}\mathrm{Ca}+^{40}\mathrm{Ca}$	1.2

<u>FAZIA-PRE Data Analysis</u>

Isospin Effects → w.r.t. Beam Energy

 \rightarrow For beam energy dependence, the difference between fragment $\langle N\rangle$ /Z from 25 and 40 MeV/A systems was taken:

$$\delta_{\rm E} \langle {\rm N} \rangle / {\rm Z} = \langle {\rm N} \rangle / {\rm Z}_{25} - \langle {\rm N} \rangle / {\rm Z}_{40}$$

 \rightarrow Expected \rightarrow interaction time and nucleon exchange between the participants reduces with increasing beam energy. Thus, an N-rich projectile should produce QP fragments with relatively higher $\langle N \rangle / Z$ at higher beam energy.

 \rightarrow Observed \rightarrow the $\delta_E \langle N \rangle / Z$ is positive for almost all Z: the fragment $\langle N \rangle / Z$ decreases with increasing beam energy.

 \rightarrow Explanation from Phys. Rev. C 55, 1900 (1997):

Rate of charged particle pre-equilibrium emission saturates at $E_B \sim 20$ MeV/A. But pre-equilibrium neutron emission increases with beam energy, thus decreasing overall N/Z of the system and consequently the fragment $\langle N \rangle / Z$.

 $\rightarrow \langle N \rangle\!/Z$ also reduces as a consequence of secondary decays following the isospin equilibration



FAZIA-PRE Data Analysis

Isospin Effects → w.r.t. Beam Energy

 \rightarrow For beam energy dependence, the difference between fragment $\langle N \rangle$ /Z from 25 and 40 MeV/A systems was taken:

$$\delta_{E} \langle N \rangle / Z = \langle N \rangle / Z_{25} - \langle N \rangle / Z_{4}$$

 \rightarrow Expected \rightarrow interaction time and nucleon exchange between the participants reduces with increasing beam energy. Thus, an N-rich projectile should produce QP fragments with relatively higher $\langle N \rangle / Z$ at higher beam energy.

 \rightarrow Observed \rightarrow the $\delta_{_F} \langle N \rangle / Z$ is positive for almost all Z: the fragment $\langle N \rangle / Z$ decreases with increasing beam energy.

 \rightarrow Explanation from Phys. Rev. C 55, 1900 (1997):

Rate of charged particle pre-equilibrium emission saturates at $E_{\rm B} \sim 20$ MeV/A. But pre-equilibrium neutron emission increases with beam energy, thus decreasing overall N/Z of the system and consequently the fragment $\langle N \rangle / Z$.

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FAZIA-PRE Data Analysis

Isospin effects

Looking for Pre-equilibrium and Neck Emission contributions

Pre-equilibrium emissions have very less contributions as the detector is far from target (1 m)

Energetic neck emissions can still reach the detector

<u>FAZIA-PRE Data Analysis</u>

Isospin effects



 $\langle N \rangle \! / Z$ vs $v_{_{||}}$ correlation shows that:

- \rightarrow Fragment $\langle N \rangle \! / Z$ is relatively higher around $v_{_{CM}}$
- \rightarrow Fragment $\langle N \rangle$ /Z is relatively higher around v_B for 25 MeV/A systems. Not seen in 40 MeV/A systems due
- to dominance of multifragmentation

FAZIA-PRE Data Analysis



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FAZIA-PRE Data Analysis



→ The fragments with v_{\parallel} comparable to the v_{B} are emitted in forward direction (FWD) and the fragments with slightly lower v_{\parallel} are emitted in the backward direction (BWD) in the QP reference frame → The "hierarchy effect" shown by neck fragments leads to emission of light fragments (Z~3-6) with lower v_{\parallel} which is comparable to the v_{CM} from the very central dilute region of the neck. → The fragments emitted from denser parts of the neck, i.e., near QP and QT have higher v_{\parallel} than v_{CM} but lower than the BWD or FWD fragments

FAZIA-PRE Data Analysis



To verify, we split the QP phase space for BWD and FWD fragments Split point slightly lower than v_B: E_B = 25 MeV/A → 6 cm/ns ; E_B = 40 MeV/A → 8 cm/ns If v|| ≤ 6/8 cm/ns → fragment is emitted in backward direction in QP phase space → BWD + neck emissions If v|| > 6/8 cm/ns → fragment is emitted in forward direction in QP phase space → FWD

<u>FAZIA-PRE Data Analysis</u>



$$\Delta \langle N \rangle / Z = \langle N \rangle / Z_{BWD} - \langle N \rangle / Z_{FWD} vs Z$$

 \rightarrow Isospin equilibration should have equal effects for BWD and FWD fragment ($\Delta \langle N \rangle / Z \cong 0$)

 $\rightarrow \Delta \langle N \rangle / Z > 0$ for fragments Z \approx 7: Confirming that these lighter fragments in the backward phase space with high $\langle N \rangle / Z$ come from neck

 \rightarrow These fragments with $v_{\parallel} \sim v_{CM}$ have high $\langle N \rangle / Z \rightarrow$ neutron enrichment of the neck region This confirms the observation of Isospin Drift !

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FAZIA-PRE Data Analysis



Isospin diffusion cannot be observed in this data because we also need a ⁴⁰Ca beam to compare the data on the basis of isospin asymmetry

<u>FAZIA-PRE Data Analysis</u>



 $\rightarrow \Delta \langle N \rangle / Z < 0$ for heavier fragments \rightarrow FWD fragments have higher $\langle N \rangle / Z$ due to presence of neutron rich PLFs and heavier fragment

 \rightarrow But the trend changes in the 40 MeV/A systems: Possible reason because of pre-equilibrium neutron emissions from projectile reducing the overall $\langle N \rangle / Z$ in PLFs \rightarrow thus reducing $\langle N \rangle / Z$ of FWD fragments than BWD

HIPSE Event Generator

HIPSE Event Generator

HIPSE (Heavy-Ion Phase-Space Exploration) is a phenomenological event generator able to simulate heavy ion collisions at the intermediate energy range (20-100 MeV/A).

It can easily perform nuclear simulations at all impact parameter ranges and thus is a valuable tool for the understanding of nuclear reaction processes.

LACROIX, VAN LAUWE, AND DURAND

not ensure a general agreement with the data collected in the intermediate energy range. Indeed, the free parameters of the model have been adjusted with help of a limited set of data and additional studies leading to the same values of these parameters are necessary to confirm the validity of our scenario.

In the near future, we plan to extend the model to mass asymmetric systems. Note also that the calculation gives the final charge to mass ratio of fragments created during the reaction and could be a valuable tool to explore the N/Z PHYSICAL REVIEW C 69, 054604 (2004)

effects in nuclear collisions in the Fermi-energy range.

Finally, it is worth noting that the technique of exploration of phase space described in this paper is flexible enough to test rapidly other (possibly thermalized) nucleon momentum distributions to describe nuclear fragmentation.

ACKNOWLEDGMENT

We thank warmly INDRA Collaboration for permission to use its data.

It says here that HIPSE "could be a valuable tool" for N/Z studies but it has never been directly checked !!

HIPSE Event Generator



Initialisation of n and p position and momentum Thomas-Fermi distribution for each nucleus in the centre-of-mass frame

distribution among fragments

energy

SIMON or GEMINI++

using

HIPSE Event Generator



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HIPSE Simulations



HIPSE Simulations



HIPSE Simulations

HIPSE full 4π simulations cannot be directly compared to the experimental data due to experimental constraints HIPSE data is filtered w.r.t. the FAZIA-PRE experimental constraints for comparison



HIPSE Simulations



HIPSE Simulations



Data Comparison

Basic reaction observables



 \rightarrow Fragment Charge (Z), Fragment Mass (A), Multiplicity of Charged Particles (M_{tot}) and Parallel Velocity (v_{||}) are reasonably well reproduced by HIPSE \rightarrow A qualitative agreement is observed for all systems between

HIPSE and experimental data (EXP)



Dominance of elastic scattering events not seen in HIPSE due to choice of b_{max} only up to semi-peripheral region

Data Comparison



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Data Comparison



Data Comparison



Data Comparison



Summary and conclusions

 \rightarrow Heavy-ion collisions at intermediate energies can help to study isospin transport in peripheral collisions.

 \rightarrow FAZIA is a charged particle multi-detector with an excellent mass resolution of up to Z~25.

 \rightarrow The research work presented here mainly focused on the analysis of the data from the FAZIA-PRE experiment performed using 6 FAZIA blocks at LNS-INFN, Catania, Italy in February 2018.

 \rightarrow ⁴⁸Ca projectile was bombarded on ¹²C, ²⁷Al and ⁴⁰Ca targets at 25 MeV/A and on ¹²C and ²⁷Al targets at 40 MeV/A.

 \rightarrow The aim of the experiment was to study the effects on the fragment N/Z due to initial neutron abundance and pre-equilibrium neutron emissions from the projectile.

 \rightarrow Mainly QP fragments detected as the detector setup had an angular acceptance of ~ 2° - 18°.

 \rightarrow Dependence on the target mass and beam energy for various basic reaction observables such as charge (Z), mass (A), multiplicity of charged particles (M_{tot}), parallel velocity (v₁) and isospin observable $\langle N \rangle / Z$ was observed.

 \rightarrow A full range of N distributions for Z=3-20 was obtained for all reaction systems from N=A-Z

 \rightarrow The fragment $\langle N \rangle \! / Z$ was found to be decreasing with increasing target mass.

 \rightarrow The fragment $\langle N \rangle \! / Z$ was found to be decreasing with increasing beam energy.

 \rightarrow Isospin drift is also observed in the FAZIA-PRE experimental data

Conclusions from comparison

 \rightarrow A detailed study of HIPSE simulations with SIMON and GEMINI++ de-excitation codes in comparison with FAZIA-PRE experimental data for 5 different systems has lead to the following conclusions:

 \rightarrow No dependence on the mass of the target is observed in HIPSE data.

 \rightarrow Effects of beam energy dependence on fragments is observed in HIPSE data.

 \rightarrow HIPSE gives an overall qualitative reproduction of the experimental data with respect to basic reaction observables (Z, A, M_{tot} and v_{\parallel}).

 \rightarrow A low quantitative agreement is observed.

 \rightarrow HIPSE is able to produce the fragment $\langle N \rangle / Z$ with a reasonable agreement to the experimental data. But it cannot give a detailed structure of the fragment $\langle N \rangle / Z$.

HIPSE event generator is a fine tool to generate the fragments from nuclear reactions at intermediate energy range with respect to the basic reaction observables and to have an idea of the overall dynamics. But due to its inability to produce a detailed structure of fragment N/Z, it is not advisable to use HIPSE to study N/Z effects. One would rather continue to use Antisymmetrized Molecular Dynamics (AMD) model or Constrained Molecular Dynamics (CoMD) model to investigate isospin and symmetry energy related studies.





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