

Heavy Ion Collision Physics

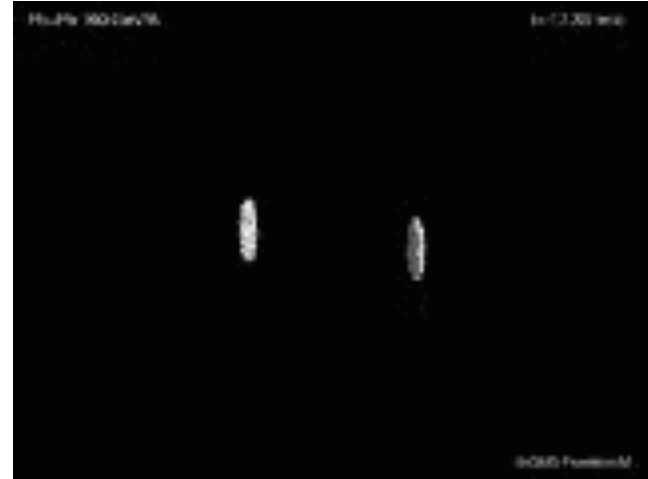
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28th January 2021

KISD, Lecture: Particle physics for specialists



- The main goal is to understand the dynamics of **dense and hot medium** created in heavy-ion collisions
- Quantum Chromo-dynamics (QCD), a non-abelian $SU(3)_{\text{color}}$ gauge theory describing strong interactions between quarks and gluons

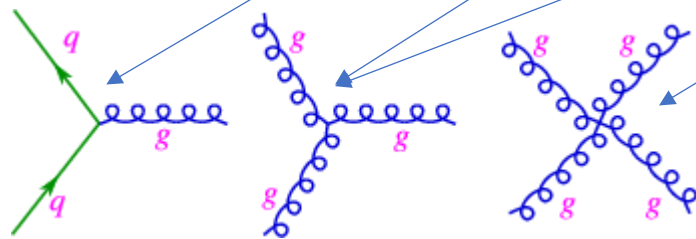
Quantum Chromo-Dynamics (QCD)

Theory describing strong interactions between quarks and gluons

- **Quarks:** $\psi_i^f(x)$ spin-1/2, Dirac fermions
 color triplet: $i = 1, 2, 3$ (green, red, blue)
 flavour: $f = u, d, s, c, b, t$ (*fractional electric charge*)
- **Gluons:** $A_{\mu,a}(x)$ spin-1
 color octet: $a = 1, \dots, 8$
- **QCD Lagrangian:**

$$\mathcal{L}_{QCD}(\psi, A) = \sum_f \bar{\psi}_i^f [(i\partial_\mu \delta_{ij} - gA_{\mu,a}(t_a)_{ij})\gamma^\mu - m_f \delta_{ij}] \psi_j^f - \frac{1}{4} [\partial_\mu A_{\nu,a} - \partial_\nu A_{\mu,a} - gC_{abc}A_{\mu,b}A_{\nu,c}]^2$$

- **Interactions:**



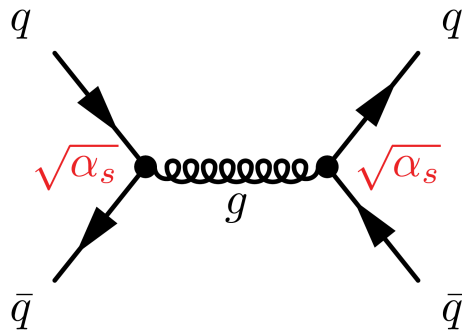
“running” coupling constant
 $\alpha_s = \frac{g^2}{4\pi}$

Asymptotic Freedom and Confinement

- High Q : **Asymptotic Freedom**

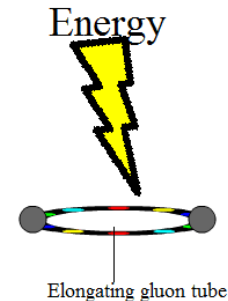
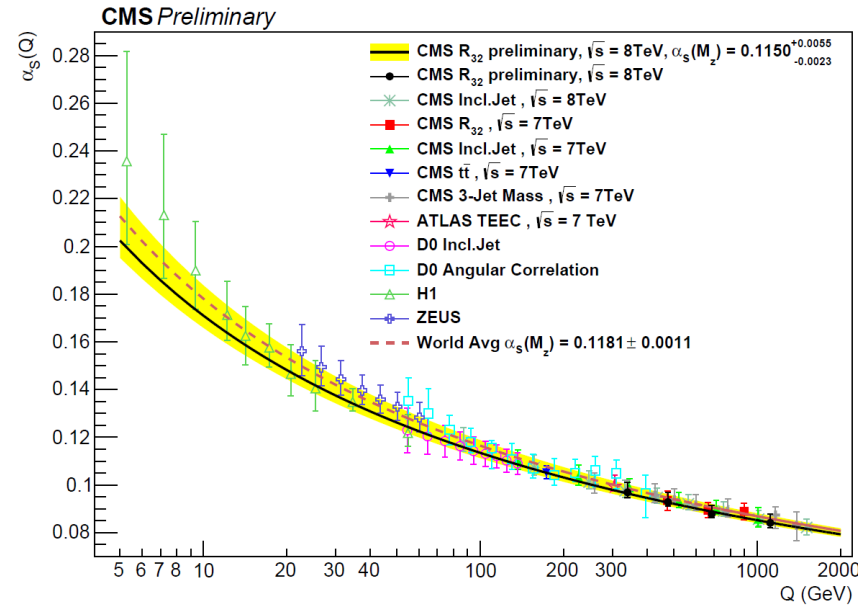
1973 – D. Gross, F. Wilczek, D. Politzer

- α_s decreases for $r \searrow 0$ (anti-screening)
- Perturbative expansion works $\alpha_s \ll 1$



- Low Q: **Confinement**

- α_s increases for $r \nearrow \text{inf}$
- Importance of non-perturbative QCD effects
 - Breakdown of perturbative calculations
- String of gluons between quarks



wikipedia

$$V(r) \sim -\alpha/r + \sigma r$$

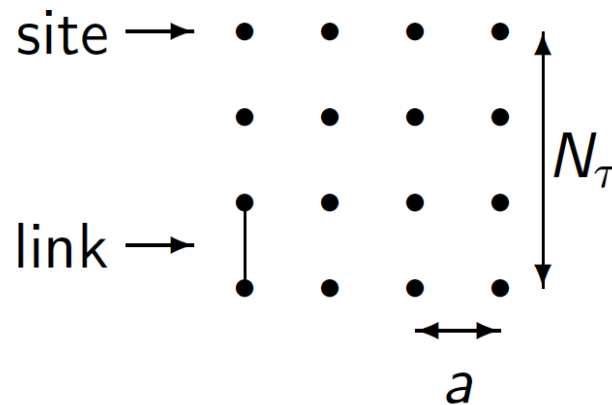
A Historical Remark

- **1973**: The beginning of QCD as a theory of strong interactions
D. Gross, F. Wilczek, D. Politzer
 - Quark Model + Yang-Mills gauge theory
- **1975**: Prediction of a deconfined phase
 - J.C. Collins and M.J. Perry, *Superdense Matter or Asymptotically Free Quarks?*, PRL **34**, 1353

The deconfined phase of quark and gluons, called "quark soup", was later called "quark-gluon plasma (QGP)" due to analogies to similar phenomena in other physics branches

QCD Lattice - Ab-initio calculations

- In QCD, the coupling constant is large at the low energy scales, hence non-perturbative approach is used
- QCD is formulated at a discrete space-time grid $N_s^3 \times N_\tau$

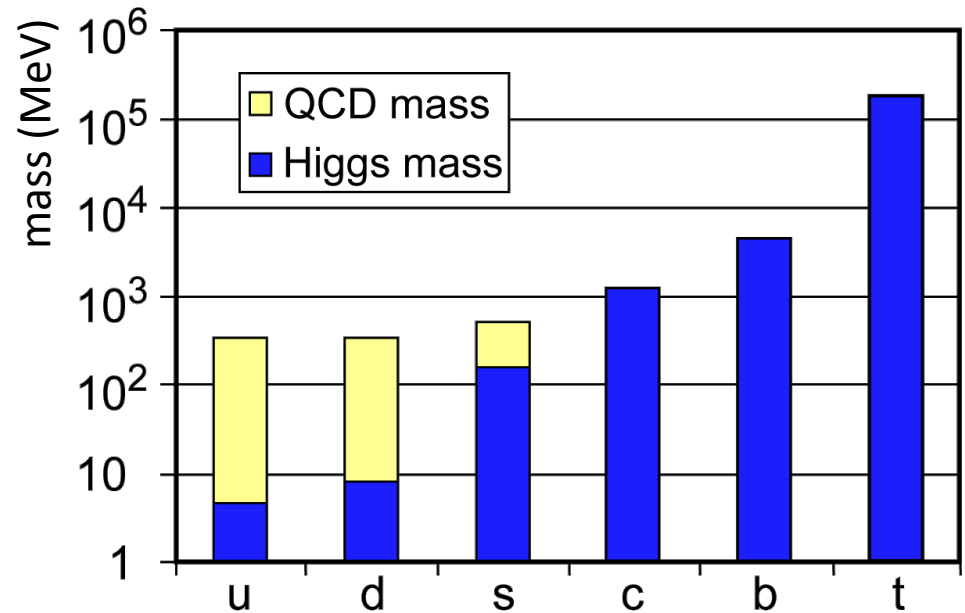
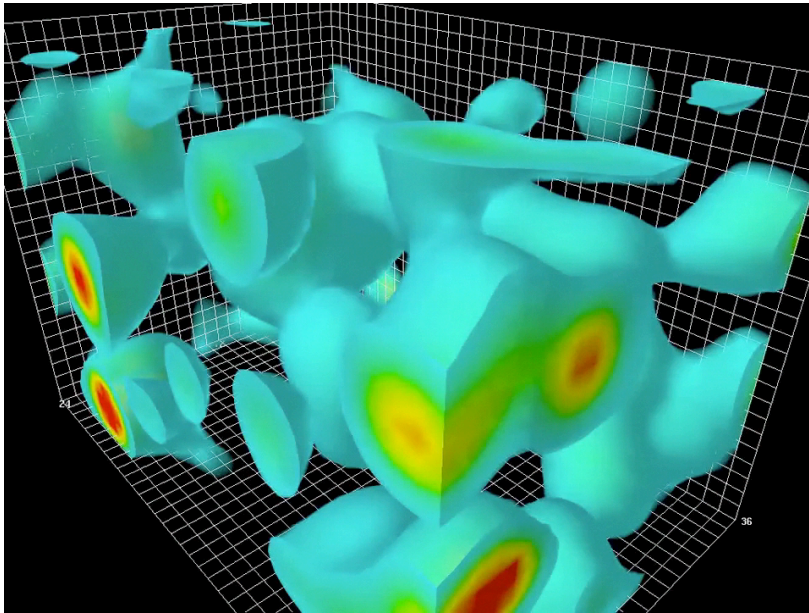


- Fermionic degrees of freedom are on the sites while links represent gauge degrees of freedom
 - Example: on lattice $32^3 \times 8$ there is $32^3 \times 8 \times 4 \times 18 \approx 2 \times 10^7$ gauge field DoF which needs to be integrated
- Physical observables are obtained via performing path integrations of partition function with a QCD action $e^{-S_{QCD}}$ as a weight, using modern MC simulation on supercomputers

QCD Vacuum

- The vacuum of QCD is composed a condensate of gluons and quarks – “empty system” unstable

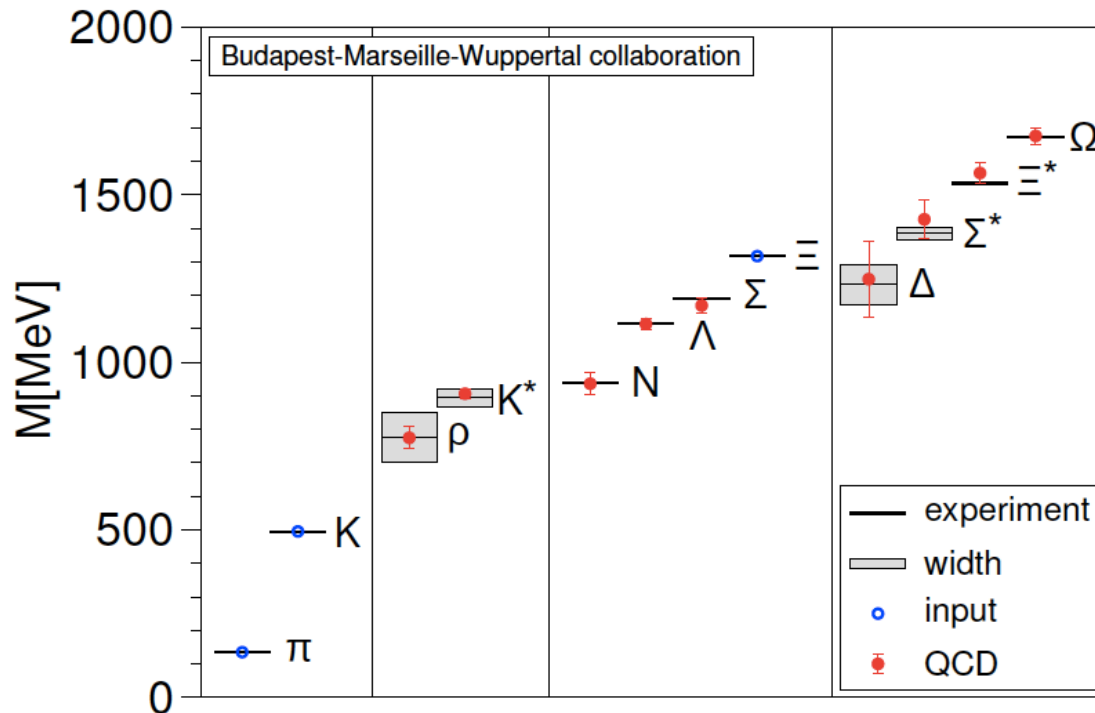
Derek Leinweber physics.adelaide.edu.au



- In the limit $m \approx 0$ \mathcal{L}_{QCD} is invariant with wrt $\psi_R \leftrightarrow \psi_L$ (chiral symmetry, $\psi_{L/R} = 1/2 (1 -/+ \gamma_5) \psi$)
- Chiral symmetry not preserved at particle level, quark bound states, e.g. $m_p \approx 938$ MeV (natural units, $c=\hbar=1$)

Exploring Non-Perturbative Regions of QCD

Using the lattice leads to predictions for the spectrum of hadrons



- Predictions from lattice-QCD agree with experimental measurements for wide variety of hadrons including light- and heavy-hadrons and proton-neutron mass difference (-1.7 MeV)

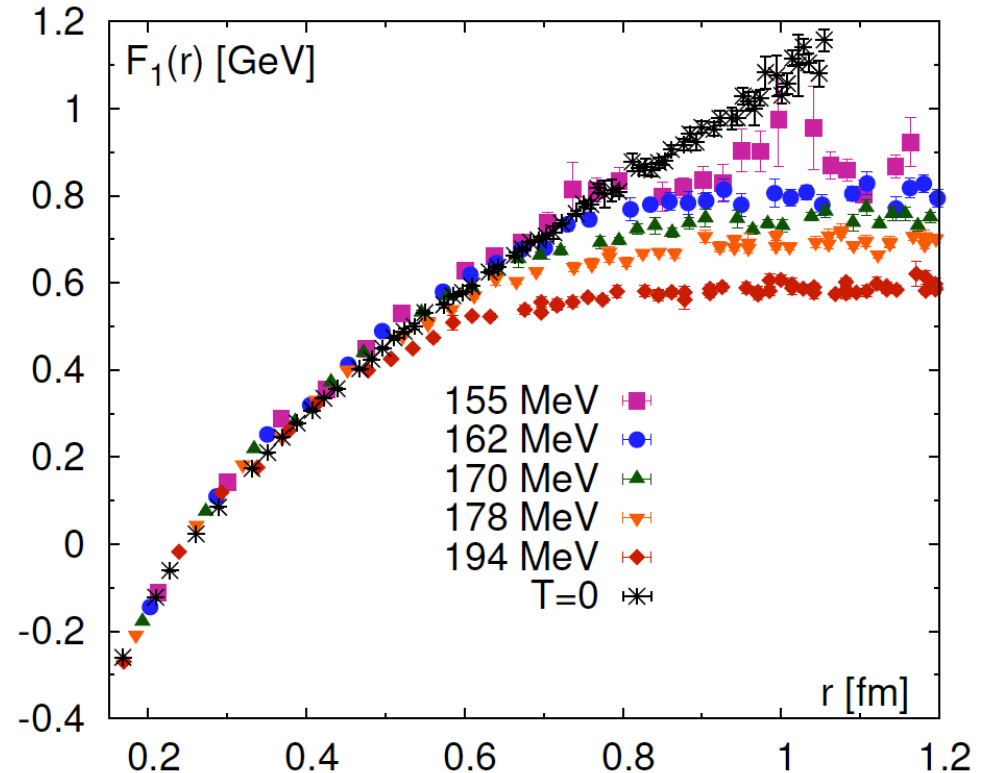
Quark Anti-quark Potential vs T - Lattice Results

- The free energy of quark anti-quark pair as function of separation r

$$V(r) \sim -\alpha/r + \sigma r$$

$T=155 \text{ MeV} \sim 10^{12} \text{ K}$

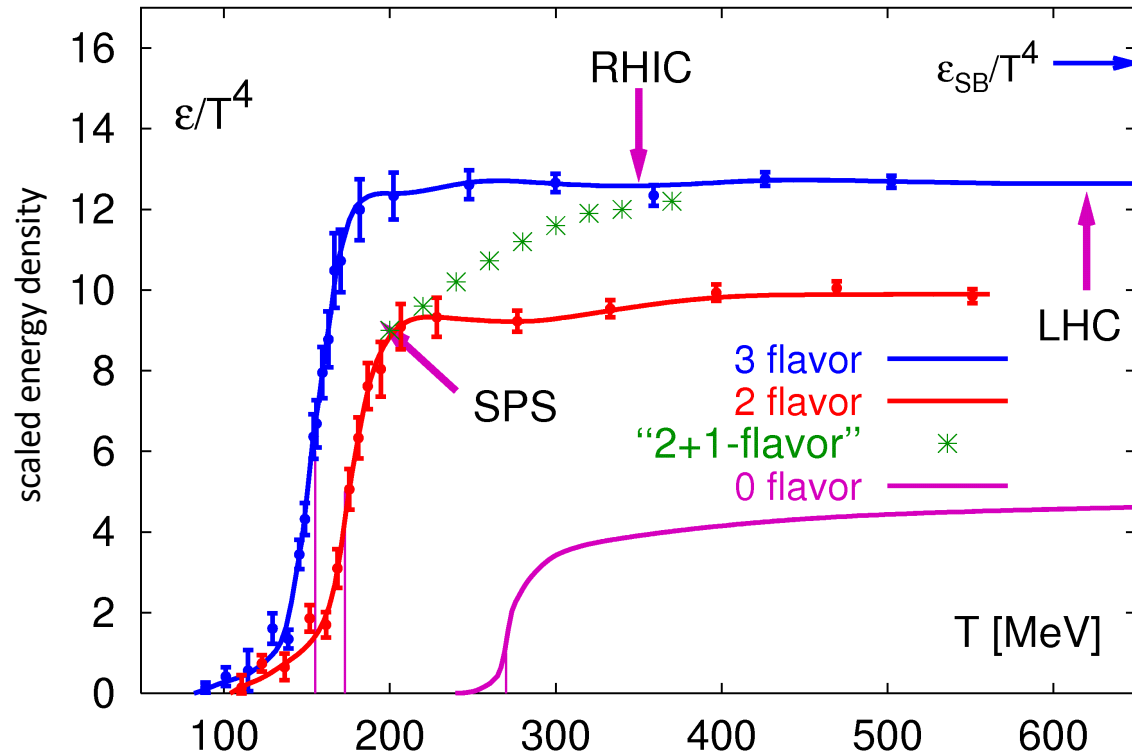
Sun core: $\sim 10^7 \text{ K}$



- Decrease of F_1 with increasing temperature at fixed r is related to the Debye screening effect in QGP

Thermal Behavior of QCD - Lattice Results

- Rapid rise in the number of degrees of freedom at $T \sim 150$ MeV
- ε below Stefan-Boltzmann limit for free quarks and gluons



[Karsch et al. (2009)]

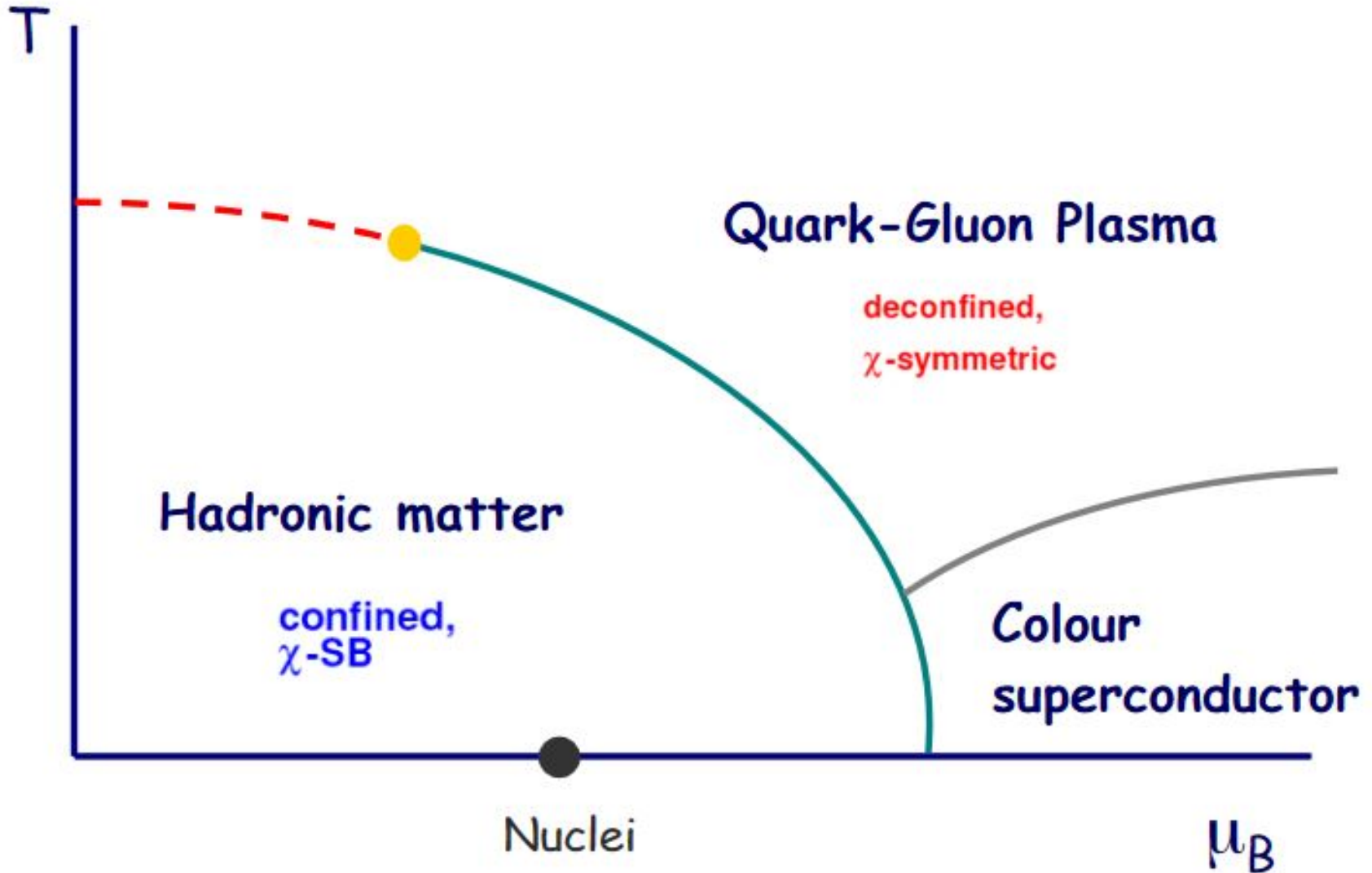
$$\begin{aligned} \varepsilon_{SB}^g &= \\ &= 16 \times 4\pi \int_0^\infty \frac{p^2 dp}{(2\pi)^3} \frac{p}{e^{p/T} - 1} \\ &= 16 \times \frac{\pi^2}{30} T^4 \end{aligned}$$

$$\begin{aligned} \varepsilon_{SB}^q + \varepsilon_{SB}^{anti-q} &= \\ &= 36 \times 4\pi \int_0^\infty \frac{p^2 dp}{(2\pi)^3} \frac{p}{e^{p/T} + 1} \\ &= 36 \times \frac{7\pi^2}{8 \cdot 30} T^4 \end{aligned}$$

$$\varepsilon_{SB} = \varepsilon_{SB}^g + \varepsilon_{SB}^q + \varepsilon_{SB}^{anti-q} = 15.6 \times T^4$$

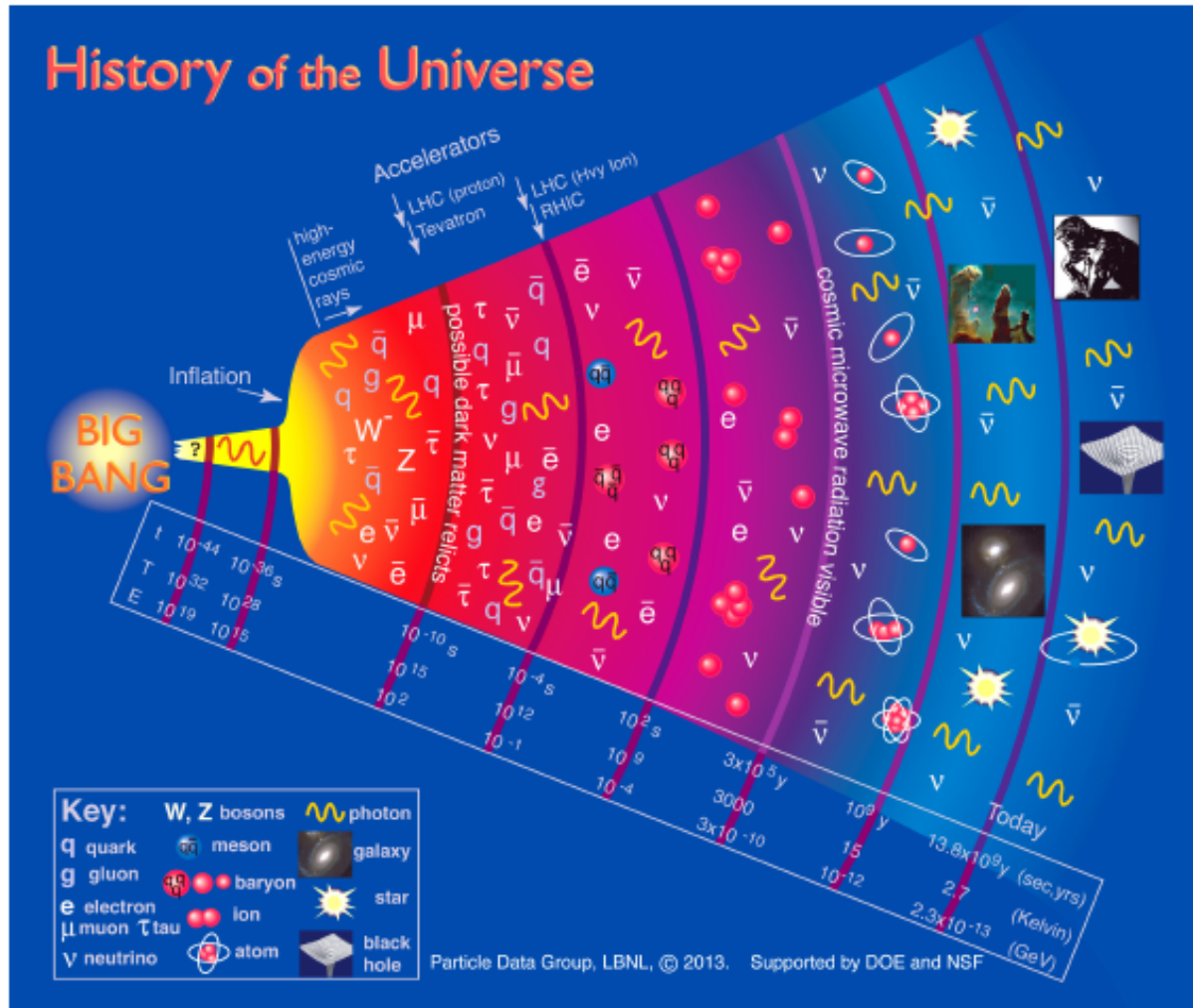
- Quantum Chromodynamics shows a rapid crossover (a smooth transition) from hadronic gas to QGP ($\mu_B=0$)
 - T_c of 150–200 MeV, $\varepsilon = 1-3$ GeV/fm³

Finite baryon density QCD matter



- Baryon chemical potential measures the imbalance between matter and antimatter

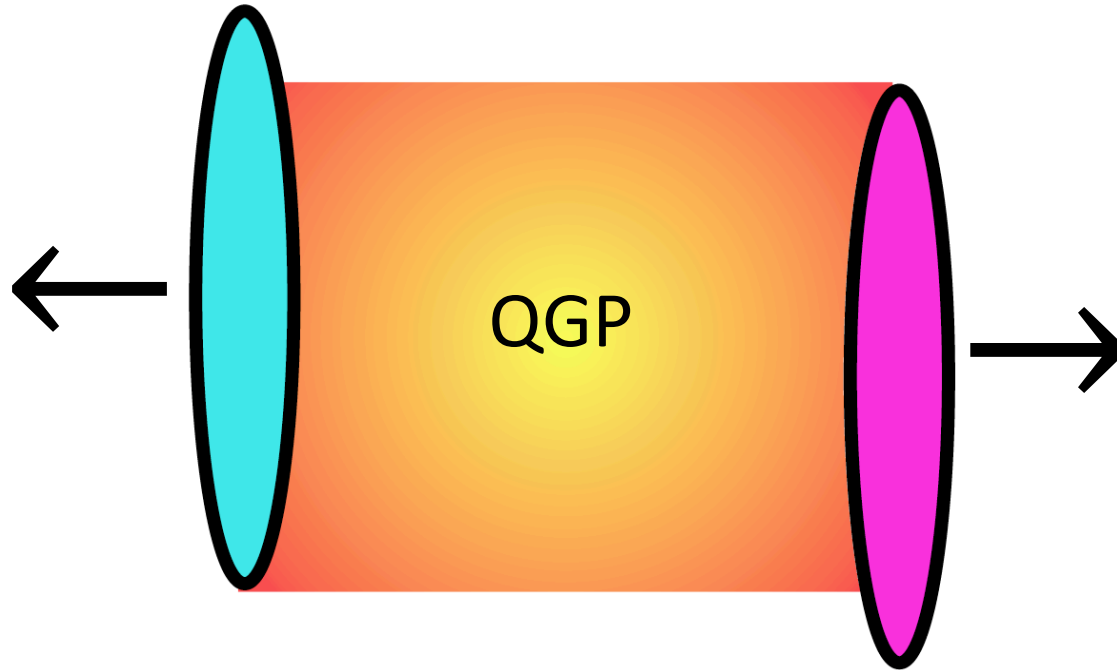
QGP in Cosmology



- QGP existed about $\sim 10 \mu\text{sec}$ after the big bang
 - Then the temperature of the universe was about 100 000 times larger than the temperatures at the center of the sun

Schematic View of a Heavy-ion Collision

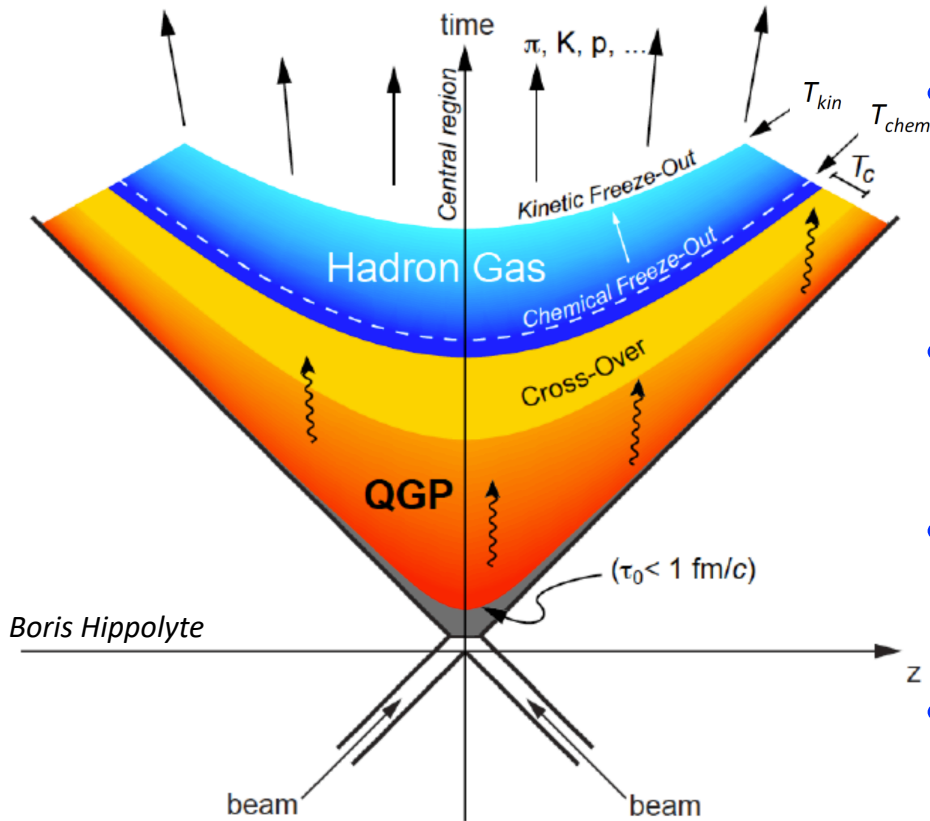
It is expected that in heavy ion collisions there are sufficient conditions to generate a “droplet” of QGP



- The life time of QGP is a few fm/c, i.e. $\sim 10^{-23}$ s

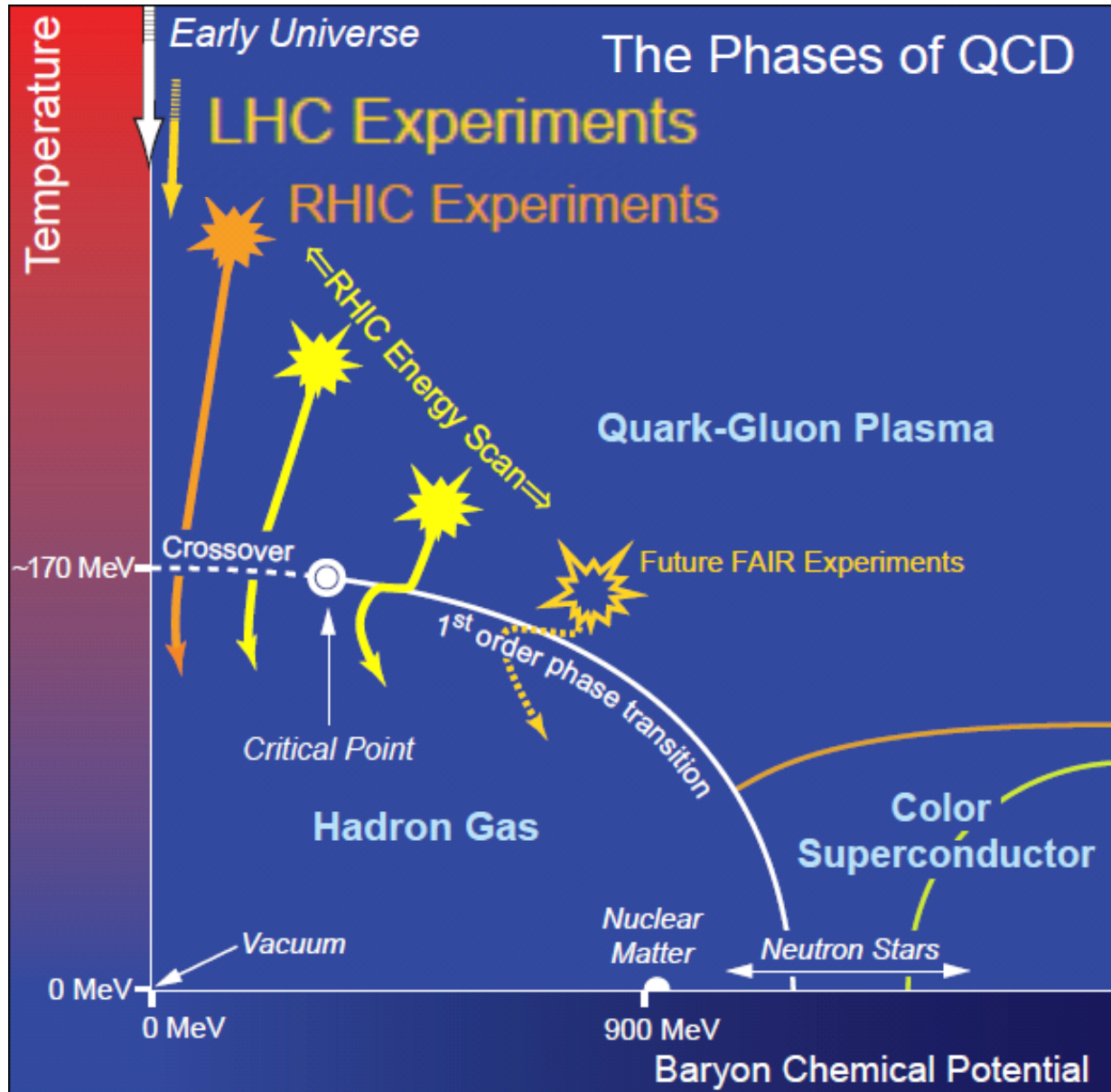
1 fm = 1 fermi = 10^{-15} m

Stages of a Heavy-ion Collision



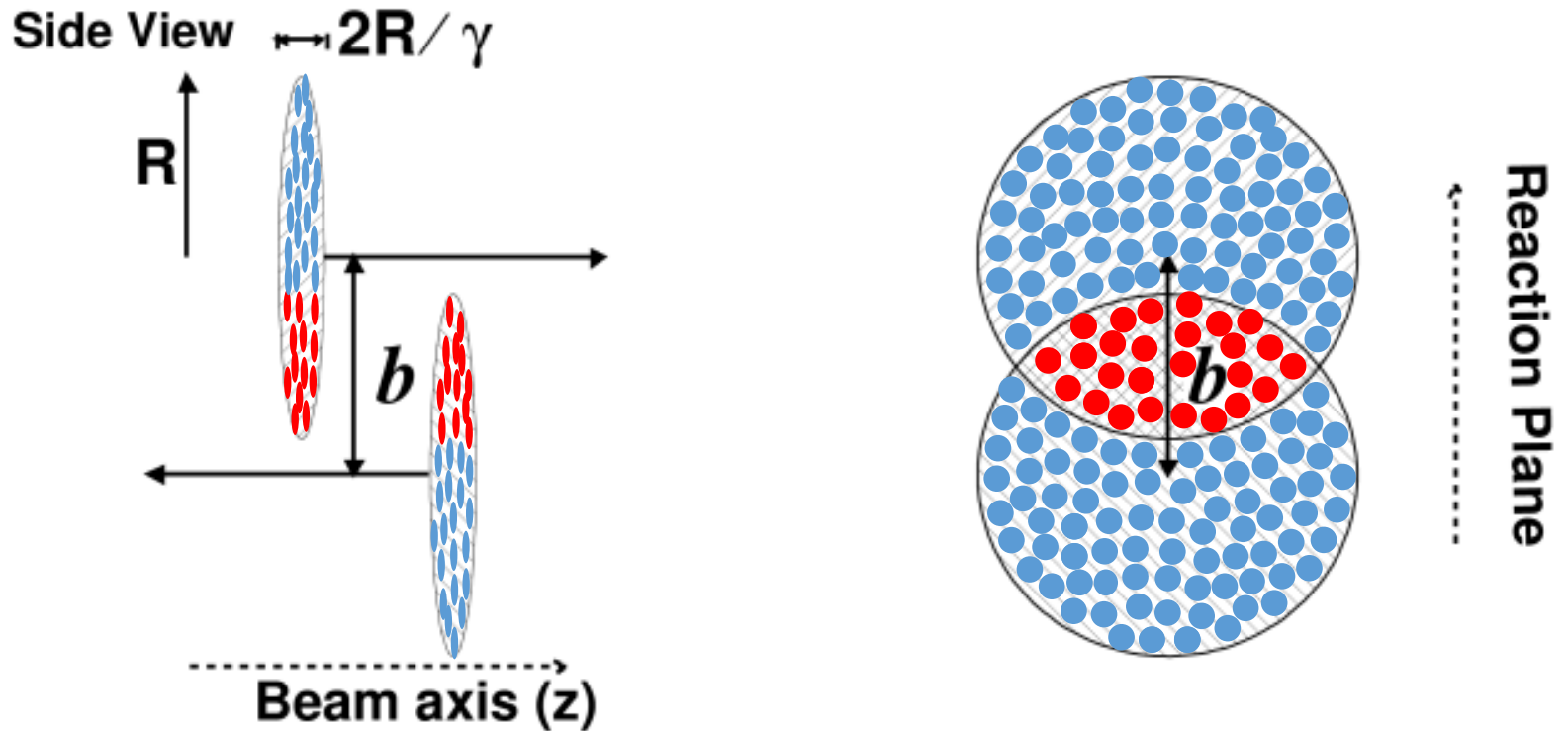
- The space time evolution starts from a hot-fireball in a pre-equilibrium phase ($\tau_0 < 1 \text{ fm}$)
- Equilibrate state, thermalization, QGP
- Cross-over phase transition to a hadron gas (T_C).
- Emitting of different kinds of particles (T_{chem}, T_{kin})

Study of the QCD Phase Diagram



- Changing beam energy leads to changes in the temperature and μ_B of the system

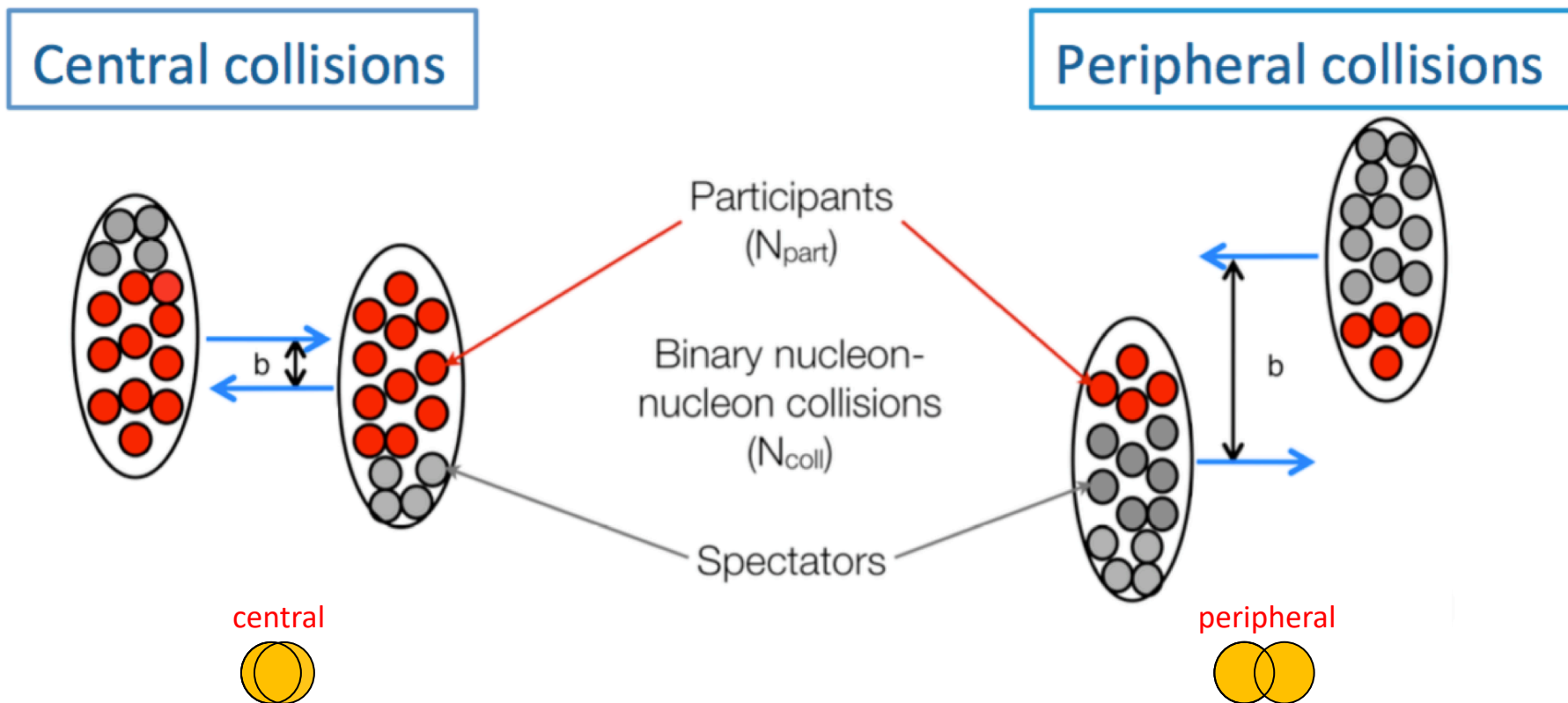
Layout of a Heavy-ion Collision



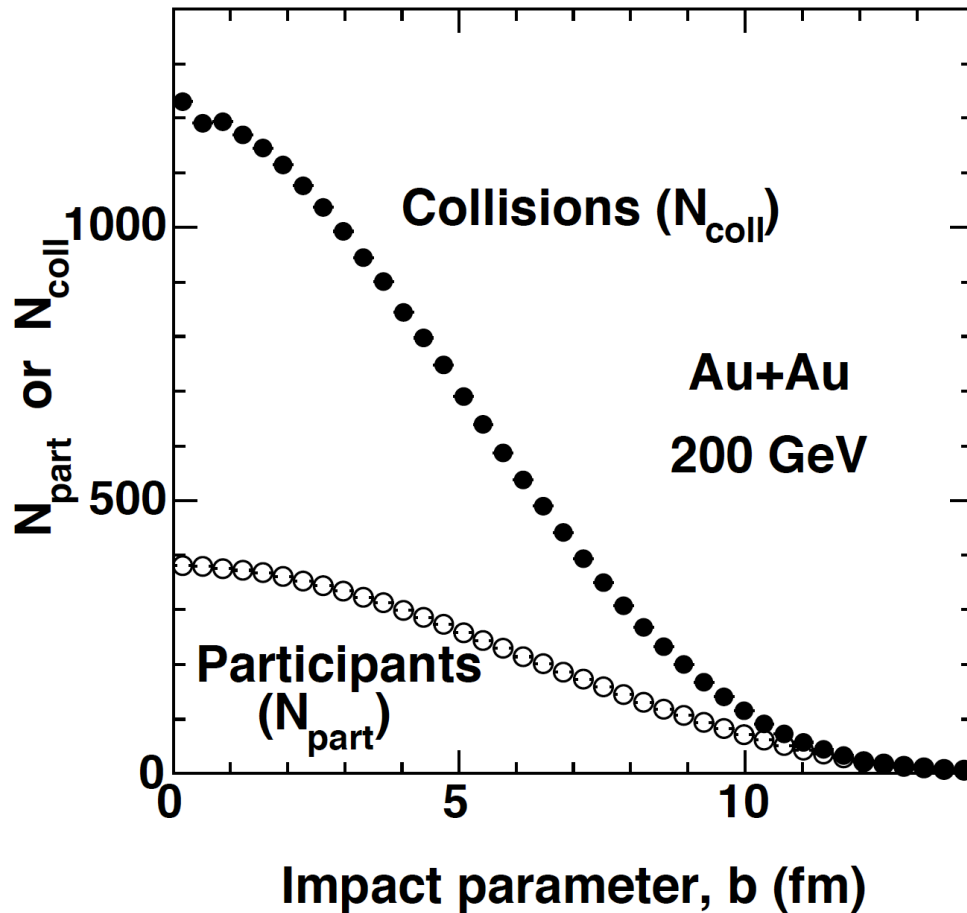
- nucleon participating in at least 1 inelastic collision, N_{part}
- spectator nucleon

N_{coll} the number of binary nucleon-nucleon collisions in a heavy ion reaction

Heavy-ion Collisions of Different Centrality



N_{coll} and N_{part} versus impact parameter



PHOBOS,
Nucl.Phys.A 757 (2005) 28-101

Woods-Saxon function

$$P(R) = R^2 \left(1 + e^{\frac{(R-r_0)}{a}} \right)^{-1}$$

$$r_0 = 6.38 \text{ fm}, a = 0.535 \text{ fm}$$

In nucleus (A)-nucleus (B) reaction

$$T_{\text{AB}}(\mathbf{b}) = \frac{N_{\text{coll}}}{AB\sigma_{\text{in}}}$$

$T_{\text{AB}}(\mathbf{b})$ - nucleus-nucleus thickness function

σ_{in} - NN inelastic cross section,
 $\approx 42 \text{ mb}$ RHIC
 $\approx 78 \text{ mb}$ LHC

arxiv1710.07098

$$1 \text{ b} = 100 \text{ fm}^2 = 10^{-28} \text{ m}^2$$