Cumulants of net-proton number fluctuations from ALICE at the LHC

Nirbhay Kumar Behera for the ALICE collaboration

23rd May 2018, WPCF 2018, IFJ PAN, Krakow, Poland







Motivation: To map the QCD phase diagram



(Y. Akiba et al, arXiv:1502.02730v1)

• At LHC energies, $\mu_B \simeq 0$.

• Lattice QCD: $T_c \simeq 154 \pm 9 \text{ MeV}$

Estimation from models with ALICE data:

- $T_{\text{freeze-out}}$: ~156 ± 3 MeV
 - [J. Stachel et al J. Phys.Conf. Ser. 509, 012019, (2014)]

• Chemical freeze-out line is close to the crossover line!

Precise determination of freeze-out parameters (T and μ_B) at the LHC can help to locate the phase boundary at $\mu_B \simeq 0$.

S. Borsányi *et al*, PRL 111, 062005 (2013), **Frithjof Karsch**, Central Eur.J.Phys. 10 (2012) 1234-1237

Motivation: Connecting experiment with lattice QCD

In lattice QCD, the cumulants (C_n) of the distributions of conserved charges (net-charge, net-baryon, net-strangeness) are related to the generalized quark-number susceptibilities (χ_n^q)



$$\chi_n^q = \frac{\partial^n \left(p/T^4 \right)}{\partial (\mu_B/T)^n}$$

A direct comparison of experimental results with lattice QCD!

Motivation: Connecting experiment with lattice QCD Freeze-out temperature from the ratio of cumulants of net-baryons





From lattice QCD, freeze-out parameters can be extracted from experimental results at $\mu_B = 0$. Net-proton fluctuations are a good proxy for net-baryon number fluctuations.

[Y. Hatta, M.A. Stephanov, PRL 91 102003 (2003)]

Experimental measurements of the ratio of cumulants of net-protons at the LHC will help to constrain the lattice QCD predictions in a model independent way.

Experimental setup and dataset ALICE detector



V0 detector: Trigger, centrality estimation V0A: $2.8 < \eta < 5.1$, V0C: $-3.7 < \eta < -1.7$ Inner Tracking System (ITS): Vertex, tracking, PID

Time Projection Chamber (TPC): Tracking, PID

Minimum-bias Pb-Pb collision data

Collision energy	Pb-Pb 2.76 TeV (Run1)	Pb-Pb 5.02 TeV (Run2)
Number of events	14×10^{6}	59 × 10 ⁶

- Kinematic cuts: $0.4 < p_T < 1.0 \text{ GeV}/c, |\eta| < 0.8$
- **Proton identification: TPC (next slide)**

Proton Identification

Energy loss (dE/dx) of particles in TPC gas



 (Anti-)Proton identified using the nσ cut around the expected value of energy loss of particles in TPC detector (this presentation).

Generalized Gaussian fit to TPC dE/dx



- Proton identification using Identity method: Assign a Bayesian probability for each particle species.
- Get the true numbers using response matrix built from the parameterization of TPC dE/dx.

M. Gazdzicki et al., PRC 83, 054907 (2011) M. I. Gorenstein, PRC 84, 024902 (2011) A. Rustamov, M. I. Gorenstein, PRC 86, 044906 (2012)

C₂ results from both the methods agree with each other

Higher moments and cumulants

- Measure net-proton numbers on event-by-event basis: $\Delta \mathbf{p} = \mathbf{p} \overline{\mathbf{p}}$
- The nth moments: $m'_n = \langle (\Delta p) \rangle^n > 1$



1) Detector efficiency correction: method 1

> A. Bzdak and V. Koch, PRC 91, 027901 (2015).

 $K_4 = N - 6K_1^4 + F_{04} + 6F_{03} + 7F_{02} - 2F_{11} - 6F_{12} - 4F_{13}$

 $-4K_1(K_1 - F_{03} - 3F_{02} + 3F_{12} + 3F_{20} - 3F_{21} + F_{30}),$

 $+12K_{1}^{2}(N+F_{02}-2F_{11}+F_{20})-3(N+F_{02}-2F_{11}+F_{20})^{2}$

 $+7F_{20}-6F_{21}+6F_{22}+6F_{30}-4F_{31}+F_{40}$

> Assumption: detector efficiency follows Binomial distribution.

Factorial

where

$$F_{i,k} = \sum_{x_1,...,x_i} \sum_{\bar{x}_1,...,\bar{x}_k} \frac{u_{i,k}(x_1,...,x_i,x_1,...,x_k)}{\epsilon(x_1)\dots\epsilon(x_i)\bar{\epsilon}(\bar{x}_1)\dots\bar{\epsilon}(\bar{x}_k)}.$$

$$a_{i,k}(x_1,...,x_i;\bar{x}_1,...,\bar{x}_k) = \langle n(x_1)[n(x_2) - \delta_{x_1,x_2}] \cdots [n(x_i) - \delta_{x_1,x_i} - \dots - \delta_{x_{i-1},x_i}] \\ \bar{n}(\bar{x}_1)[\bar{n}(\bar{x}_2) - \delta_{\bar{x}_1,\bar{x}_2}] \cdots [\bar{n}(\bar{x}_k) - \delta_{\bar{x}_1,\bar{x}_k} - \dots - \delta_{\bar{x}_{k-1},\bar{x}_k}] \rangle.$$

$$Cumulants from factorial moments$$

$$K_1 = \langle N_1 \rangle - \langle N_2 \rangle, \\ K_2 = N - K_1^2 + F_{02} - 2F_{11} + F_{20}, \\ K_3 = K_1 + 2K_1^3 - F_{03} - 3F_{02} + 3F_{12} + 3F_{20} - 3F_{21} + F_{30} \\ - 3K_1(N + F_{02} - 2F_{11} + F_{20}), \end{cases}$$
where
$$N \equiv \langle N_1 \rangle + \langle N_2 \rangle = F_{10} + F_{01}.$$

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1) Detector efficiency correction: method 2 (Used here)

Recent method: T Nonaka et al ,PRC 95, 064912, (2017)

> Assumption: detector efficiency follows Binomial distribution.

$$\langle Q \rangle_{c} = \langle q_{(1,1)} \rangle_{c}, \langle Q^{2} \rangle_{c} = \langle q_{(1,1)}^{2} \rangle_{c} + \langle q_{(2,1)} \rangle_{c} - \langle q_{(2,2)} \rangle_{c}, \langle Q^{3} \rangle_{c} = \langle q_{(1,1)}^{3} \rangle_{c} + 3 \langle q_{(1,1)} q_{(2,1)} \rangle_{c} - 3 \langle q_{(1,1)} q_{(2,2)} \rangle_{c} + \langle q_{(3,1)} \rangle_{c} - 3 \langle q_{(3,2)} \rangle_{c} + 2 \langle q_{(3,3)} \rangle_{c}, \langle Q^{4} \rangle_{c} = \langle q_{(1,1)}^{4} \rangle_{c} + 6 \langle q_{(1,1)}^{2} q_{(2,1)} \rangle_{c} - 6 \langle q_{(1,1)}^{2} q_{(2,2)} \rangle_{c} + 4 \langle q_{(1,1)} q_{(3,1)} \rangle_{c} + 3 \langle q_{(2,1)}^{2} \rangle_{c} + 3 \langle q_{(2,2)}^{2} \rangle_{c} - 12 \langle q_{(1,1)} q_{(3,2)} \rangle_{c} + 8 \langle q_{(1,1)} q_{(3,3)} \rangle_{c} - 6 \langle q_{(2,1)} q_{(2,2)} \rangle_{c} + \langle q_{(4,1)} \rangle_{c} - 7 \langle q_{(4,2)} \rangle_{c} + 12 \langle q_{(4,3)} \rangle_{c} - 6 \langle q_{(4,4)} \rangle_{c},$$
where
$$\sum_{n=1}^{M} \langle q_{n} q_{n}$$

Moments

$$q_{(r,s)} = q_{(a^r/p^s)} = \sum_{i=1}^{M} (a_i^r/p_i^s) n_i.$$

 $a_i = 1$ for proton, -1 for anti-proton $p_i = Efficiency in i^{th} p_T bin.$ $n_i = Number of observed (anti-) proton$ $M = Total Number of p_T bins$

Simpler and faster than method 1.
Method 1 and 2 give the same results!

2) Centrality bin-width correction (CBWC):

To minimize the impact parameter (or volume) variations due to the finite centrality bin. [X. Luo et al J. Phys. G 40, 105104 (2013)]



$$\mathbf{C}_n = \frac{\sum_{i=1}^k m_i \mathbf{C}_{n,i}}{\sum_{i=1}^k m_i}$$

 $C_{n,i} = n^{th}$ order cumulant measured in ith centrality bin $m_i =$ number of events in ith centrality bin k = numbers of centrality bins

With or without CBWC, the results are consistent for ALICE measurements.
 Cumulants up to 3rd order are not sensitive to volume fluctuations!

[P. Braun-Munzinger, A. Rustamov, J. Stachel, NPA 960, 114 (2017)]

3) Statistical error estimation:

 \succ Subgroup method: Divide the data sample into 'k' subsamples randomly.

$$C_n = \frac{\sum_{i=1}^{k} C_n^i}{k}$$

Statistical error
$$\delta C_n = \frac{\sigma}{\sqrt{k}}$$
 where $\sigma = \sqrt{\frac{\sum_{i=1}^k (C_n^i - \langle C_n \rangle)^2}{k-1}}$

> Number of subsamples used : 30

> Statistical errors estimated from Bootstrap and Subgroup method agree with each other.

Results: proton and anti-proton C₁



• For both energies, proton and anti-proton numbers are similar within systematic errors.

Results: Cumulants of net-proton distributions



- Cumulants of net-proton distributions for both energies are the same within statistical and systematic errors.
- C_2 results are consistent with Identity method.

[A. Rustamov NPA, 967 (2017) 453-456 (QM 2017)]

Baseline estimation

Assume proton and anti-proton distributions are independent Poissonian distributions => net-proton distribution is a Skellam distribution.

- $\succ C_n(\text{Skellam}) = C_1(p) + (-1)^n C_1(\overline{p})$
- $\succ C_1(\Delta p) = C_3(\Delta p) = C_1(p) C_1(\overline{p})$
- > $C_2(\Delta p) = C_4(\Delta p) = C_1(p) + C_1(\overline{p}) => C_4/C_2 = 1.$

Used in Hadron Resonance Gas (HRG) model calculations.
 [P. Braun-Munzinger *et al* PLB 747, 292 (2015), P. Garg *et al* PLB 726, 691 (2013)]

Results: Ratio of cumulants of net-proton distributions



- C_3/C_2 and C_4/C_2 results for 2.76 TeV and 5.02 TeV are the same within the statistical error bars.
- C_3/C_2 and C_4/C_2 of central events agree with Skellam expectations within statistical errors.

Results: RHIC to LHC



- C_3/C_2 follows a decreasing trend as a function of collision energy.
- Within experimental uncertainties C_4/C_2 at the LHC agrees with the Skellam baseline.

Freeze-out temperature at the LHC



Implied freeze-out temperature from C_4/C_2 consistent with thermal fits to particle yields. (assuming net-protons are good proxy to net-baryons on the lattice)

> Need more precise data to further constrain the freeze-out temperature.

Summary

- The net-proton higher order cumulants up to 4th order and their ratios in Pb-Pb collisions at 2.76 and 5.02 TeV are presented as a function of centrality.
- \succ C_{3,} C₄ and their ratios with respect to C₂ generally agree with Skellam expectations within the uncertainties.
- > From RHIC to LHC the ratios of cumulants approach the Skellam baseline.
- The results may be used in conjunction with lattice calculations to estimate the freeze-out temperature.
- News: Second cumulants of net-Λ fluctuations are measured in ALICE (see A. Ohlson talk in QM 2018)

Outlook

- > The analysis will be extended for higher p_T (up to 2 GeV/c) and for other particle species in different pseudo-rapidity windows.
- > Measurement of net-proton higher order cumulants with Identity method.
- Off-diagonal cumulants of net-particles.
- The upcoming dedicated Pb-Pb run at 5.02 TeV will increase the statistics to further constrain the freeze-out parameters.

THANK YOU

Backup slides

Efficiency correction method

$$N_{rec.} = N_{rec.primary} + N_{cont}.$$

where
$$N_{cont.} = N_{material} + N_{secondary} + N_{mis.Identified}$$

Now $N_{rec.primary} = N_{rec} - N_{cont.}$

$$\Rightarrow N_{gen.primary} = N_{rec.} \frac{(1-\delta)}{\varepsilon}$$

where contamination factor (
$$\delta$$
) = $\frac{N_{cont.}}{N_{rec.}}$

Efficiency (
$$\varepsilon$$
) = $\frac{N_{rec.primary}}{N_{gen.primary}}$

$$\Rightarrow N_{gen.primary} = \frac{N_{rec.}}{\omega}$$

correction factor
$$\boldsymbol{\omega} = \frac{1}{(1-1)^2}$$

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