



Neutrino Physics

Lecture 2

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Outline

- Neutrino oscillations - experimental aspects
- Reactor neutrinos
- Accelerator neutrinos
- Future experiments

Oscillations - experimental aspects

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} - 4 \sum_{i < j} \Re(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \cdot \sin^2 \Phi_{ij} + 2 \sum_{i < j} \Im(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \cdot \sin^2 \Phi_{ij}$$
$$\Phi_{ij} = \Delta m_{ij}^2 \frac{L}{4E_\nu} = 1.27 \cdot \Delta m_{ij}^2 [eV^2] \cdot \frac{L [km]}{E_\nu [GeV]}$$

- In the 3 neutrino case - transition probabilities $P(\nu_\alpha \rightarrow \nu_\beta)$ depend on:
 - Mixing angles: θ_{23} , θ_{13} , θ_{12}
 - Complex phase: δ_{CP}
 - Independent mass splittings: Δm_{32}^2 , Δm_{12}^2
 - Detector-source distance (L), neutrino energy (E) - chosen experimentally
- Formulas above are valid for vacuum oscillations. Matter effects related to the interactions of electron neutrinos with matter modify the mixing angles and mass splittings giving effective angles and splittings (not discussed here).

Oscillations - experimental aspects

- Neutrino oscillation experiments:
 - Appearance probability: $P(\nu_\alpha \rightarrow \nu_\beta) = ?$
 - Disappearance (survival) probability: $P(\nu_\alpha \rightarrow \nu_\alpha) = ?$
- Number of observed (detected) neutrino interactions is proportional to the neutrino-target (eg. neutrino-nucleus) interaction cross section, neutrino flux and number of interacting targets (eg. number of target nuclei)

$$N_{Obs} \approx \sigma * \Phi * T$$

Neutrino oscillation experiments

$$\sin^2(\theta_{12}) = 0.307 \pm 0.013$$

$$\Delta m_{21}^2 = (7.53 \pm 0.18) \times 10^{-5} \text{ eV}^2$$

$$\sin^2(\theta_{23}) = 0.547 \pm 0.021 \quad (\text{Inverted order})$$

$$\sin^2(\theta_{23}) = 0.545 \pm 0.021 \quad (\text{Normal order})$$

$$\Delta m_{32}^2 = (-2.546^{+0.034}_{-0.040}) \times 10^{-3} \text{ eV}^2 \quad (\text{Inverted order})$$

$$\Delta m_{32}^2 = (2.453 \pm 0.034) \times 10^{-3} \text{ eV}^2 \quad (\text{Normal order})$$

$$\sin^2(\theta_{13}) = (2.18 \pm 0.07) \times 10^{-2}$$

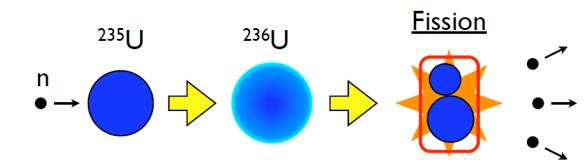
δ , CP violating phase = $1.36 \pm 0.17 \pi$ rad

- Experiments that contributed to these results:

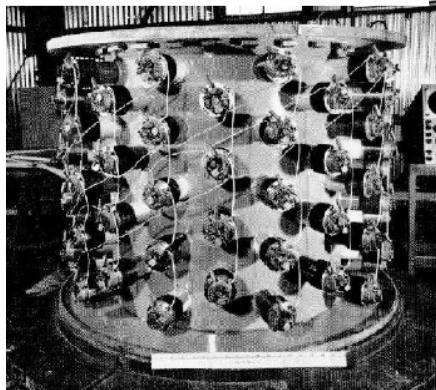
Experiment	Dominant	Important
Solar Experiments		
Reactor LBL (KamLAND)	θ_{12}	$\Delta m_{21}^2, \theta_{13}$
Reactor MBL (Daya-Bay, Reno, D-Chooz)	Δm_{21}^2	θ_{12}, θ_{13}
Atmospheric Experiments (SK, IC-DC)	$\theta_{13}, \Delta m_{31,32}^2 $	$\theta_{23}, \Delta m_{31,32}^2 , \theta_{13}, \delta_{\text{CP}}$
Accel LBL $\nu_\mu, \bar{\nu}_\mu$, Disapp (K2K, MINOS, T2K, NO ν A)	$ \Delta m_{31,32}^2 , \theta_{23}$	
Accel LBL $\nu_e, \bar{\nu}_e$ App (MINOS, T2K, NO ν A)	δ_{CP}	θ_{13}, θ_{23}

Reactor neutrinos

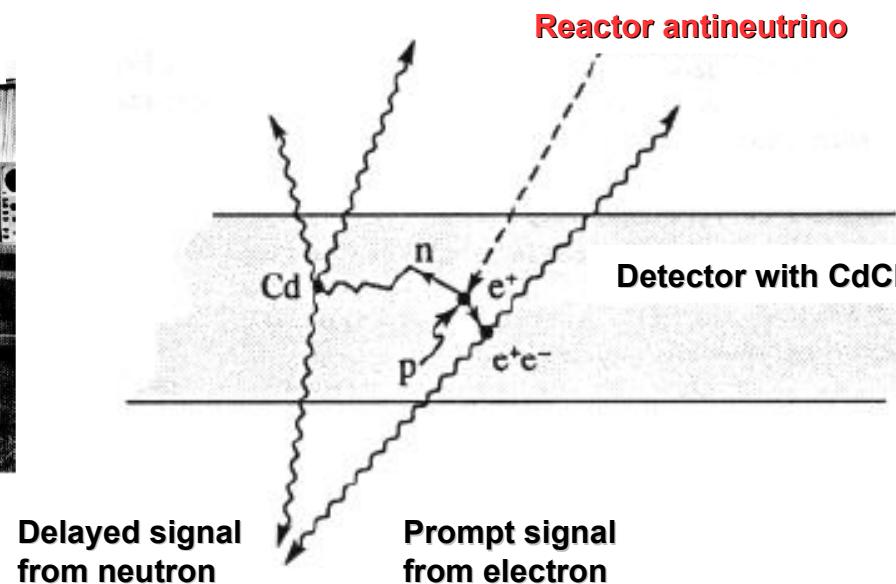
- In the nuclear reactors antineutrinos are produced as a result of the nuclear fission of heavy isotopes ^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu
- 6 antineutrinos are produced per one fission reaction
- Typical nuclear powerplant (1 GW reactor) produces 2×10^{20} antineutrinos per second
- Number of detected antineutrinos is proportional to the reactor power
- There's a long tradition of using reactor antineutrinos to study neutrino properties. Started in 1953 by Reines & Cowan (detector next to the Savannah River powerplant)



Reines & Cowan detector



Kraków, 25.02.2021



Frederick Reines & Clyde Cowan



Nobel prize for Reines in 1995

Reactor neutrinos

- Reactor experiments are measuring the probability of disappearance of electron antineutrinos from the reactor.

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = ?$$

- A precise knowledge of the antineutrino flux from the nuclear reactor is crucial because we need to compare the measured neutrino spectrum with the predicted one.
- Detailed calculation of the antineutrino flux from the nuclear reactor is challenging (summing up the spectra of beta decays)
 - Fission of four main isotopes involves thousands(!) of beta decay branches → The main problem of reactor neutrino experiments.
 - Many improvements in the reactor neutrino flux calculations recently.
- Currently operating reactor neutrino experiments:

Name	Reactor power (GW _{th})	Baseline (km)	Detector mass (t)	Year
KamLAND	various	180 (ave.)	1,000	2001–
Double Chooz	4.25×2	1.05	8.3	2011–2018
Daya Bay	2.9×6	1.65	20×4	2011–
RENO	2.8×6	1.38	16	2011–
JUNO	26.6 (total)	53	20,000	

Reactor neutrinos - KamLAND result

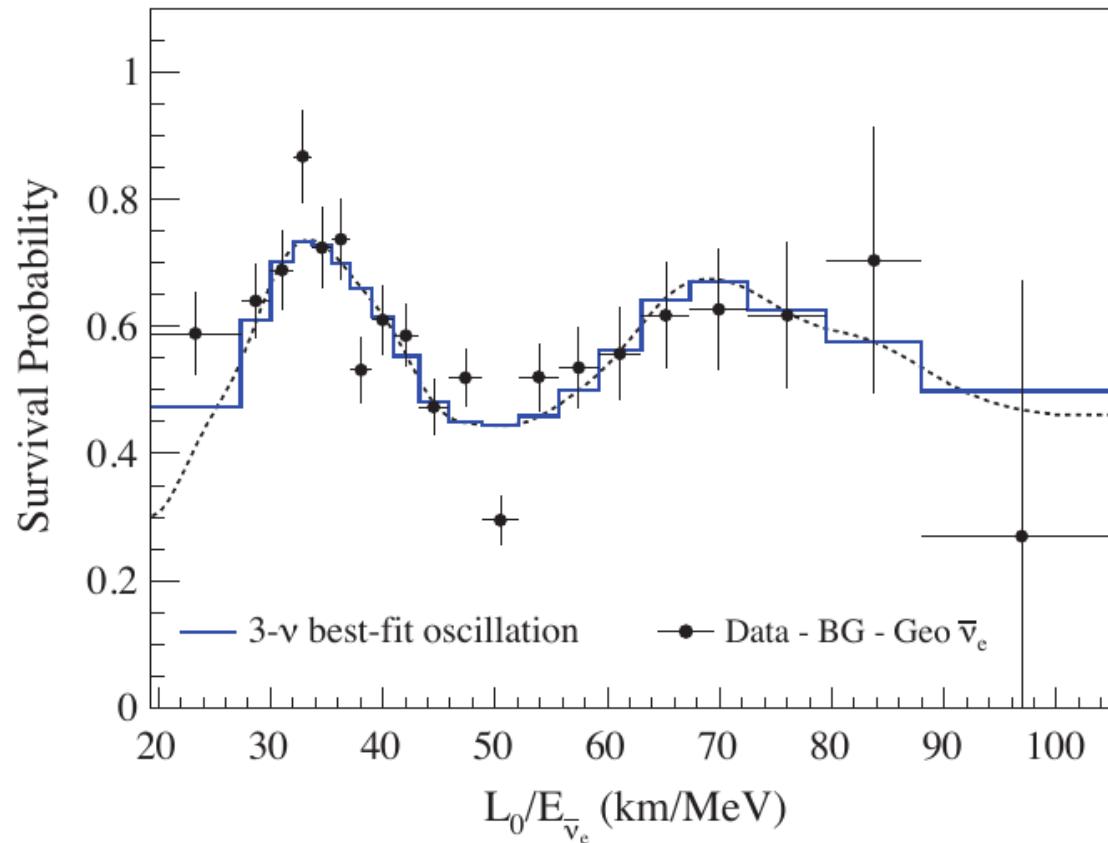
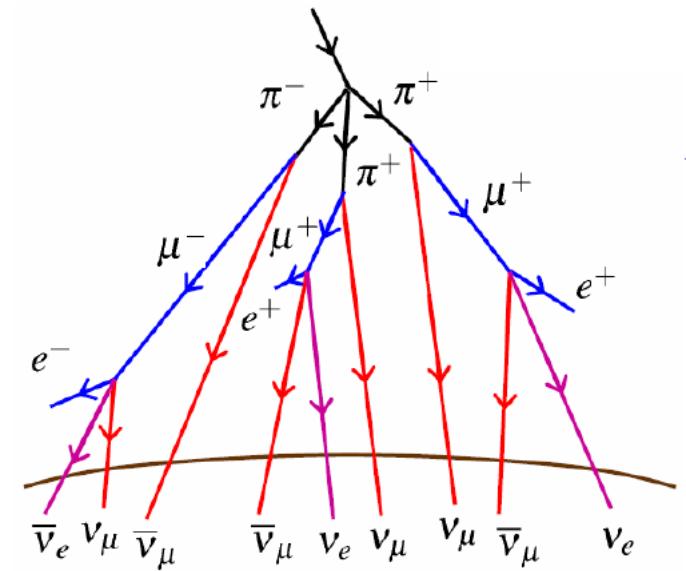


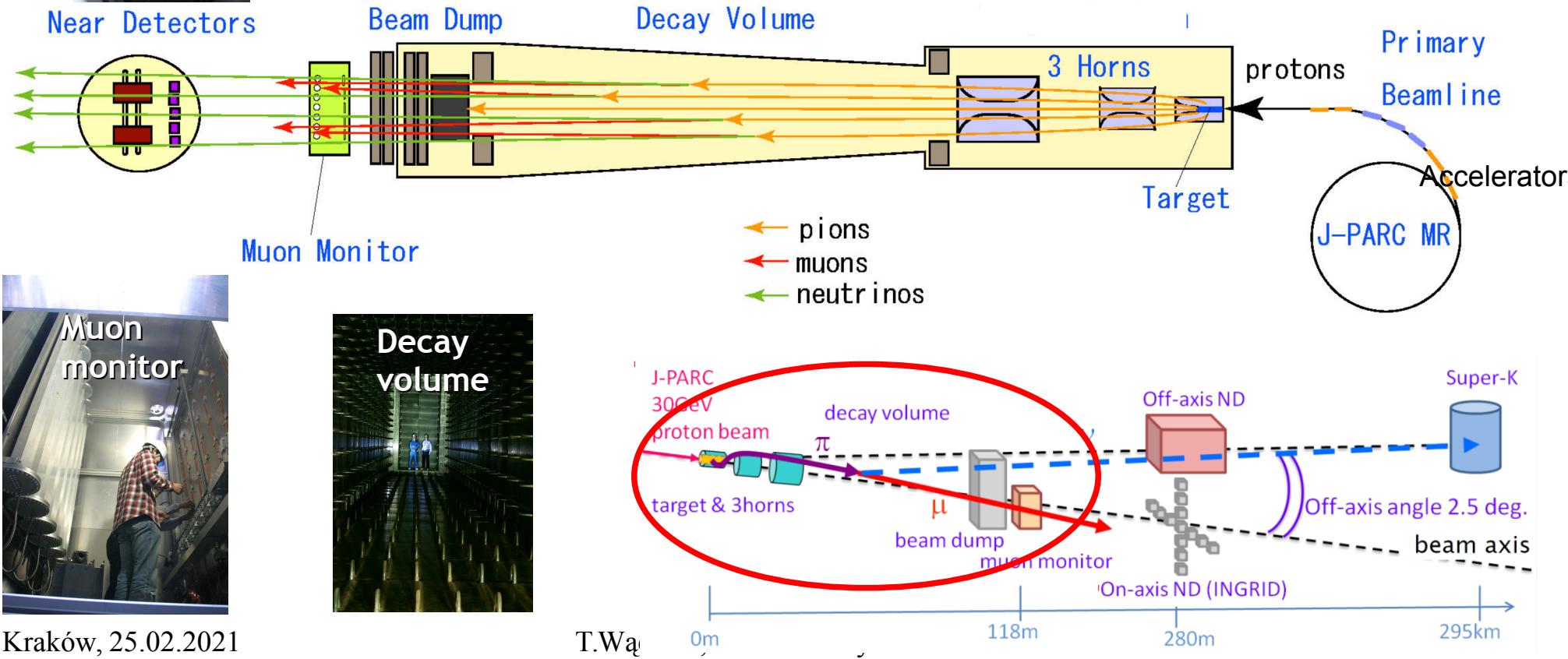
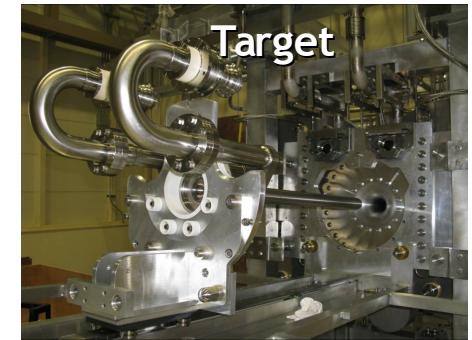
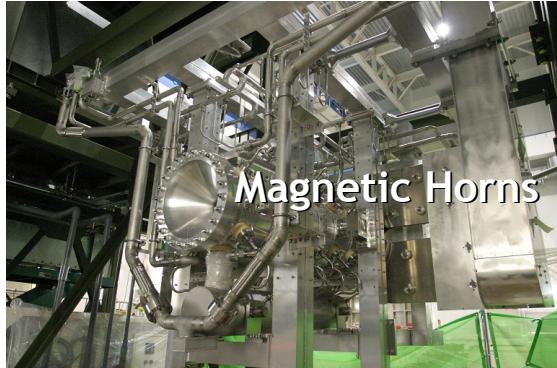
Figure 14.7: Ratio of the observed $\bar{\nu}_e$ spectrum to the expectation for no-oscillation versus L_0/E for the KamLAND data. $L_0 = 180$ km is the flux-weighted average reactor baseline. The 3- ν histogram is the best-fit survival probability curve from the three-flavour unbinned maximum-likelihood analysis using only the KamLAND data. This figure is taken from [150].

Accelerator neutrinos

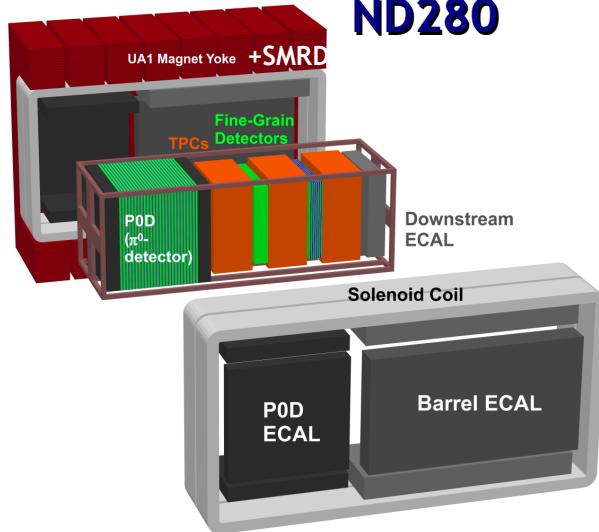
- The main principle behind the production accelerator neutrinos is very similar to atmospheric neutrino production mechanism
 - Accelerate protons to high energies with the accelerator (~50 GeV)
 - Smash protons into the target (eg. graphite) and produce secondary particles, mainly pions
 - Let the pions decay and produce neutrinos
- Extra ideas in producing accelerator neutrinos:
 - Focus charged pions with the same sign (and deflect pions with opposite sign) with magnetic horns (toroidal magnetic field).
 - ✓ Ultimately able to get either neutrino (positive pions are focused) or antineutrino (negative pions are focused) beam.
 - ✓ Can change the polarity of the horns → one experiment can operate in two modes: neutrino or antineutrino mode
 - Stop muons produced in the charged pion decay using a block of graphite and iron → beam dump.



Accelerator neutrinos



Near detector example (T2K)

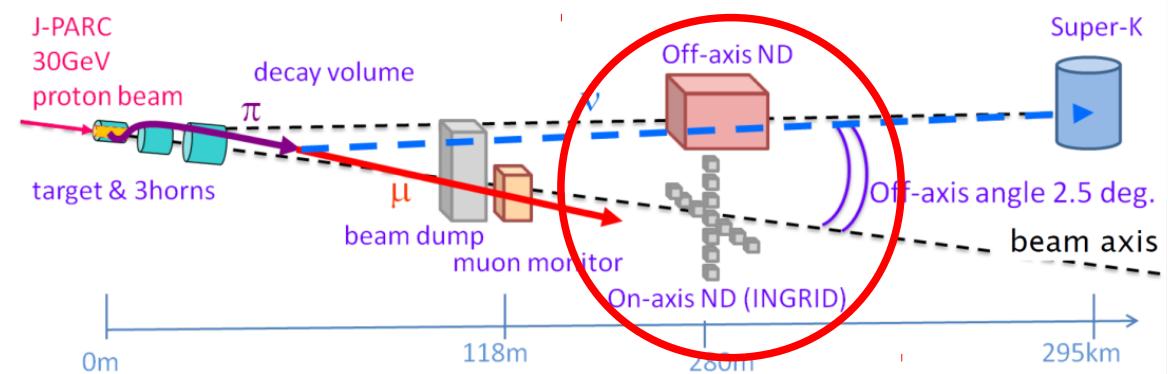
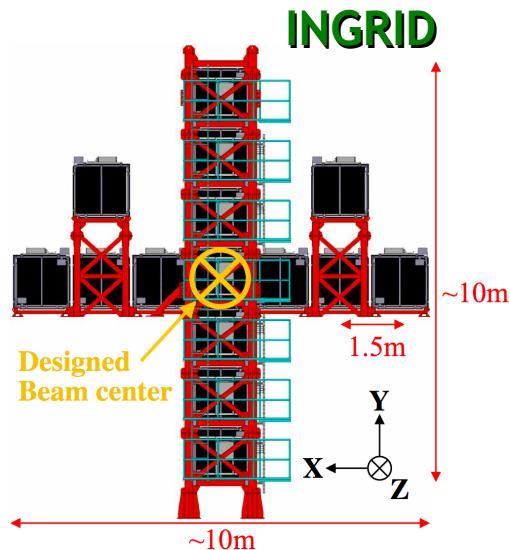


- Off-axis detector:

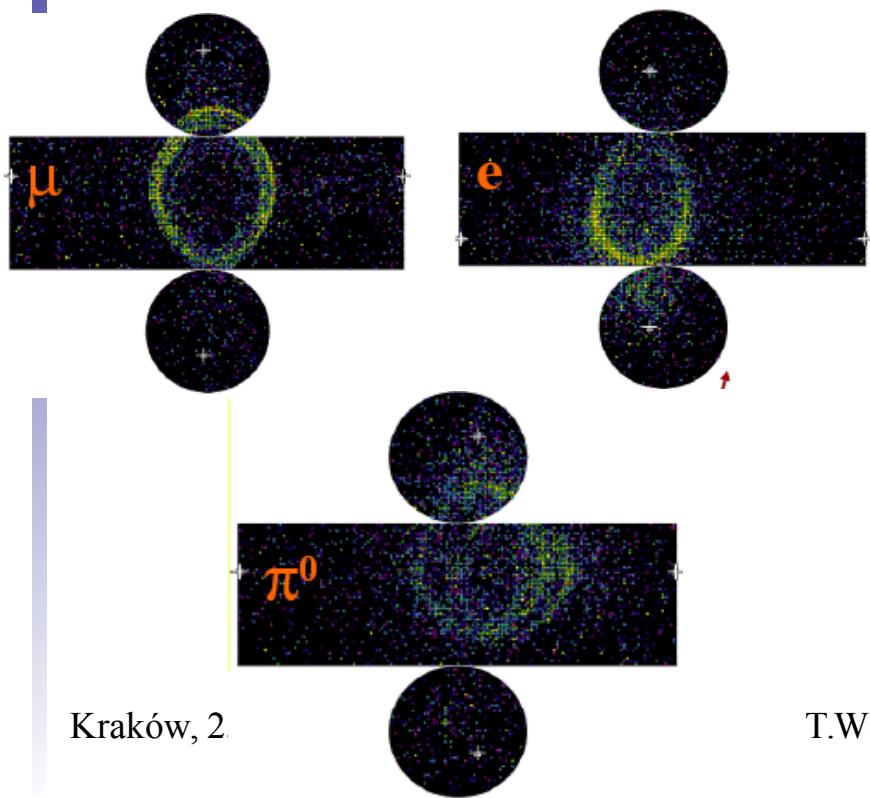
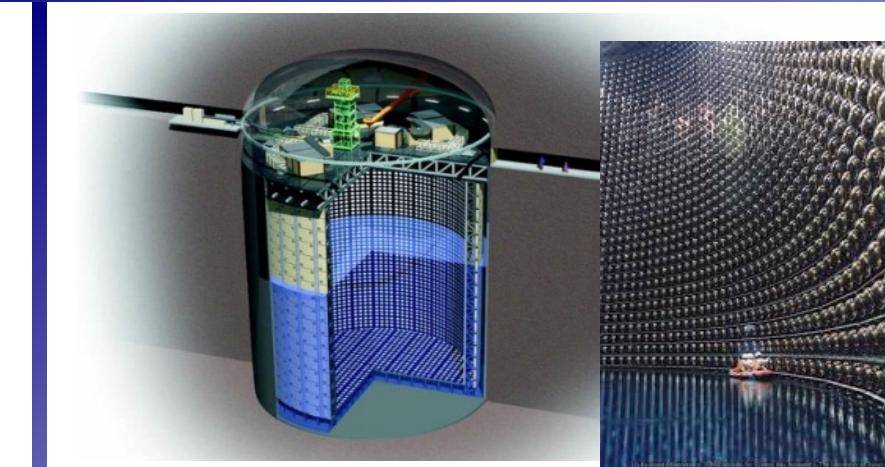
- Several sub-detectors in 0.2T magnetic field:
 - ✓ Tracker (TPC + FGD), pizero detector (POD), electromagnetic calorimenter (ECAL), muon ranger (SMRD)
- Measures the neutrino flux before the oscillations occur
- Measures intrinsic ν_e contamination
- Measures neutrino interaction cross sections

- On-axis detector (INGRID):

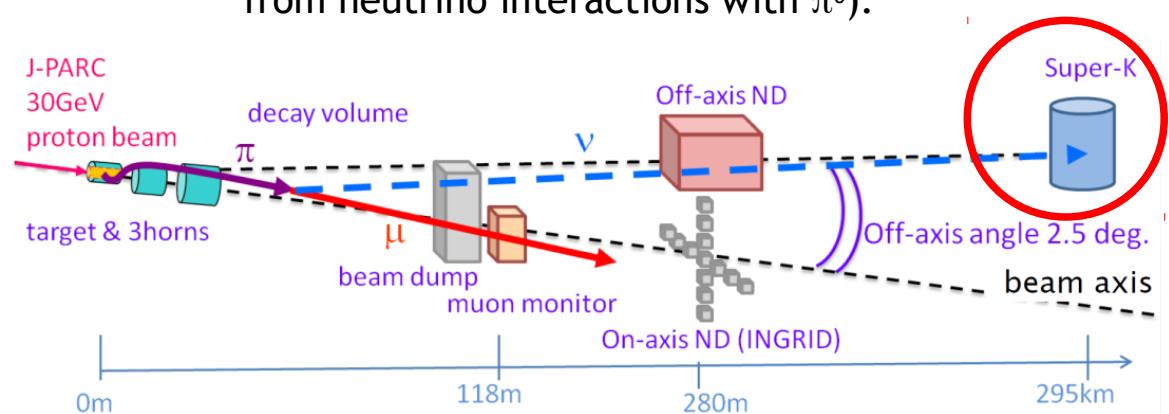
- 16 iron-scintillator modules form the cross
- Monitoring flux, direction and stability of the neutrino beam
- Neutrino cross section measurements



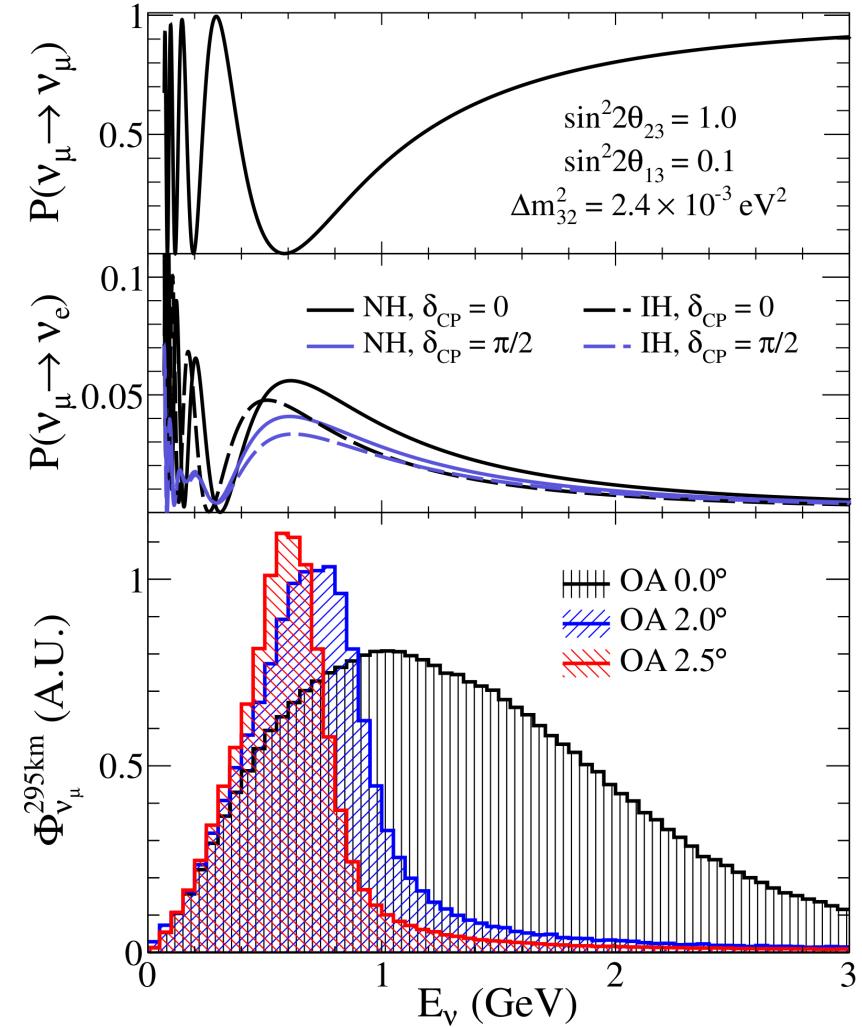
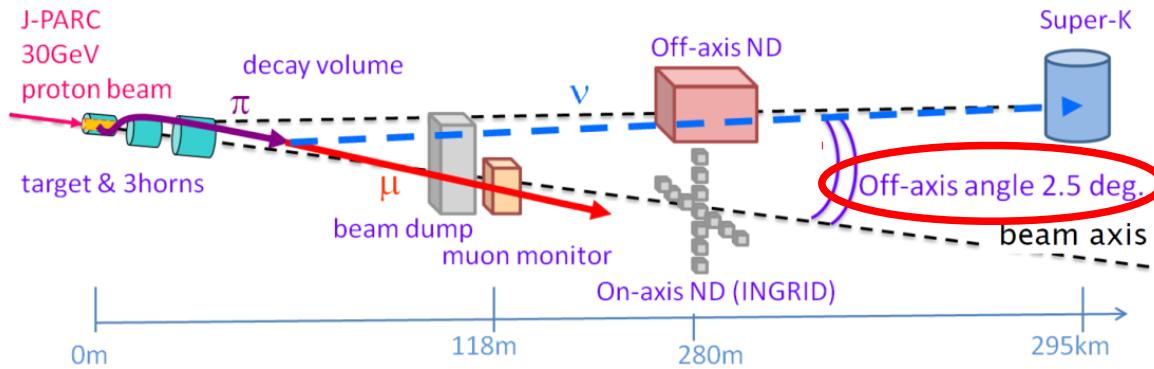
Far detector (T2K)



- Super-Kamiokande (operating since 1996):
 - Water Cherenkov (50 kt, 22.5 kt fiducial volume).
 - 11 000 (inner) + 2000 (outer) photomultipliers
 - Neutrino energy resolution ~10%
 - Particle identification:
 - ✓ Good electron-muon discrimination (<1% muons identified as electrons)
 - ✓ Neutral pion detection (rejecting background from neutrino interactions with π^0).



Off-axis beam



- Currently two world leading accelerator experiments (T2K, NovA) use 'off-axis' beam idea:
 - Pion decay kinematic effect
 - Thin energy spectrum with the mean energy tuned to the neutrino oscillation probability
 - Lower background from high energy interactions that are difficult to reconstruct

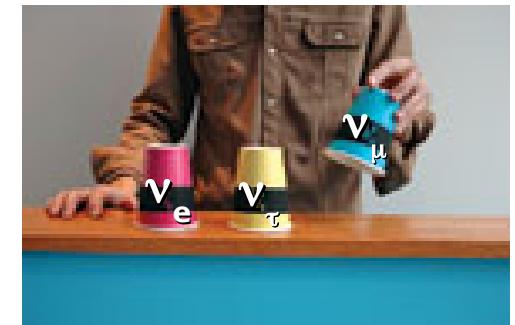
Accelerator neutrino oscillations

Disappearance of muon neutrinos/antineutrinos from the beam

(ν_μ /anti- ν_μ disappearance)

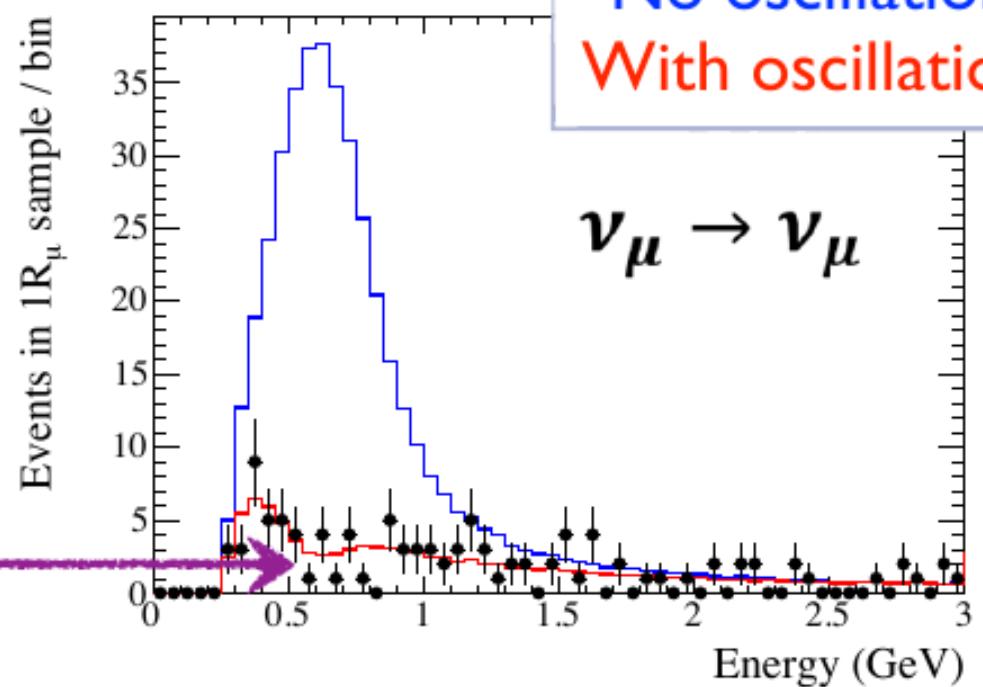
$$P(\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu) \simeq 1 - 4 \cos^2 \theta_{13} \sin^2 \theta_{23}$$
$$\times [1 - \cos^2 \theta_{13} \sin^2 \theta_{23}] \sin^2 \frac{\Delta m_{32}^2 L}{4E}$$

+ (solar, matter effect terms)



No oscillation
With oscillation

$\nu_\mu \rightarrow \nu_\mu$



Location of min: Δm_{32}^2
Depth of min: $\sin^2 2\theta_{23}$

Acelerator neutrino oscillations

Electron neutrino/antineutrino appearance in the muon neutrino/antineutrino beam (ν_e /anti- ν_e appearance)

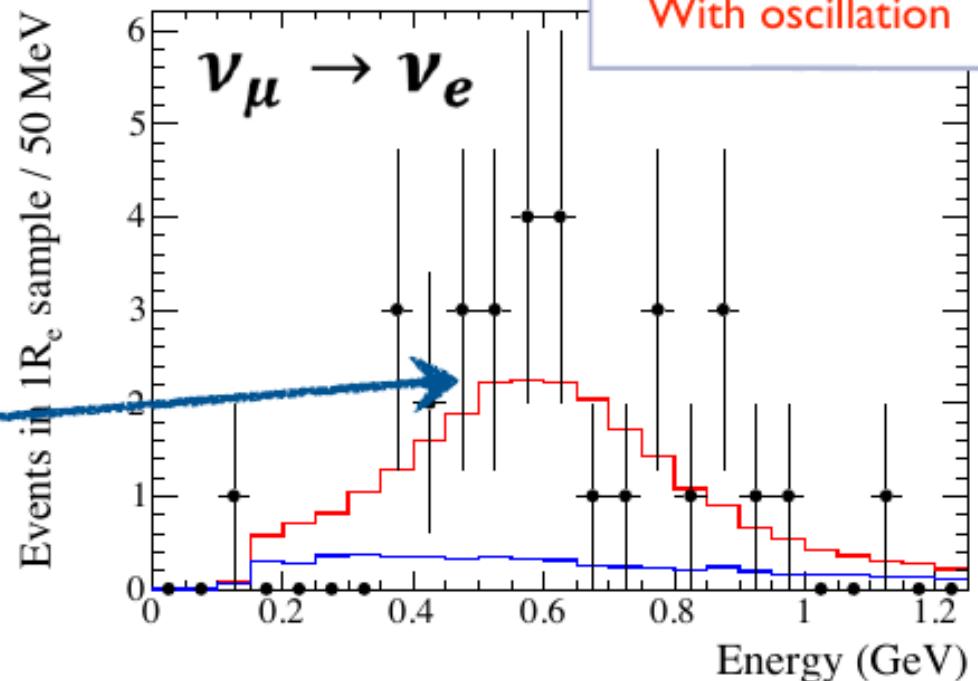
$$P(\overleftarrow{\nu_\mu} \rightarrow \overleftarrow{\nu_e}) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{32}^2 L}{4E}$$

$$(+)- \left[\sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \right.$$

$$\times \sin \frac{\Delta m_{21}^2 L}{4E} \sin^2 \frac{\Delta m_{32}^2 L}{4E} \sin \delta_{CP} \left. \right]$$

+ (CP-even, solar, matter effect terms)

Magnitude of the peak
 $\sin^2 \theta_{23}, \sin^2 2\theta_{13}, \delta_{CP}$



Accelerator neutrino oscillations

CPV: $P(\nu_\mu \rightarrow \nu_e) \neq P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$

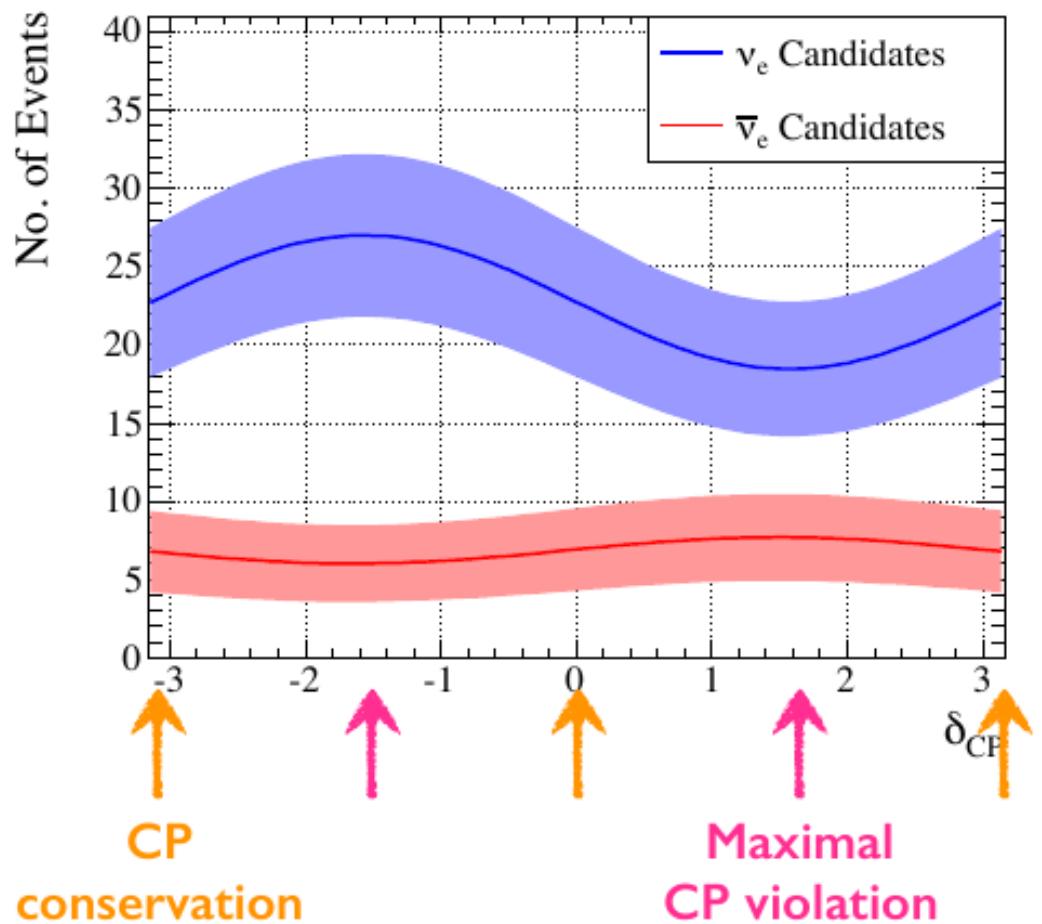
$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{32}^2 L}{4E}$$
$$(+)- \left[\sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \right.$$
$$\times \sin \frac{\Delta m_{21}^2 L}{4E} \sin^2 \frac{\Delta m_{32}^2 L}{4E} \sin \delta_{CP}]$$

+ (CP-even, solar, matter effect terms)

$\sin \delta_{CP}$ occurs in ν_e and $\bar{\nu}_e$ appearance probability with opposite sign

Complicated measurement because the sensitivity to measure δ_{CP} depends on:

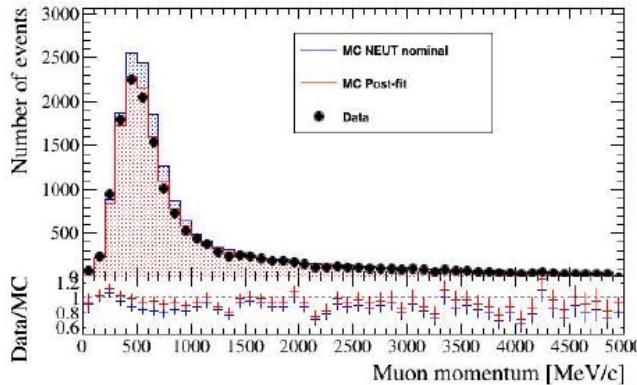
δ_{CP} true value, θ_{23} true value, mass ordering



Accelerator neutrino oscillation analysis

Neutrino flux model:
Monte Carlo simulations + external data from other experiments (eg. NA61)

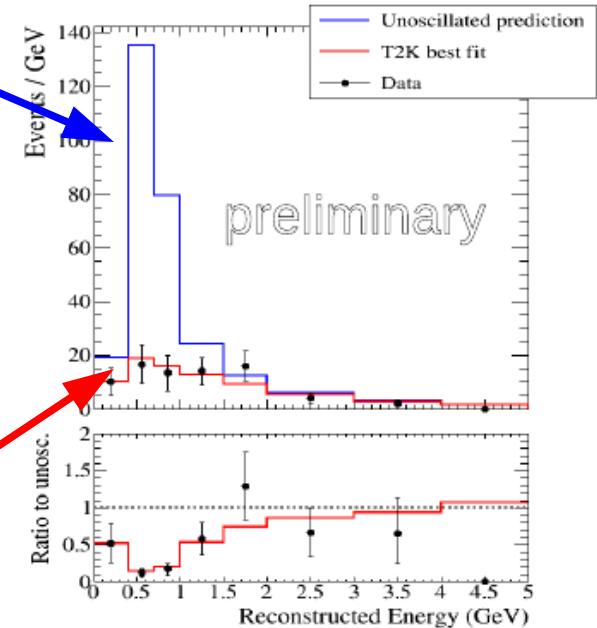
Neutrino interaction models and their uncertainties:
Monte Carlo simulations + external data (MINERvA, MiniBooNE experiments)



From near detector fit:
Flux model parameters,
Neutrino interaction model parameters,
Backgrounds in far detector

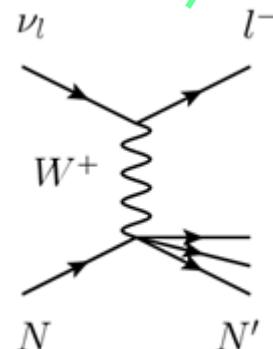
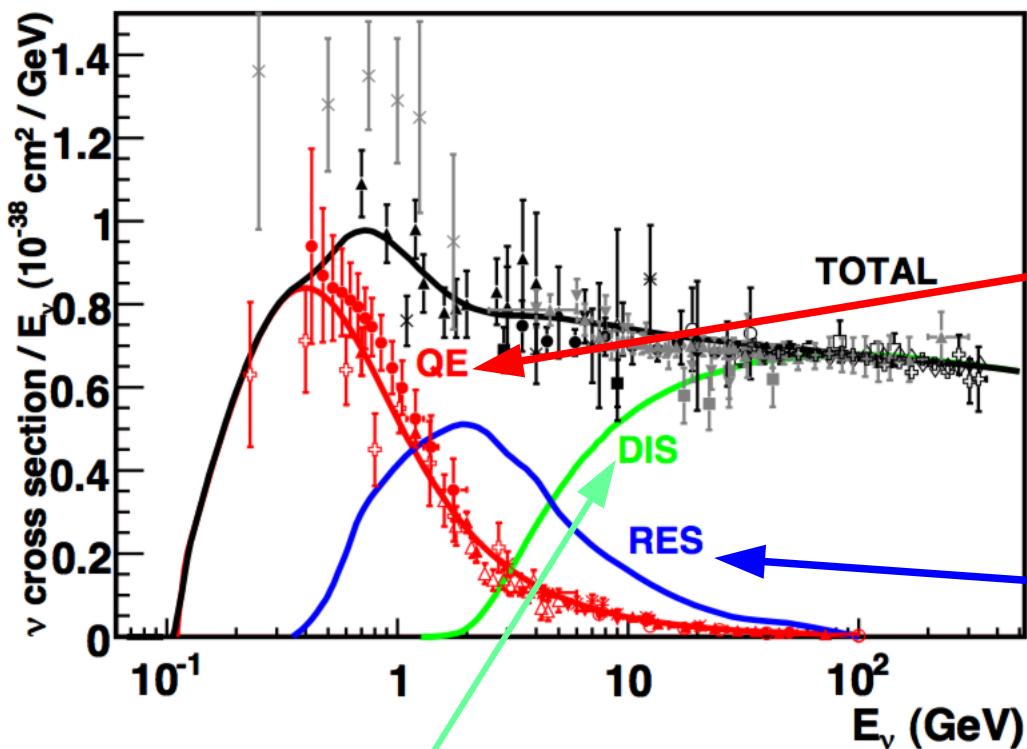
Fit to the data
from the near detector

Expected neutrino
spectrum
In far detector



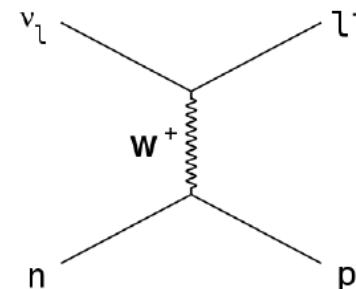
Neutrino interactions (accelerator neutrinos)

Charged-current interactions

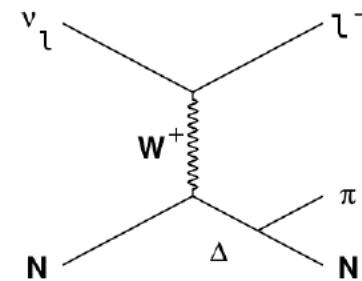


Deep Inelastic Scattering (DIS)
(NOvA)

Quasi-elastic scattering - CCQE
(dominant in T2K)



Resonance pion production - RES
(dominant in NOvA)



- + additional complications:
 - Nuclear 'initial-state' models
 - Nuclear re-interactions models (Final State Interactions)
 - ...

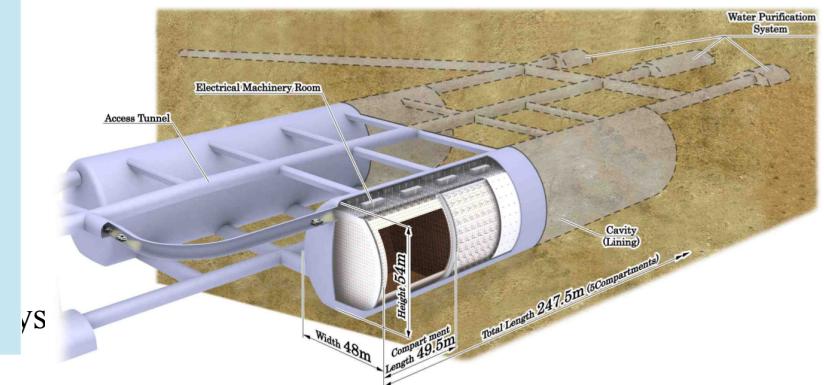
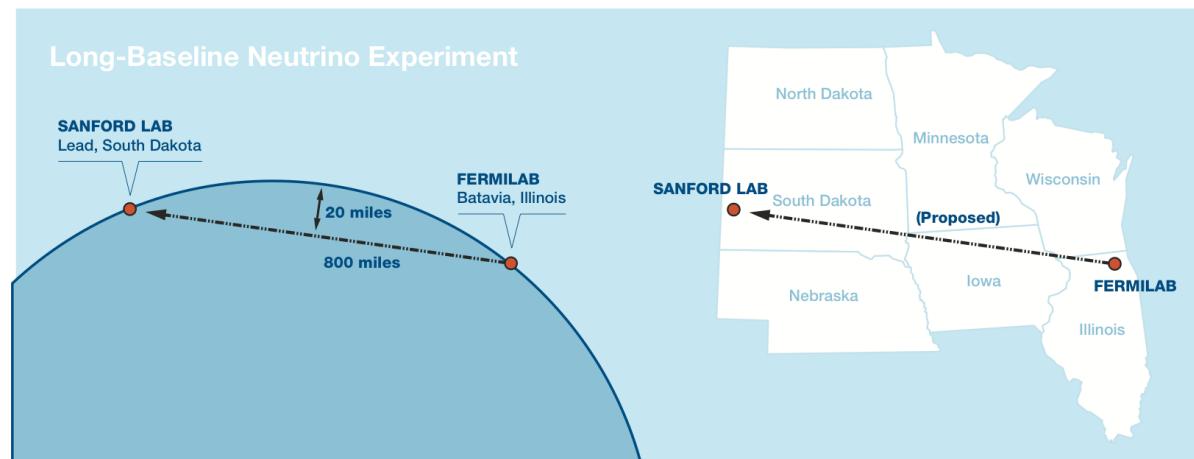
Accelerator neutrino experiments

Name	Beamline	Far Detector	L (km)	E_ν (GeV)	Year
K2K	KEK-PS	Water Cherenkov	250	1.3	1999–2004
MINOS	NuMI	Iron-scintillator	735	3	2005–2013
MINOS+	NuMI	Iron-scintillator	735	7	2013–2016
OPERA	CNGS	Emulsion	730	17	2008–2012
ICARUS	CNGS	Liquid argon TPC	730	17	2010–2012
T2K	J-PARC	Water Cherenkov	295	0.6	2010–
NOvA	NuMI	Liquid scint. tracking calorimeter	810	2	2014–

- Accelerator neutrino experiments and their major results:
 - K2K: confirming atmospheric neutrino oscillations reported by Super-Kamiokande
 - MINOS, MINOS+ : measuring atmospheric neutrino oscillations m223, 23
 - OPERA: confirming oscillations
 - ICARUS: first detector in liquid argon TPC technique
 - T2K: 13 angle measurement, m223 measurement, first measurement of CP phase
 - NOvA: 13 angle measurement, CP phase measurement

Future accelerator neutrino experiments

- LBNE (USA)
 - US flagship accelerator project.
 - Neutrino beam from Fermilab (Chicago)
 - 10 kton liquid argon detector in South Dakota (1300 km from Fermilab).
 - Should be able to measure δ_{CP} with 20 degree accuracy (10 degrees after beam upgrade).
- Hyper-Kamiokande (Japonia)
 - T2K experiment extension
 - High-intensity neutrino beam (1.7 MW) from J-PARC complex
 - Far detector - 0.5 kton water Cherenkov 300 km from J-PARC.
 - Measuring δ_{CP} with accuracy of 18 degrees but depends on mass ordering measurements.



Summary & prospects

- Covered only a small part of neutrino physics:
 - Neutrino sources
 - ✓ Nautural
 - ✓ Artificial
 - Oscillations
- For many other interesting topics such as:
 - Sterile neutrinos
 - Neutrino masses (theory)
 - Absolute neutrino mass measurements
 - Neutrinoless double-beta decay
 - Neutrino interactions
 - History of neutrino discoveries
 - ...