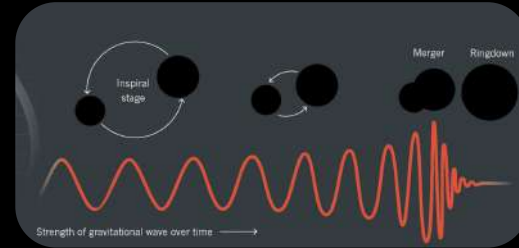


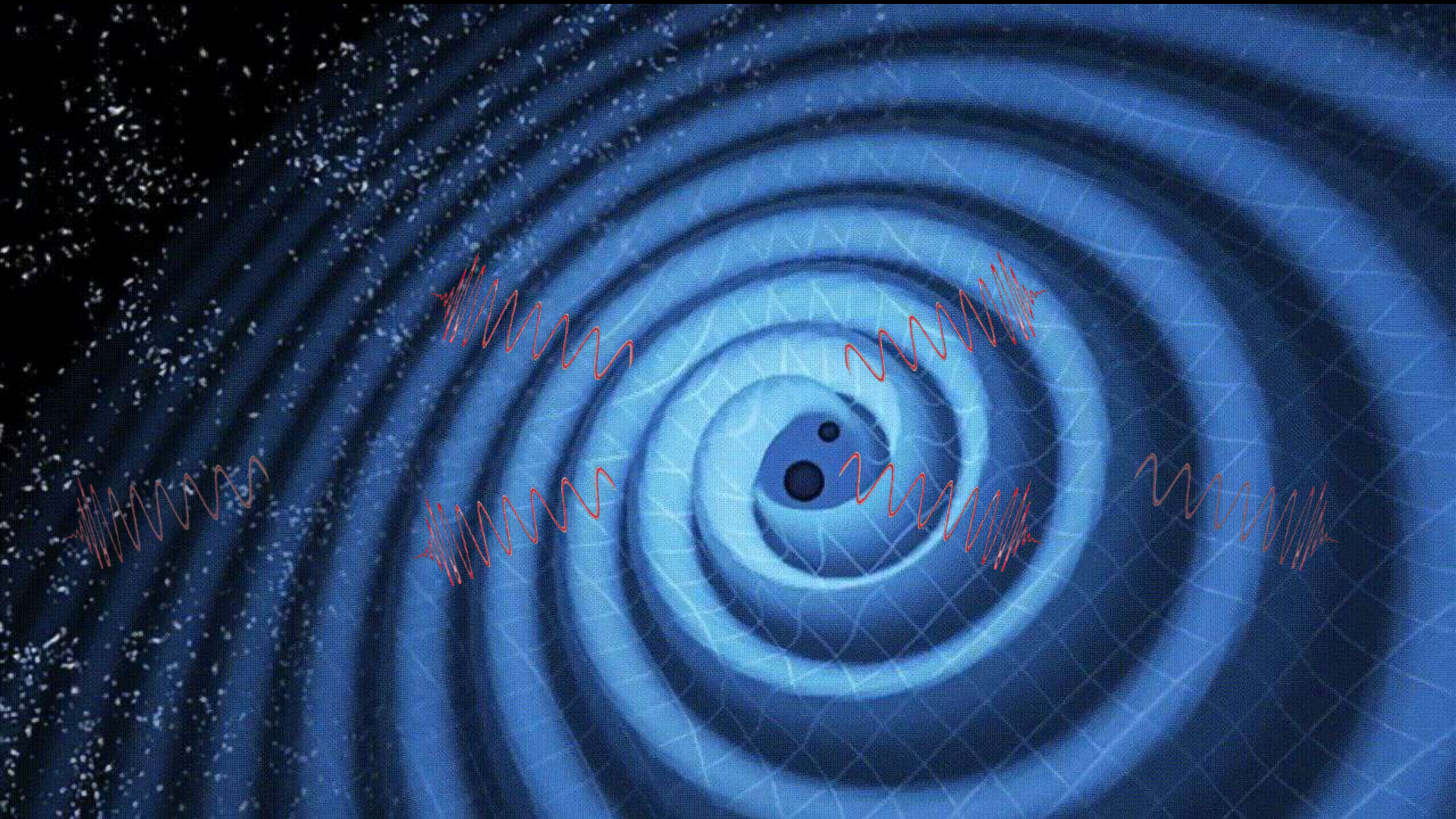
# Gravitational Waves: Principles, Detection, and Recent Discoveries

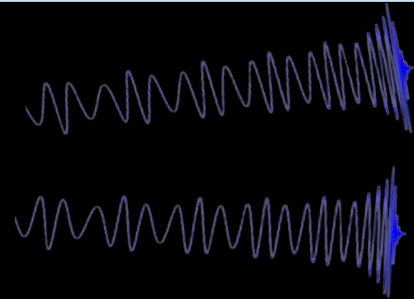
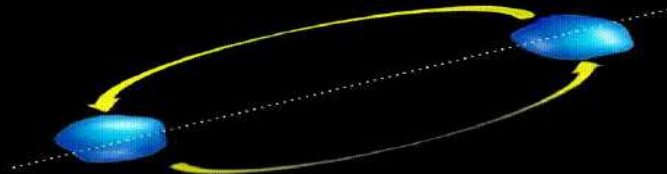
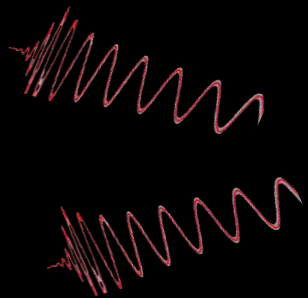
$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$



$$h = Ae^{i\Psi}$$

**Syed Naqvi, Post-Doc, NZ-15**  
**Einstein Telescope member (David E. Alvarez-Castillo group)**

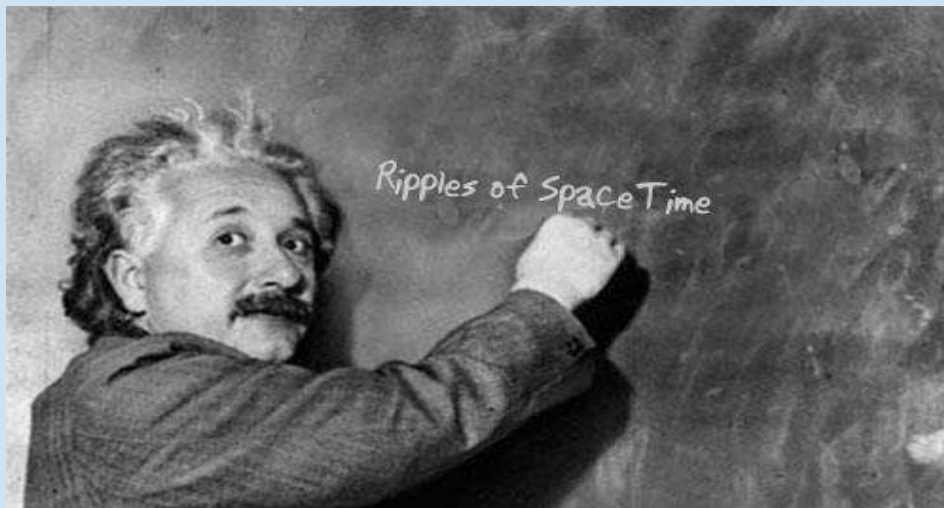




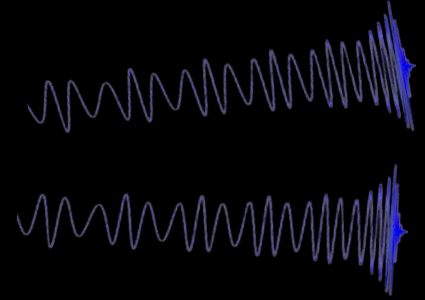
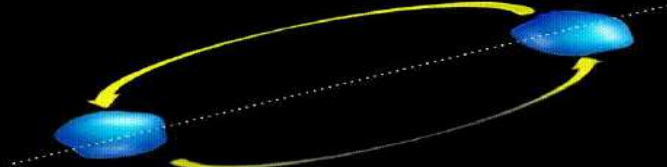
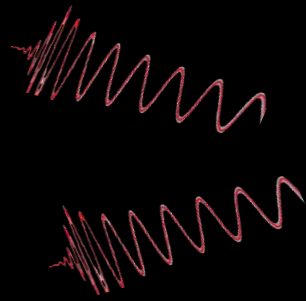
## Contents

:

- Einstein's journey to the famous quadrupole formula



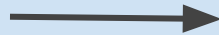
$$h \sim \ddot{Q}$$



## Contents

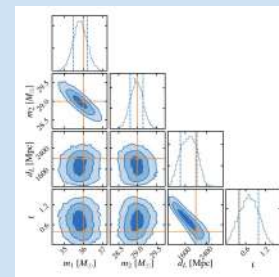
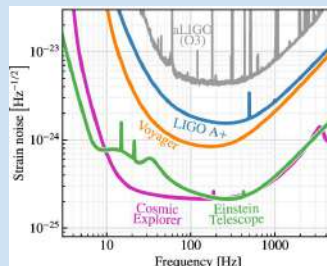
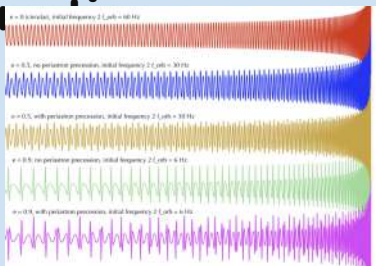
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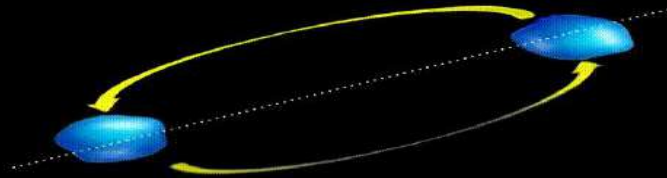
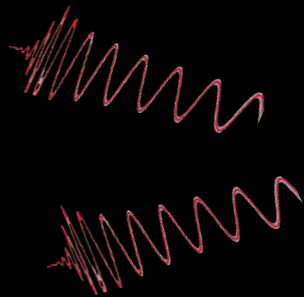
- Einstein's journey to the famous quadrupole formula
- Weber bar's to current Interferometric detectors



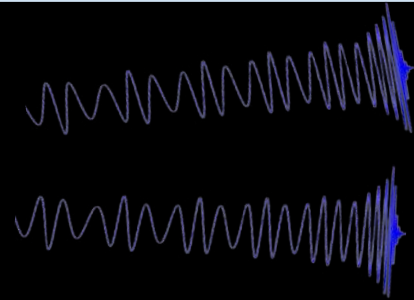
# Contents

- Einstein's journey to the famous quadrupole formula
- Weber bar's to current Interferometric detectors
- Gravitational Waves: Template banks + Parameter Estimation





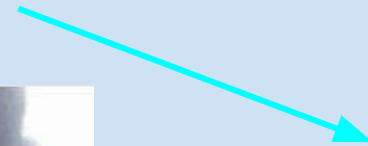
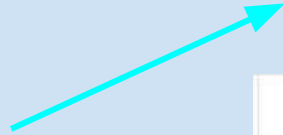
## Contents



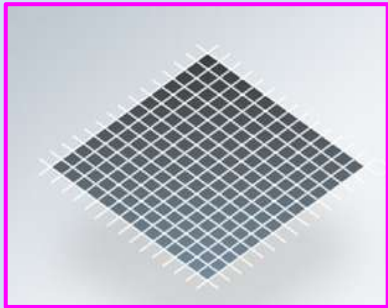
:

- **Einstein's journey to the famous quadrupole formula**
- **Weber bar's to current Interferometric detectors**
- **Gravitational Waves: Template banks + Parameter Estimation**
- **Recent results : O3 i O4 Catalogue**

# From Special Relativity to General Relativity



$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$$



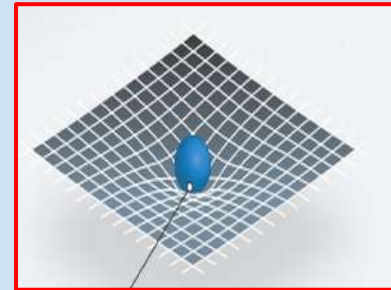
Gravity according  
to Newton.



Gravity According  
to Einstein.



$$ds^2 = g_{\mu\nu}(x) dx^\mu dx^\nu$$



Space + Time  
=>  
Spacetime

Gravity=>  
Bending of  
spacetime

# Einstein Field Equations : Meme Version

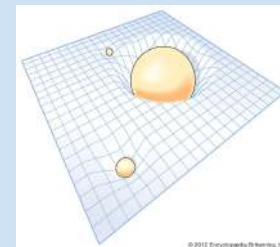


- Non-linear nature, many types of spacetimes(compact objects, universe expanding etc...)

# Different spacetimes represent different settings...



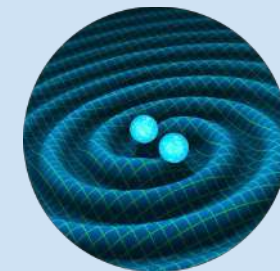
Pierogi Ruskie: Spacetimes of heavy objects Black Holes, NS etc....



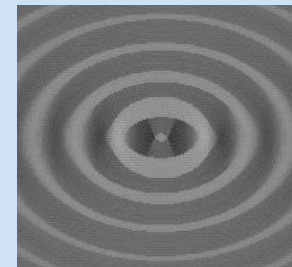
$$ds^2 = g_{\mu\nu}(x)dx^\mu dx^\nu$$



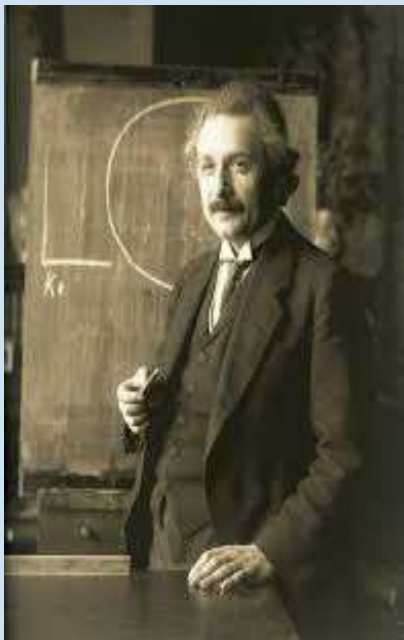
Pierogi ze szpinakiem i serem feta:  
Spacetimes having radiation/Grav Waves from astrophysical sources



Indyjskie pierogi: Spacetimes having  
Standing Grav. Waves



# From doubt...



Einstein



Rosen

- ❖ 1905 : Henri Poincare
- ❖ 1915-16: Einstein linearized gravity
- ❖ Issues :
  - > Plane GWs in full theory?
  - > Do full Einstein Equation have solution which can be interpreted a GW?
  - > Do GWs carry energy?
  - > .....
- ❖ 1937 : Einstein- Rosen metric, exact solution but with singularities

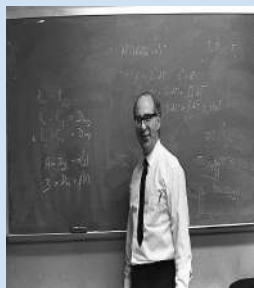
# To belief...



Bondi



Pirani



Robinson



Trautmann

- ❖ **1959, Bondi-Pirani-Robinson :**
  - > the plane wave in the full theory is defined ✓
  - > found as a class of solutions of Einstein fields equations ✓
  - > they carry energy in a form of a sandwich wave which affects test particles ✓
  
- ❖ **1958, Andrzej Trautman:**
  - > Radiation is nonlocal
  - > defining GW in full Einstein theory = boundary conditions at infinity

# GWs = Ripples of spacetime = Perturbations

Einstein Field Equations in LINEARIZED form : *weak field*

$$||h_{ab}|| \ll 1$$

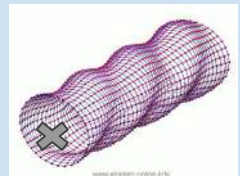
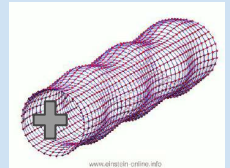
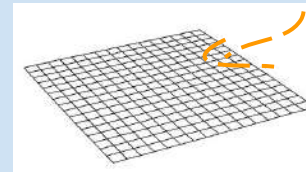
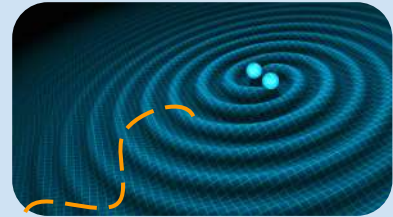
Metric Perturbations

$$g_{ab} = \eta_{ab} + h_{ab}$$

Wave Equation  
(TT gauge)

$$\square \bar{h}_{ab} = 0$$

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$



# Hartle Book: Rough GW and EM analogy

|                                                                             | Linearized<br>Gravitation                                                                             | Electromagnetism                                                                                          |
|-----------------------------------------------------------------------------|-------------------------------------------------------------------------------------------------------|-----------------------------------------------------------------------------------------------------------|
| Basic "potentials"                                                          | Linearized metric<br>perturbation<br>$h_{\alpha\beta}(x)$                                             | Vector and scalar<br>potentials<br>$(\Phi(t, \vec{x}), \vec{A}(t, \vec{x}))$                              |
| Field quantities                                                            | Linearized Riemann<br>curvature<br>$\delta R_{\alpha\beta\gamma\delta}(x)$                            | Electric and magnetic<br>fields<br>$\vec{E}(t, \vec{x}), \vec{B}(t, \vec{x})$                             |
| Gauge transformation<br>leading to new<br>potentials but<br>the same fields | $h_{\alpha\beta} \rightarrow h_{\alpha\beta} - \partial_\alpha \xi_\beta - \partial_\beta \xi_\alpha$ | $\vec{A} \rightarrow \vec{A} + \nabla \Lambda$<br>$\Phi \rightarrow \Phi - \partial \Lambda / \partial t$ |
| Example of a gauge<br>condition                                             | Lorentz gauge<br>$\partial_\beta h^\beta_\alpha - \frac{1}{2} \partial_\alpha h^\beta_\beta = 0$      | Lorentz condition<br>$\vec{\nabla} \cdot \vec{A} + \partial \Phi / \partial t = 0$                        |
| Field equations<br>simplified by the<br>gauge condition                     | $\square h_{\alpha\beta} = 0$                                                                         | Maxwell's equations<br>$\square \vec{A} = 0$<br>$\square \Phi = 0$                                        |

# Generation of G.Ws = Quadrupole radiation

Including matter source term

$$\square \bar{h}_{ab} = -16\pi T_{ab} \quad \longrightarrow \quad \bar{h}_{ab}(t, \mathbf{x}) = 4 \int d^3x' \frac{T_{ab}(t - |\mathbf{x} - \mathbf{x}'|, \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|}$$

Quadrupole Formula  
(in units c,G)

$$\bar{h}^{ij}(t, r) = \frac{2G}{r c^4} \frac{d^2}{dt^2} Q^{ij} \left( t - \frac{r}{c} \right)$$

Reduced Quadrupole Moment

$$Q^{jk} = I^{jk} - \frac{1}{3} \delta^{jk} \delta_{lm} I^{lm}$$

$$I^{jk} = \int d^3x \rho(t, \vec{x}) x^j x^k.$$

Moment of  
Inertia tensor

$$I_p = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix}$$

# Generation of G.Ws = Quadrupole radiation

Including matter source term

$$\square \bar{h}_{ab} = -16\pi T_{ab} \quad \longrightarrow \quad \bar{h}_{ab}(t, \mathbf{x}) = 4 \int d^3x' \frac{T_{ab}(t - |\mathbf{x} - \mathbf{x}'|, \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|}$$

Quadrupole Formula  
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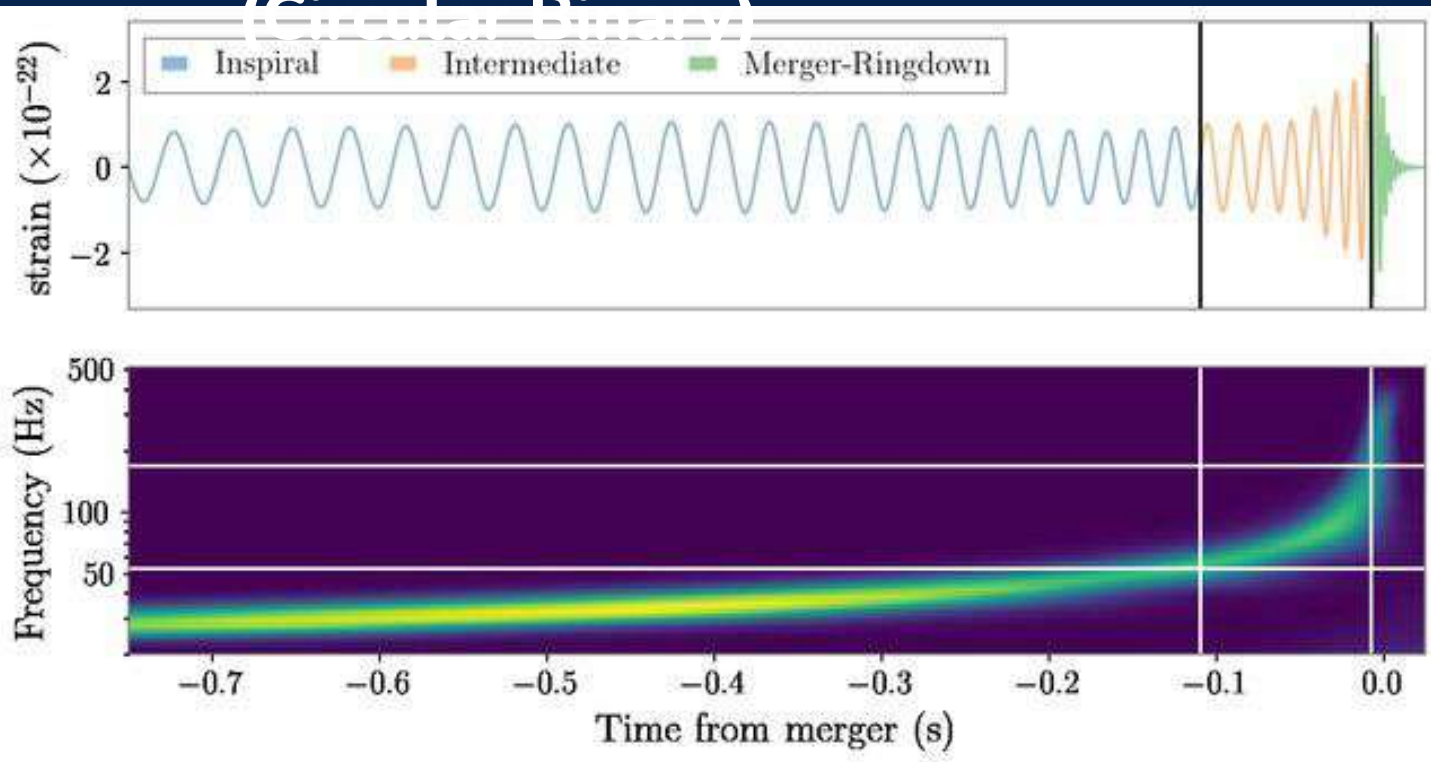
GW Luminosity

$$L_{GW} = -\frac{dE}{dt} = \frac{1}{5} \boxed{\frac{G}{c^5}} \left\langle \frac{\partial^3 Q_{ij}}{\partial t^3} \frac{\partial^3 Q_{ij}}{\partial t^3} \right\rangle$$

Energy carried away by GWs per second ~ Peak luminosity @ BH merger :  **$10^{50}$  watts!!!**

# Gravitational Waveform: Time Domain vs Freq Domain

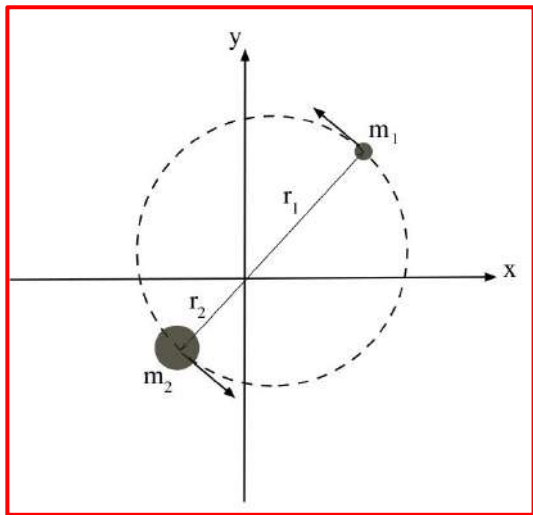
(Circular Binary)



“Chirp/ćwierkanie”  
Signal



# GWs from binary system in **Circular** orbit



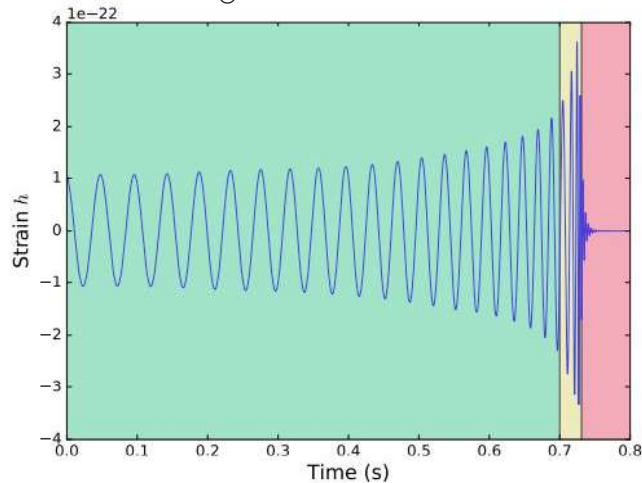
Radiated power:  
Evaluated from Reduced Quadrupole Moment

$$\langle E_{GW} \rangle \approx \langle \dot{h}_+^2 + \dot{h}_\times^2 \rangle$$

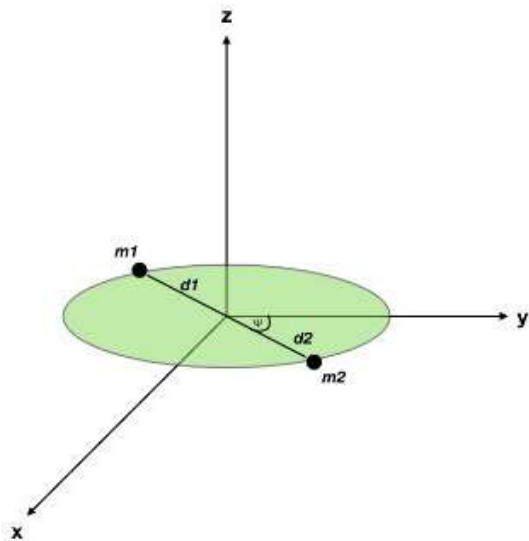


$$\langle E_{GW} \rangle = \frac{32}{5} \frac{G^4}{c^5} \frac{m_1^2 m_2^2 (m_1 + m_2)}{a}$$

35-30  $M_\odot$  : quasi circular case



# GWs from binary system in **Eccentric** orbit



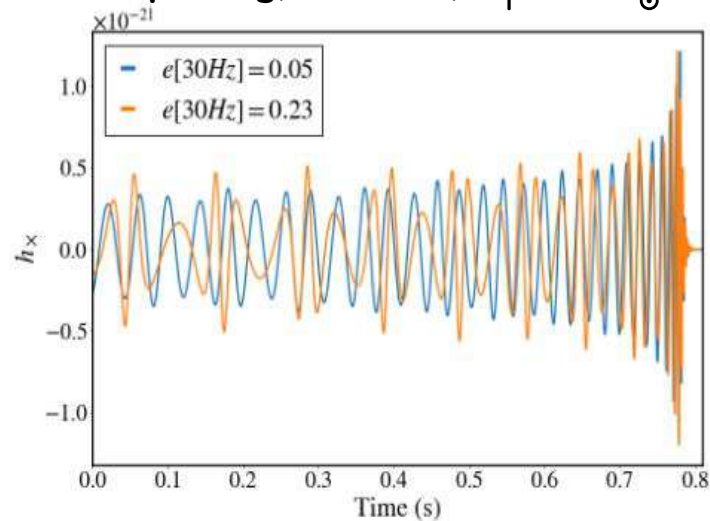
Radiated power:  
Evaluated from Reduced Quadrupole Moment

$$\langle E_{GW} \rangle \approx \langle \dot{h}_+^2 + \dot{h}_\times^2 \rangle$$



$$\langle E_{GW} \rangle = \frac{32}{5} \frac{G^4}{c^5} \frac{m_1^2 m_2^2 (m_1 + m_2)}{a} f(e)$$

Non-spinning, eccentric,  $M_T = 50M_\odot$



# First Indirect GW detection: Hulse–Taylor binary pulsar ( PSR 1913+16)

- Discovered in 1974; binary neutron-star system: one star is a millisecond pulsar
- Orbital period: ~7.75 hours
- Eccentric orbit → strong relativistic effects
- Provided the first evidence that gravitational waves carry energy

## PSR 1913+16

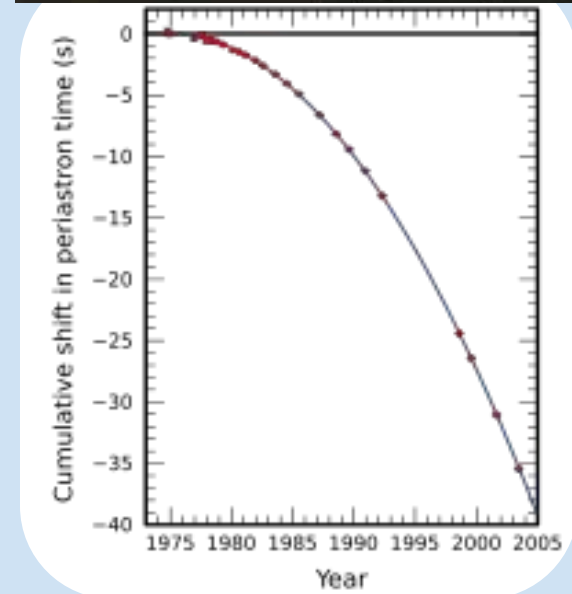
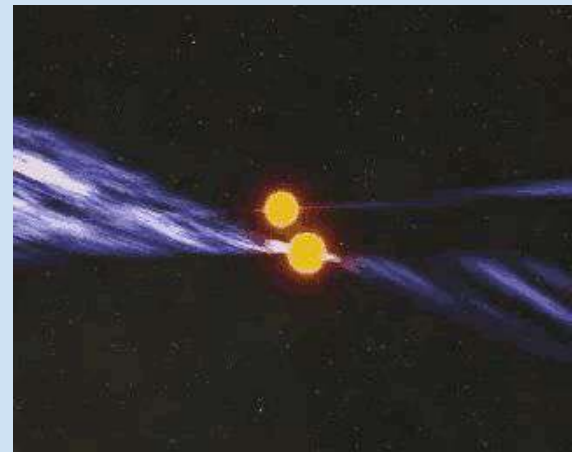
$$l_0 \sim 10^{11} \text{ cm}$$

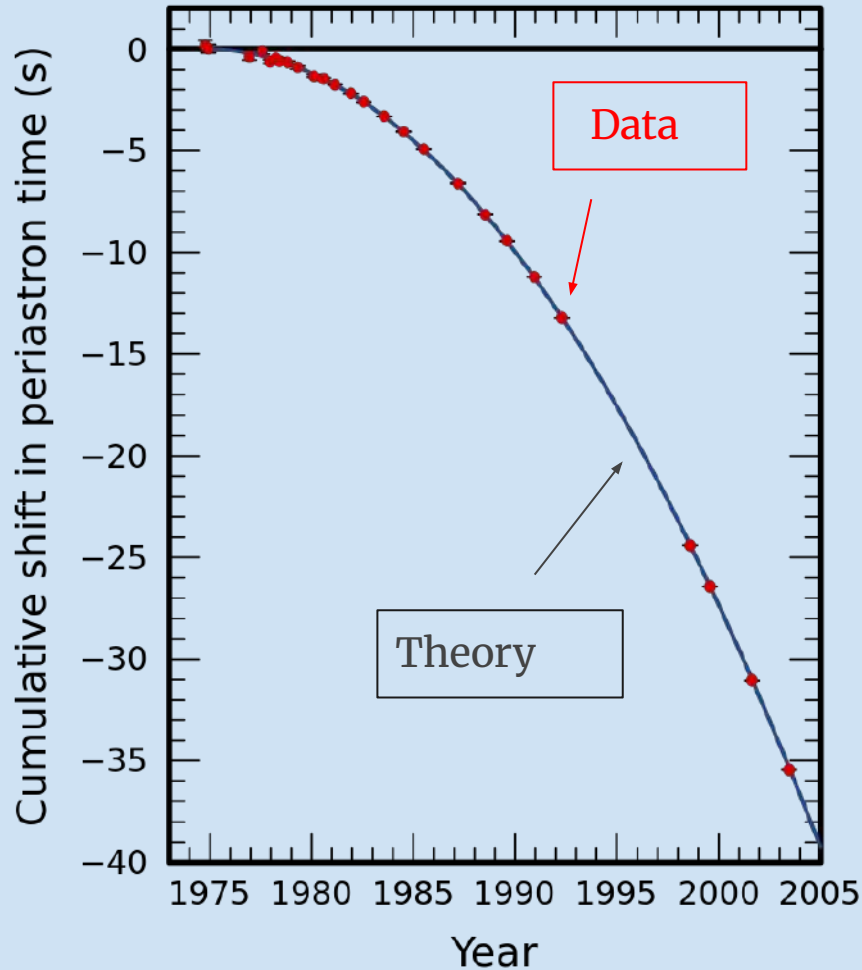
$$\nu_{GW} \sim 10^5 \text{ Hz}$$

$$\lambda_{GW} = \frac{c}{\nu_{GM}} \sim 10^{14} \text{ cm}$$

$$(\lambda_{GW} \gg l_0)$$

(slow-motion approximation,  
on which the quadrupole formalism is based)





Hulse-Taylor binary pulsar  
 $e_o \sim 0.617$  and  $f(e_o) \sim 0.184$ ,

Time to coalescence  
shorter by a factor  $1/f(e_o) \sim 5.4$ , compared  
to a binary on a circular  
orbit with the same period.

# How to detect them? Past

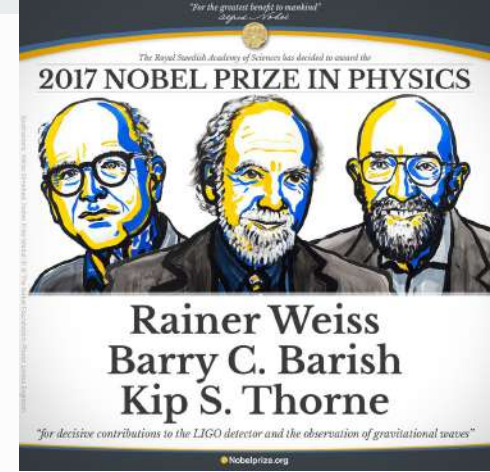
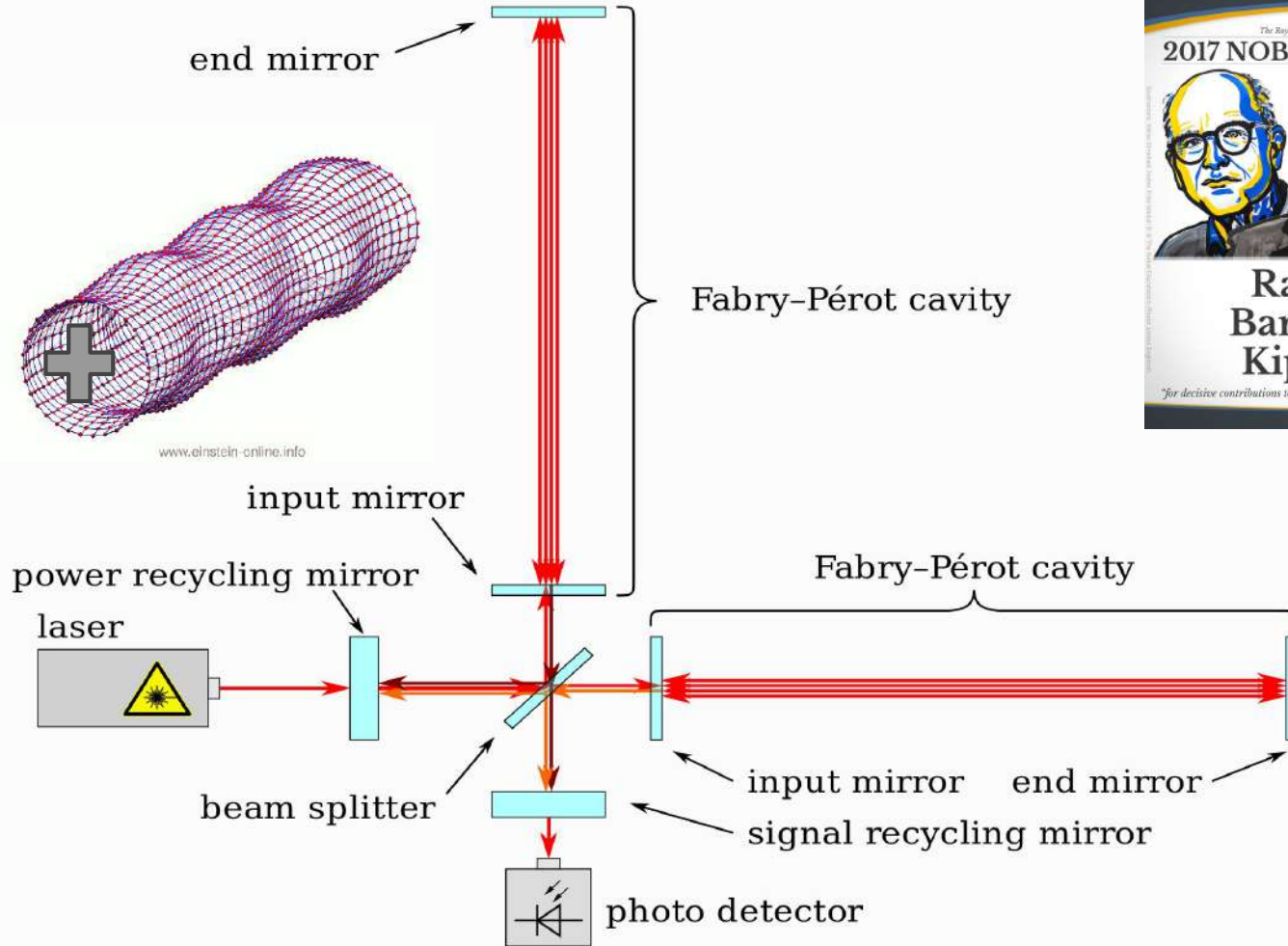
-Joseph Weber and his resonant bars in 1960s



-By the early 2000s, a set of initial detectors was completed, including TAMA 300 in Japan, GEO 600 in Germany, the Laser Interferometer Gravitational-Wave Observatory (LIGO) in the United States, and Virgo in Italy.

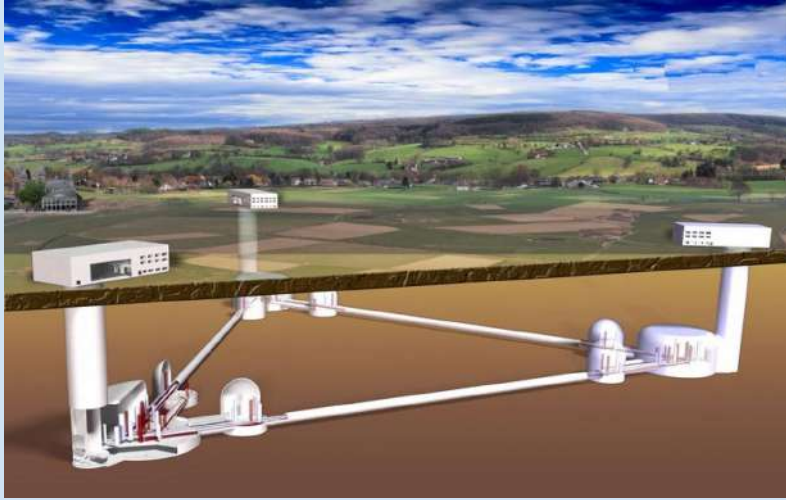


# How It works?



# How to detect them? Future Ground Based

## Einstein Telescope (~2035)

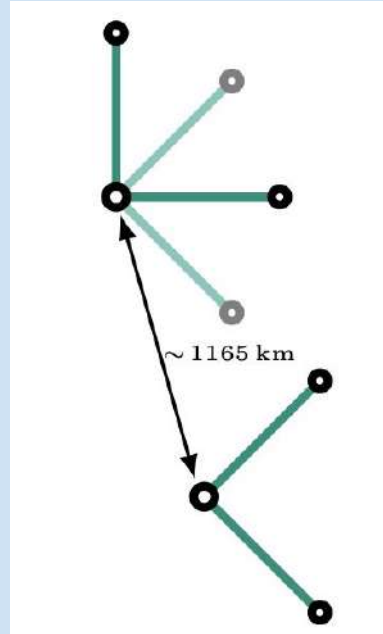


2 proposed config: L-shape(15km) vs Triangle(10km)

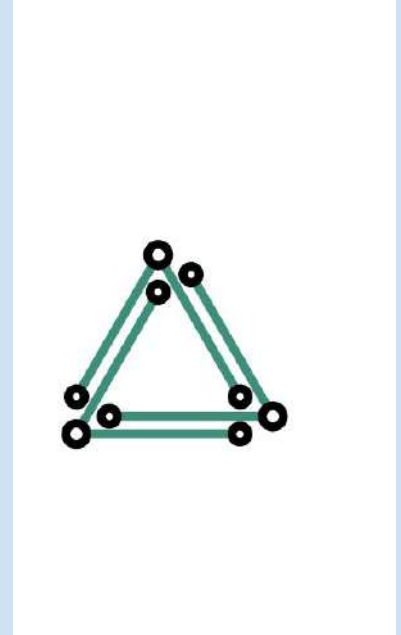
Lower frequencies ( $\sim 1$  Hz)  $\rightarrow$  much earlier inspiral information

Three candidate sites: i) Italy (Sardina) ii) Meuse-Rhine Euroregion(Maastricht, Liège and Aachen)  
iii) Germany (Saxony)

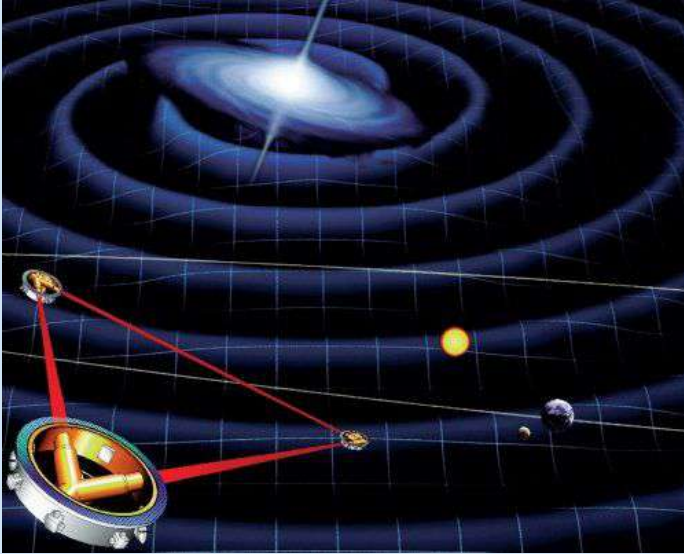
## 2L Config.



## $\Delta$ Config.



# How to detect them? Future Space based



**USA: LISA ;**

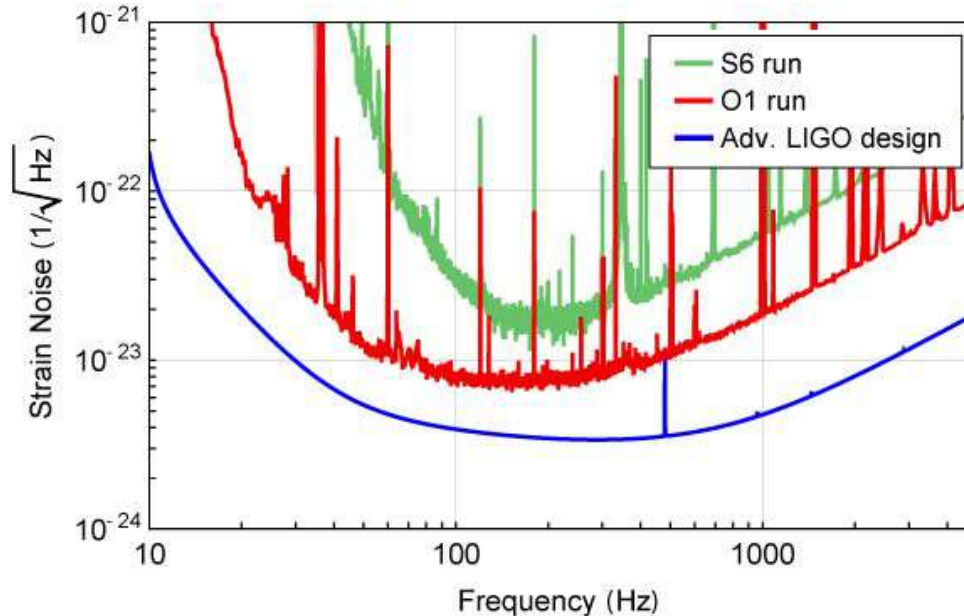
**China: Tian-Qin**

**Operate in the milli-Hz to 1  
Hz band**

# The Detector Noise : PSD $S_n(f)$

-2015 Advanced LIGO: the multi-stage optic suspensions and their attached internal seismic isolation platforms .

These subsystems produced a significant improvement on the low-frequency side of LIGO's noise spectrum

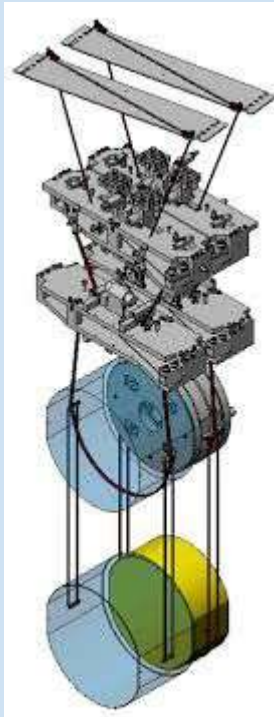


Noise Sources:

seismic & Newtonian noise at **low frequencies**,

suspension and thermal noise at **mid frequencies**,

Quantum (shot + radiation-pressure) noise at **high frequencies** — all combining to shape the detector's sensitivity curve.



# The Detector Noise : PSD $S_n(f)$

$$S_n(f) = \frac{1}{L^2} [S_{x,\text{seismic}}(f) + S_{x,\text{suspension}}(f) + S_{x,\text{thermal}}(f) + S_{x,\text{quantum}}(f) + S_{x,\text{coating}}(f) + \dots]$$

## Mirror Thermal Noise

PSD of displacement

$$S_x(f) = \frac{4k_B T}{\pi^2 f Y} \frac{d}{r_0^2} \left( \frac{Y'}{Y} \phi_{\parallel} + \frac{Y}{Y'} \phi_{\perp} \right)$$

Temperature

Boltzmann constant

Geometrical coating thickness

Loss angle of coating

Young's modulus of mirror substrate

laser beam radius

Young's modulus of coating

Harry et al, CQG 19, 897-917, 2002

## Observatories & experiments

Ground-based experiment



Space-based observatory



Pulsar timing array



Cosmic microwave background polarisation



## Cosmic sources



Supernova



Pulsar



Compact object falling onto a supermassive black hole



Merging supermassive black holes



Merging neutron stars in other galaxies



Merging stellar-mass black holes in other galaxies



Merging white dwarfs in our Galaxy

Cosmic fluctuations in the early Universe

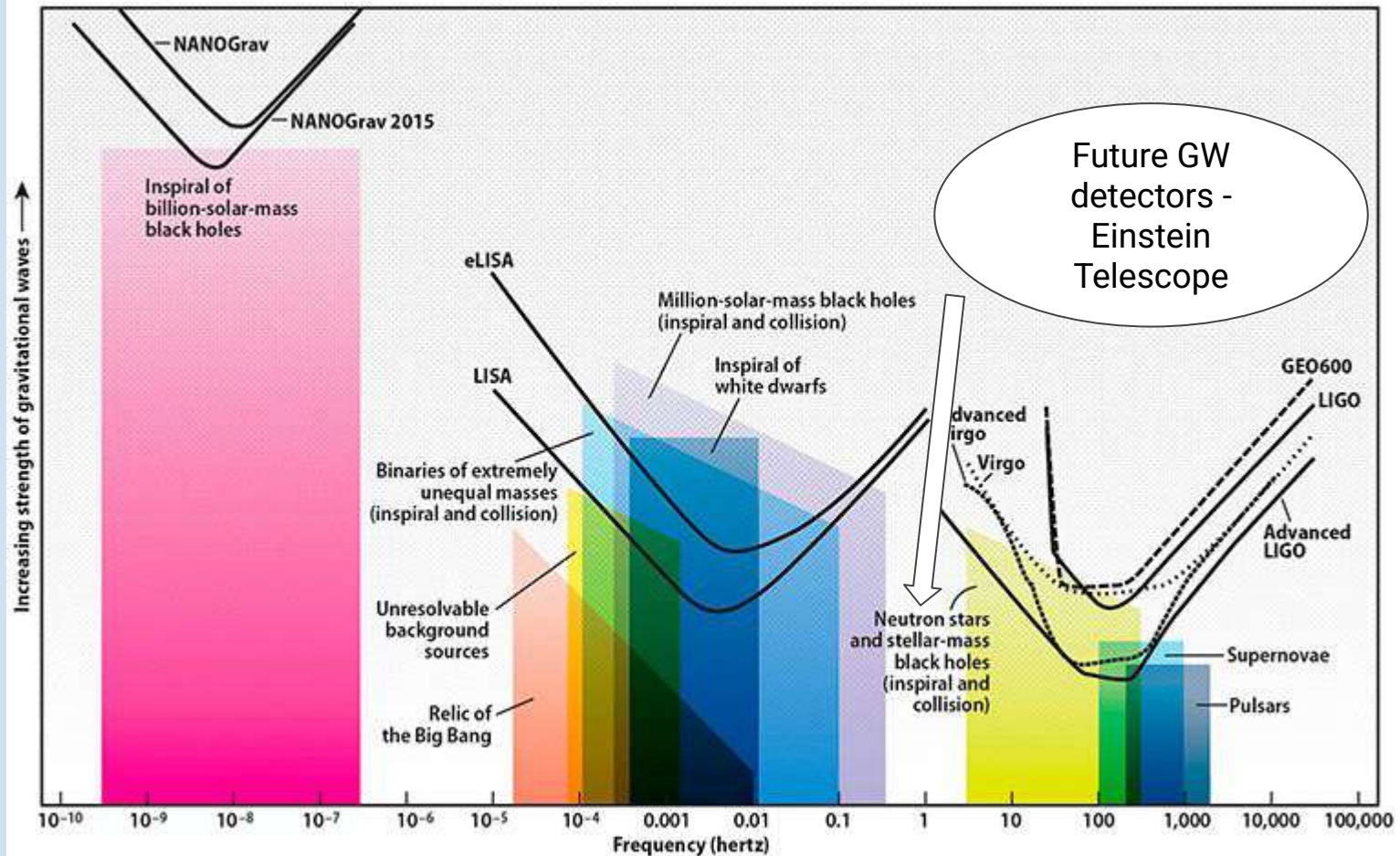
# GW Astronomy

Credit: [ESA](#)

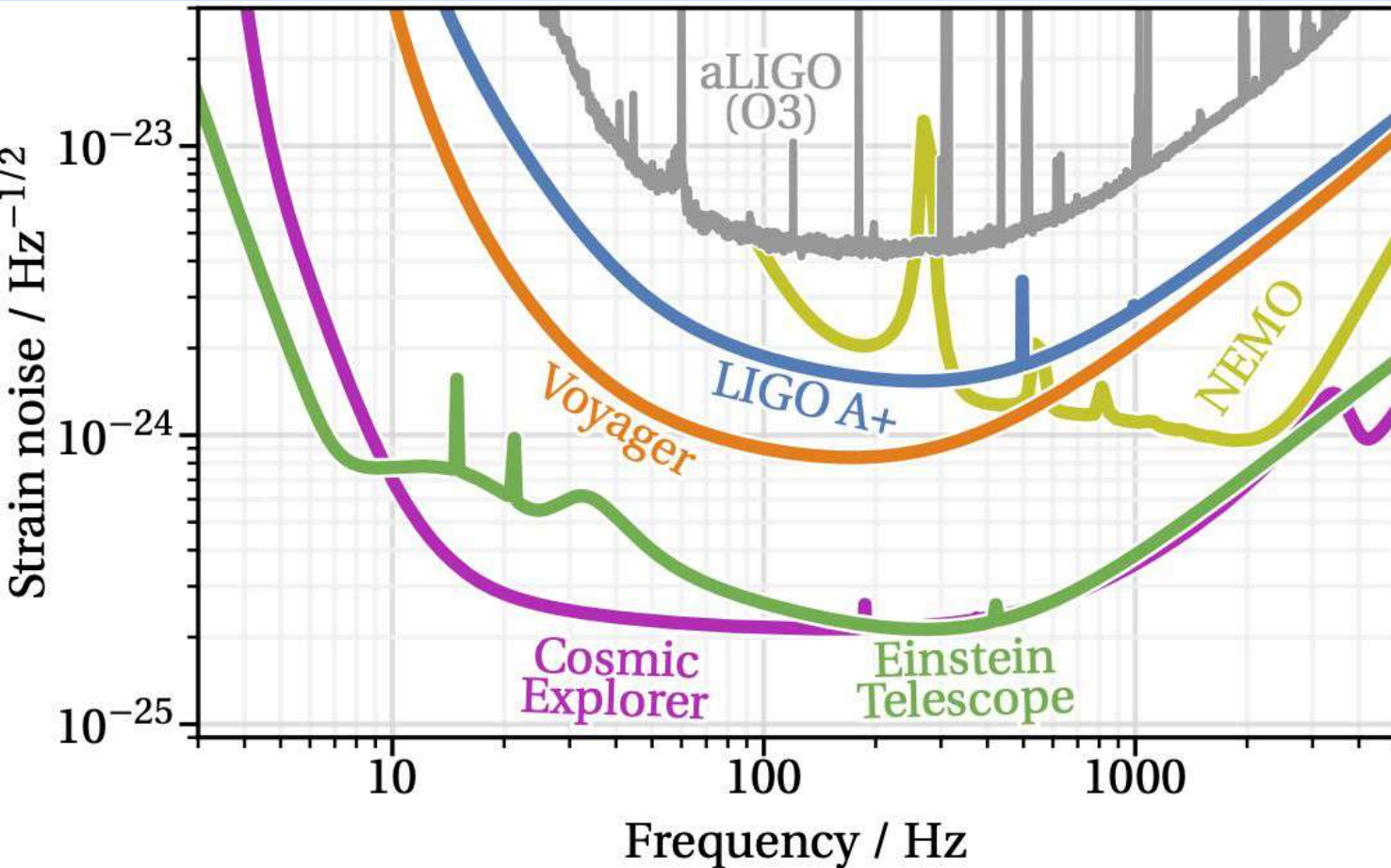


#LISA

# Future of GWs : A new window to the universe



# Future of GWs : A new window to the universe



Lower freq  
sensitivity (1Hz)=>  
 $O(10^5\text{-}10^6/\text{yr})$  +  
longer-duration  
signals

Improved reach vs. 2G:

- SNR  $\rightarrow 10\times$
- Redshift horizon  $z \rightarrow 20\times$
- Detection rate  $\rightarrow 100\times$

## Our Group Goal

- Neutron Stars
- oscillations of compact stars
- stochastic GWs

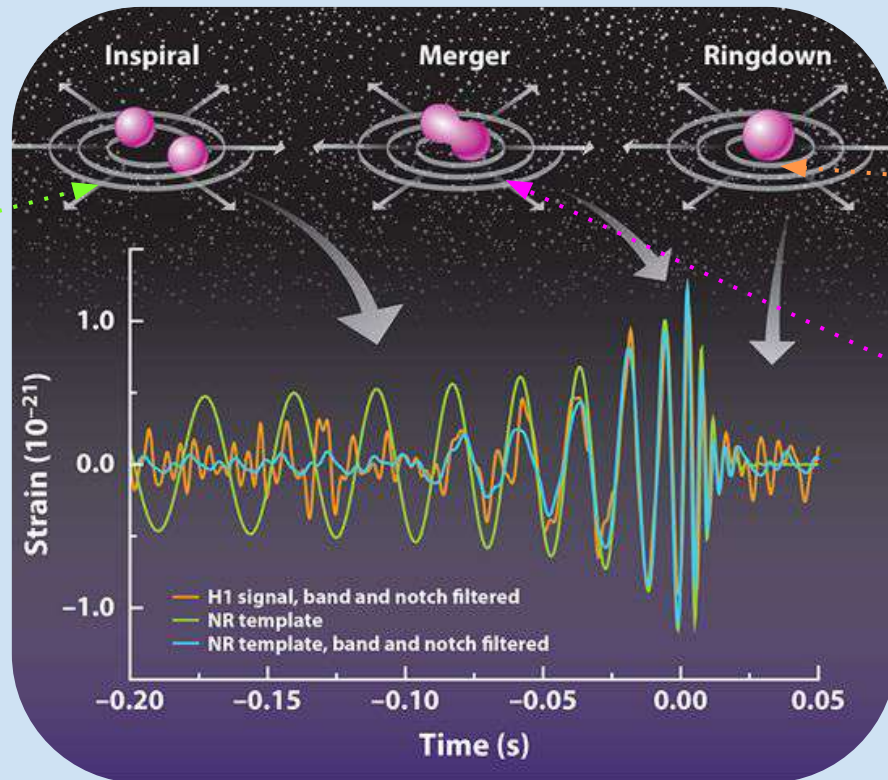
# Detecting Gravitational

$$\frac{dE_{orb}}{dt} + L_{GW} = 0$$

**Theory:**  $h = Ae^{i\Psi}$

**Data:**  $h = \frac{\Delta L}{L}$

# GWs : Methods to compute templates



Post-Newtonian  
Theory

BH Perturbation  
Theory

Numerical Rel.

$$h = Ae^{i\Psi}$$

the correct template/needle!

[Gravitational Waves and the Effort to Detect them By Peter Shawhan](#)

# GWs : Data-Analysis



**Parameter  
Estimation**



$$h = Ae^{i\Psi}$$

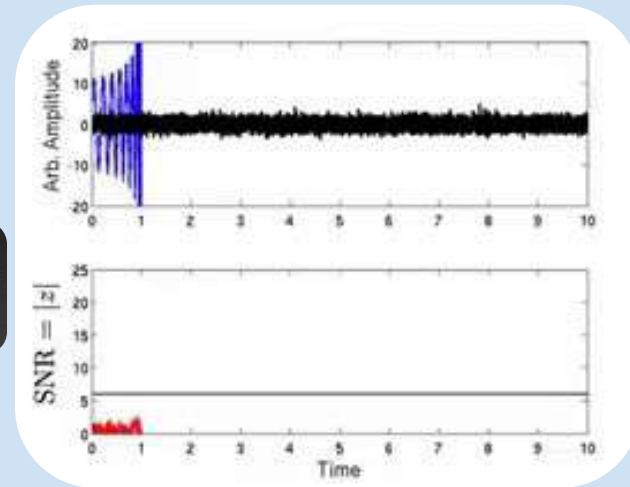
# GW Data-analysis => Matched Filtering



$$d(t) = h(t) + n(t)$$

$$z = \langle d | \hat{h} \rangle = 4 \int_{f_{\text{low}}}^{f_{\text{high}}} \frac{\tilde{d}(f) \hat{h}^*(f)}{S_n(f)} df$$

- Cross-correlate with detector data:  $d(t)$  using the noise-weighted inner product  $\rightarrow$  Signal-to-Noise Ratio (SNR)
- Template bank: Covers a grid of masses, spins, orientations  
Search = “scan over templates” to find the best match
- Noise is non-Gaussian and non-stationary



**Need ranking statistics  
& likelihood**

## GW Parameter Estimation



**Bayesian Analysis:**  
Comprehensive, reliable for realistic data,  
especially when the likelihood is non-Gaussian.

Given the detector data,:  
What is the probability that the  
observed strain was produced by an  
astrophysical signal rather than by  
noise?

**Fisher Analysis:**  
Quick, approximate, best for high SNR and gaussian  
noise

### \*References:

- *Post-Newtonian theory for gravitational waves*, *Living Reviews in Relativity*, Luc Blanchet 2024
- *Sources of Gravitational Waves: Theory and Observations*, Alessandra Buonanno and B.S. Sathyaprakash

# Bayesian Analysis: Given the data, inferring parameters

**Posterior:** prob. distn. of param ( $\theta$ )

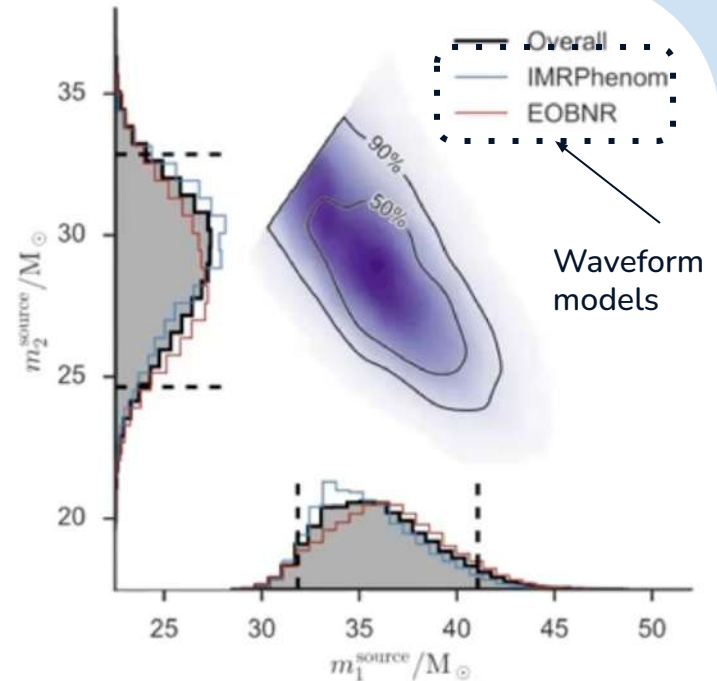
$$p(\theta|\vec{d}, M) = \frac{p(\vec{d}|\theta, M)p(\vec{\theta}, M)}{p(\vec{d}|M)}$$

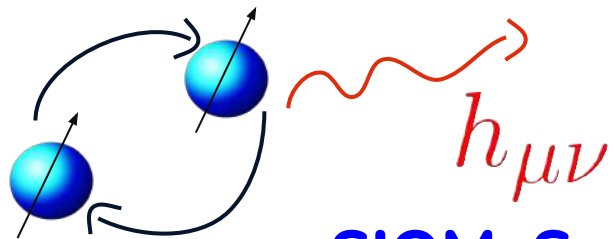
**Likelihood:** measures how well the data fits param

$$\mathcal{L}(\mathcal{H}_1|s) = \frac{p(s|\mathcal{H}_1)}{p(s|\mathcal{H}_0)} = \boxed{e^{(s,h)}} e^{-(h,h)/2}$$

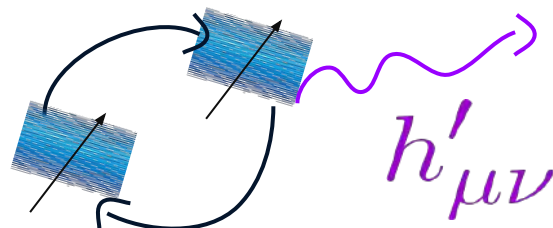
**Prior:** encodes our knowledge about param

**Evidence:** normalization



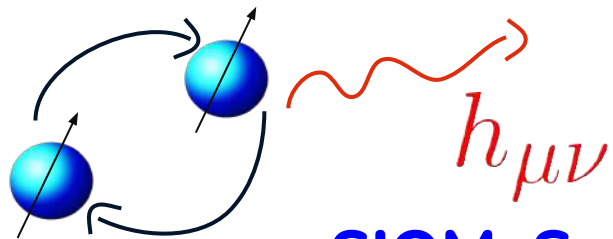


Testing the  
Compact binary  
nature

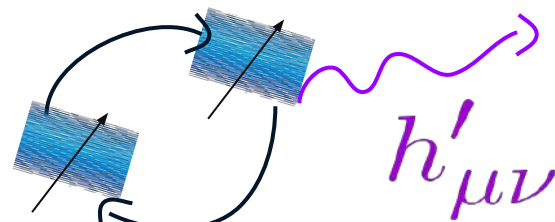


**SIQM: Spin-Induced Quadrupole Moment**

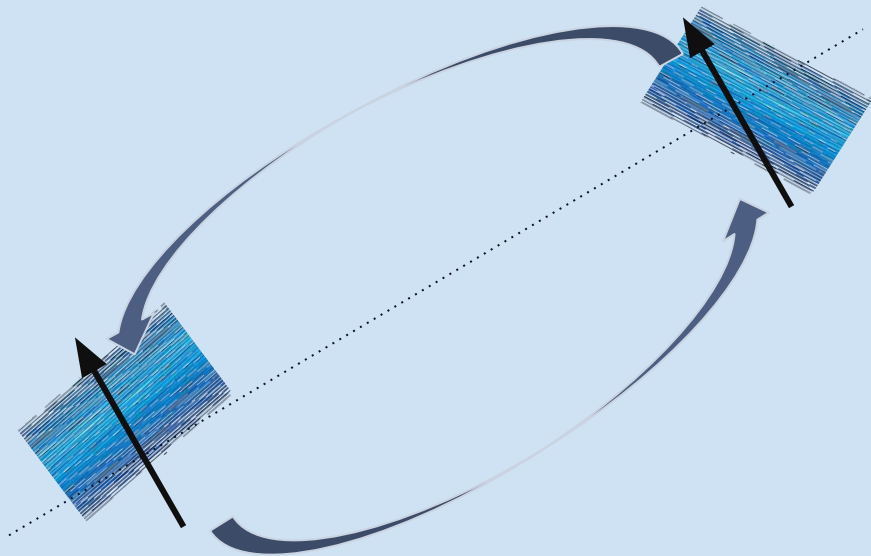
$$\left\{ \begin{array}{l} \kappa = 1, \text{ BH} \\ \kappa \neq 1, \text{ N.S, Exotic Compact Objects} \end{array} \right\}$$



## Testing the Compact binary nature



## SIQM: Spin-Induced Quadrupole Moment



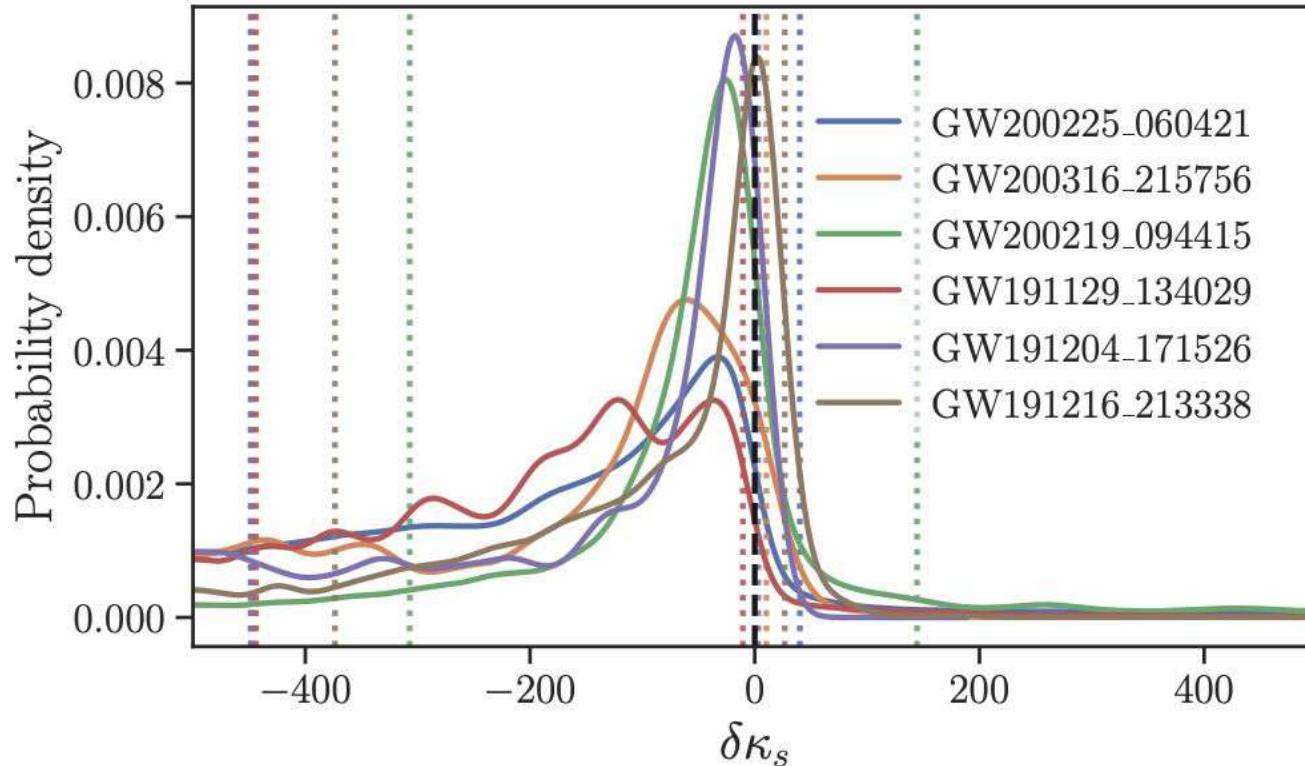
No-hair theorem test ➡ Is the object truly a Kerr BH?

Astrophysical identity test ➡ Neutron star vs boson star vs BH

Modified gravity/GR test ➡ Deviations from GR predictions

# Constrain on kappa GWTC 3 : Waveform IMRPhenomPv2

Parameterised deviations  $\kappa_s = 1 + \delta\kappa_s$  assuming  $\kappa_a = 0$ .



Get to know

# GW230529

Full name GW230529\_181500

~ 650 million light years away



## Detectors



- Offline OR not operational
- Online BUT not used for analysis\*
- Online AND used for analysis

Discovered on 29 May 2023 at 18h15 UTC

most likely a merger between a  
Neutron Star & Black Hole (NSBH)



~1.4  $M_{\odot}$



~3.6  $M_{\odot}$

**Most symmetric NSBH event so far**

more likely than prior GW NSBHs to have the neutron star  
ripped apart by the black hole

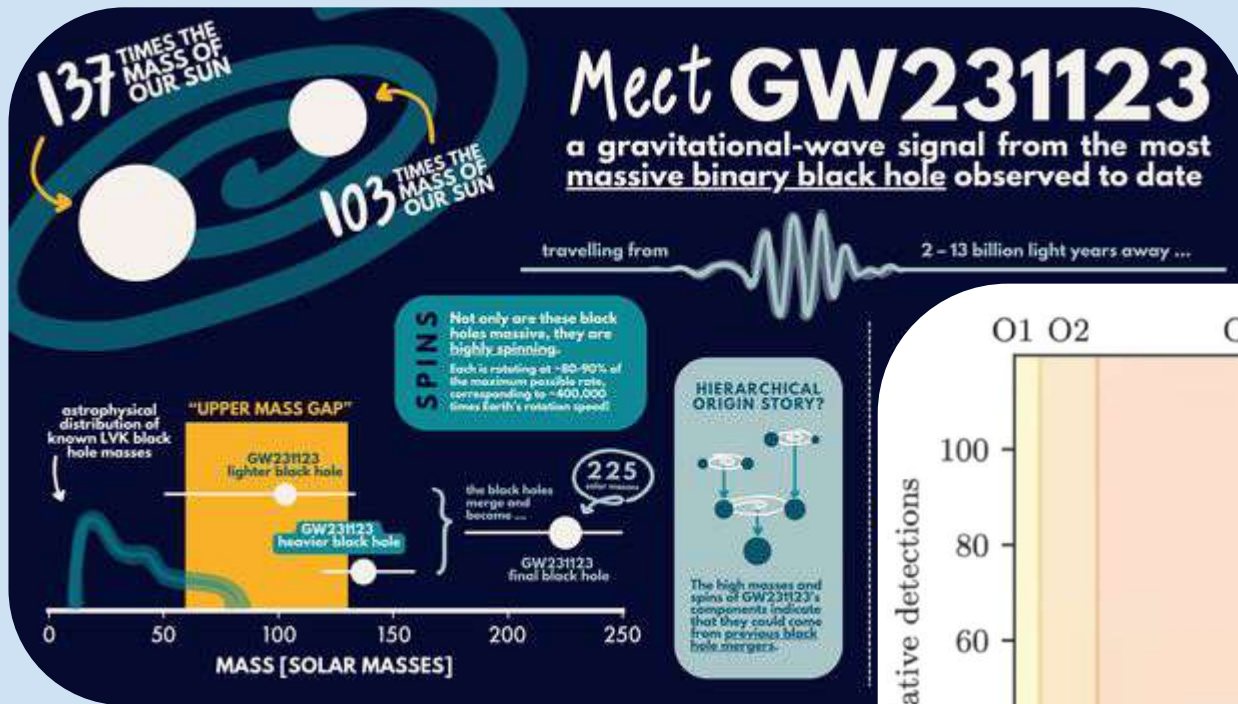
## Primary object in lower mass gap

further supports that this region is not empty

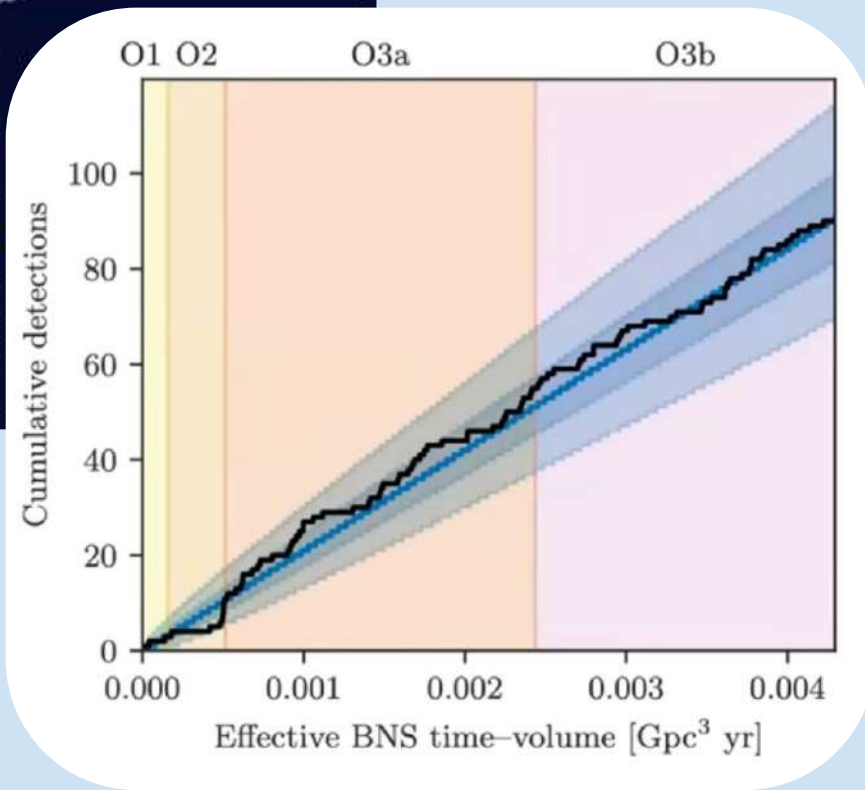


\* Although the KAGRA detector was in observing mode, its sensitivity was insufficient to impact the analysis of GW230529

@astronerdika



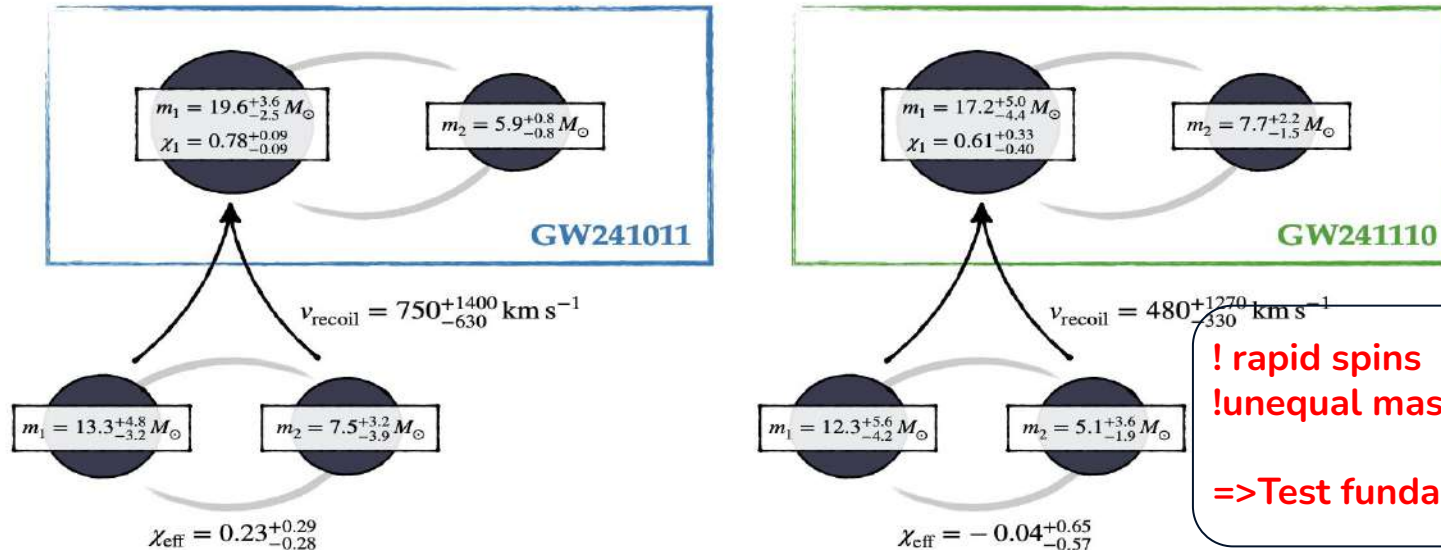
Now O4a: 128 new significant GW signal candidates!!!





# GW241011 and GW241110: Exploring Binary Formation and Fundamental Physics with Asymmetric, High-spin Black Hole Coalescences

Involve the mergers of second-generation black holes in dense stellar environments, such as globular, nuclear, and young massive star clusters



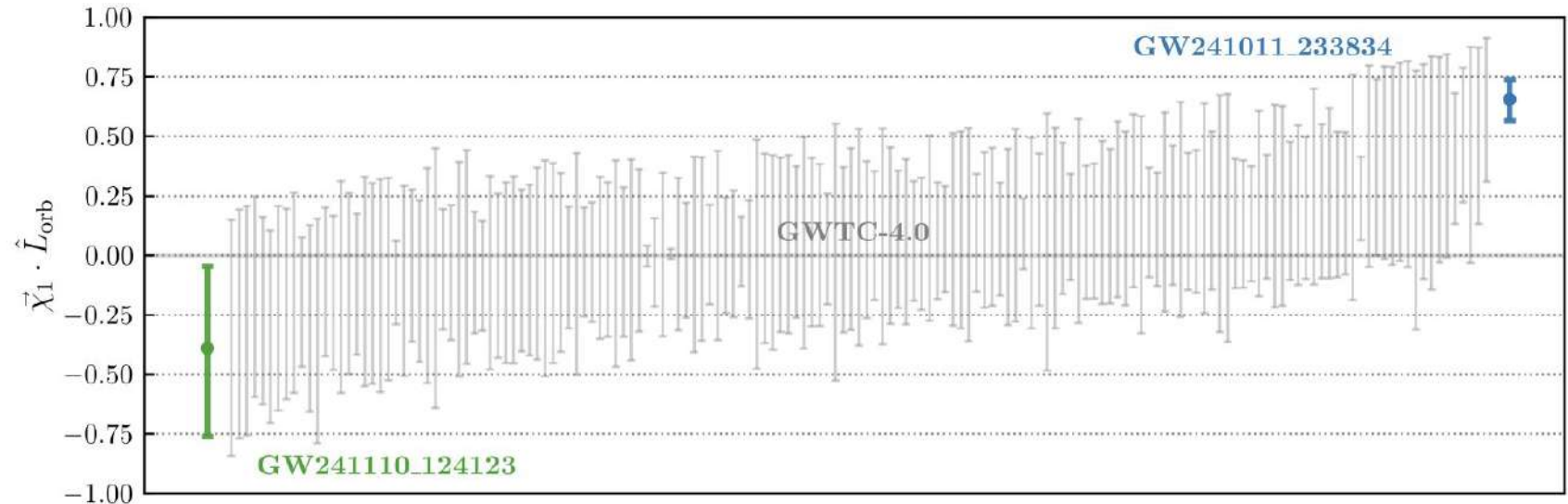
**! rapid spins**  
**! unequal mass ratio**

**=> Test fundamental physics**



CrossMark

# GW241011 and GW241110: Exploring Binary Formation and Fundamental Physics with Asymmetric, High-spin Black Hole Coalescences



Compact Binary Detections



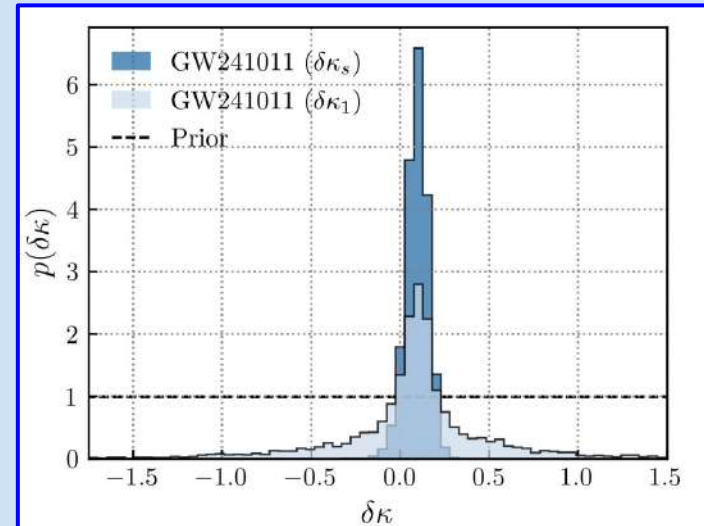
# GW241011 and GW241110: Exploring Binary Formation and Fundamental Physics with Asymmetric, High-spin Black Hole Coalescences

## Test fundamental physics

Define  $\kappa_1 = 1 + \delta\kappa_1$  and  $\kappa_2 = 1 + \delta\kappa_2$  as the spin-induced quadrupole coefficients of each compact object in GW241011's source binary.

Repeat parameter estimation with a modified IMRPHENOMXPHM waveform model allowing for nonzero  $\delta\kappa_1$  and  $\delta\kappa_2$

## Black Hole Spin-induced Quadrupole Moment





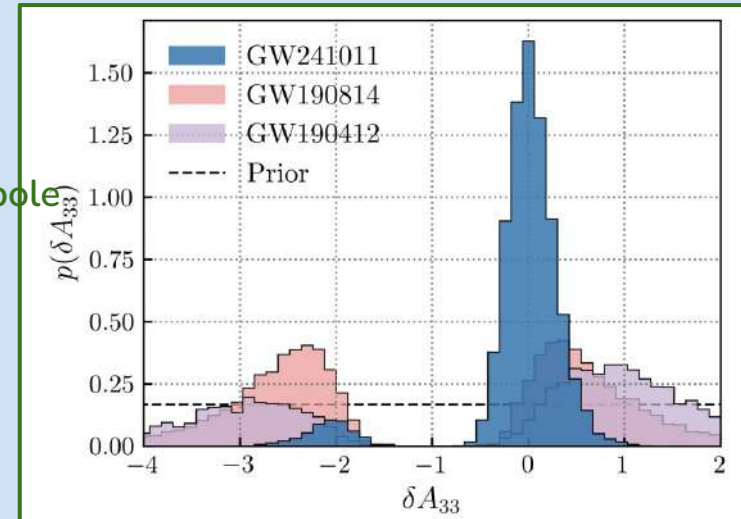
CrossMark

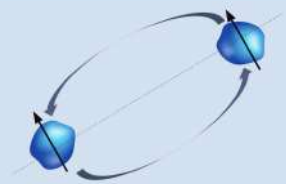
# GW241011 and GW241110: Exploring Binary Formation and Fundamental Physics with Asymmetric, High-spin Black Hole Coalescences

## Test fundamental physics

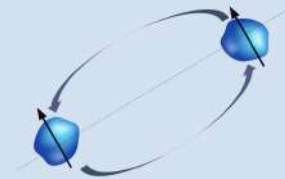
- Typically, GW signals are dominated by the quadrupole  $(\ell, m) = (2, \pm 2)$  multipole.
- Higher-order multipoles such as  $(2, \pm 1)$  and  $(3, \pm 3)$  become increasingly relevant for systems with unequal masses

## Radiation beyond the Quadrupole Approximation





# Conclusion



## Current Detections

### ➤ High-mass, asymmetric events enable higher-mode measurements

- GW241011 and GW241110 have mass ratios and orientations that make the (3,3) mode measurable, allowing precision tests of GR beyond the (2,2) mode.

### ➤ Isolated binary evolution struggles to explain the spins

- Standard stellar evolution predicts low and aligned spins → tension with observed primary spins.
- Requires exotic channels.

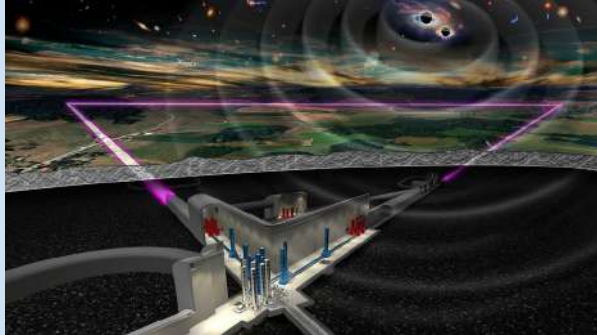
## Implications of GW241011 for rotating exotic compact objects

arXiv:2511.1734

1

N. V. Krishnendu, Tamara Evstafyeva, Aditya Vijaykumar, William E. East, Rimo Das, Sayantani Datta, Nils Siemonsen, Nami Uchikata, Poulami Dutta Roy, Anuradha Gupta, Ish Gupta, Syed U. Naqvi, Manuel Piarulli, Muhammed Saleem, Elise M. Sängner, Pratyusava Baral, Sajad A. Bhat, Thomas A. Callister, Mattia Emma, Carl-Johan Haster

# Conclusion



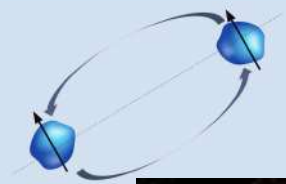
## ➤ A Leap Beyond Current Detectors

- ET's underground, cryogenic,  $10\times$  improved sensitivity will reveal signals far below the LIGO–Virgo noise floor.
- Access to 1–10 Hz band opens an entirely new discovery space (IMBH inspirals, early inspiral phases, massive binaries).

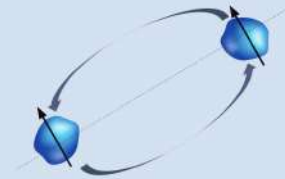
## ➤ A Complete Census of Compact Objects

- Detects stellar-mass merger in the universe up to redshift 20+.
- Allows population studies with orders-of-magnitude larger statistics, enabling:
  - BH formation channels
  - Cosmic evolution of spins, masses, eccentricities

**Future  
Detections  
:  
Einstein  
Telescope**



# Conclusion



# Thank You!

## ➤ Key Challenges

Ahead Data volumes  $\times 100$  larger  $\rightarrow$   
need new data analysis pipelines (fast inference, ML tools).

Theory and waveform modeling must improve to  
match ET's precision.

Environmental noise control  
(seismic, Newtonian noise) is critical.

