



# Accelerating hydrodynamic description of pseudorapidity distributions and the initial energy density at RHIC and LHC

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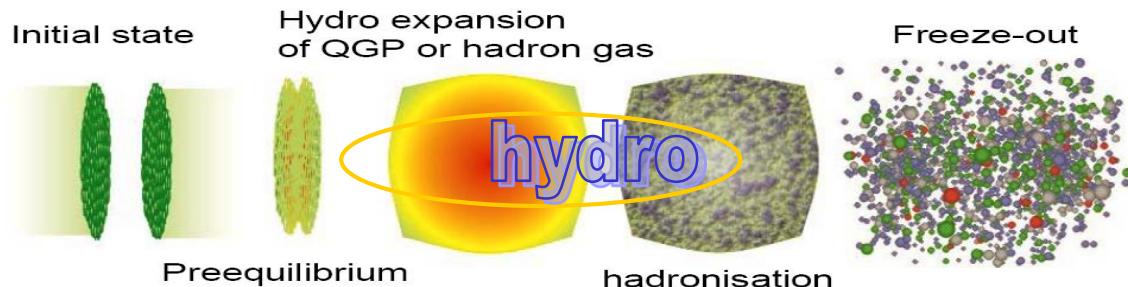
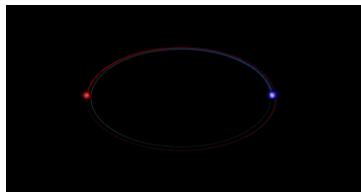
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Many thanks to:

Xin-Nian Wang, T.S. Biró, M. I. Nagy.

# Motivation



- *What is the initial temperature and thermal evolution of the produced matter?*
- *What is the viscosity of the produced matter? ... <http://www.bnl.gov/physics/rhiciiscience/>*

## Outline

- Accelerating ideal solution of hydrodynamics and results (perfect flow for fitting p+p, Cu+Cu, Au+Au, Pb+Pb data).  
Csörgő, Nagy, Csanád (CNC) arXiv: 0605070, 0710.0327, 0805.1562,  
Csanád, et. arXiv:1609.07176. Z. F. Jiang, et. arXiv: 1711.10740.
- New perturbative solutions of viscous hydrodynamics.
- Summary and outlook.

# Motivation

The next phase ... will focus on detailed investigations of the QGP, “both to quantify its properties and to understand precisely how they emerge from the fundamental properties of QCD”

-- The frontiers of nuclear science, a long range plan

- *What is the initial temperature and thermal evolution of the produced matter?*
- *What is the viscosity of the produced matter?* ... <http://www.bnl.gov/physics/rhiciiscience/>

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# CNC solutions of relativistic hydrodynamics

Hydrodynamics can be a universal tool to study the QGP. In Rindler coordinate, 5 different sets of the parameters  $\lambda$ ,  $\kappa$ ,  $d$  and  $K$  for 5 possible cases as follows  
( $\lambda$  is the accelerate parameter.):

[Csörgő, Nagy, Csanád\(CNC\)](#)  
arXiv: 0605070,  
0710.0327, 0805.1562,....

Case	$\lambda$	$d$	$\kappa$	$\phi$	
a.)	2	R	d	0	→ accelerating, $d$ dimension
b.)	1/2	R	1	$(\kappa+1)/\kappa$	← $d$ dimensional ( T. S. Biró)
c.)	3/2	R	$(4d-1)/3$	$(\kappa+1)/\kappa$	
d.)	1	R	R	0	→ Hwa-Bjorken, Buda-Lund type
e.)	R	R	1	0	→ Special EoS, but general velocity

In all ideal cases, the **velocity field** and the **pressure** is expressed as :

$$v = \tanh \lambda \eta_s, \quad p = p_0 \left( \frac{\tau_0}{\tau} \right)^{\lambda d \frac{(\kappa+1)}{\kappa}} \left( \cosh \frac{\eta_s}{2} \right)^{-(d-1)\phi}$$

# The initial energy density estimation

Rapidity distribution:

$$\frac{dN}{dy} \approx \left. \frac{dN}{dy} \right|_{y=0} \cosh^{\frac{\pm\alpha}{2}-1} \left( \frac{y}{\alpha} \right) e^{-\frac{m}{T_f} [\cosh^\alpha (\frac{y}{\alpha}) - 1]}$$



Pseudorapidity distribution:

$$\frac{dN}{d\eta} \approx \frac{\bar{p}}{E} \frac{dn}{dy} = \frac{\bar{p}_T \cosh \eta}{\sqrt{m^2 + \bar{p}_T^2}} \frac{dN}{dy}$$

$$\alpha = \frac{2\lambda - 1}{\lambda - 1}.$$

# The initial energy density estimation

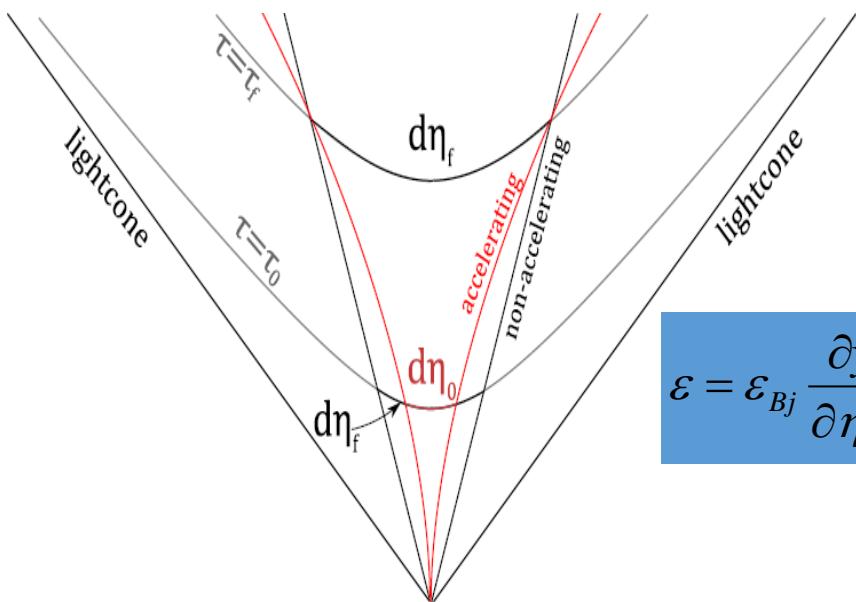
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Flow element's motion path

For more recent, exact results that include the work done by pressure, see [G.Kasza's talk at WPCF2018](#)

$$\varepsilon_{Bj} = \frac{1}{R^2 \pi \tau_0} \frac{dE}{d\eta} = \frac{\langle E \rangle}{R^2 \pi \tau_0} \frac{dn}{d\eta}$$

J. D. Bjorken, Phys. Rev. D 27, 140 (1983)

For an accelerating flow,  
Two main modifications for **Initial conditions**:  $y \neq \eta$ ;  $\eta_{\text{final}} \neq \eta_{\text{initial}}$ .

$$\varepsilon = \varepsilon_{Bj} \frac{\partial y}{\partial \eta_f} \frac{\partial \eta_f}{\partial \eta_i}$$

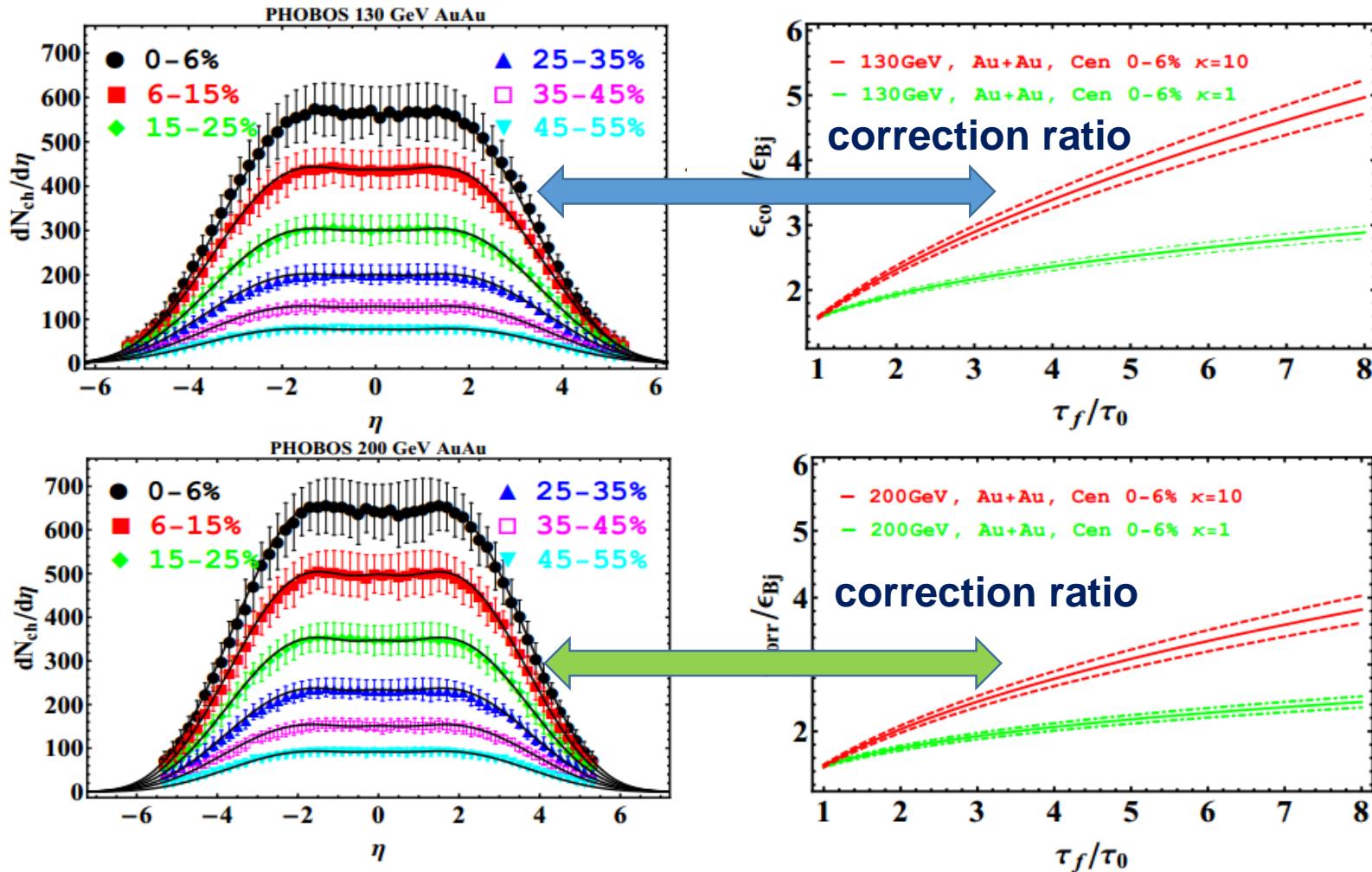
$$\lambda = 1.18 \pm 0.01, \quad \tau_f / \tau_0 = 8 \pm 2,$$

$$\varepsilon_{corr} = (2.0 \pm 0.1) \varepsilon_{Bj} = 10.0 \pm 0.5 \text{ GeV/fm}^3.$$

BRAHMS 200 GeV Au+Au data  
arXiv: 0805.1562

# The pseudorapidity distributions

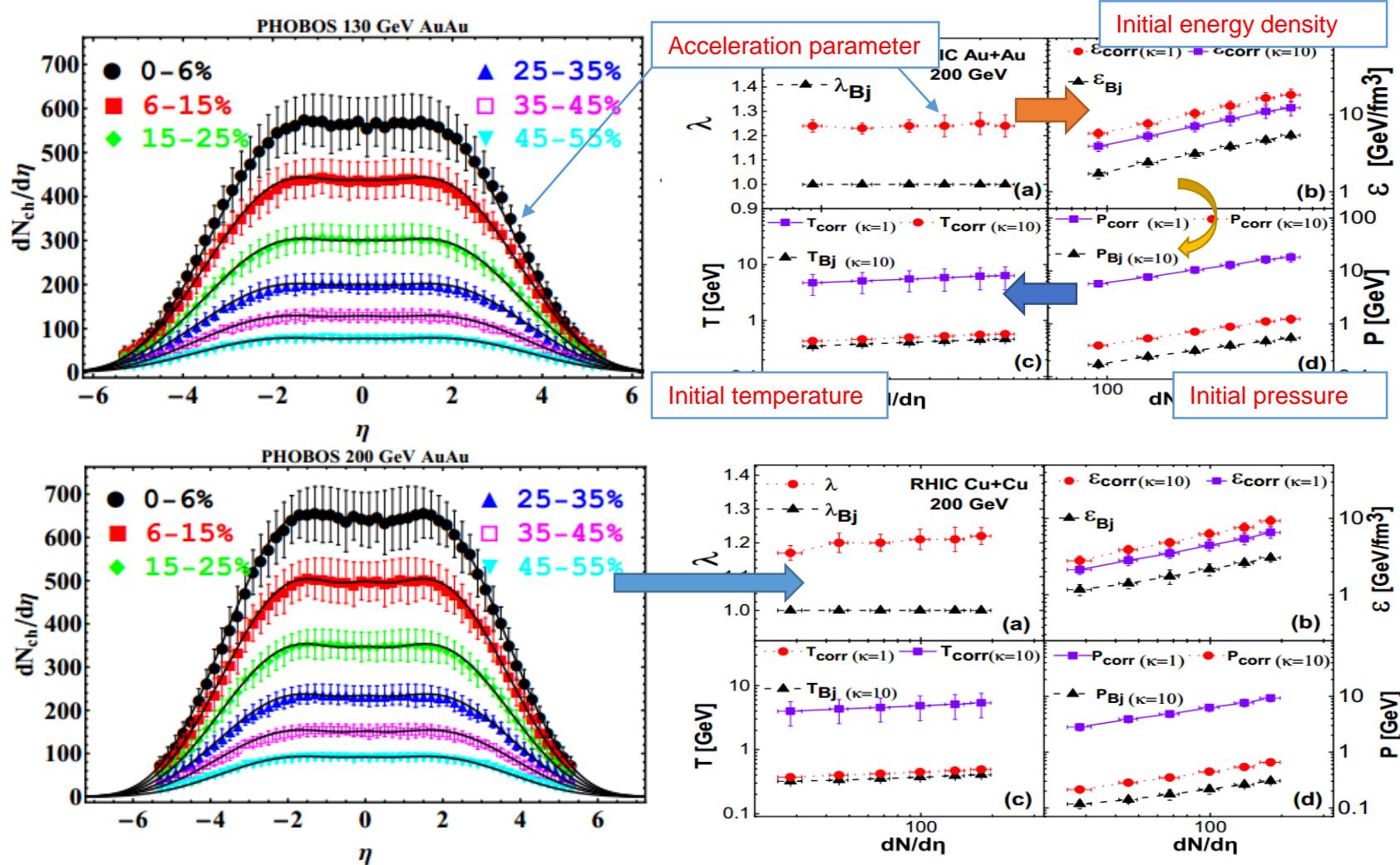
Results for 130 GeV Au+Au and 200 GeV Au+Au collisions



(See Z.F. Jiang, C.B. Yang, M. Csanad, T. Csorgo, arXiv:1711.10740)

# The initial energy density estimation

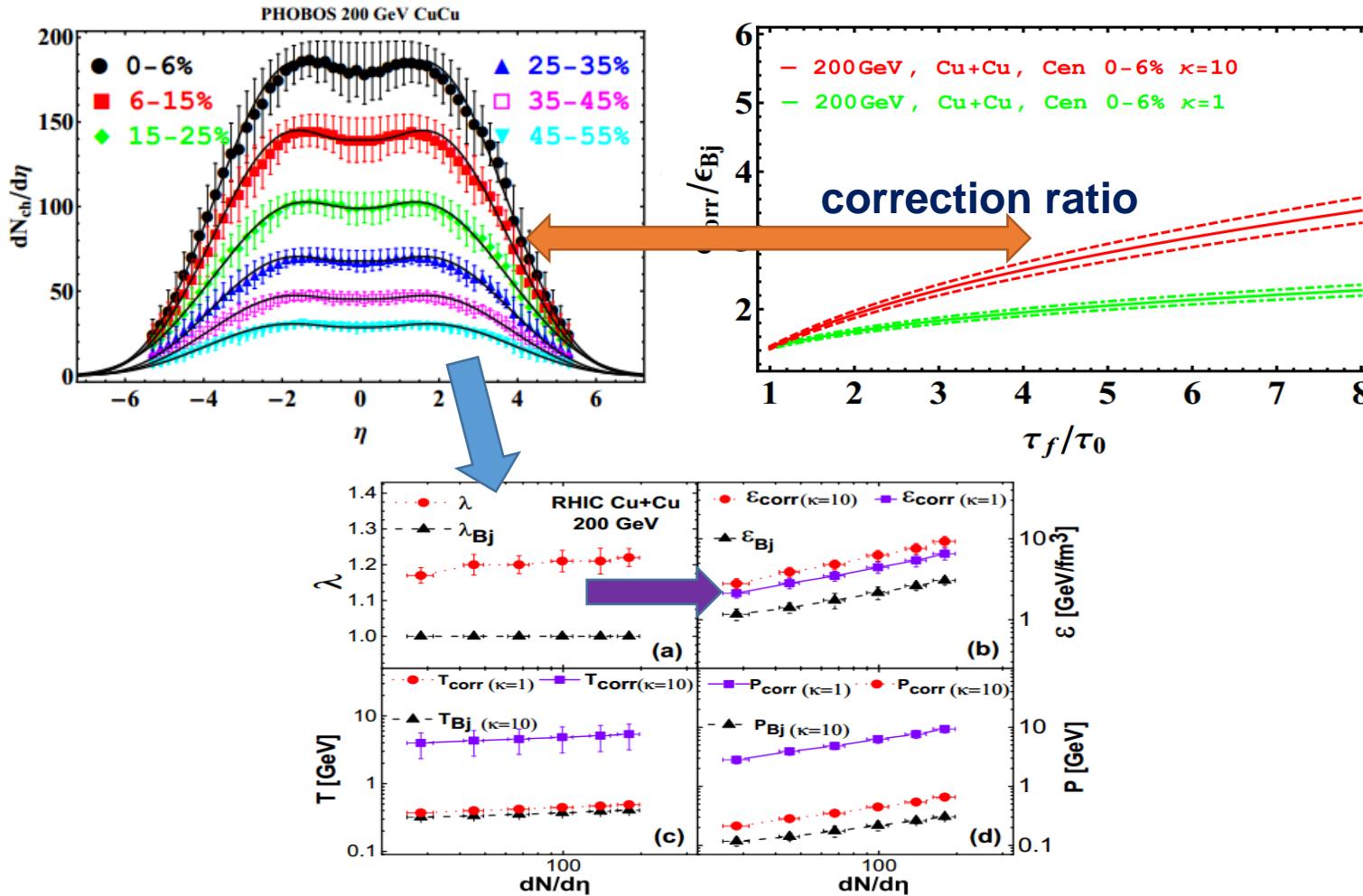
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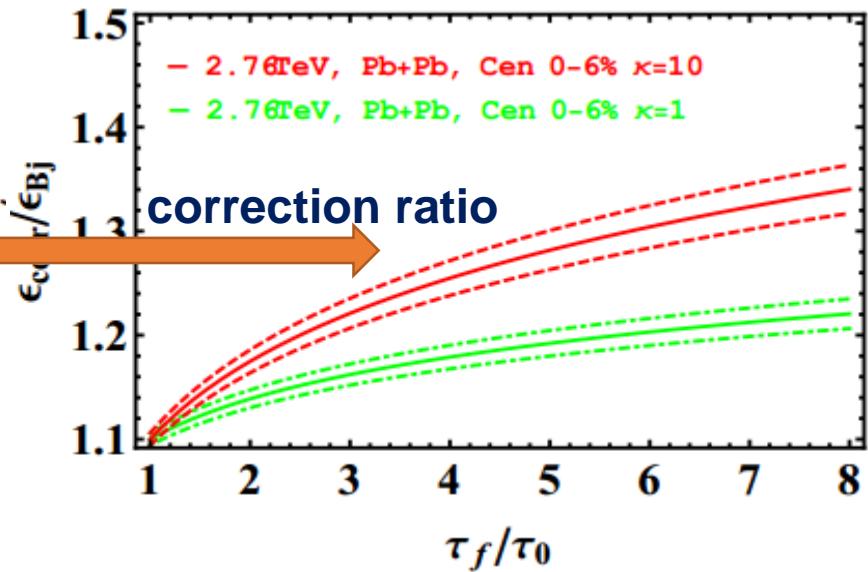
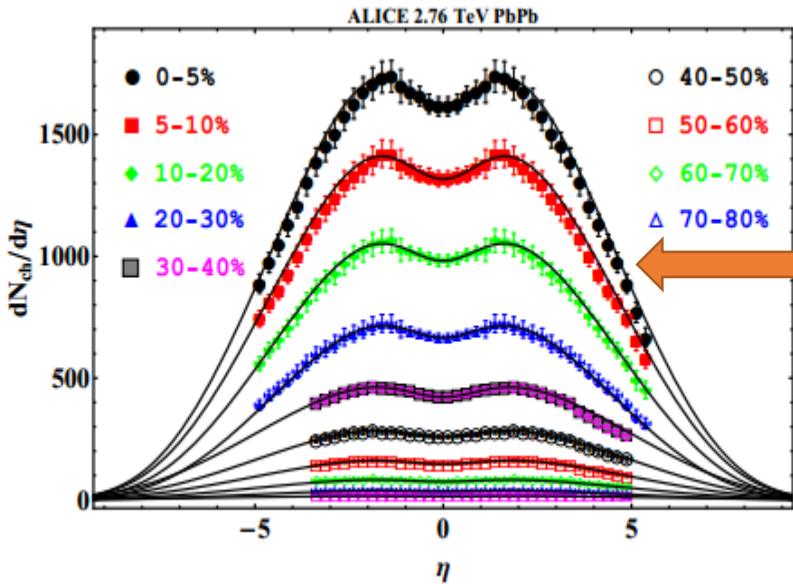
## Results for 200 GeV Cu+Cu collisions



(See Z.F. Jiang, C.B. Yang, M. Csanad, T. Csorgo, arXiv:1711.10740)

# The initial energy density estimation

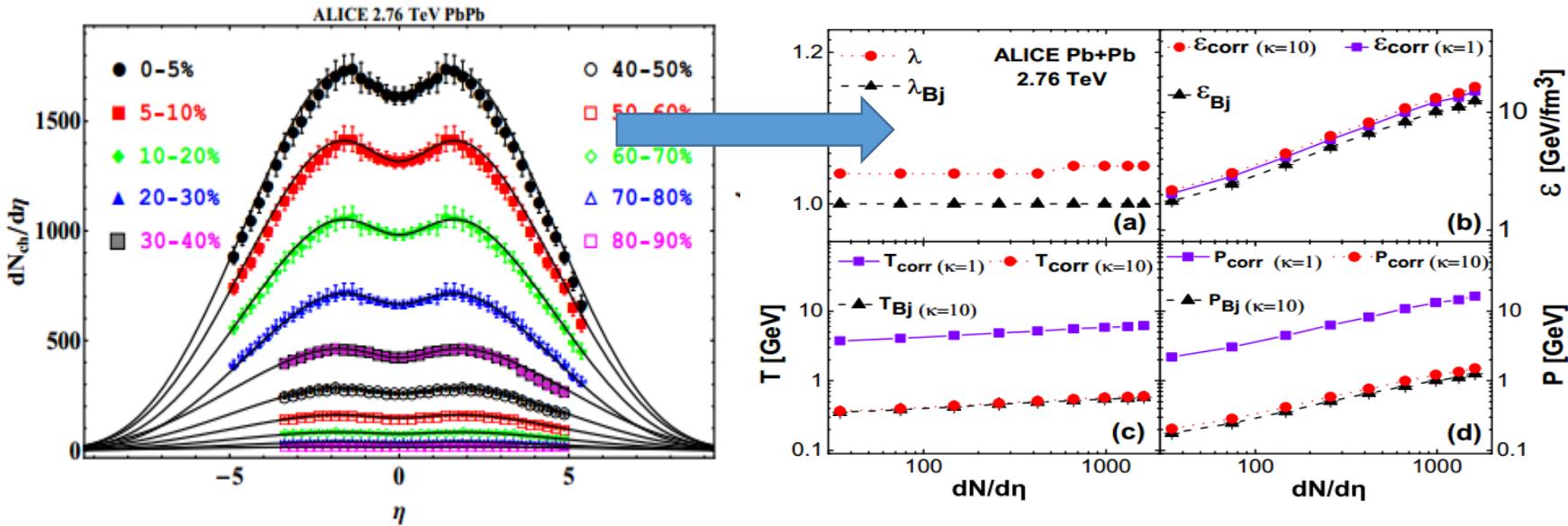
## Results for 2.76 TeV Pb+Pb collisions



(See Z.F. Jiang, C.B. Yang, M. Csanad, T. Csorgo, arXiv:1711.10740)

# The initial energy density estimation

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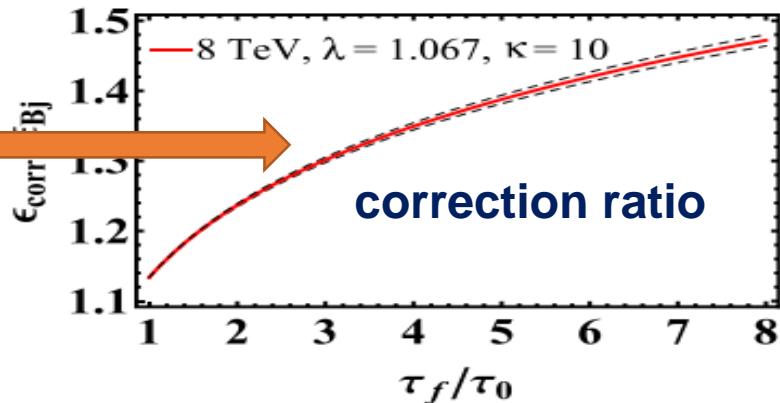
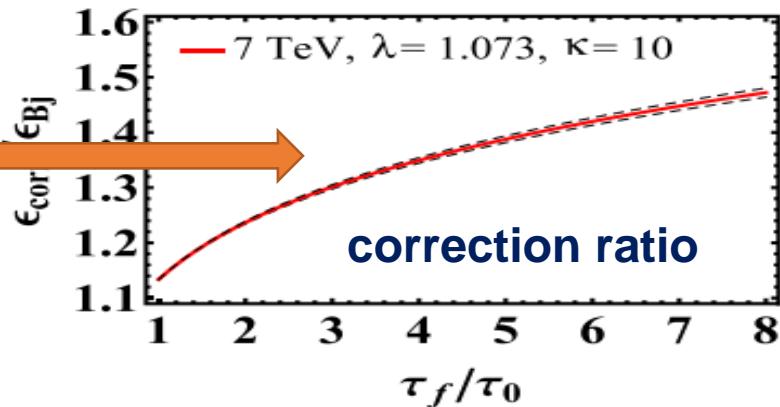
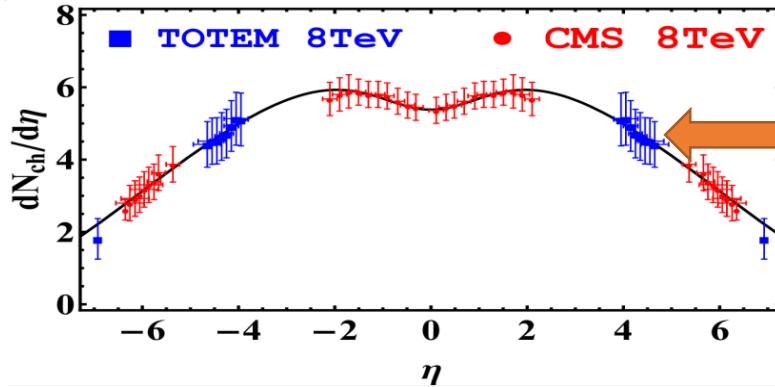
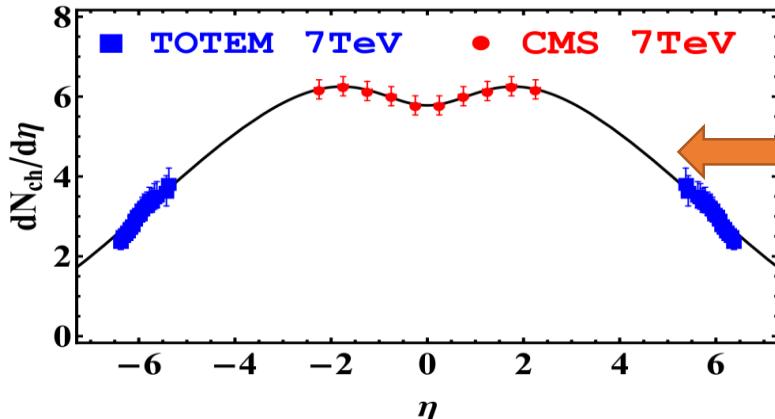


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- Hydro can be used to estimate the initial  $\epsilon_0$ . (J. D. Bjorken, Phys. Rev. D27 (1983) 140-151. Gyulassy M, Matsui T. 1984. Phys. Rev. D29: 419)
- The larger the center of mass energy, the smaller the acceleration parameter  $\lambda$ . The bigger  $\langle N_{part} \rangle$ , the bigger  $\lambda$ .
- From the fits, the initial energy density, initial temperature and initial pressure are reconstructed.

# The pseudorapidity distributions

Results for 7 TeV & 8 TeV p+p collisions:

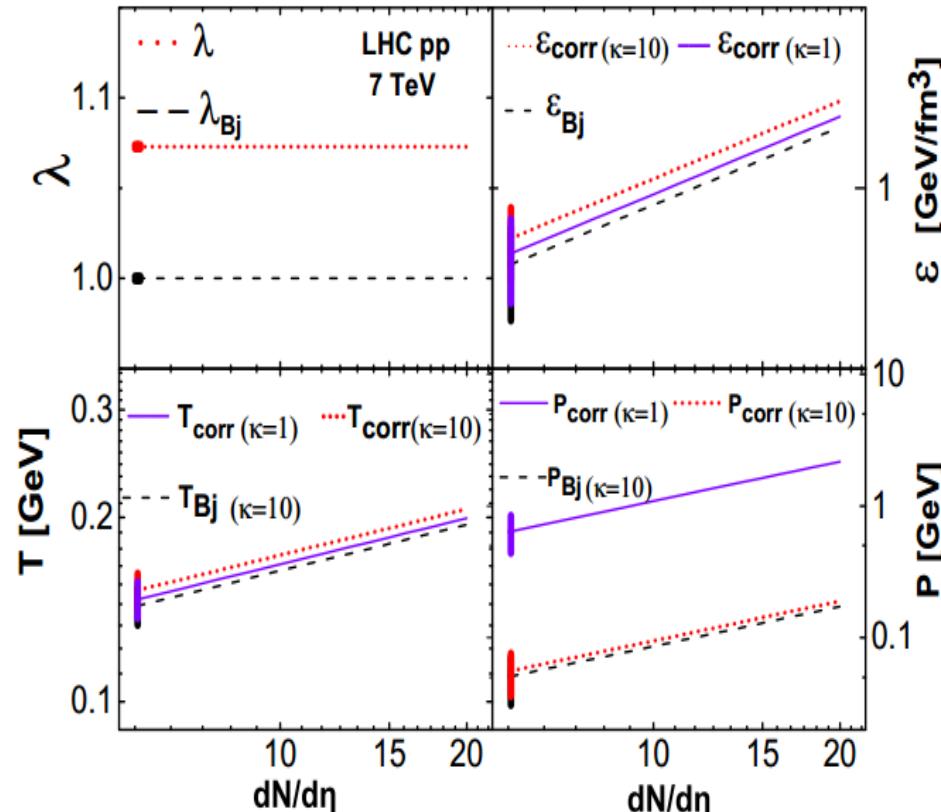
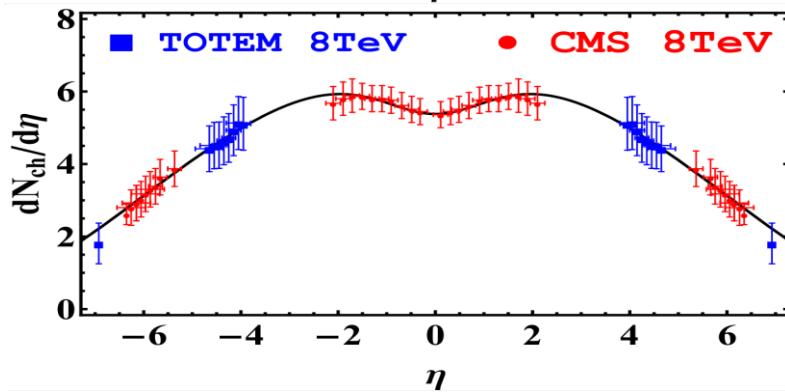
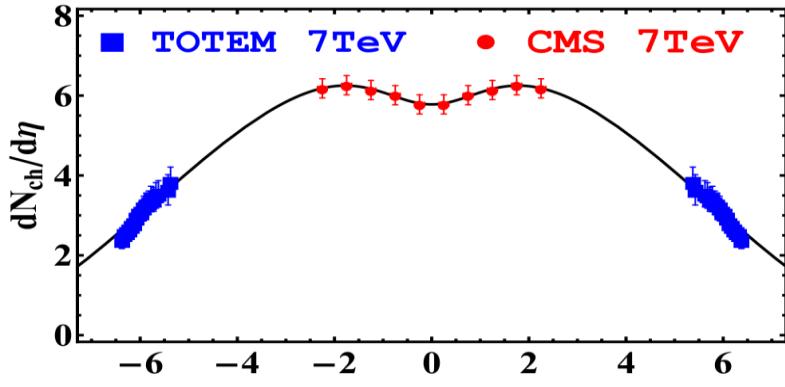


(M. Csanad, T. Csorgo, Z.F. Jiang, C.B. Yang, Universe. vol 3(2017). pp 1-9.)

$\sqrt{S}$	$\mathcal{E}_{Bj}$	$f_c$	$\mathcal{E}_{corr}$	$\lambda$	$c_s^2$	$dN / d\eta  _{\eta=0}$
7 TeV	0.507	1.262	0.640	1.073	0.10	5.895(NSD)
8 TeV	0.500	1.240	0.644	1.067	0.10	5.38(Inelastic)

# The initial energy density estimation

Results for 7 TeV & 8 TeV p+p collisions:

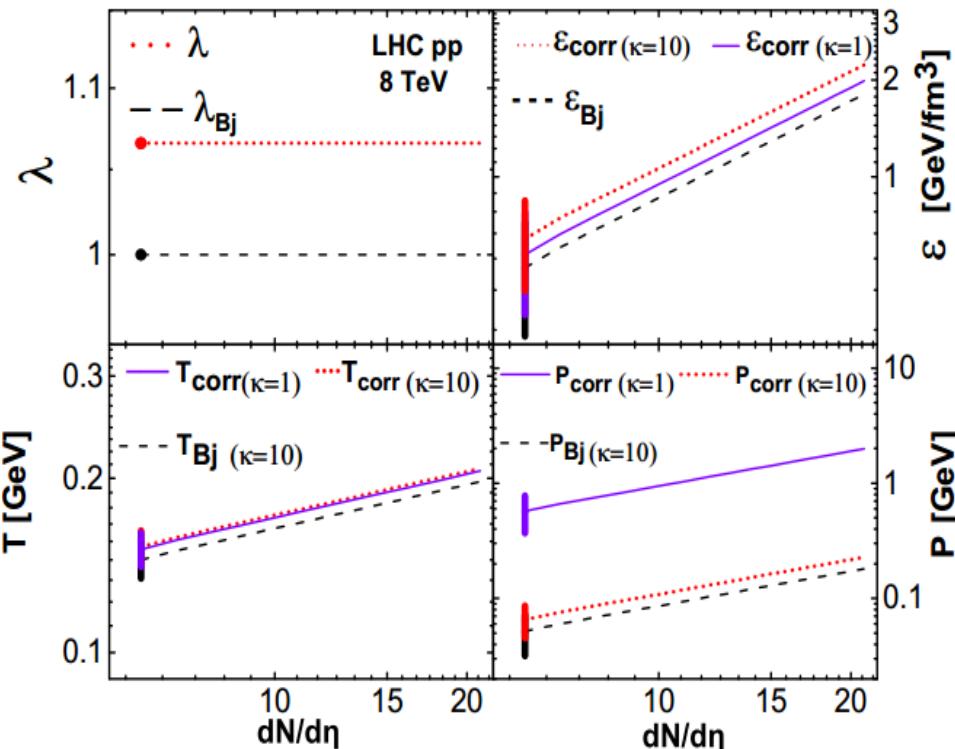
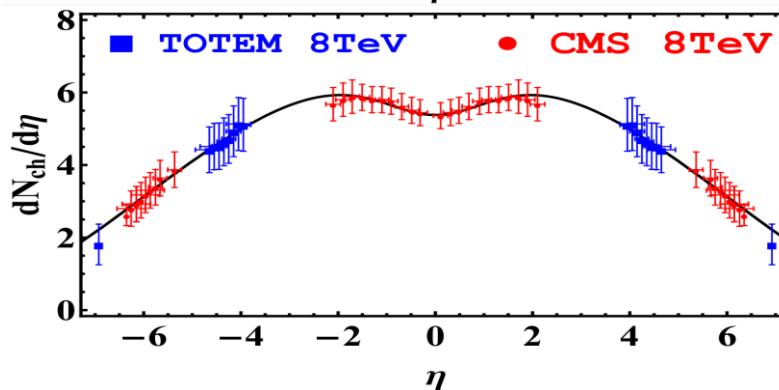
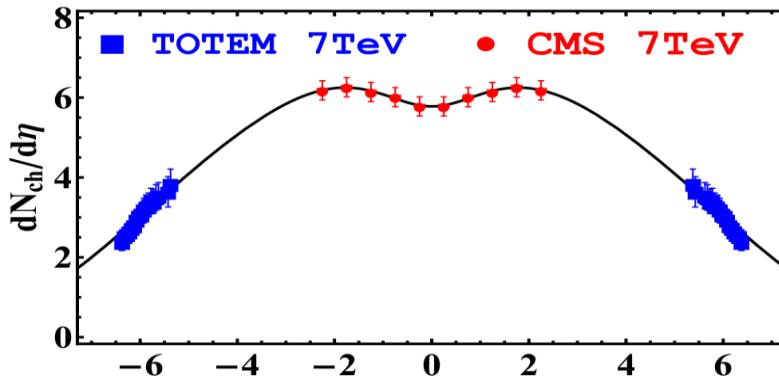


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# The CKCJ solutions

## The new CKCJ family of exact hydro solutions

$$\begin{aligned}\eta_x(H) &= \Omega(H) - H, \\ \Omega(H) &= \frac{\lambda}{\sqrt{\lambda-1}\sqrt{\kappa-\lambda}} \arctan\left(\sqrt{\frac{\kappa-\lambda}{\lambda-1}} \tanh(H)\right), \\ \sigma(\tau, H) &= \sigma_0 \left(\frac{\tau_0}{\tau}\right)^\lambda \mathcal{V}_\sigma(s) \left[1 + \frac{\kappa-1}{\lambda-1} \sinh^2(H)\right]^{-\frac{\lambda}{2}}, \\ T(\tau, H) &= T_0 \left(\frac{\tau_0}{\tau}\right)^{\frac{\lambda}{\kappa}} \mathcal{T}(s) \left[1 + \frac{\kappa-1}{\lambda-1} \sinh^2(H)\right]^{-\frac{\lambda}{2\kappa}}, \\ \mathcal{T}(s) &= \frac{1}{\mathcal{V}_\sigma(s)}, \\ s(\tau, H) &= \left(\frac{\tau_0}{\tau}\right)^{\lambda-1} \sinh(H) \left[1 + \frac{\kappa-1}{\lambda-1} \sinh^2(H)\right]^{-\lambda/2}.\end{aligned}$$

T. Csörgő, G. Kasza, M. Csanád, Z.F. Jiang, CKCJ solutions:  
arXiv: [1805.01427](https://arxiv.org/abs/1805.01427)

For more details, [see T. Csörgő's talk](#) at WPCF2018 on a new family of exact solutions of relativistic hydrodynamics.

# The CKCJ solutions

Final state observables:

$$\frac{dn}{dy} \approx \left. \frac{dn}{dy} \right|_{y=0} \cosh^{-\frac{1}{2}\alpha(\kappa)-1} \left( \frac{y}{\alpha(1)} \right) \exp \left( -\frac{m}{T_f} \left[ \cosh^{\alpha(\kappa)} \left( \frac{y}{\alpha(1)} \right) - 1 \right] \right),$$

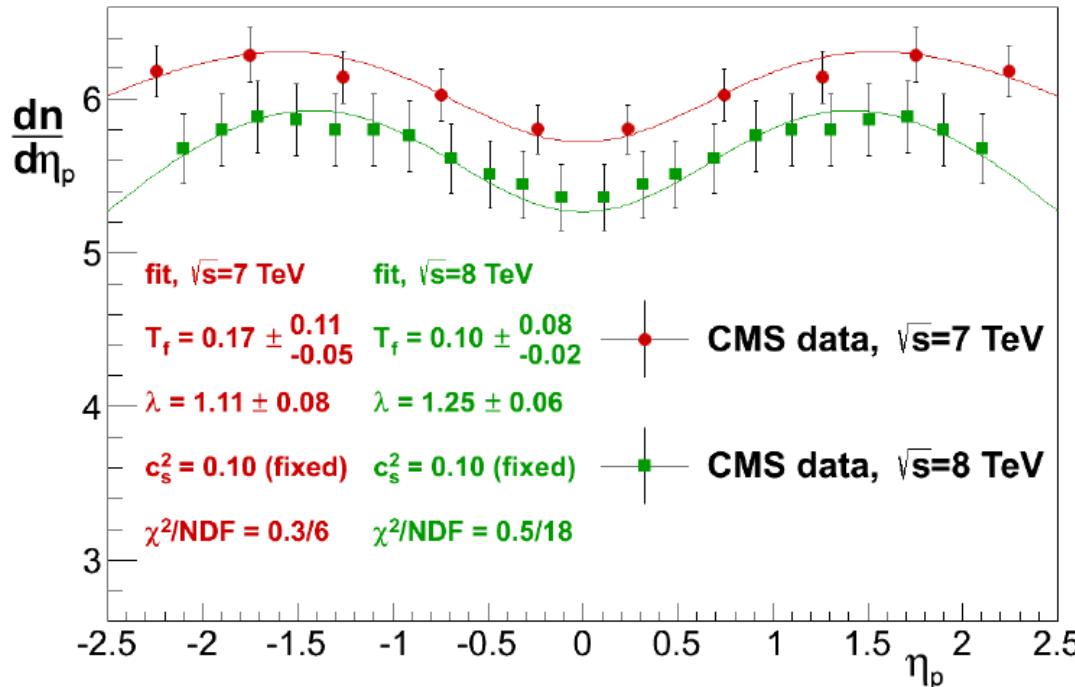
$$\frac{dn}{d\eta_p} \approx \left. \frac{dn}{dy} \right|_{y=0} \frac{\langle p_T(y) \rangle \cosh(\eta_p)}{\sqrt{m^2 + \langle p_T(y) \rangle^2 \cosh(\eta_p)}} \cosh^{-\frac{1}{2}\alpha(\kappa)-1} \left( \frac{y}{\alpha(1)} \right) \exp \left( -\frac{m}{T_f} \left[ \cosh^{\alpha(\kappa)} \left( \frac{y}{\alpha(1)} \right) - 1 \right] \right),$$

T. Csörgő, G. Kasza, M. Csanád, Z.F. Jiang, CKCJ solutions:  
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For more details, [see T. Csörgő's talk](#) at WPCF2018 on a new family of exact solutions of relativistic hydrodynamics.

# The CKCJ solutions

Results for 7 TeV & 8 TeV p+p collisions at CMS:



For a new method to estimate the initial energy density from the CKCJ solution: see [G.Kasza's talk](#) at WPCF 2018.

See May 22th, [G.Kasza's talk](#) talk: Initial energy density from new, exact solutions of relativistic hydrodynamics.

# Relativistic accelerating viscous hydrodynamics

Longitudinal acceleration effect makes the fluid cooling **faster**,

The viscosity will creating heat and makes the fluid cooling **slower**.

$$T^{\mu\nu} = e u^\mu u^\nu - (p + \boxed{\Pi}) \Delta^{\mu\nu} + \boxed{\pi^{\mu\nu}}$$

$$u^\mu = (\cosh \Omega, 0, 0, \sinh \Omega) \quad \Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$$

**Shear viscosity tensor:**  $\pi^{\mu\nu}$

**Bulk viscosity:**  $\Pi$ .

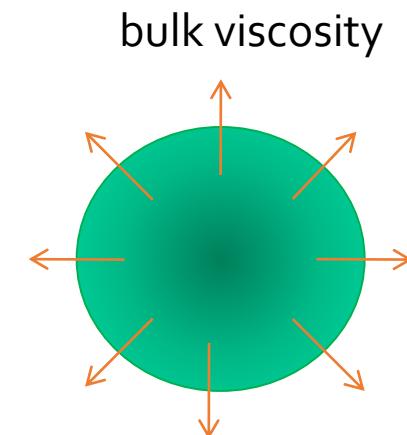
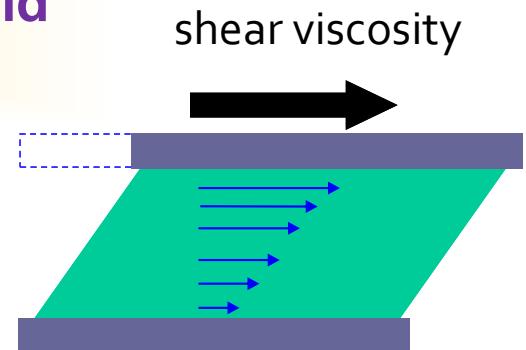
**Shear tensor:**

$$\sigma^{\mu\nu} \equiv \partial^{\langle\mu} u^{\nu\rangle} \equiv \left( \frac{1}{2} (\Delta_\alpha^\mu \Delta_\beta^\nu + \Delta_\beta^\mu \Delta_\alpha^\nu) - \frac{1}{d} \Delta^{\mu\nu} \Delta_{\alpha\beta} \right) \partial^\alpha u^\beta.$$

**The fundamental equations of viscous fluid:**

$$e = \kappa p, \quad \partial_v T^{\mu\nu} = 0.$$

**Assume:**  $n \approx 0, \quad \partial(\text{nu}^\mu) = 0.$



# Equations of viscosity hydrodynamic

The second law of thermodynamics:  $\partial_\mu S^\mu \geq 0$

$$\tau_\pi \Delta^{\alpha\mu} \Delta^{\beta\nu} \dot{\pi}_{\alpha\beta} + [\pi^{\mu\nu}] = 2\eta \sigma^{\mu\nu} - \frac{1}{2} \pi^{\mu\nu} \frac{\eta T}{\tau_\pi} \partial_\lambda \left( \frac{\tau_\pi}{\eta T} u^\lambda \right)$$

Israel-Stewart

$$\tau_\Pi \dot{\Pi} + [\Pi] = -\zeta (\partial \cdot u) - \frac{1}{2} \Pi \frac{\zeta T}{\tau_\Pi} \partial_\lambda \left( \frac{\tau_\Pi}{\zeta T} u^\lambda \right)$$

equations.

viscous hydro: near-equilibrium system

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equations.

viscous hydro: near-equilibrium system

The Navier-Stokes approximation,

$$\boxed{\pi^{\mu\nu} = 2\eta \sigma^{\mu\nu}}$$

$$\boxed{\Pi = -\zeta (\partial_\rho \mu^\rho)}$$

The shear viscosity and bulk viscosity,

Strongly coupled AdS/CFT prediction:  $\eta / s \geq 1/4\pi \approx 0.08$

D.T. Son, et.al. 05

Via lattice calculation:

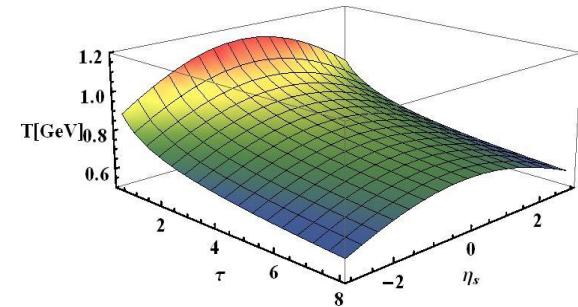
$$\boxed{\zeta / s \leq 0.015}$$

H.B. Meyer, et.al. 07 10.3717

# Solutions form viscous hydrodynamics

## The temperature profile:

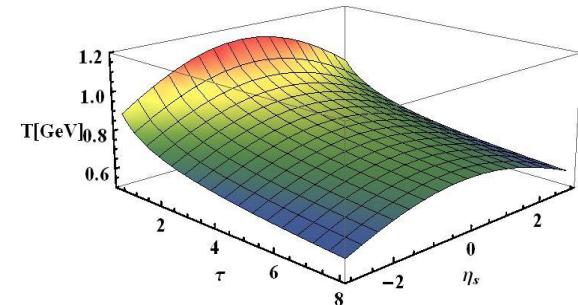
$$T(\tau, \eta_s) = T_0 \left( \frac{\tau_0}{\tau} \right)^{\frac{1+\epsilon}{\kappa}} \times \left[ \exp\left[-\frac{1}{2}\epsilon\left(1 - \frac{1}{\kappa}\right)\eta_s^2\right] + \frac{R_0^{-1}}{\kappa-1} \left( 2\epsilon + \exp\left[-\frac{1}{2}\epsilon\left(1 - \frac{1}{\kappa}\right)\eta_s^2\right] - (2\epsilon+1) \left( \frac{\tau_0}{\tau} \right)^{\frac{\kappa-\epsilon-1}{\kappa}} \right) \right]$$



# Solutions form hydrodynamic equations

## The temperature profile:

$$T(\tau, \eta_s) = T_0 \left( \frac{\tau_0}{\tau} \right)^{\frac{1+\epsilon}{\kappa}} \times \left[ \underbrace{\exp[-\frac{1}{2}\epsilon(1-\frac{1}{\kappa})\eta_s^2] + \frac{R_0^{-1}}{\kappa-1} \left( 2\epsilon + \exp[-\frac{1}{2}\epsilon(1-\frac{1}{\kappa})\eta_s^2] - (2\epsilon+1) \left( \frac{\tau_0}{\tau} \right)^{\frac{\kappa-\epsilon-1}{\kappa}} \right)}_{\text{Contribution from ideal terms.}} \right]$$



Contribution from ideal terms.  
One special case of our recent work.

*Contribution from viscous effect*

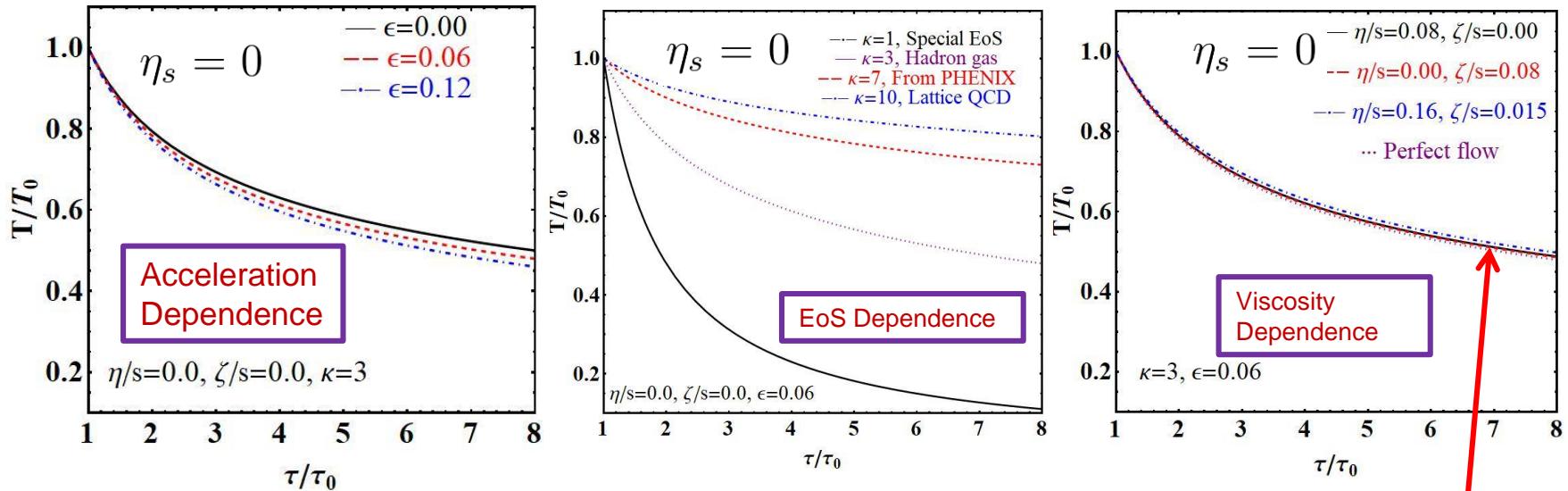
$$R_0^{-1} = \frac{\Pi_d}{T_0 \tau_0}$$

Reynolds number  
[A. Muronga, arXiv: 0309055 ]

→ { A non-zero Reynolds numbers  $R_0^{-1}$  makes cooling rate smaller,  
A non-vanishing acceleration  $\epsilon$  makes the cooling rate larger.

Open question: setting viscosity as the perturbative term.

# Temperature evolution



- Acceleration effect comes from the pressure gradient, makes the cooling ratio **larger** than non-acceleration flow.
- [**M. Nagy, T. Csorgo, M. Csand: arXiv:0709.3677v1**]
- EoS is an important modified factor.  
 $\kappa=1$  a very special case, CNC solution.  
 $\kappa=7$  comes from [**PHENIX, arXiv:nucl-ex/0608033v1**].
- Viscosity effect make the cooling rate smaller. [**H. Song, S. Bass, U. Heinz. et, PRL2011**]

The acceleration effect is almost fully compensated by the viscous contribution.

# A brief summary and outlook

## Summary:

1. By using exact accelerating solutions of perfect flow, the **initial thermodynamics quantities** are estimated.
2. The **perturbative solution** with viscous correction are obtained.

## Outlook:

1. 2rd I-S problem, rotation, CLVisc 3+1D code;
2. Jet/heavy quarkonium disturbance evolution on accelerating medium background...



# Thank you for your attention



[arXiv: 1609.07176, 1711.10740, 1805.01427...](#)