

Interferometry correlations in central p+Pb collisions

based on: <https://doi.org/10.1140/epjc/s10052-017-5482-5> (with P. Bożek)

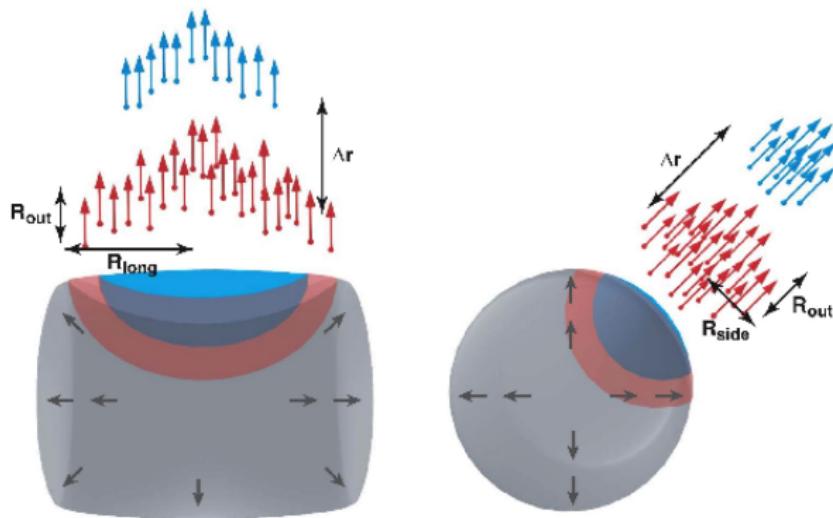
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XIII Workshop on Particle Correlations and Femtoscopy,
Kraków 2018

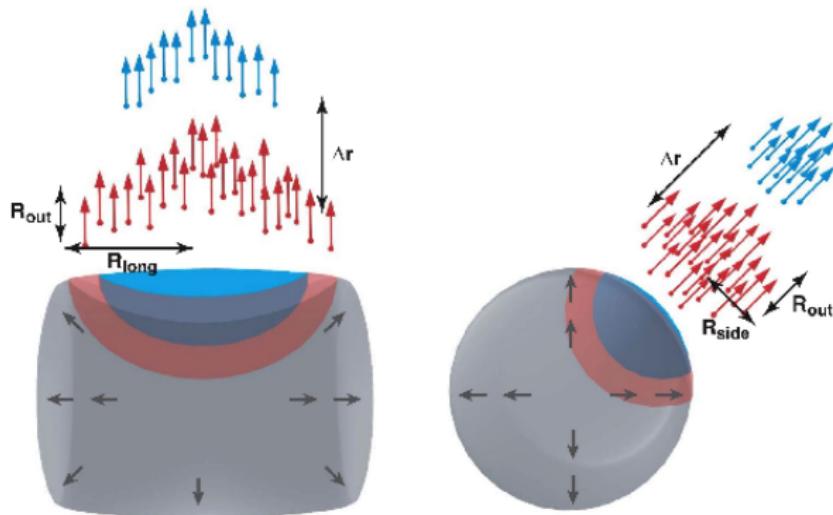
HBT radii and Bersch – Pratt coordinate system

M. A. Lisa and S. Pratt (2008), arXiv:0811.1352



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$$k_T = \frac{p_1 + p_2}{2} \quad q = p_1 - p_2$$

$$C(q) = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2)$$

Form of correlation function:

$$C(q, k) = \frac{\int d^4x_1 d^4x_2 S(x_1, p_1) S(x_2, p_2) |\Psi(q, x_1 - x_2)|^2}{\int d^4x_1 S(x_1, p_1) \int d^4x_2 S(x_2, p_2)}$$

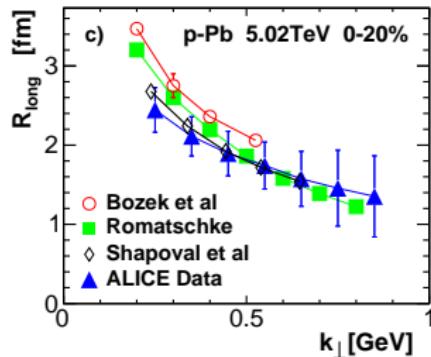
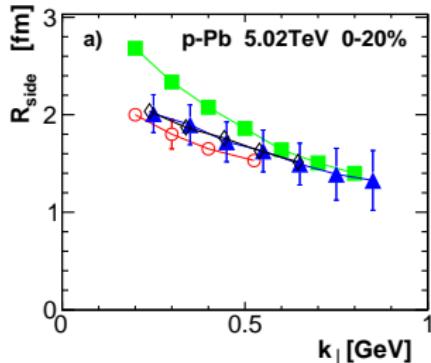
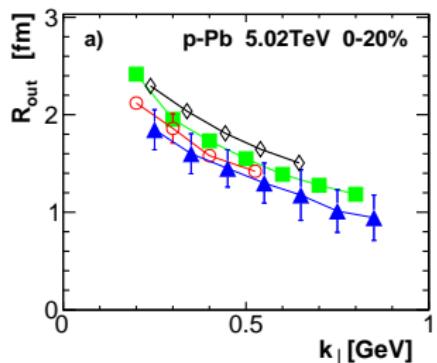
$$C(p^a, p^b) = \frac{\frac{dN^{ab}}{d^3p^a d^3p^b}}{\left(\frac{dN^a}{d^3p^a}\right)\left(\frac{dN^b}{d^3p^b}\right)}$$

$$C(q_a, k_b) = \frac{\frac{1}{N_{pairs,num}} \sum_{j=1}^{N_h} \sum_{l \neq m=1}^{N_e} \sum_{s=1}^{M_I} \sum_{f=1}^{M_m} \delta_{q_a} \delta_{k_b} |\Psi(q, x_1 - x_2)|^2}{\frac{1}{N_{pairs,den}} \sum_{i \neq j=1}^{N_h} \sum_{l,m=1}^{N_e} \sum_{s=1}^{M_I} \sum_{f=1}^{M_m} \delta_{q_a} \delta_{k_b}}$$

numerator: pion pairs from events generated from **the same** hypersurface
denominator: pion pairs from events generated from **different** hypersurfaces ($i \neq j$) – combinatorical background

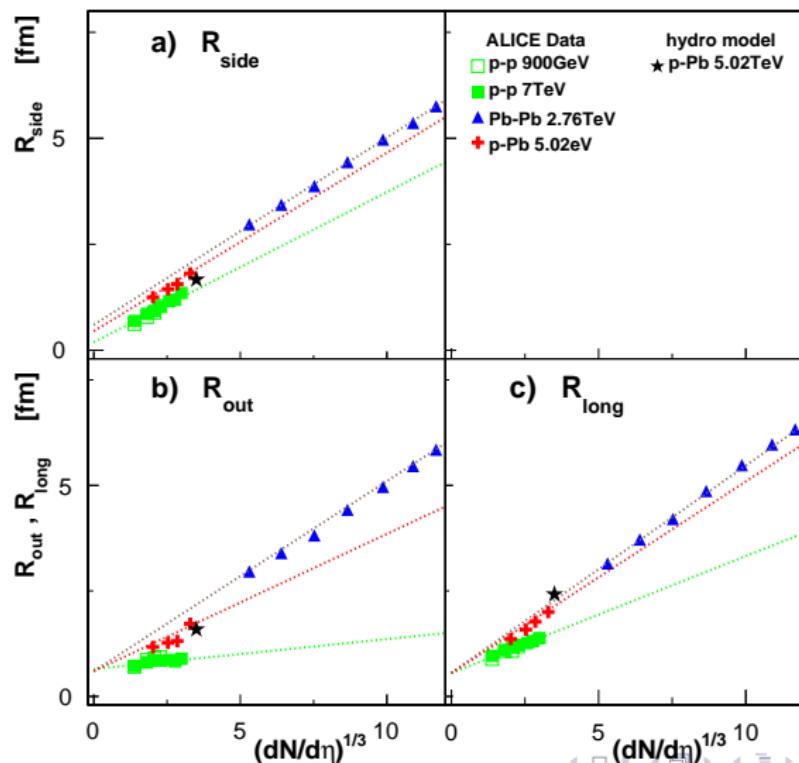
3 standard HBT radii in hydro

Hydro works in small systems



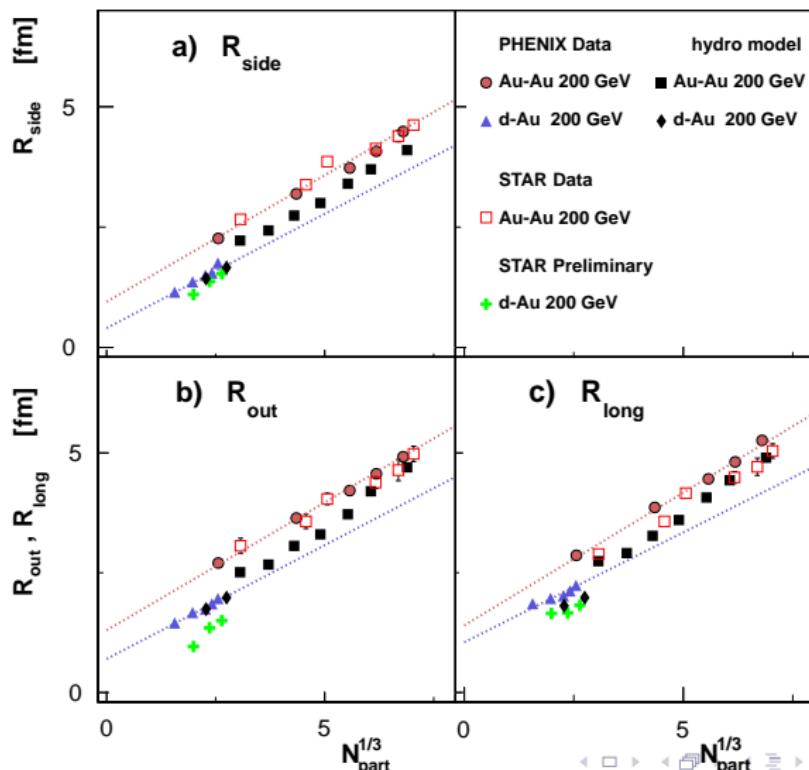
3 standard HBT radii in hydro – various systems

R_{out} , R_{side} , R_{long} are fairly well reproduced by hydro models.



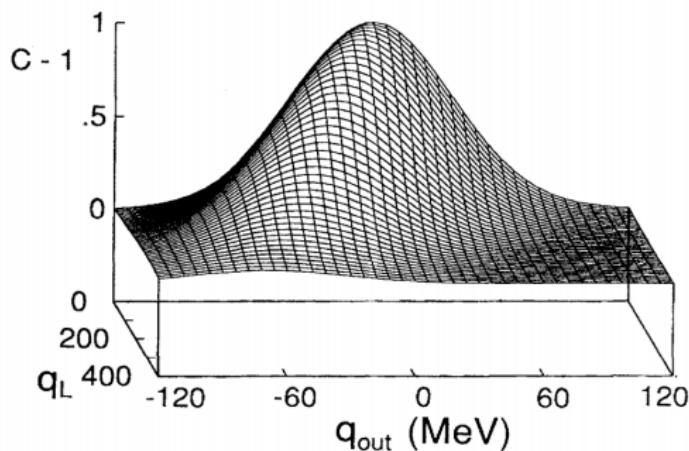
3 standard HBT radii in hydro – various systems

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$R_{out-long}$ cross term

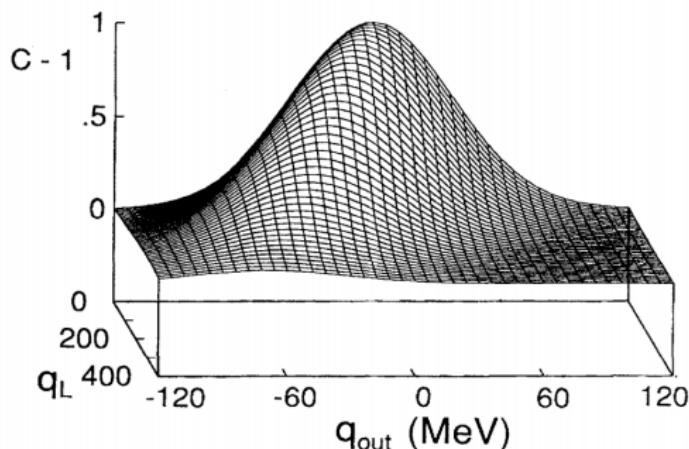
Proposed in *S. Chapman, et al, Phys. Rev. Lett. 74, 4400 (1995)*.



$$C(\mathbf{q}) = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2 - 2R_{ol}^2 q_o q_l)$$

$R_{out-long}$ cross term

Proposed in *S. Chapman, et al, Phys. Rev. Lett. 74, 4400 (1995)*.



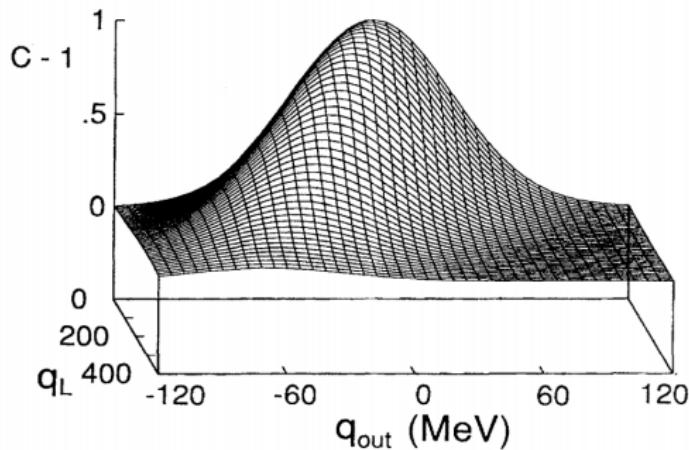
$$C(\mathbf{q}) = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2 - 2R_{ol}^2 q_o q_l)$$

In symmetric systems: $R_{ol} = 0$ in central rapidity.

$R_{ol} \neq 0$ in non-central rapidity.

$R_{out-long}$ cross term

Proposed in *S. Chapman, et al, Phys. Rev. Lett. 74,4400 (1995)*.



$$C(\mathbf{q}) = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2 - 2R_{ol}^2 q_o q_l)$$

$R_{ol} \neq 0$ also for asymmetric systems.

Cross terms

In azimuthally sensitive analyses (*arXiv:1408.1264*) :

$$R_{ol}^2 = H_1 + I_1 - G_0 \beta_{\perp} + (I_1 + I_3 - H_1 + H_3) \cos(2\Phi)$$

where, G_i, H_i, I_i :

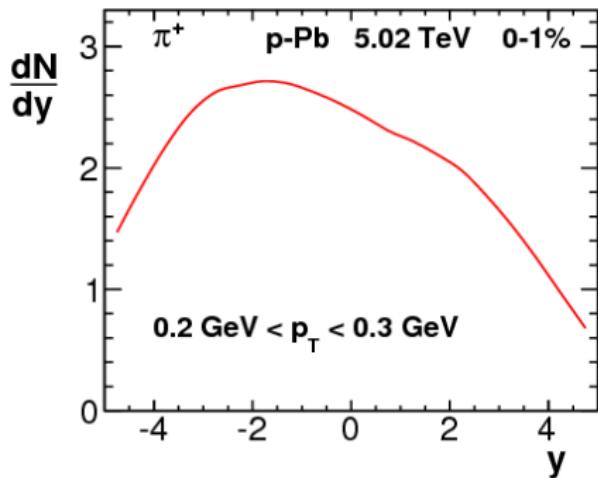
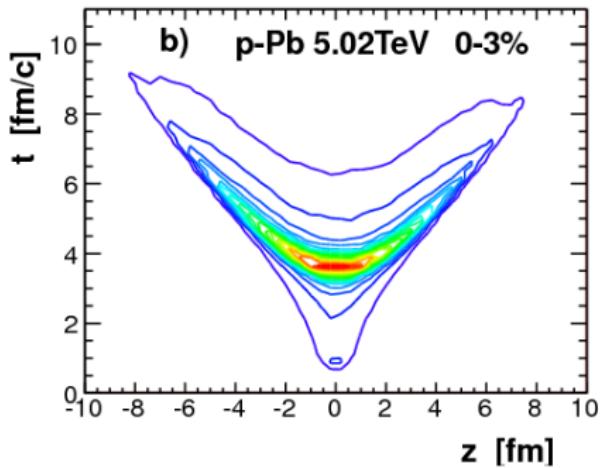
$$\langle tz \rangle - \langle t \rangle \langle z \rangle = \textcolor{red}{G_0} + 2 \sum_{n=2,4,\dots} G_n \cos(n\Phi)$$

$$\langle xz \rangle - \langle x \rangle \langle z \rangle = 2 \sum_{n=1,3,\dots} H_n \cos(n\Phi)$$

$$\langle yz \rangle - \langle y \rangle \langle z \rangle = 2 \sum_{n=1,3,\dots} I_n \sin(n\Phi)$$

Most significant among angle independent terms is G_0 , appearing in Fourier expansion of time – beam-along moment of the emission function $S(x, k)$

$R_{out-long}$ – consequence of asymmetry in time evolution of source \Rightarrow it results in not boost invariant y distribution

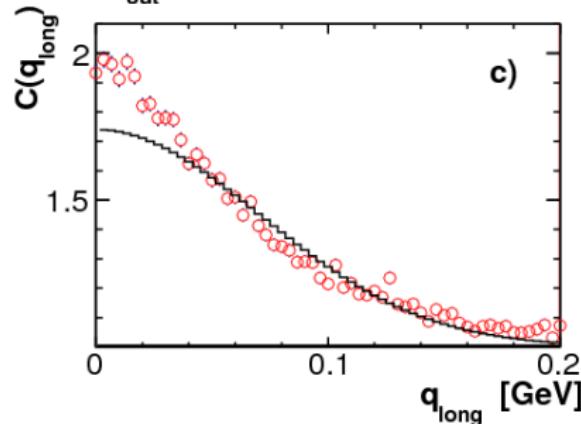
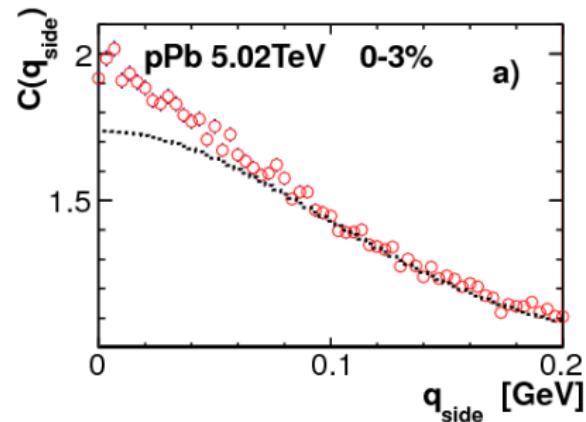
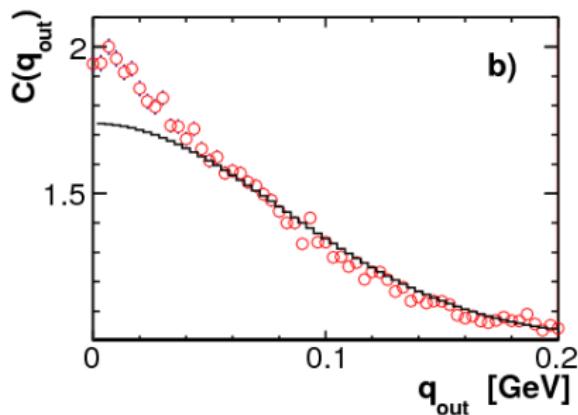


hydrodynamic emission source
in asymmetric $p+Pb$ collisions

$$\langle tz \rangle - \langle t \rangle \langle z \rangle \neq 0$$

Gaussian paramterization:

$$C(q, k) = 1 + \lambda \exp(-R_o^2 q_o^2 - R_s^2 q_s^2 - R_l^2 q_l^2 - 2R_{ol}^2 q_o q_l)$$



Non – gaussian parametrization used by ATLAS

- First measurements including out-long coupling
- First time radii as a function of rapidity

Non – gaussian parametrization used by ATLAS

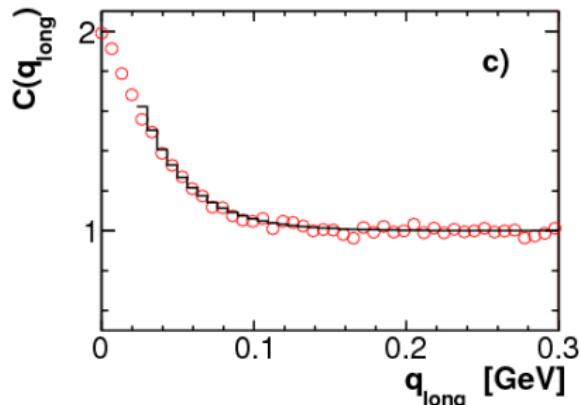
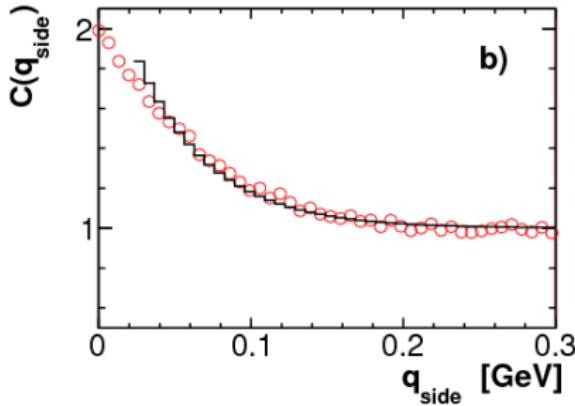
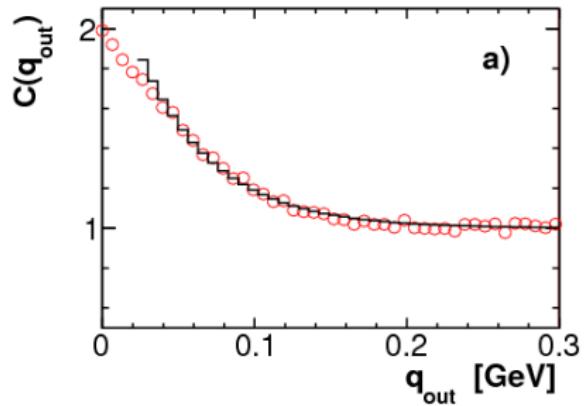
- First measurements including out-long coupling
- First time radii as a function of rapidity

$$C(\mathbf{q}) = 1 + \lambda \exp(-\|\mathbf{R}\mathbf{q}\|)$$

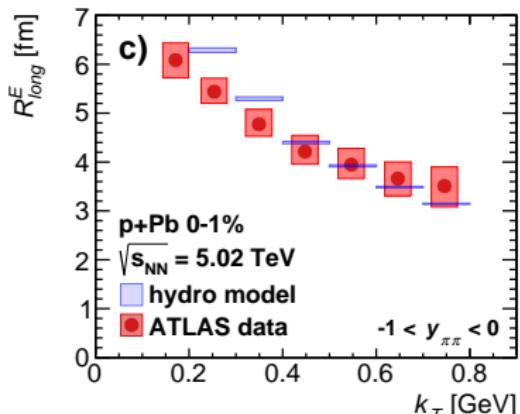
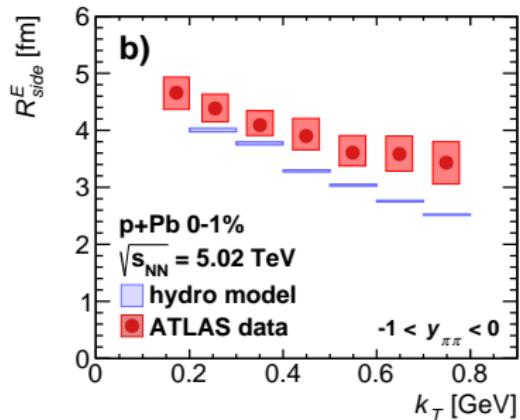
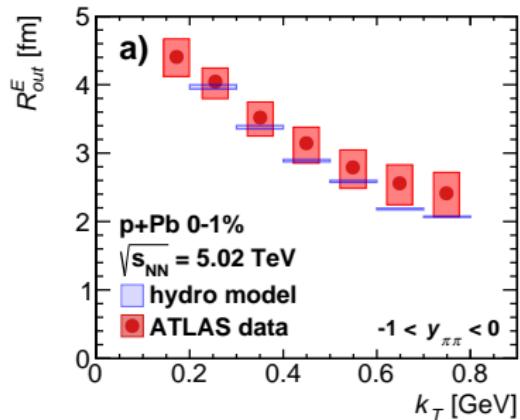
$$\begin{aligned} \|\mathbf{R}\mathbf{q}\| = & \left[\left(R_{out}^E q_{out} + R_{mix}^E q_{long} \right)^2 + \left(R_{side}^E q_{side} \right)^2 \right. \\ & \left. + \left(R_{long}^E q_{long} + R_{mix}^E q_{out} \right)^2 \right]^{1/2} \end{aligned}$$

$$\begin{aligned} \|\mathbf{R}\mathbf{q}\| = & \left[\left(R_{out}^E {}^2 + R_{mix}^E {}^2 \right) q_{out}^2 + R_{side}^E {}^2 q_{side}^2 + \left(R_{long}^E {}^2 + R_{mix}^E {}^2 \right) q_{long}^2 \right. \\ & \left. + 2 \left(R_{out}^E + R_{long}^E \right) R_{mix}^E q_{out} q_{long} \right]^{1/2} \end{aligned}$$

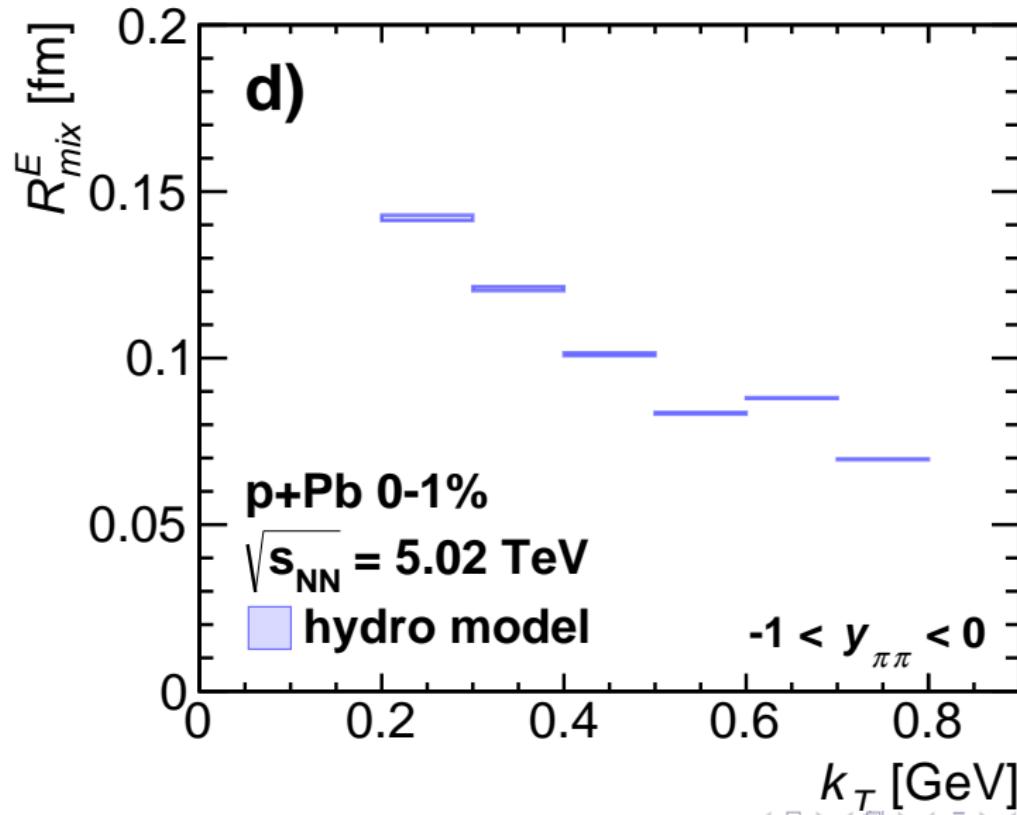
Non – gaussian parametrization used by ATLAS



Results – $R_{o,s,l}^E(k_T)$

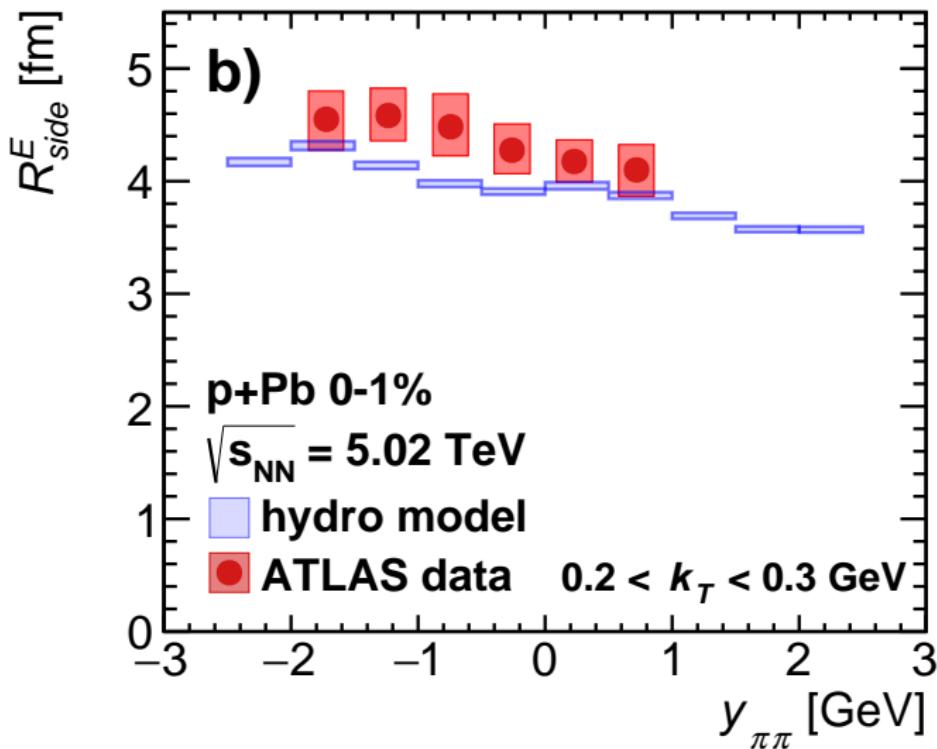


Results – $R_{mix}^E(k_T)$



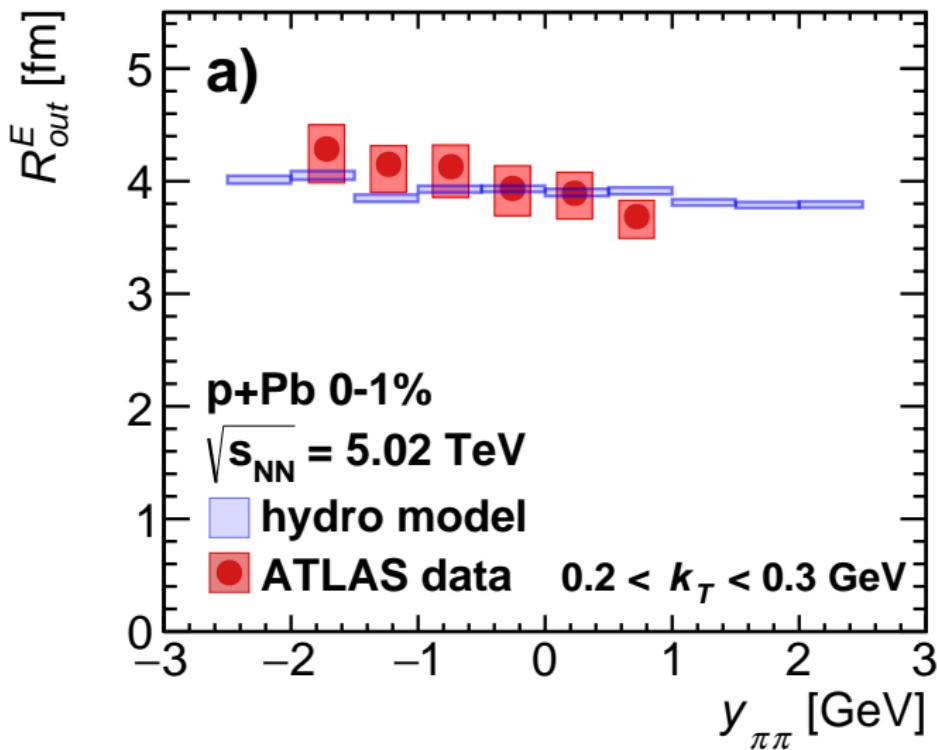
Results – $R_{side}^E(y)$

first measurement of y -dependent HBT radii



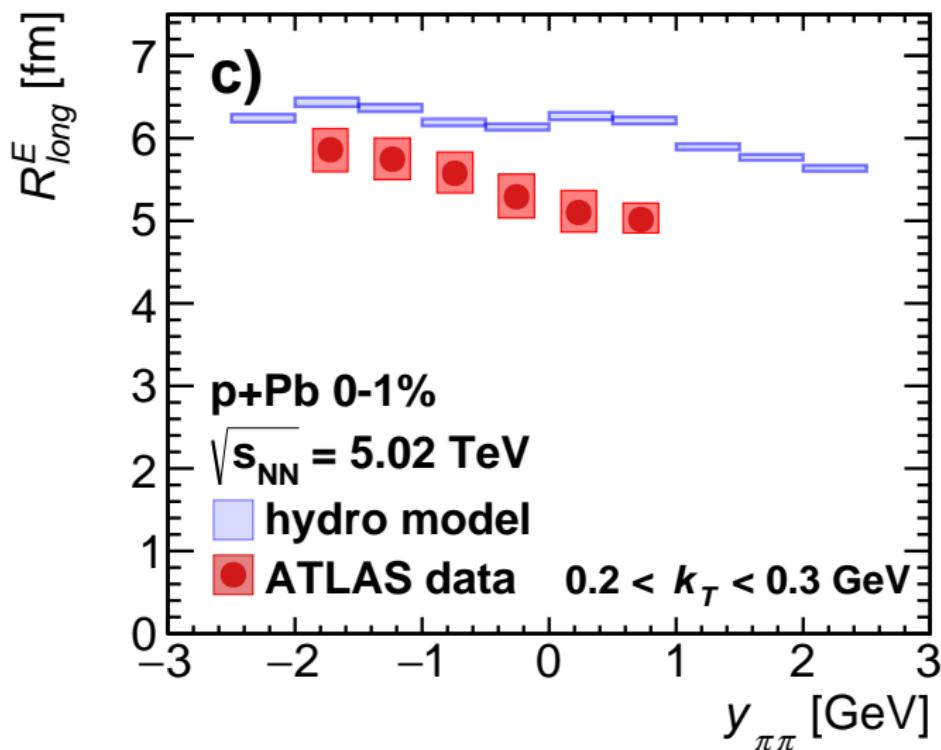
Results – $R_{out}^E(y)$

first measurement of y -dependent HBT radii

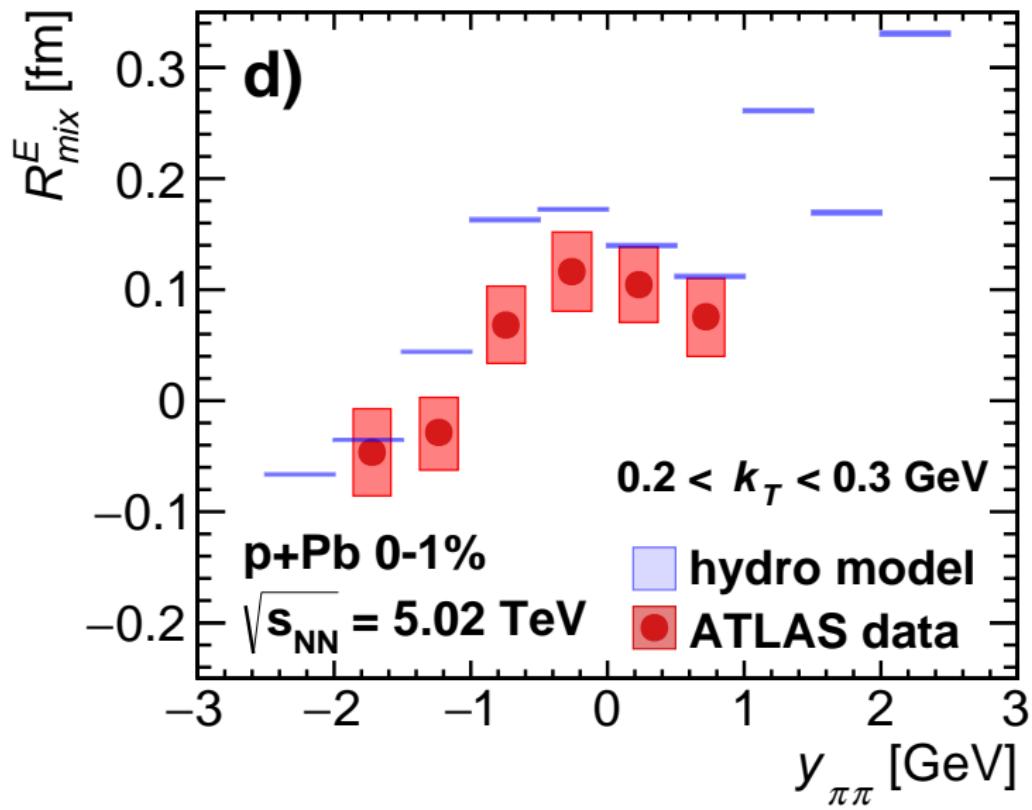


Results – $R_{long}^E(y)$:

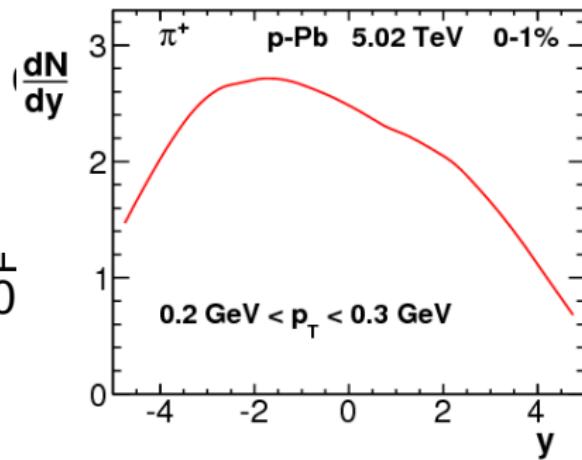
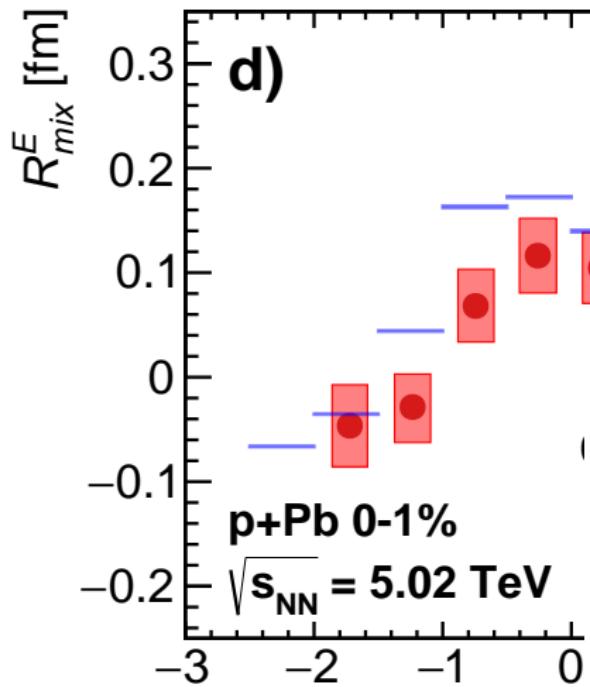
first measurement of y -dependent HBT radii



Results – $R_{mix}^E(y)$



Results – $R_{mix}^E(y)$



Conclusions

- hydrodynamic models can describe (semi-)quantitatively HBT radii dependence on k_T in pPb collisions – both in gaussian and exponential parametrization
- rapidity dependence of HBT radii can be reproduced partly-qualitatively in simulations
- hydro models can explain out-long cross term in asymmetric collisions