

ISOSPIN MIXING VIA GDR DECAY: AN OVERVIEW AND RECENT RESULTS

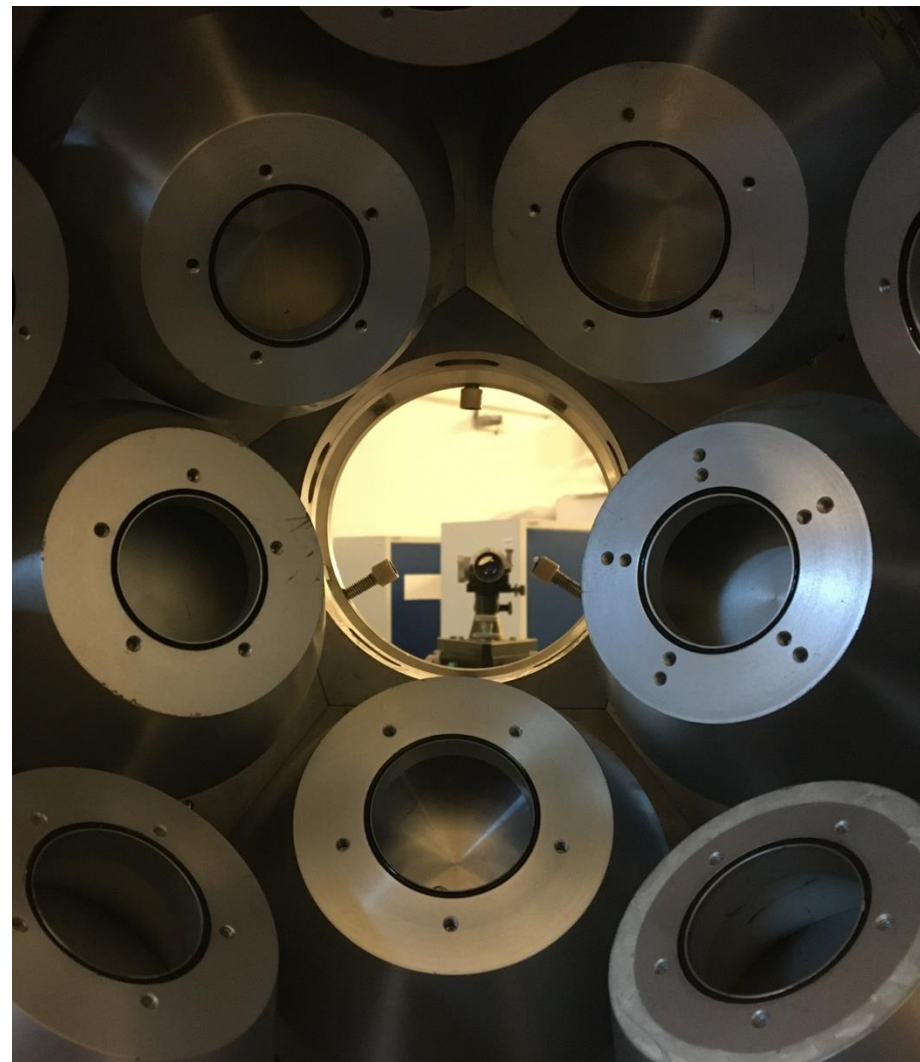
A. Giaz

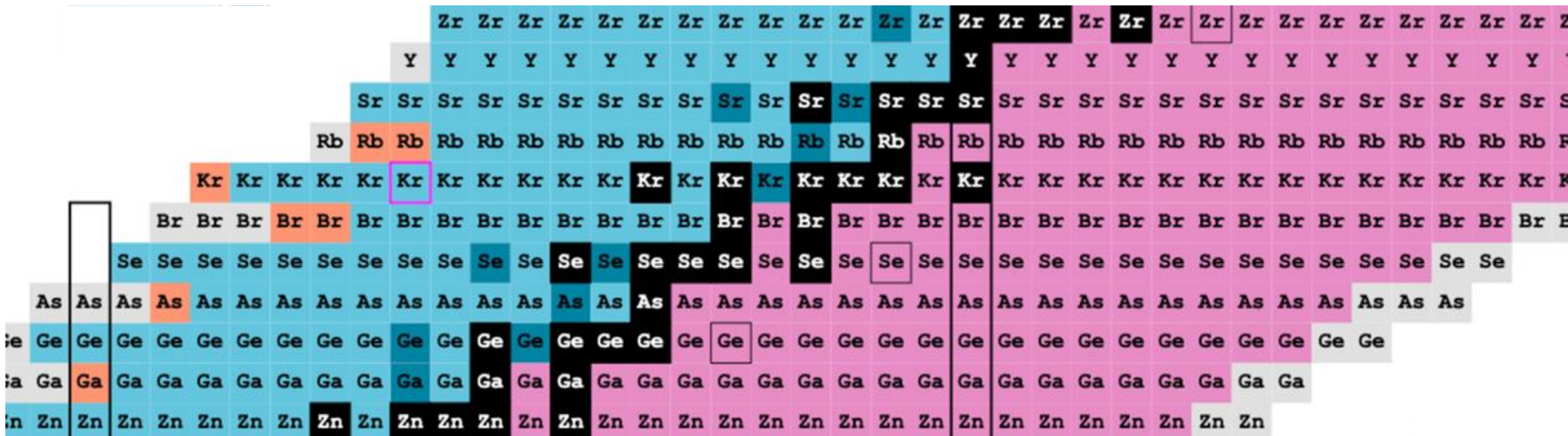
INFN Sezione di Milano



OUTLINE

- Motivations
 - Isospin mixing
 - Isospin mixing via GDR
 - Why ^{72}Kr ?
- Experimental Technique & Setup
- Results
 - The isospin mixing value α^2 and the correction term δ_c
- Summary & Perspectives





MOTIVATIONS

^{72}Kr AT LOWEST POSSIBLE TEMPERATURE

ISOSPIN SYMMETRY & ISOSPIN MIXING

Isospin symmetry plays a key role in nuclear structure and nuclear reaction.

$$\bullet |p\rangle \equiv |I = \frac{1}{2}, I_z = +\frac{1}{2}\rangle \quad \bullet |n\rangle \equiv |I = \frac{1}{2}, I_z = -\frac{1}{2}\rangle$$

ISOSPIN SYMMETRY & ISOSPIN MIXING

Isospin symmetry plays a key role in nuclear structure and nuclear nuclear reaction.

$$\bullet |p\rangle \equiv |I = \frac{1}{2}, I_z = +\frac{1}{2}\rangle \quad \bullet |n\rangle \equiv |I = \frac{1}{2}, I_z = -\frac{1}{2}\rangle$$

The Coulomb interaction (weaker than the nuclear interaction) breaks the Isospin symmetry inducing a mixing between state with different Isospin:

$$\alpha^2 = \frac{|\langle I = 1 | H_c | I = 0 \rangle|^2}{\Delta E^2}$$

Mixing Probability in the g.s.

ISOSPIN SYMMETRY & ISOSPIN MIXING

Isospin symmetry plays a key role in nuclear structure and nuclear reaction.

$$\bullet |p\rangle \equiv |I = \frac{1}{2}, I_z = +\frac{1}{2}\rangle \quad \bullet |n\rangle \equiv |I = \frac{1}{2}, I_z = -\frac{1}{2}\rangle$$

The Coulomb interaction (weaker than the nuclear interaction) breaks the Isospin symmetry inducing a mixing between state with different Isospin:

$$\alpha^2 = \frac{|\langle I = 1 | H_c | I = 0 \rangle|^2}{\Delta E^2}$$

Mixing Probability in the g.s.

Isospin mixing beyond nuclear structure: The CKM matrix


$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

ISOSPIN SYMMETRY & ISOSPIN MIXING

Isospin symmetry plays a key role in nuclear structure and nuclear reaction.

$$\bullet |p\rangle \equiv |I = \frac{1}{2}, I_z = +\frac{1}{2}\rangle \quad \bullet |n\rangle \equiv |I = \frac{1}{2}, I_z = -\frac{1}{2}\rangle$$

The Coulomb interaction (weaker than the nuclear interaction) breaks the Isospin symmetry inducing a mixing between state with different Isospin:

$$\alpha^2 = \frac{|\langle I = 1 | H_c | I = 0 \rangle|^2}{\Delta E^2}$$

Mixing Probability in the g.s.

Isospin mixing beyond nuclear structure: The CKM matrix


$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

the V_{ud} term is obtained from the ft values of $0^+ \rightarrow 0^+$ superallowed Fermi β transition.

ISOSPIN SYMMETRY & ISOSPIN MIXING

Isospin symmetry plays a key role in nuclear structure and nuclear reaction.

$$\bullet |p\rangle \equiv |I = \frac{1}{2}, I_z = +\frac{1}{2}\rangle \quad \bullet |n\rangle \equiv |I = \frac{1}{2}, I_z = -\frac{1}{2}\rangle$$

The Coulomb interaction (weaker than the nuclear interaction) breaks the Isospin symmetry inducing a mixing between state with different Isospin:

$$\alpha^2 = \frac{|\langle I = 1 | H_c | I = 0 \rangle|^2}{\Delta E^2}$$

Mixing Probability in the g.s.

Isospin mixing beyond nuclear structure: The CKM matrix

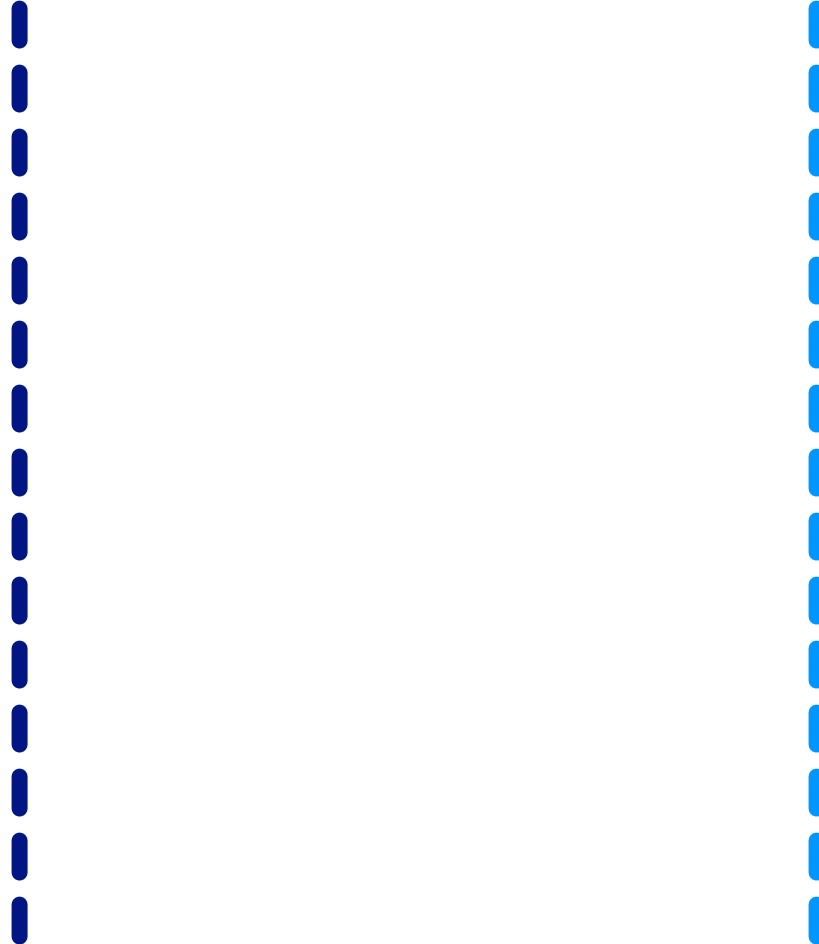
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

the V_{ud} term is obtained from the ft values of $0^+ \rightarrow 0^+$ superallowed Fermi β transition.

The isospin-symmetry-breaking corrected $\mathcal{F}t = (1 + \delta_r)(1 + \delta_c)$

HOW WE CAN MEASURE ISOSPIN MIXING?

Measuring transition strictly forbidden by isospin conservation.



HOW WE CAN MEASURE ISOSPIN MIXING?

Measuring transition strictly forbidden by isospin conservation.

Study of Fermi β transitions

$(\Delta J = 0, \Delta I = 0, \pi_i, \pi_f = +)$

- Measurement of the β -particles anisotropy emission
- For β^+ transition $\rightarrow$ initial state, for $\beta^- \rightarrow$ final state.
- Fermi operator is not nuclear structure dependent
- **Possible for all nuclei**
- **Difficult for unstable nuclei**

HOW WE CAN MEASURE ISOSPIN MIXING?

Measuring transition strictly forbidden by isospin conservation.

Study of Fermi β transitions

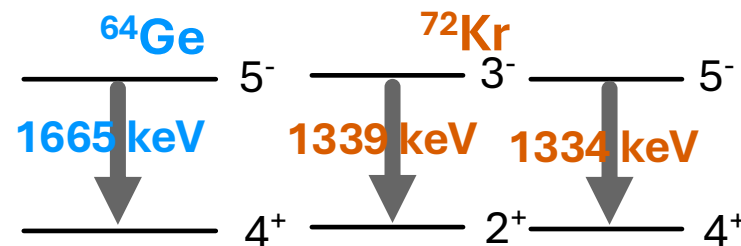
$(\Delta J = 0, \Delta I = 0, \pi_i, \pi_f = +)$

- Measurement of the β -particles anisotropy emission
- For β^+ transition $\rightarrow$ initial state, for $\beta^- \rightarrow$ final state.
- Fermi operator is not nuclear structure dependent
- **Possible for all nuclei**
- **Difficult for unstable nuclei**

Study of a forbidden E1 decay

(polarization & angular distribution)

- **Possible only in N=Z nuclei**
- **Unstable nuclei**
- It cannot discern if the mixing occur in the initial or final state



E. Farnea et al.,
PLB, **551**, (2003),
56-62

A. Ertoprak, et al.,
proposal for AGATA @ LNL

HOW WE CAN MEASURE ISOSPIN MIXING?

Measuring transition strictly forbidden by isospin conservation.

Study of Fermi β transitions

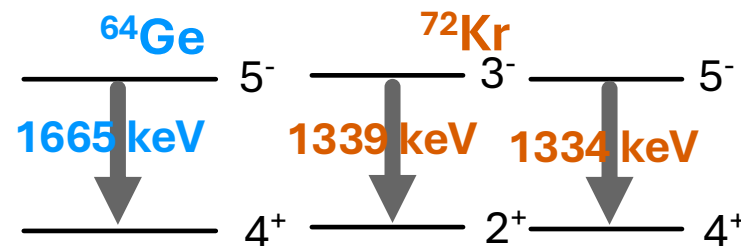
$(\Delta J = 0, \Delta I = 0, \pi_i, \pi_f = +)$

- Measurement of the β -particles anisotropy emission
- For β^+ transition $\rightarrow$ initial state, for $\beta^- \rightarrow$ final state.
- Fermi operator is not nuclear structure dependent
- **Possible for all nuclei**
- **Difficult for unstable nuclei**

Study of a forbidden E1 decay

(polarization & angular distribution)

- **Possible only in N=Z nuclei**
- **Unstable nuclei**
- It cannot discern if the mixing occur in the initial or final state

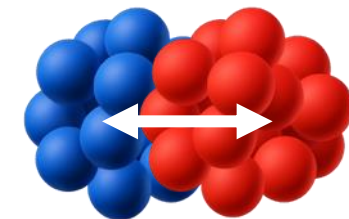


E. Farnea et al.,
PLB, **551**, (2003),
56-62

A. Ertoprak, et al.,
proposal for AGATA @ LNL

Study of E1 Giant Dipole Resonance Strength

- **Possible only in N=Z nuclei**
- **Unstable nuclei**
- It cannot discern if the mixing occur in the initial or final state

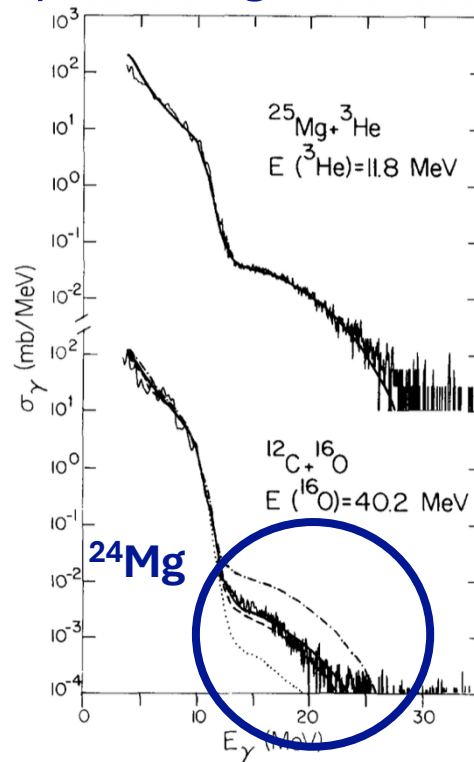


^{72}Kr this work

PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (I)

Previous measurements in light nuclei

First investigation of the role of isospin mixing in GDR decay

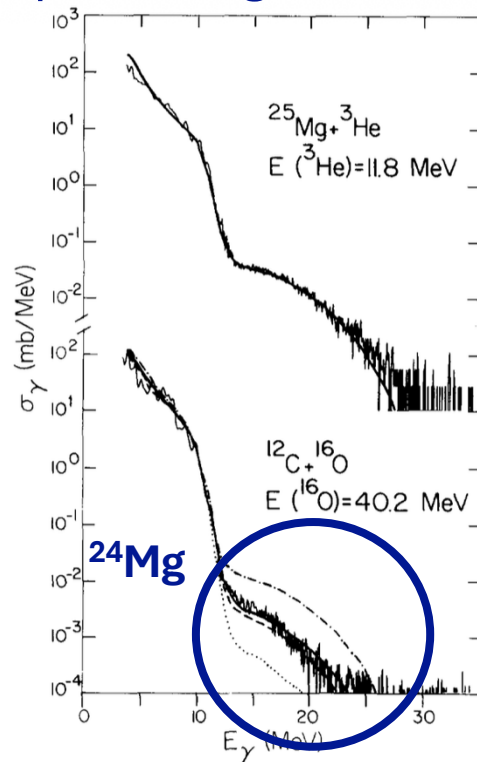


M. Harakeh, et al., Phys. Lett. B 176 (1986) 297

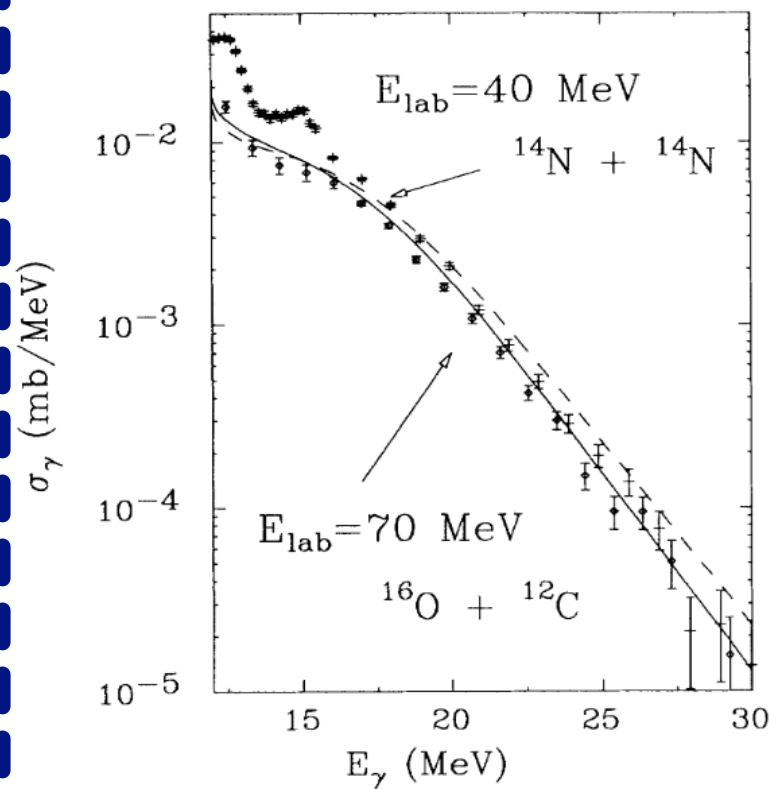
PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (I)

Previous measurements in light nuclei

First investigation of the role of isospin mixing in GDR decay



Isospin mixing in ^{28}Si and ^{30}Si



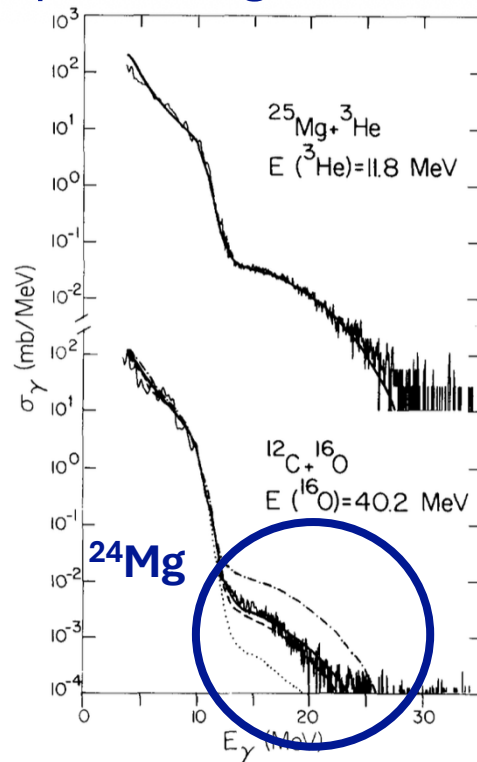
M. Harakeh, et al., Phys. Lett. B 176 (1986) 297

J. Behr, et al., Phys. Rev. Lett. 70 (1993) 3201.

PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (I)

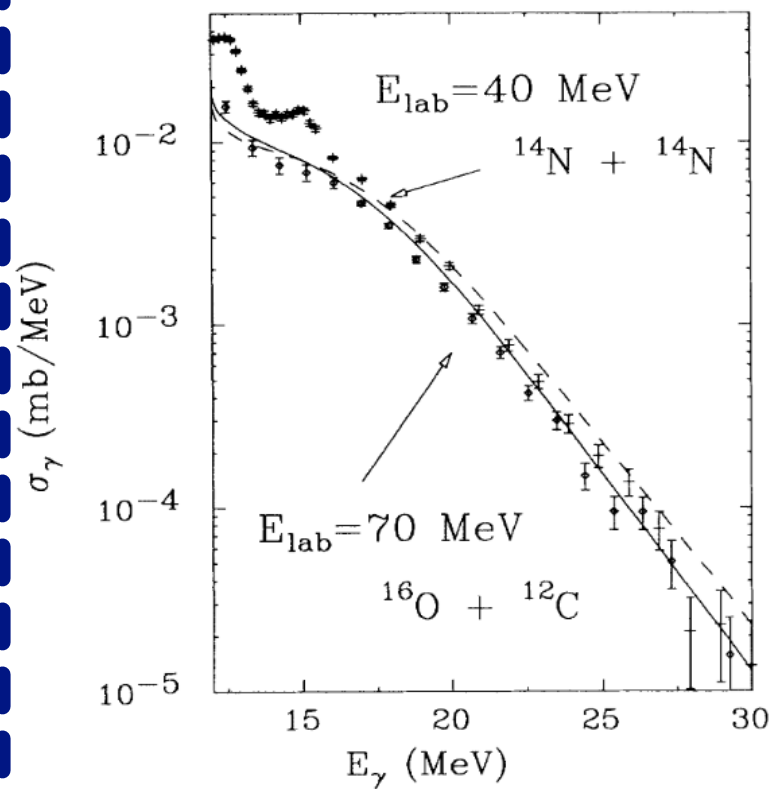
Previous measurements in light nuclei

First investigation of the role of isospin mixing in GDR decay



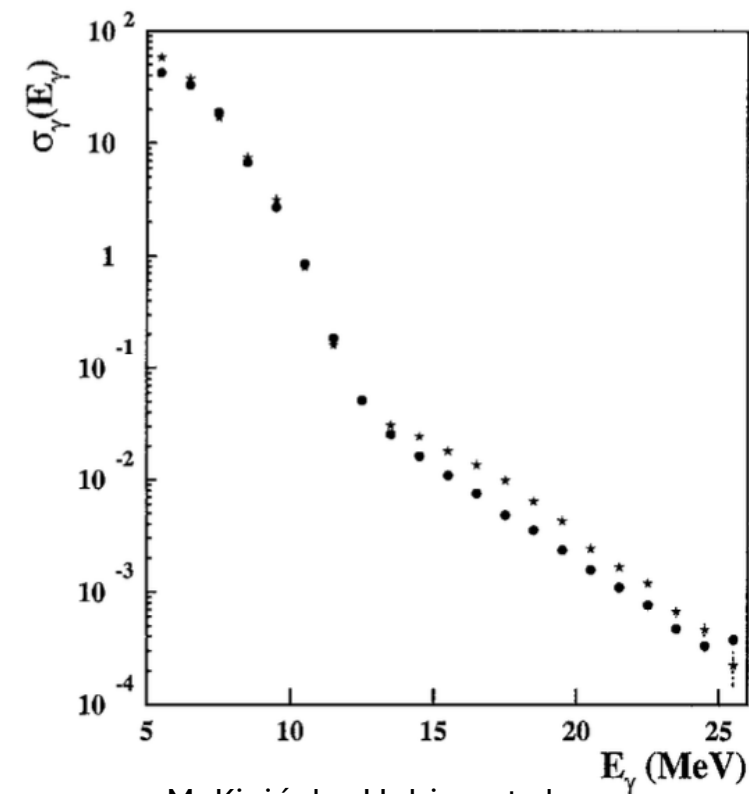
M. Harakeh, et al., Phys. Lett. B 176 (1986) 297

Isospin mixing in ^{28}Si and ^{30}Si



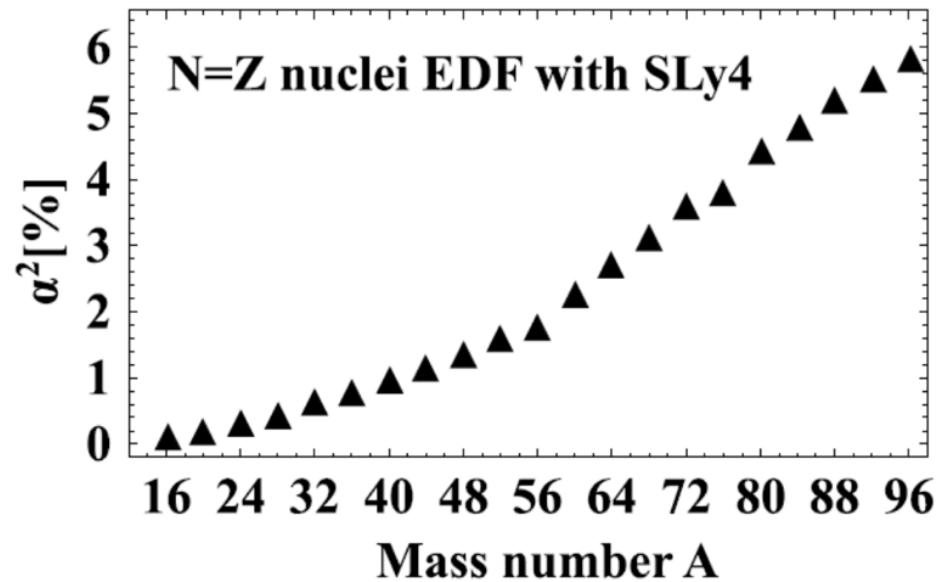
J. Behr, et al., Phys. Rev. Lett. 70 (1993) 3201.

Isospin mixing in ^{32}S



M. Kicińska-Habior, et al.,
Nucl. Phys. A 731 (2004) 138.

PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (II)



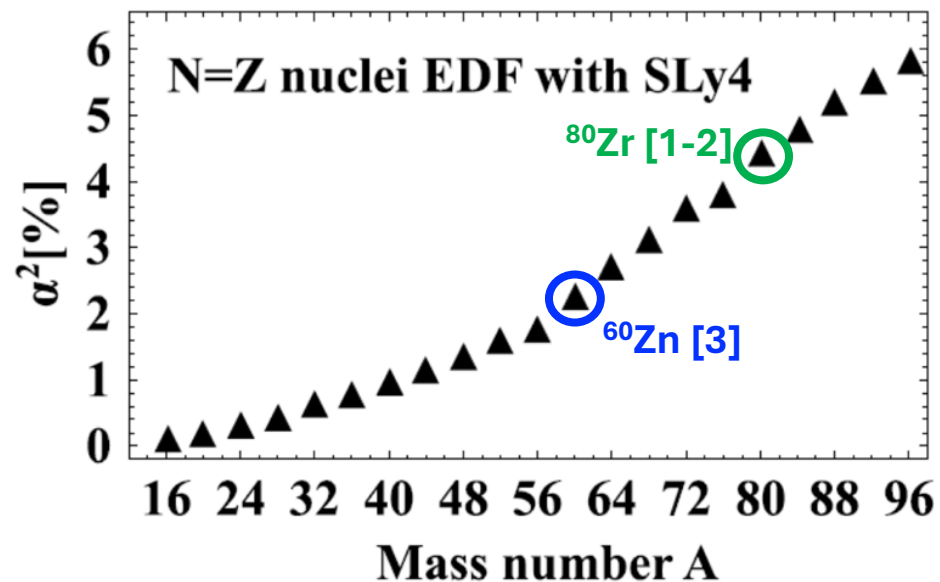
W. Satula, et al., Phys. Rev. Lett., **103**,(2009) 012502

[1] A. Corsi et al., PRC 84, 041304(R) (2011)

[2] S. Ceruti et al., PRL 115, 222502 (2015)

[3] G. Gosta et al., PRC 103, L041302 (2021)

PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (II)



W. Satula, et al., Phys. Rev. Lett., **103**,(2009) 012502

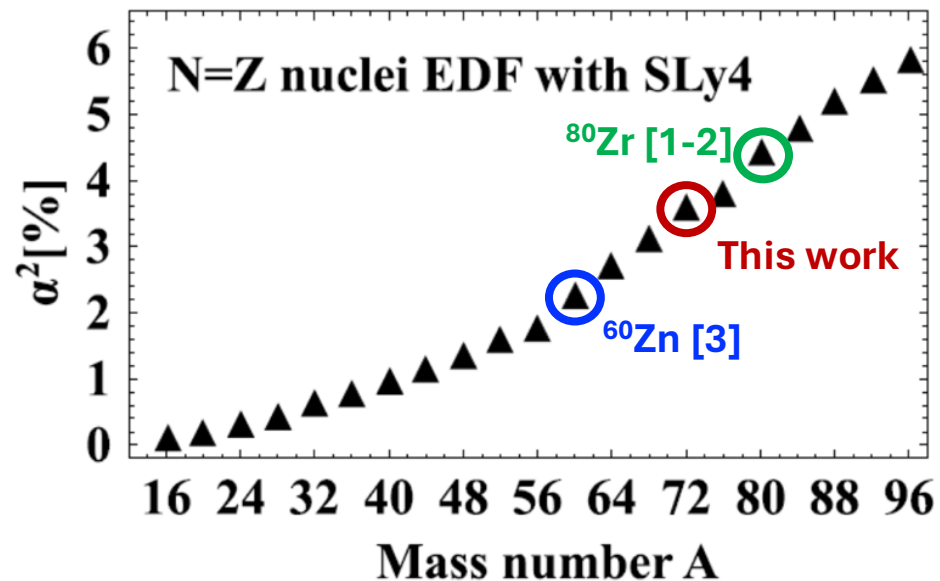
Nucleus	Temperature	Setup
^{80}Zr	2.8 MeV	HECTOR – GARFIELD [1]
^{80}Zr	1.8 MeV	AGATA – HECTOR+ [2]
^{60}Zn	2.0 MeV	GALILEO – 3''x3'' LaBr ₃ :Ce [3]
^{60}Zn	2.4 MeV	GALILEO – 3''x3'' LaBr ₃ :Ce [3]

[1] A. Corsi et al., PRC 84, 041304(R) (2011)

[2] S. Ceruti et al., PRL 115, 222502 (2015)

[3] G. Gosta et al., PRC 103, L041302 (2021)

PREVIOUS EXPERIMENTS BASED ON GDR TECHNIQUE (II)



W. Satula, et al., Phys. Rev. Lett., **103**, (2009) 012502

Nucleus	Temperature	Setup
^{80}Zr	2.8 MeV	HECTOR – GARFIELD [1]
^{80}Zr	1.8 MeV	AGATA – HECTOR+ [2]
^{60}Zn	2.0 MeV	GALILEO – 3''x3'' LaBr ₃ :Ce [3]
^{60}Zn	2.4 MeV	GALILEO – 3''x3'' LaBr ₃ :Ce [3]

The new experiment:

Isospin Mixing in ^{72}Kr at the lowest possible temperature at IFIN-HH tandem

[1] A. Corsi et al., PRC 84, 041304(R) (2011)

[2] S. Ceruti et al., PRL 115, 222502 (2015)

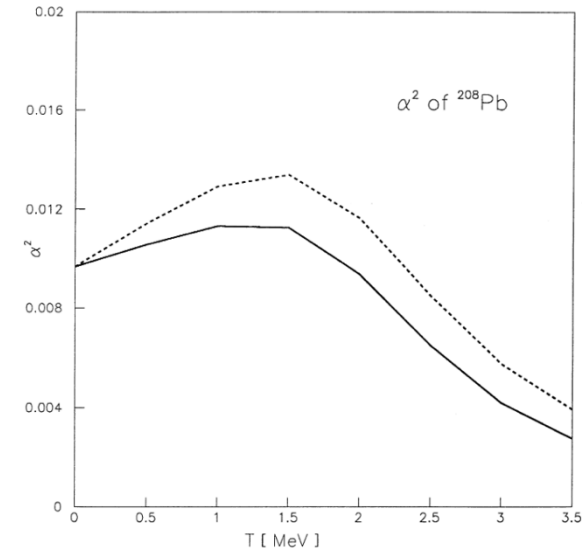
[3] G. Gosta et al., PRC 103, L041302 (2021)

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measure with two techniques.

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with two techniques.

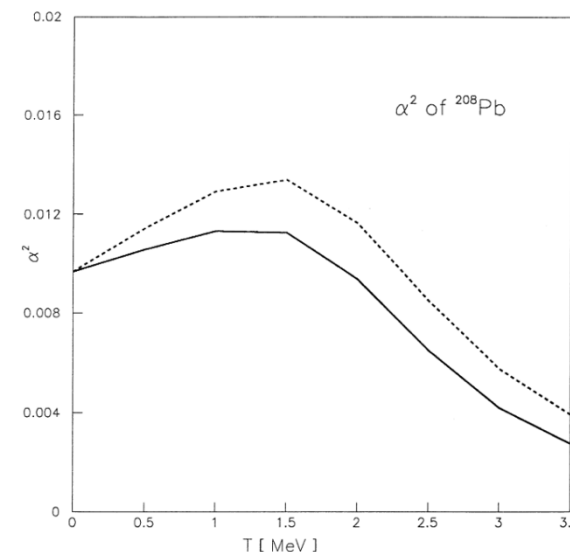
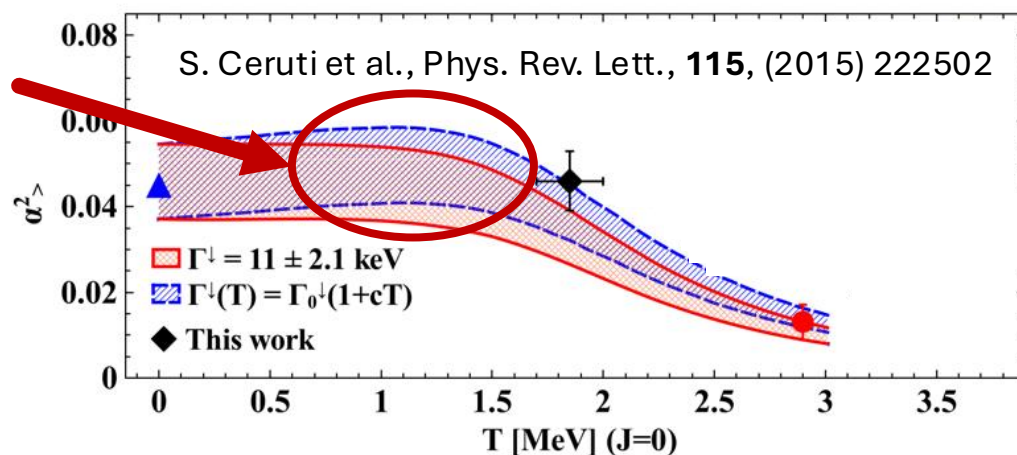


H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with the two techniques.

Add a point
in the
plateau
region

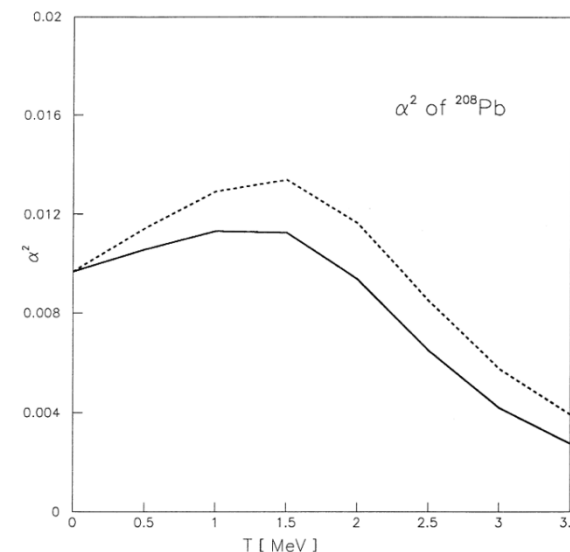
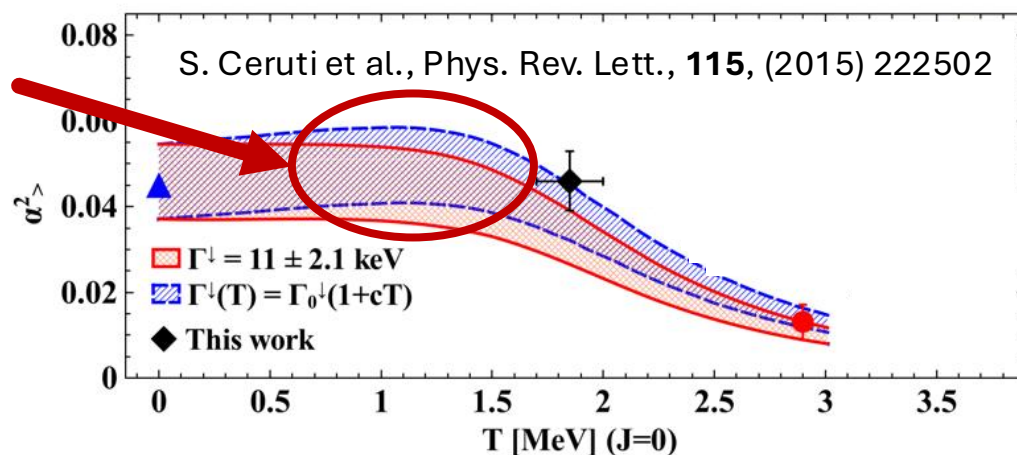


H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with two techniques.

Add a point
in the
plateau
region



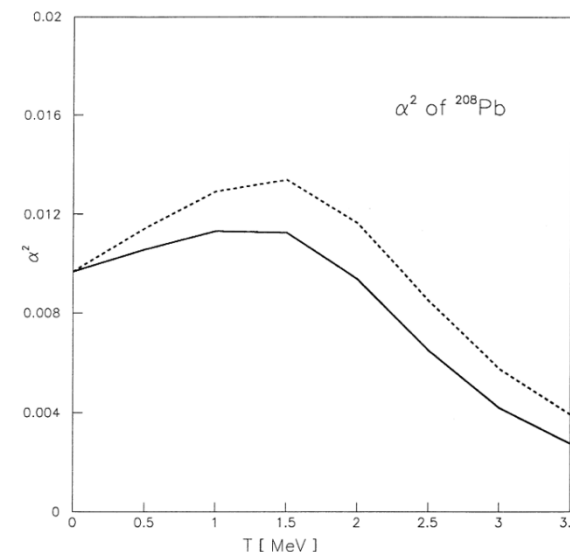
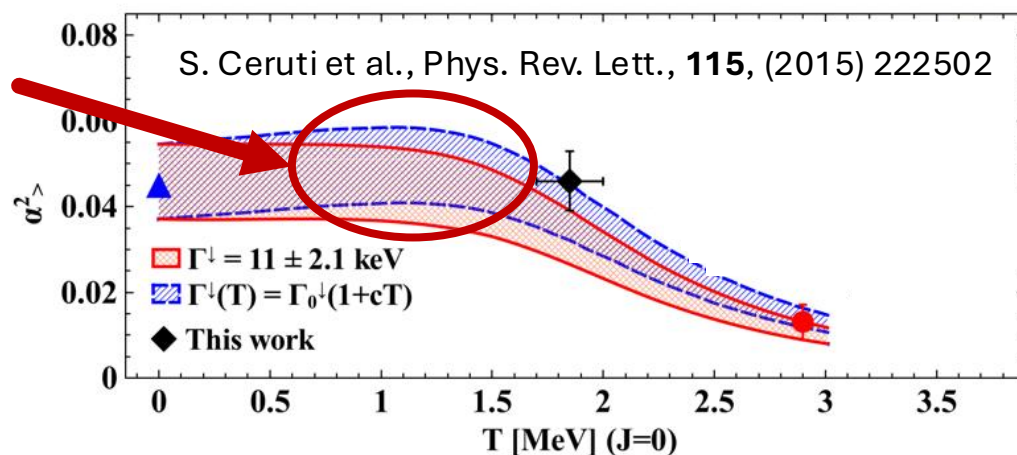
H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

ADVATAGE: In the plateau region, only one data point is needed to extract the isospin-mixing value at $T=0$.

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with two techniques.

Add a point
in the
plateau
region



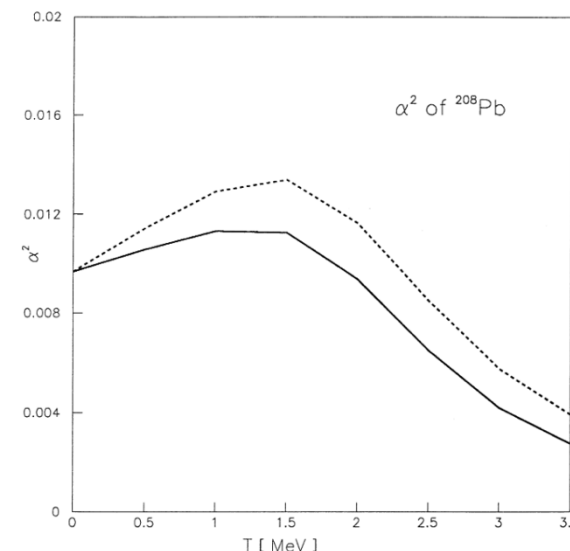
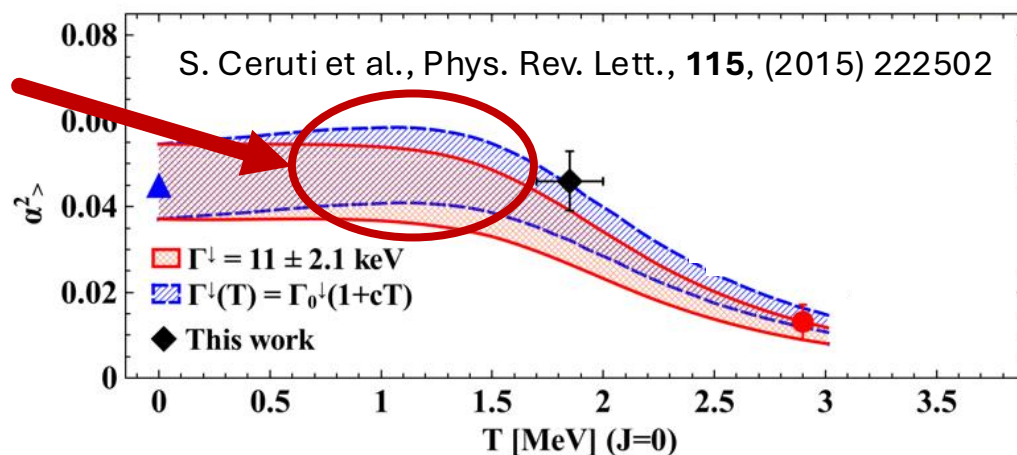
H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

ADVANTAGE: In the plateau region, only one data point is needed to extract the isospin-mixing value at $T=0$. ✓

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with two techniques.

Add a point
in the
plateau
region



H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

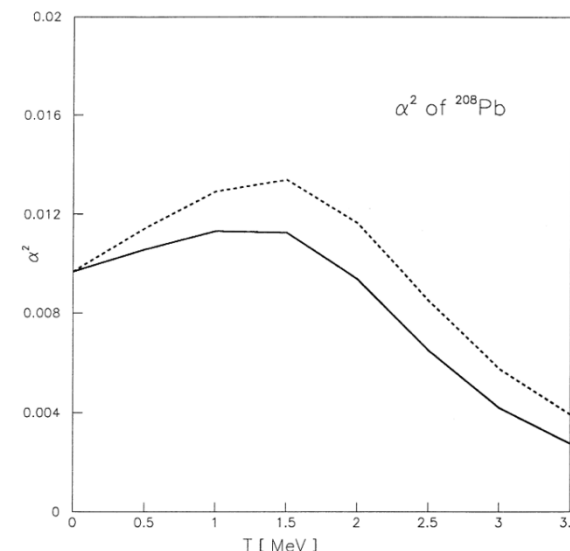
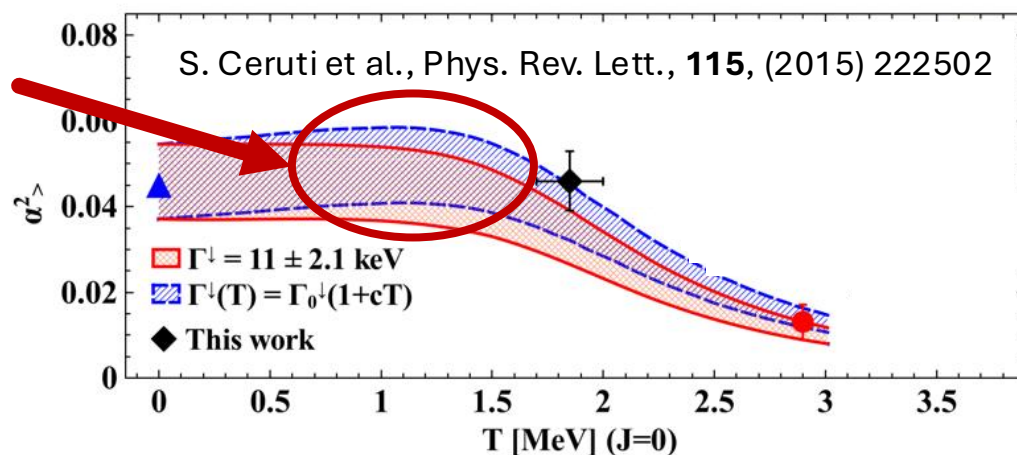
ADVATAGE: In the plateau region, only one data point is needed to extract the isospin-mixing value at $T=0$. ✓

DISADVATAGE: Low reaction cross-section

ISOSPIN MIXING IN ^{72}Kr

- ^{72}Kr is a N=Z nucleus, produced by fusion-evaporation reactions.
- It has an intermediate mass between ^{60}Zn and ^{80}Zr .
- It has a E1 discrete transition so it can be measured with two techniques.

Add a point
in the
plateau
region



H. Sagawa
et al. Phys.
Lett. B, **444**,
(1998), 1–6

ADVATAGE: In the plateau region, only one data point is needed to extract the isospin-mixing value at $T=0$. ✓

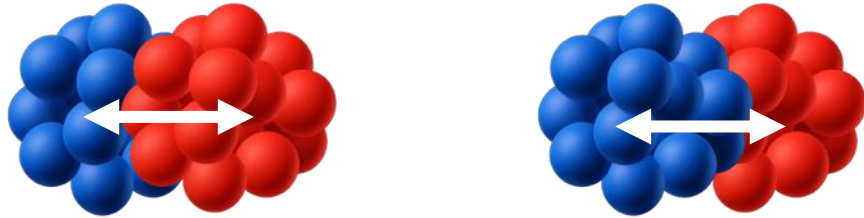
DISADVATAGE: Low reaction cross-section ✗



EXPERIMENTAL TECHNIQUE AND SETUP

ISOSPIN MIXING VIA GDR

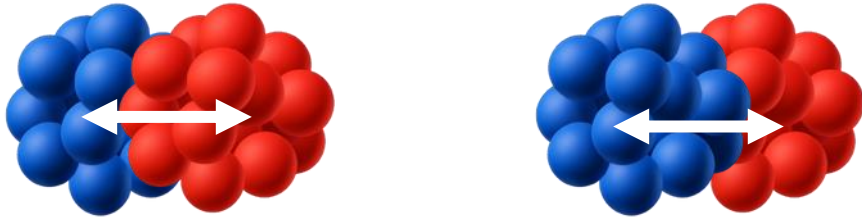
EXPERIMENTAL TECHNIQUE: GDR (I)



GDR can be described as an oscillation of the protons against the neutrons.

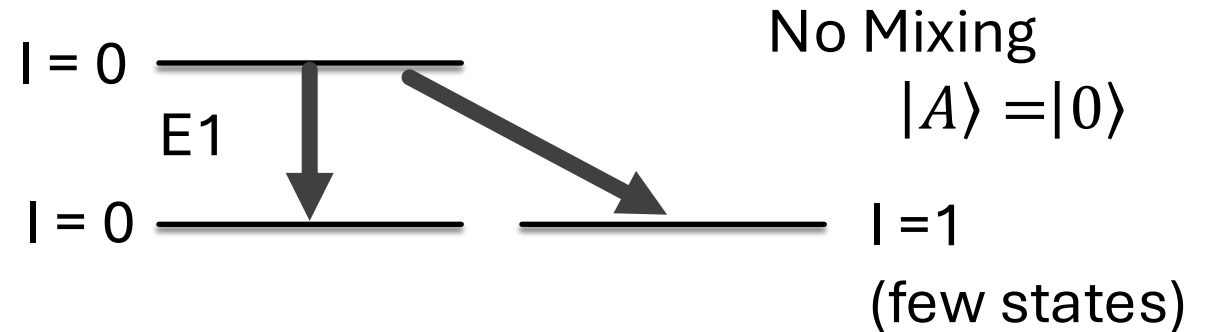
It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.

EXPERIMENTAL TECHNIQUE: GDR (I)

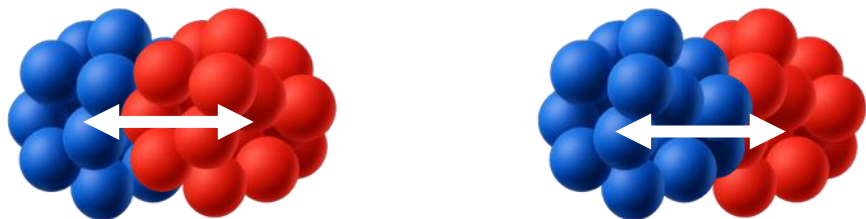


GDR can be described as an oscillation of the protons against the neutrons.

It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.

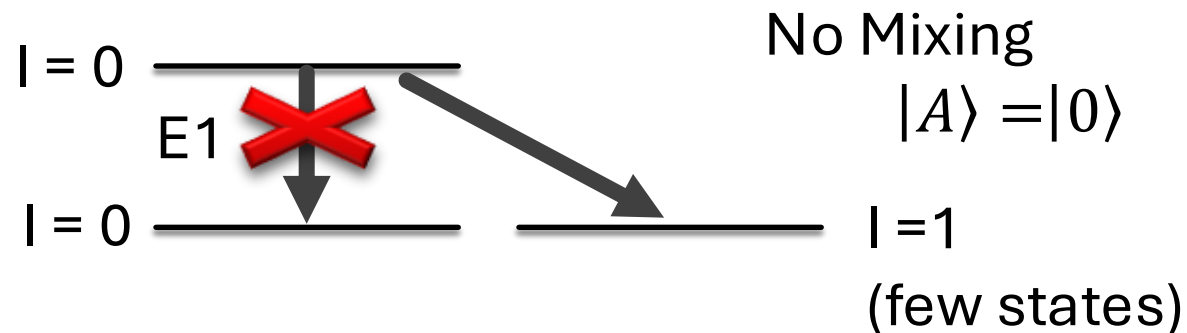


EXPERIMENTAL TECHNIQUE: GDR (I)

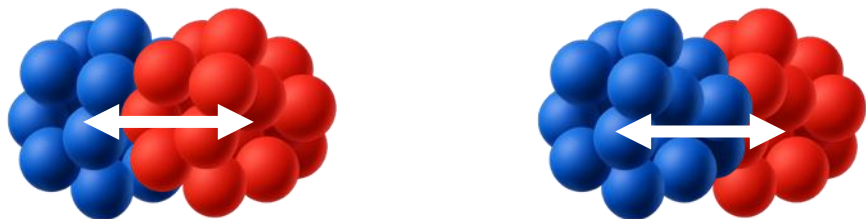


GDR can be described as an oscillation of the protons against the neutrons.

It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.

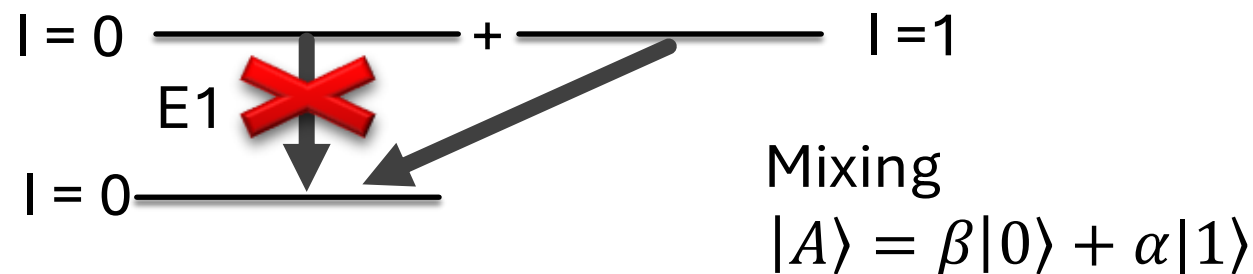
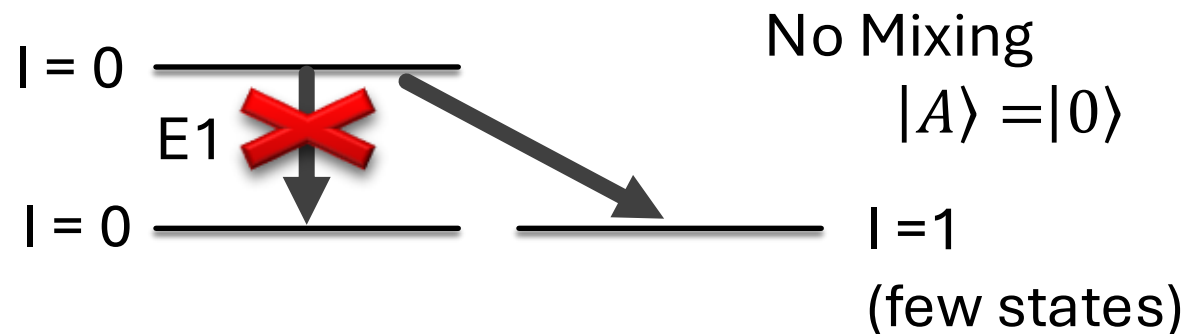


EXPERIMENTAL TECHNIQUE: GDR (I)

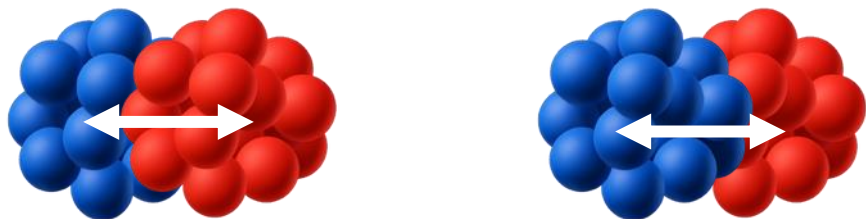


GDR can be described as an oscillation of the protons against the neutrons.

It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.



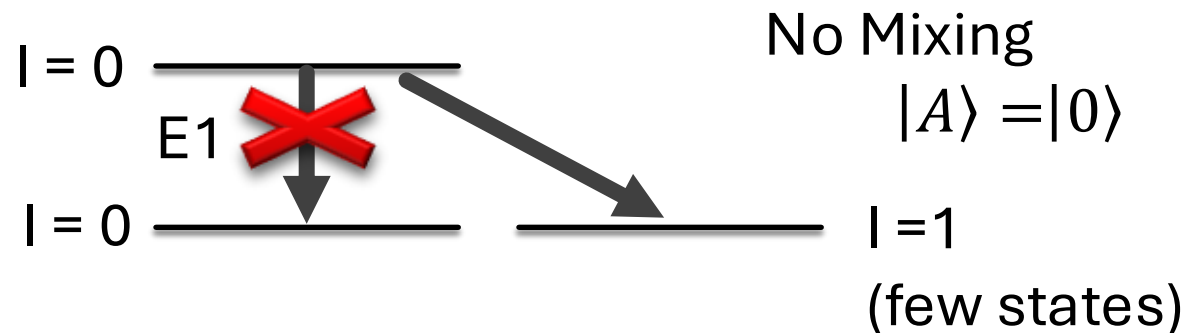
EXPERIMENTAL TECHNIQUE: GDR (I)



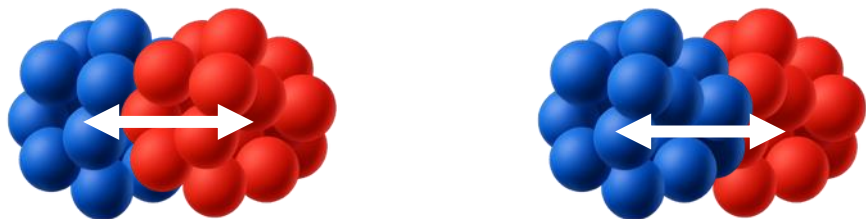
GDR can be described as an oscillation of the protons against the neutrons.

It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.

Study of the E1 GDR decay at finite temperature to find the mixing.



EXPERIMENTAL TECHNIQUE: GDR (I)

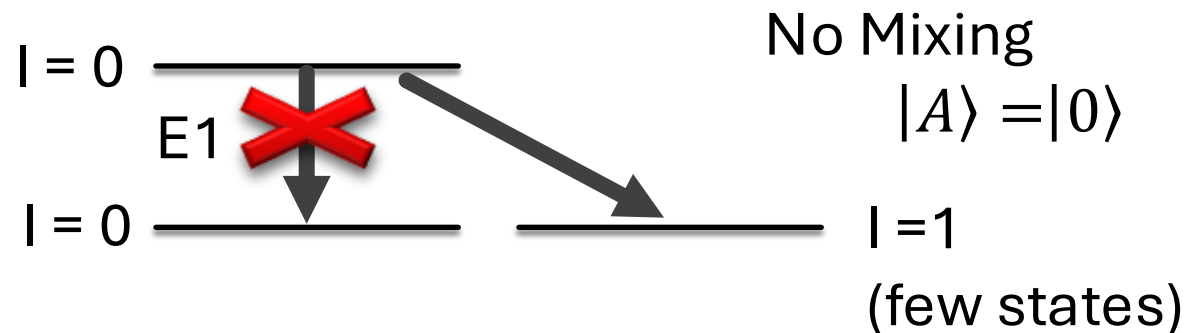


GDR can be described as an oscillation of the protons against the neutrons.

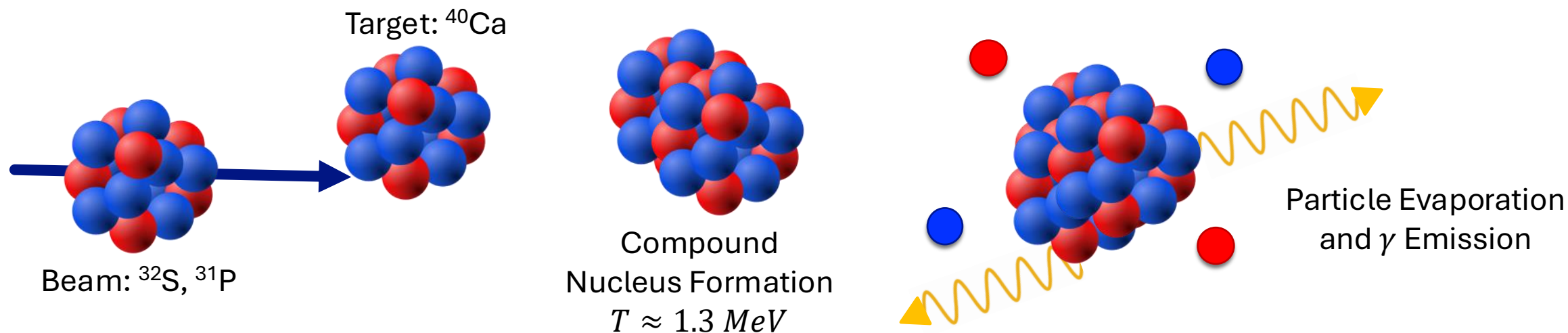
It is an electric ($\Delta S = 0$) and isovector ($\Delta I = 1$) transition.

Study of the E1 GDR decay at finite temperature to find the mixing.

To form a compound nucleus at $T > 0$ we used a fusion-evaporation reaction.



EXPERIMENTAL TECHNIQUE: GDR (II)



$I = 0$ compound nucleus reaction:



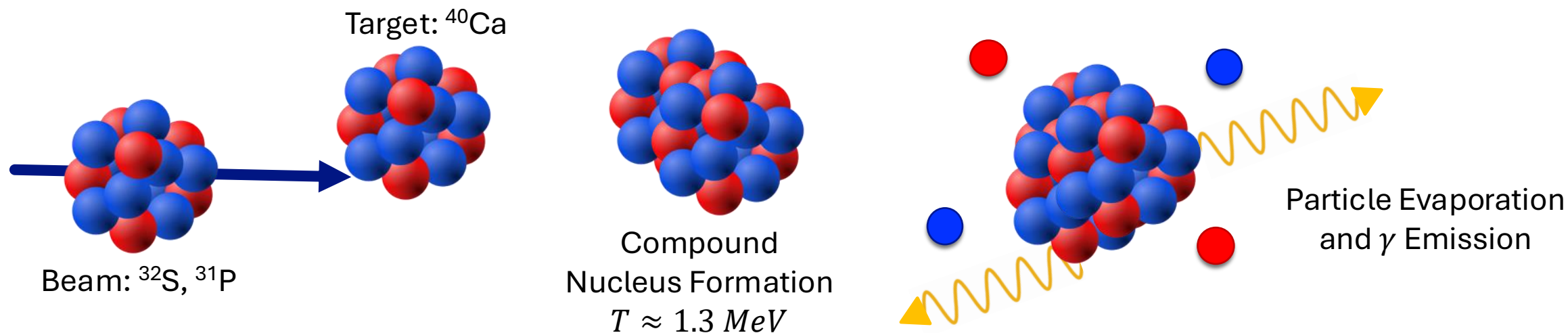
The observed E1 strength is a signature of the mixing.

$I \neq 0$ compound nucleus reaction:



Full mixing. No E1 strength reduction is expected.

EXPERIMENTAL TECHNIQUE: GDR (II)



$I = 0$ compound nucleus reaction:



The observed E1 strength is a signature of the mixing.

$I \neq 0$ compound nucleus reaction:

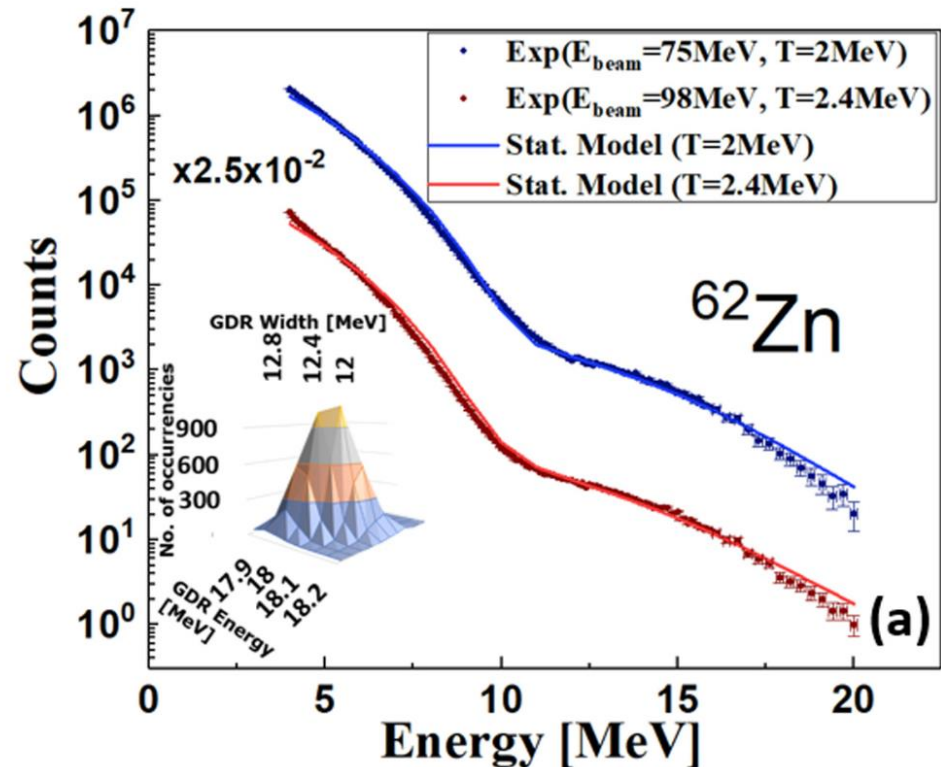


Full mixing. No E1 strength reduction is expected.

GDR manifest itself as E1 Strength $\rightarrow$ GDR is a probe of the isospin mixing

DATA ANALYSIS TECHNIQUE

Reaction with $I \neq 0$, no GDR decay suppression



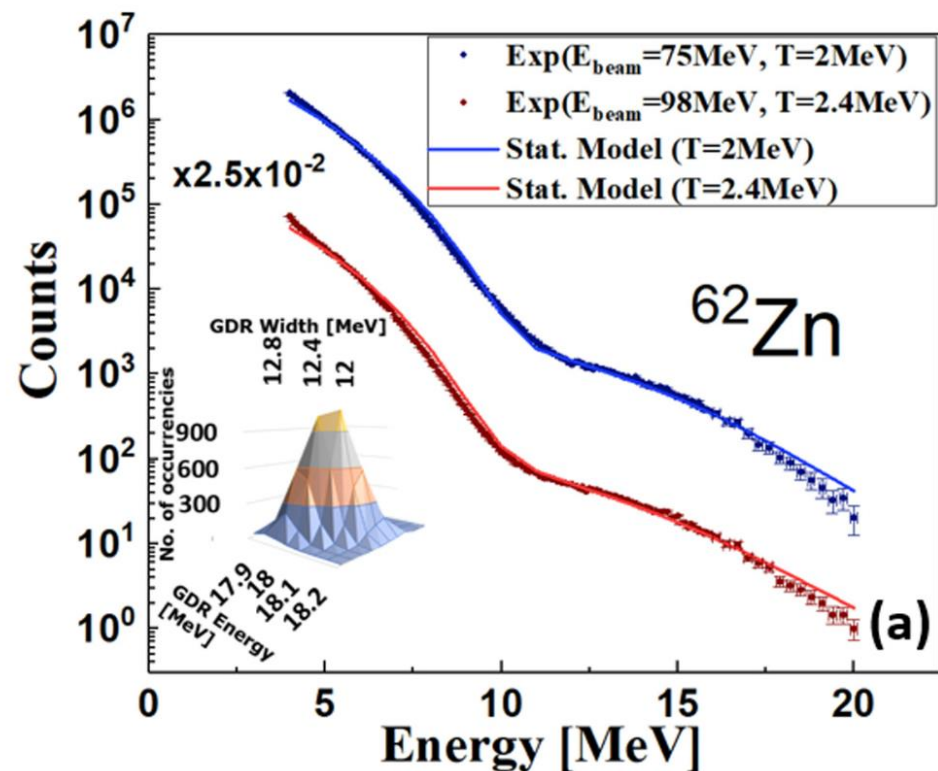
G. Gosta et al., PRC 103, L041302 (2021)

Statistical model to extract of GDR parameters:

centroid, width and strenght

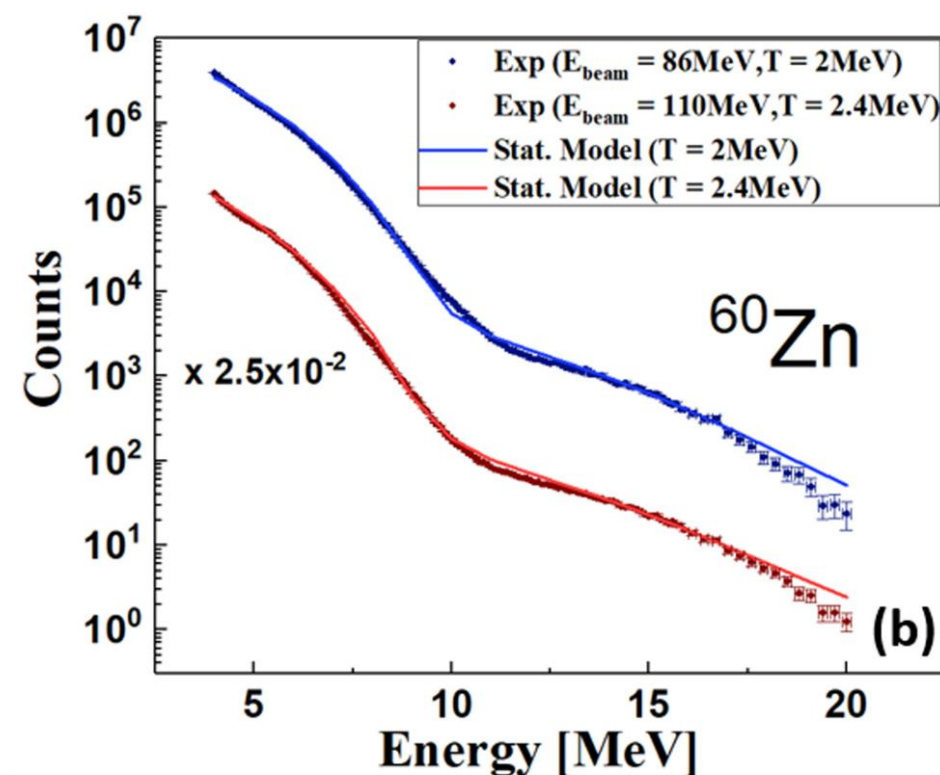
DATA ANALYSIS TECHNIQUE

Reaction with $I \neq 0$, no GDR decay suppression



G. Gosta et al., PRC 103, L041302 (2021)

Reaction with $I = 0$, GDR decay possible only in case of isospin mixing



G. Gosta et al., PRC 103, L041302 (2021)

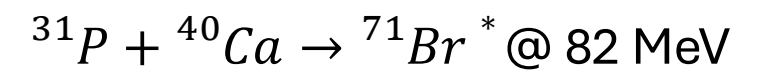
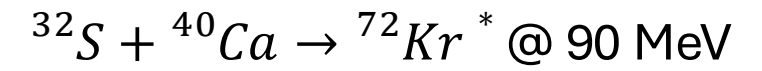
Statistical model to extract of GDR parameters:

centroid, width and strenght

Statistical model to extract the

Coulomb Spreading Width $\Gamma_c^\downarrow$

EXPERIMENTAL SETUP



Low cross-section



High efficiency

EXPERIMENTAL SETUP

ELIFANT-GG Array:

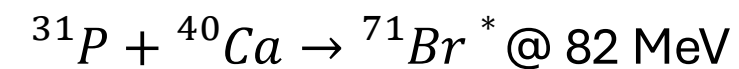
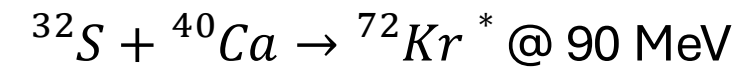
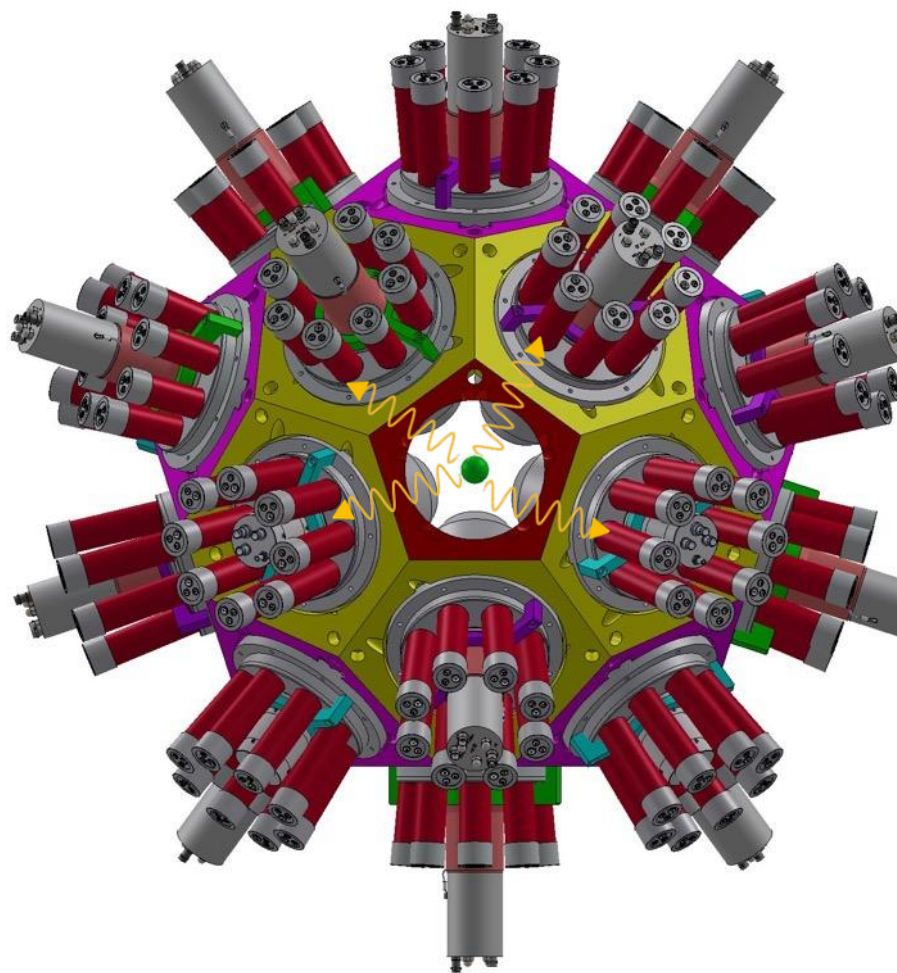
High-energy γ rays:

21 Bromide Scintillator

Detector Array: 11 3''x3''

LaBr₃:Ce + 10 CeBr₃ (3''x3'')

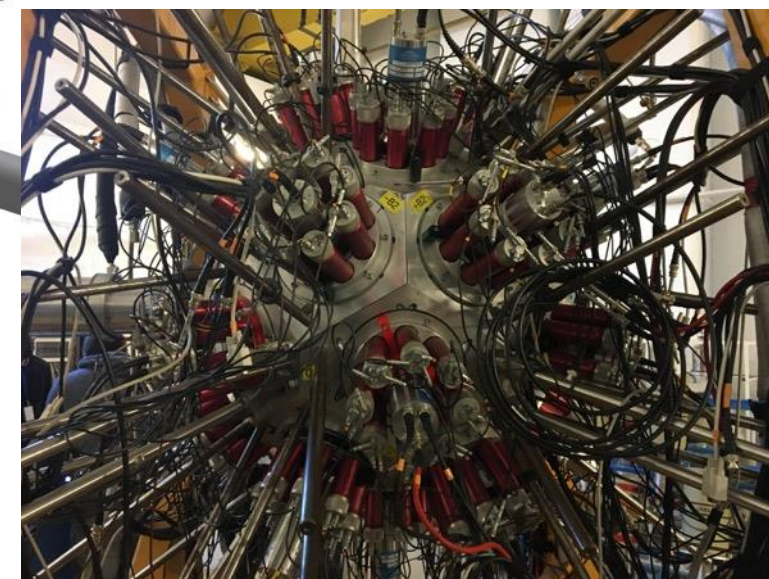
(AC shield)



Low cross-section



High efficiency



EXPERIMENTAL SETUP

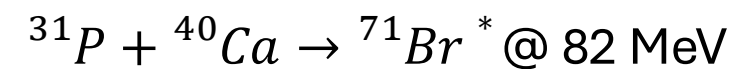
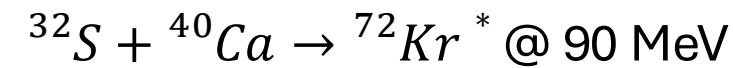
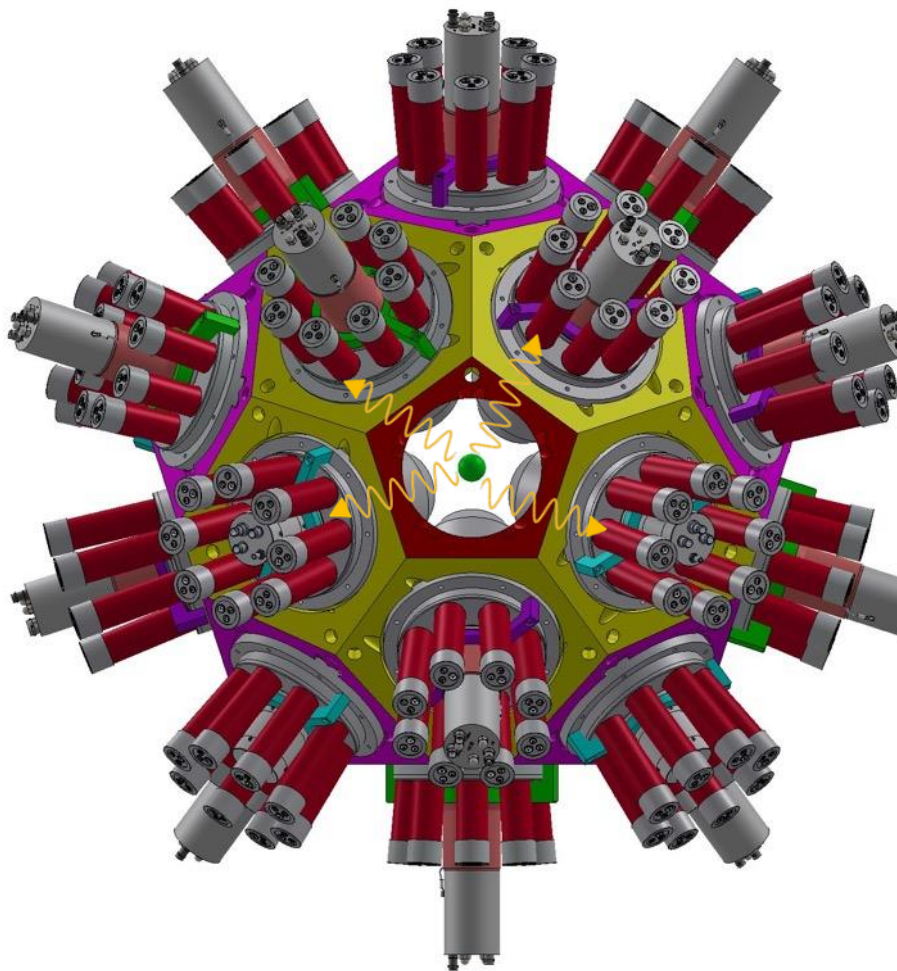
ELIFANT-GG Array:

High-energy γ rays:

21 Bromide Scintillator
Detector Array: 11 3''x3''
LaBr₃:Ce + 10 CeBr₃ (3''x3'')
(AC shield)

Fusion-evaporation residues (low-energy γ rays):

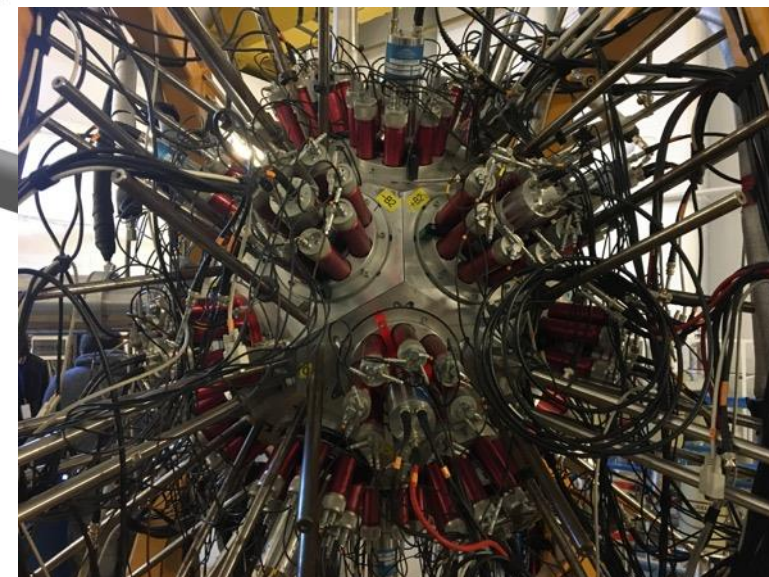
4 AC-shield HPGe detectors



Low cross-section



High efficiency





RESULTS

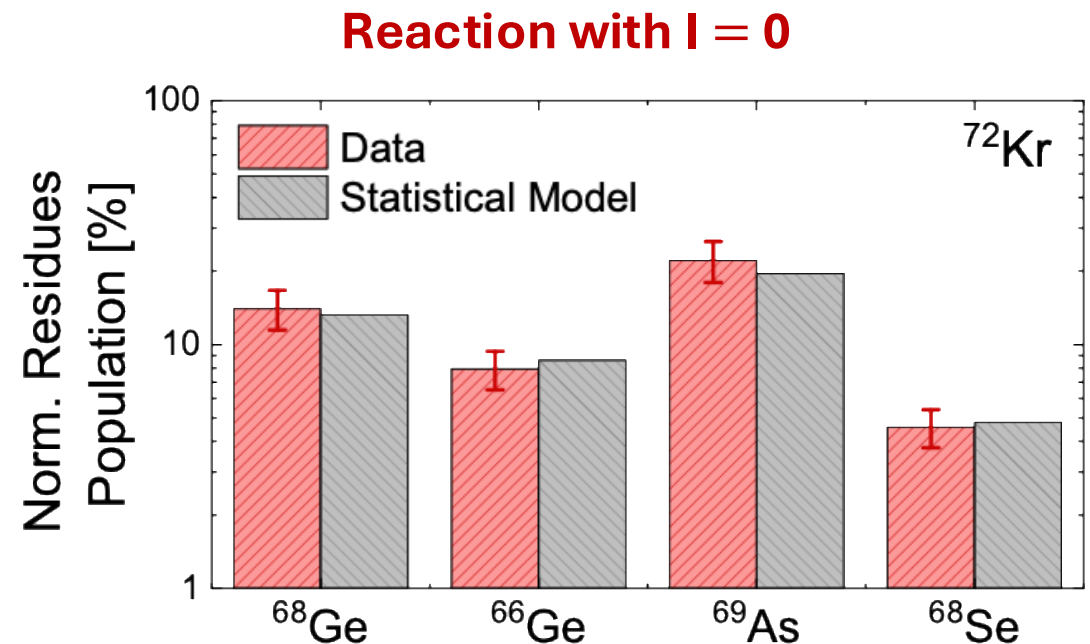
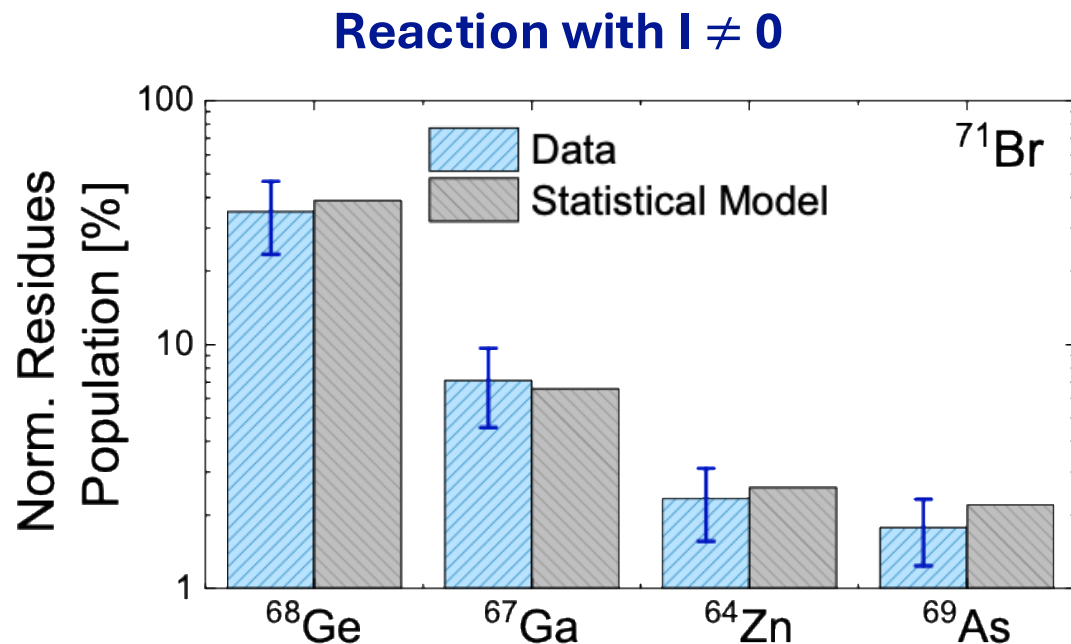
ISOSPIN MIXING PARAMETRES: α^2 & δ_c

RESULTS: EVAPORATION RESIDUES

Fusion-evaporation residues (low-energy γ rays): 4 AC-shield HPGe detectors

Experimental population of the residues compared with the statistical model prediction.

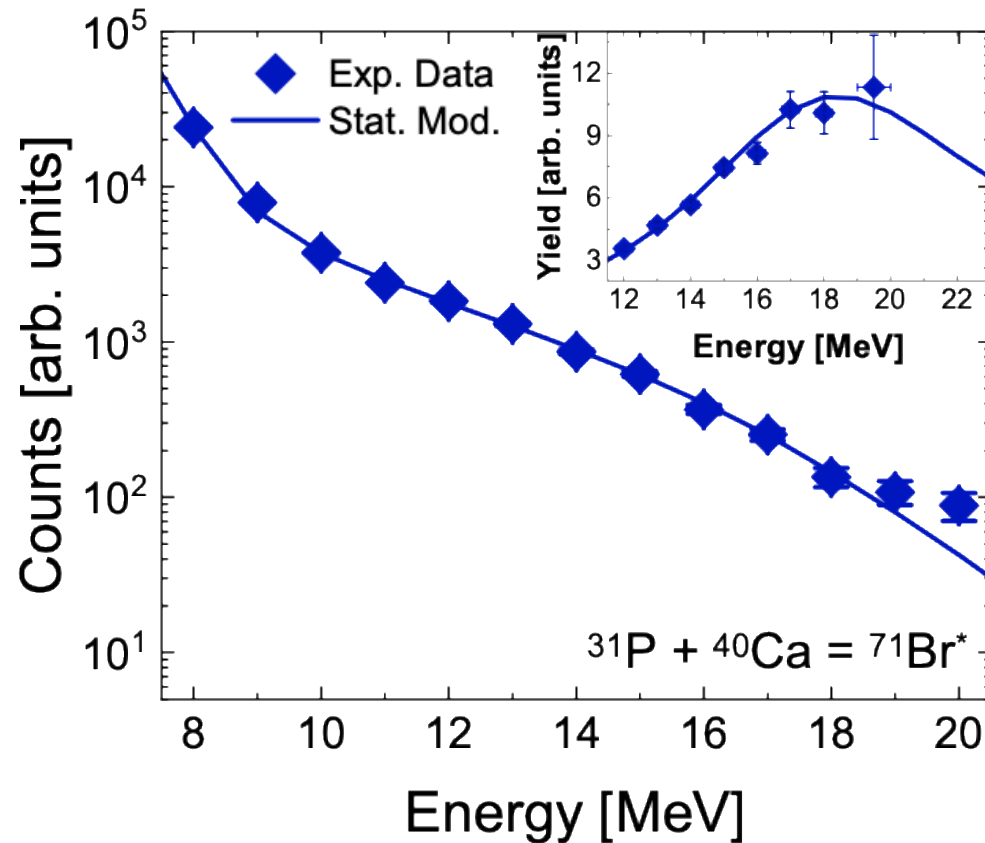
The most populated residues are reproduced within the error bars.



A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

RESULTS: HIGH-ENERGY γ -RAY SPECTRA

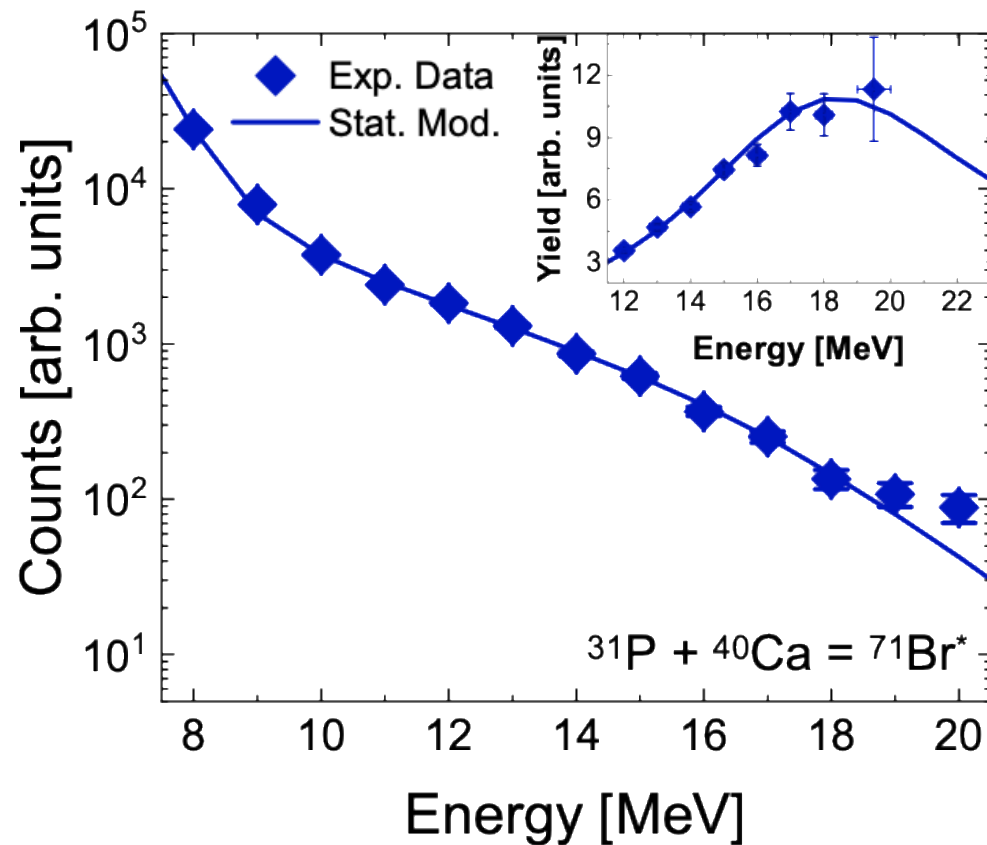
$$I \neq 0 - E_{GDR} = 18.4 \pm 0.5 \text{ MeV} \text{ \& } \Gamma_{GDR} = 11 \pm 1 \text{ MeV}$$



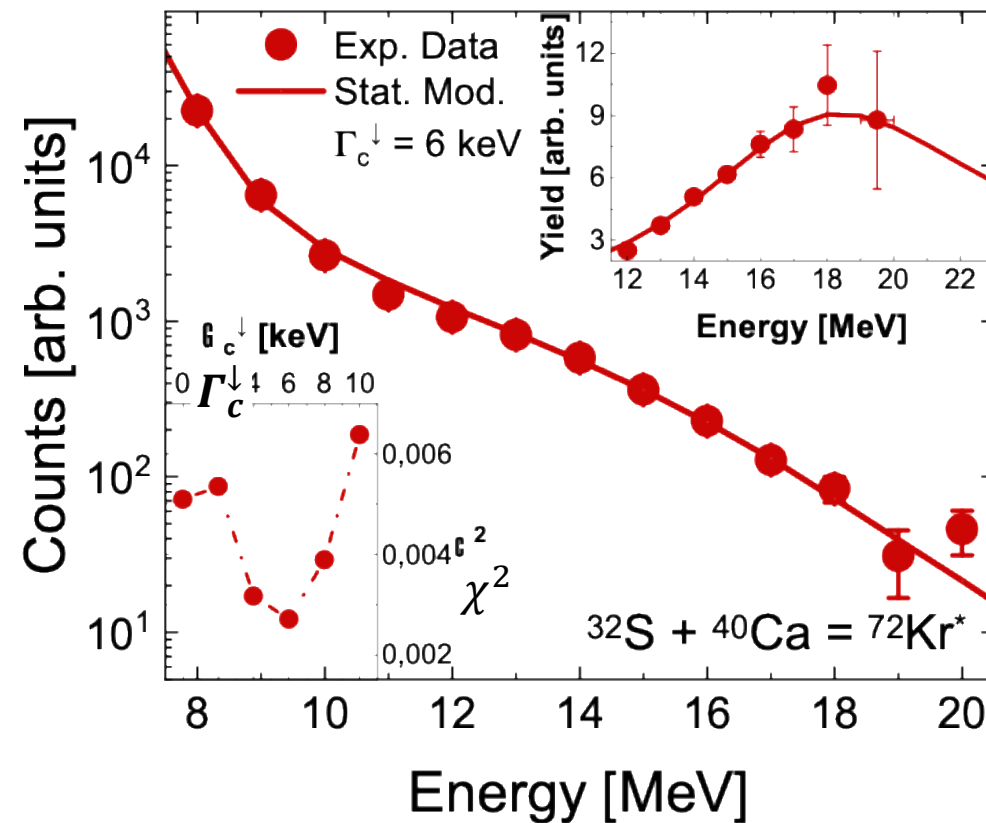
A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

RESULTS: HIGH-ENERGY γ -RAY SPECTRA

$$I \neq 0 - E_{GDR} = 18.4 \pm 0.5 \text{ MeV} \text{ \& } \Gamma_{GDR} = 11 \pm 1 \text{ MeV}$$



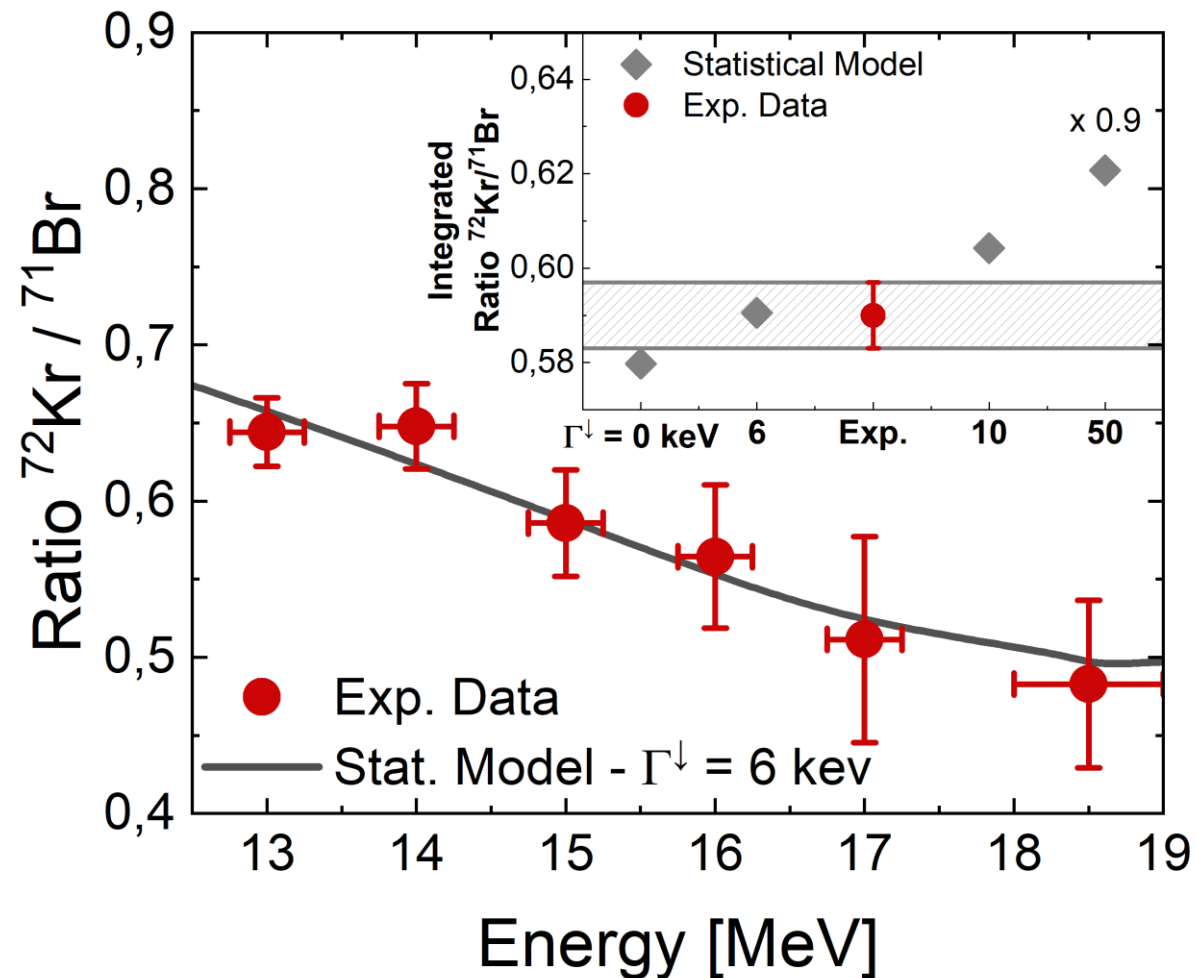
$$I = 0 - \text{Mixing } \Gamma_c^\downarrow = 6 \text{ keV}$$



A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

RESULTS: MIXING ESTIMATION

Statistical model with different values of $\Gamma_c^\downarrow$: 0 (no mixing), 6, 10 and 50 keV (full mixing).



A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

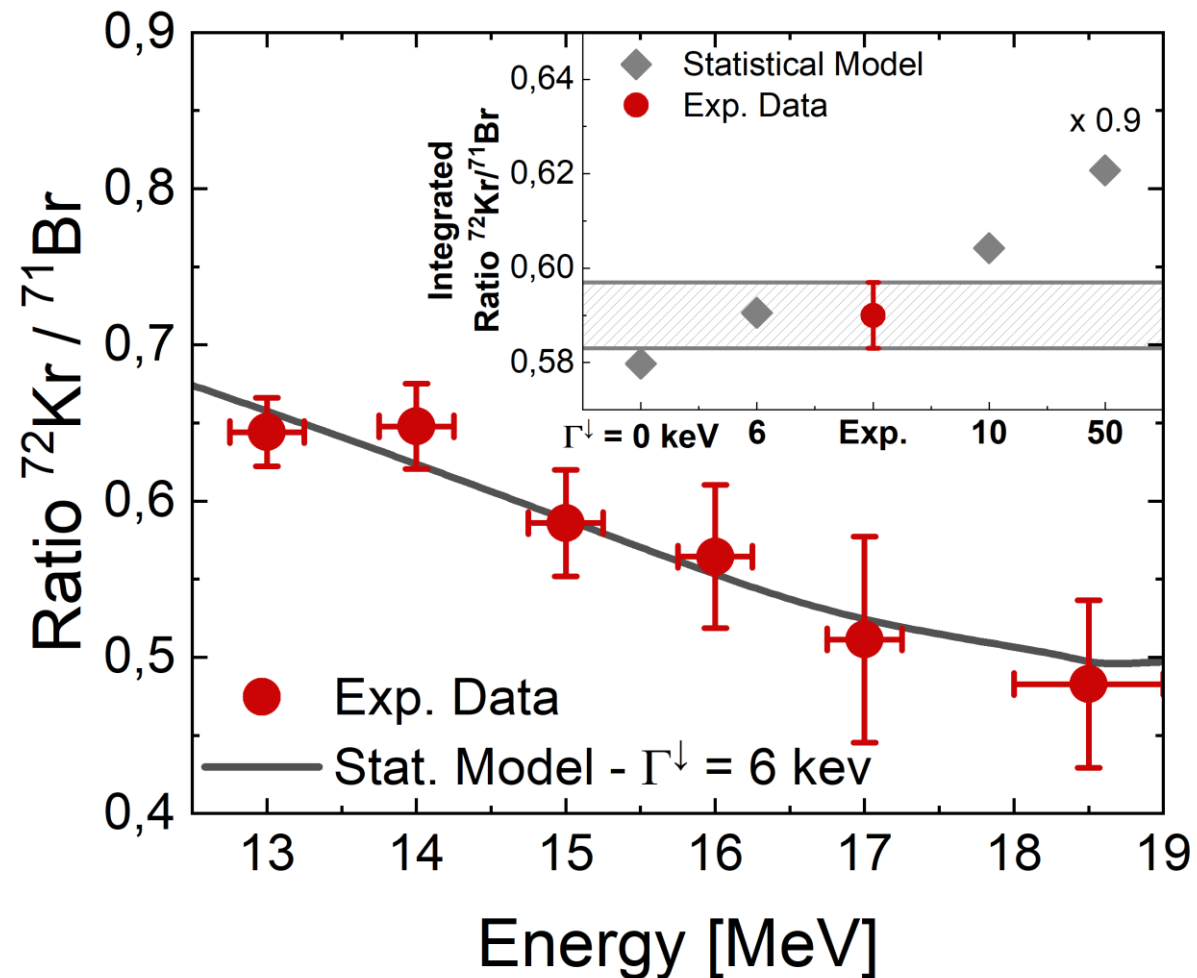
RESULTS: MIXING ESTIMATION

Statistical model with different values of $\Gamma_c^\downarrow$: 0 (no mixing), 6, 10 and 50 keV (full mixing).

Integrated Ratio ^{72}Kr and ^{71}Br from 13 MeV to 18 MeV:

- Experimental data:

$$0.590 \pm 0.007$$



RESULTS: MIXING ESTIMATION

Statistical model with different values of $\Gamma_c^\downarrow$: 0 (no mixing), 6, 10 and 50 keV (full mixing).

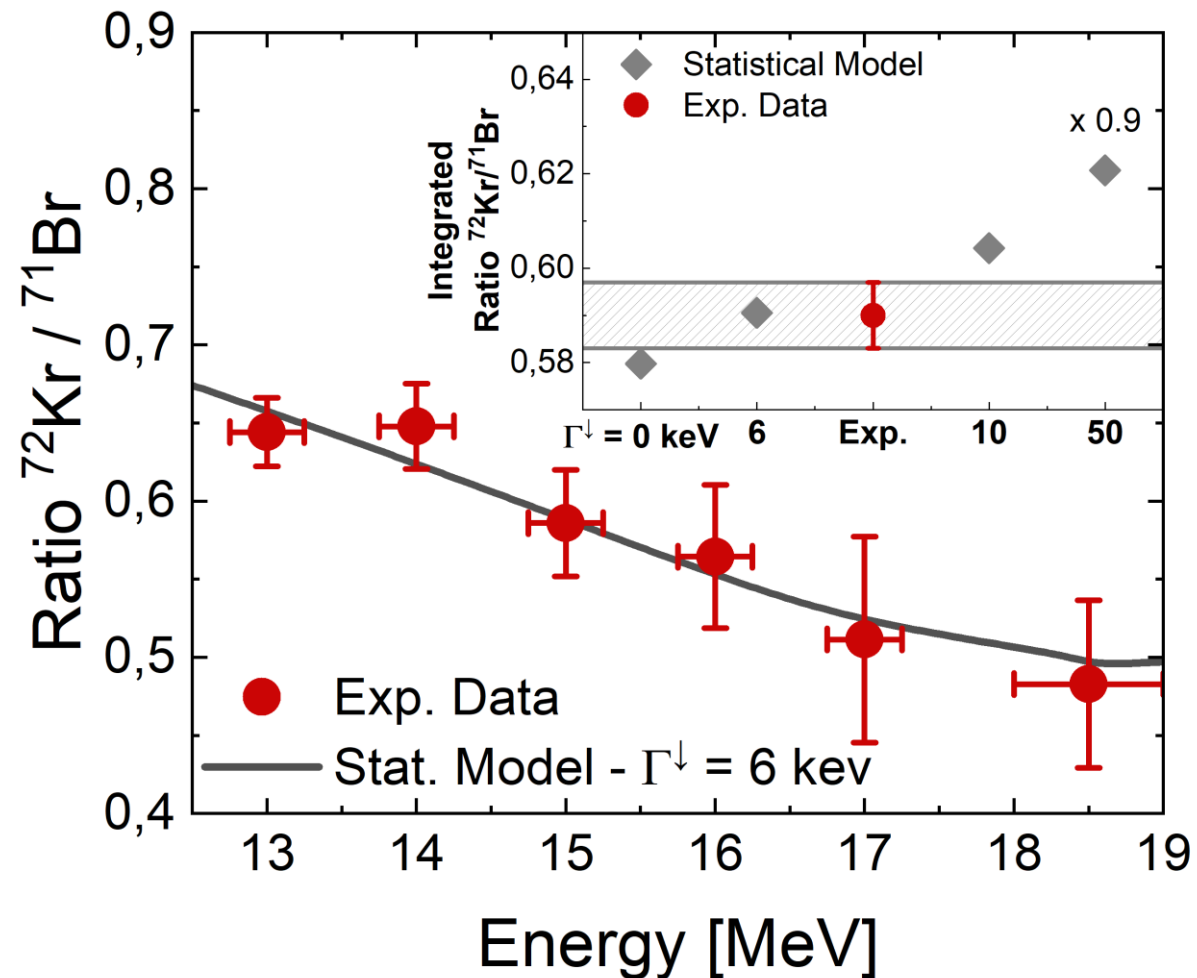
Integrated Ratio ^{72}Kr and ^{71}Br from 13 MeV to 18 MeV:

- Experimental data:

$$0.590 \pm 0.007$$

- Statistical model $\Gamma_c^\downarrow = 6 \text{ keV}$:

$$0.591$$



RESULTS: α^2 MIXING PARAMETER (I)

From $T>0$ to $T=0$

α^2 trend as a function of temperature:

$$\alpha^2(T) = \frac{1}{I_0 + 1} \frac{\Gamma_{IAS}^{\downarrow}}{\Gamma_{CN}(T) + \Gamma_{IVM}^{(IAS)}} \quad [1]$$

[1] G.C.H. Sagawa, P.F. Bortignon, Phys. Lett. B, 444, (1998), 1-6

RESULTS: α^2 MIXING PARAMETER (I)

From $T>0$ to $T=0$

α^2 trend as a function of temperature:

$$\alpha^2(T) = \frac{1}{I_0 + 1} \frac{\Gamma_{IAS}^{\downarrow}}{\Gamma_{CN}(T) + \Gamma_{IVM}^{(IAS)}} \quad [1]$$

 200 keV [1]

- $\Gamma_{IVM}^{(IAS)}$ width of IVM resonance at excitation energy of IAS

[1] G.C.H. Sagawa, P.F. Bortignon, Phys. Lett. B, 444, (1998), 1-6

RESULTS: α^2 MIXING PARAMETER (I)

From $T>0$ to $T=0$

α^2 trend as a function of temperature:

$$\alpha^2(T) = \frac{1}{I_0 + 1} \frac{\Gamma_{IAS}^\downarrow}{\Gamma_{CN}(T) + \Gamma_{IVM}^{(IAS)}} \quad [1]$$

From Stat. Model 200 keV [1]

- $\Gamma_{IVM}^{(IAS)}$ width of IVM resonance at excitation energy of IAS
- $\Gamma_{CN}(T)$ is the compound nucleus decay width

[1] G.C.H. Sagawa, P.F. Bortignon, Phys. Lett. B, 444, (1998), 1-6

RESULTS: α^2 MIXING PARAMETER (I)

From $T>0$ to $T=0$

α^2 trend as a function of temperature:

$$\alpha^2(T) = \frac{1}{I_0 + 1} \frac{\Gamma_{IAS}^\downarrow}{\Gamma_{CN}(T) + \Gamma_{IVM}^{(IAS)}} \quad [1]$$

$6.0 \pm 2.0 \text{ keV}$
 200 keV [1]

From Stat. Model

- $\Gamma_{IVM}^{(IAS)}$ width of IVM resonance at excitation energy of IAS
- $\Gamma_{CN}(T)$ is the compound nucleus decay width
- $\Gamma_{IAS}^\downarrow$ is $\Gamma_c^\downarrow$

[1] G.C.H. Sagawa, P.F. Bortignon, Phys. Lett. B, 444, (1998), 1-6

RESULTS: α^2 MIXING PARAMETER (I)

From $T > 0$ to $T = 0$

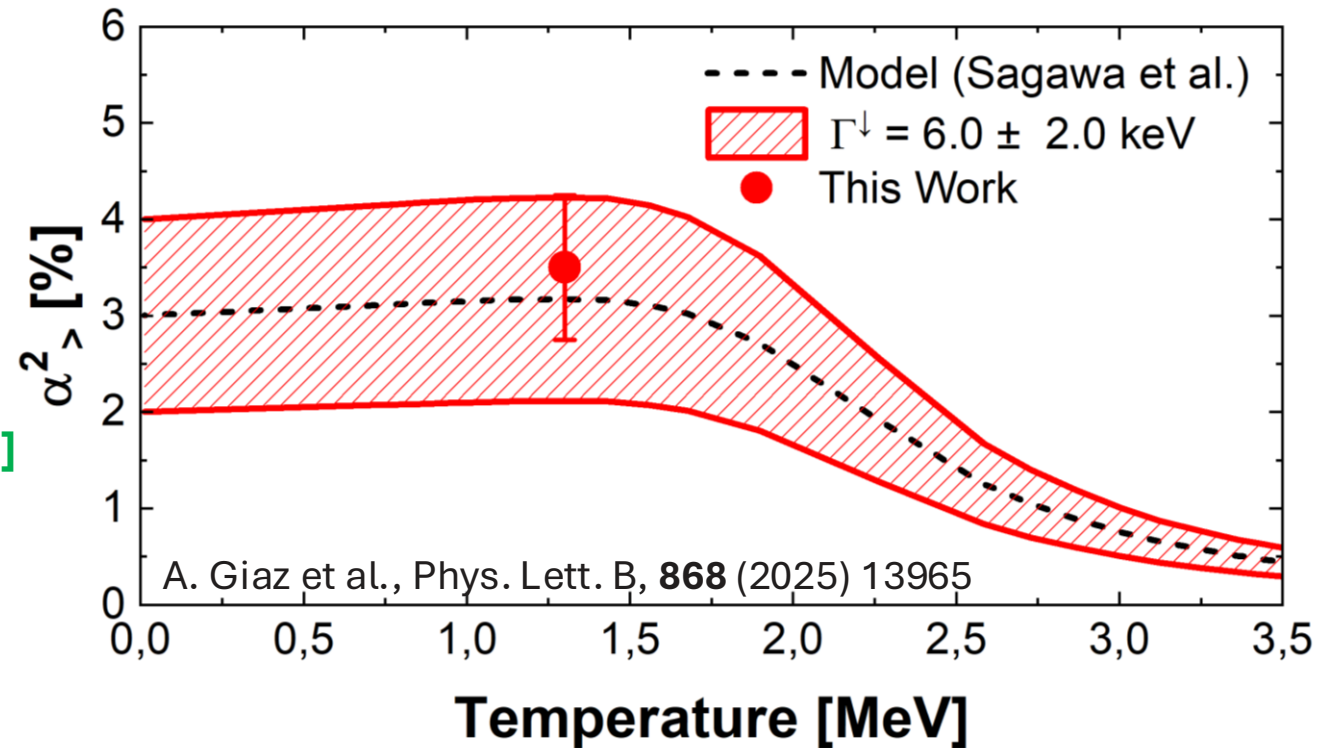
α^2 trend as a function of temperature:

$$\alpha^2(T) = \frac{1}{I_0 + 1} \frac{\Gamma_{IAS}^\downarrow}{\Gamma_{CN}(T) + \Gamma_{IVM}^{(IAS)}} \quad [1]$$

$6.0 \pm 2.0 \text{ keV}$
 200 keV [1]

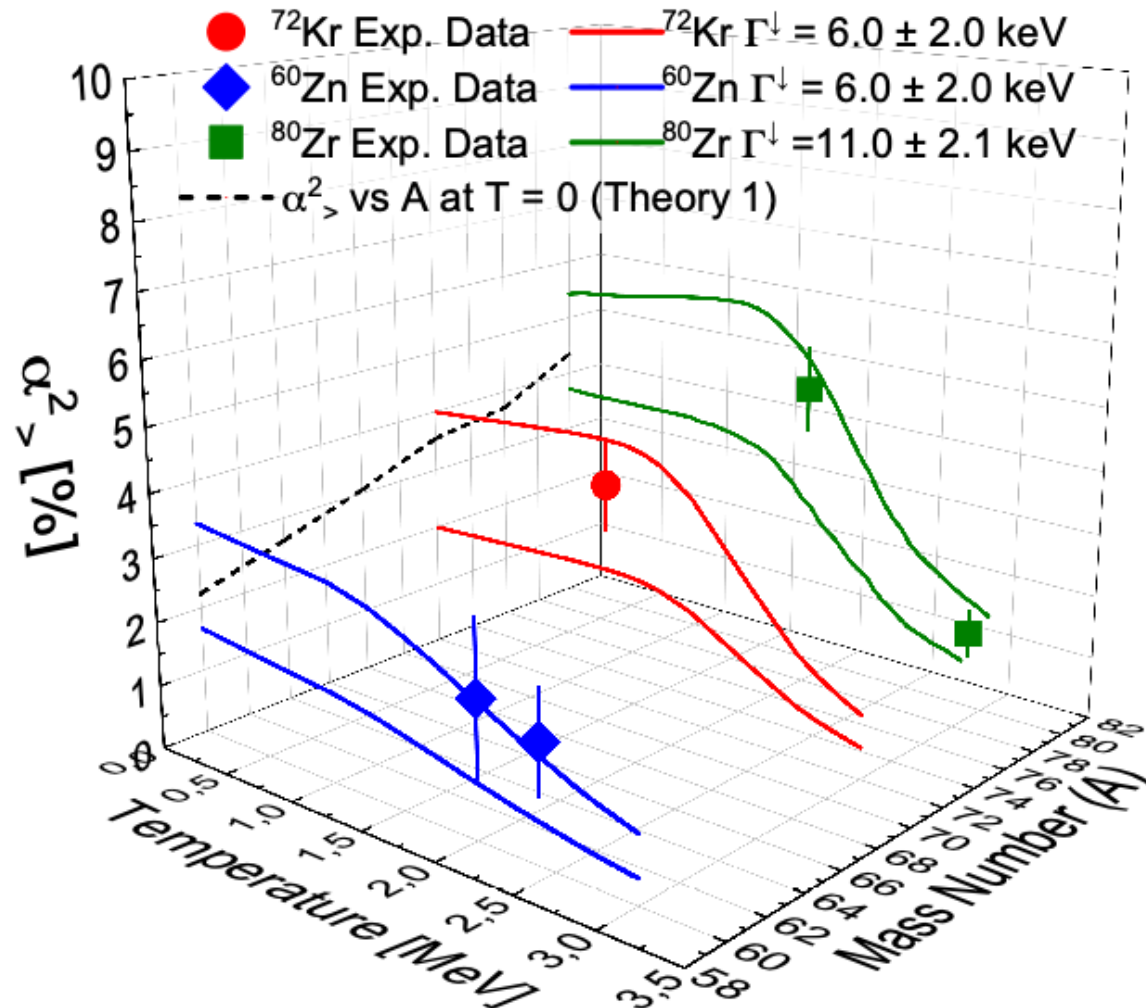
From Stat. Model

- $\Gamma_{IVM}^{(IAS)}$ width of IVM resonance at excitation energy of IAS
- $\Gamma_{CN}(T)$ is the compound nucleus decay width
- $\Gamma_{IAS}^\downarrow$ is $\Gamma_c^\downarrow$



[1] G.C.H. Sagawa, P.F. Bortignon, Phys. Lett. B, 444, (1998), 1-6

RESULTS: α^2 MIXING PARAMETER (II)

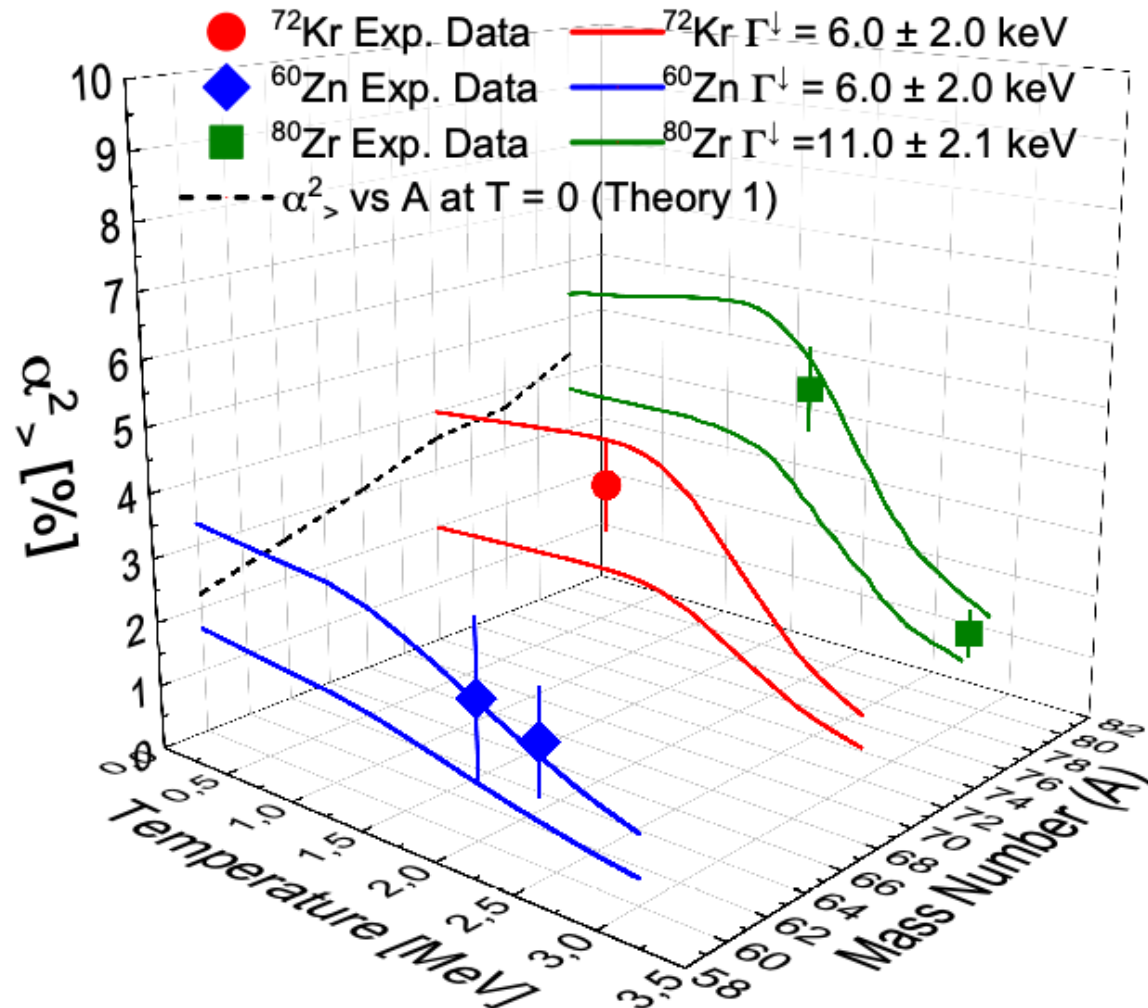


From finite temperature to zero temperature:

- ^{60}Zn : $(2.5 \pm 0.8) \%$ @ $T = 0$ [1]
- ^{80}Zr : $(4.6 \pm 0.9) \%$ @ $T = 0$ [2]
- ^{72}Kr : $(3.0 \pm 0.8) \%$ @ $T = 0$ and $(3.5 \pm 0.8) \%$ @ $T = 1.3 \text{ MeV}$ [3]

- [1] G. Gosta et al., Phys. Rev. C **103**, (2021) L041302
 [2] S. Ceruti et al., Phys. Rev. Lett. **115**, (2015) 222502
 [3] A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

RESULTS: α^2 MIXING PARAMETER (II)



From finite temperature to zero temperature:

- ^{60}Zn : $(2.5 \pm 0.8) \%$ @ $T = 0$ [1]
- ^{80}Zr : $(4.6 \pm 0.9) \%$ @ $T = 0$ [2]
- ^{72}Kr : $(3.0 \pm 0.8) \%$ @ $T = 0$ and $(3.5 \pm 0.8) \%$ @ $T = 1.3 \text{ MeV}$ [3]

Mixing Probability α^2 follows the expected trend both as a function of **Mass Number (A)** and **Temperature**

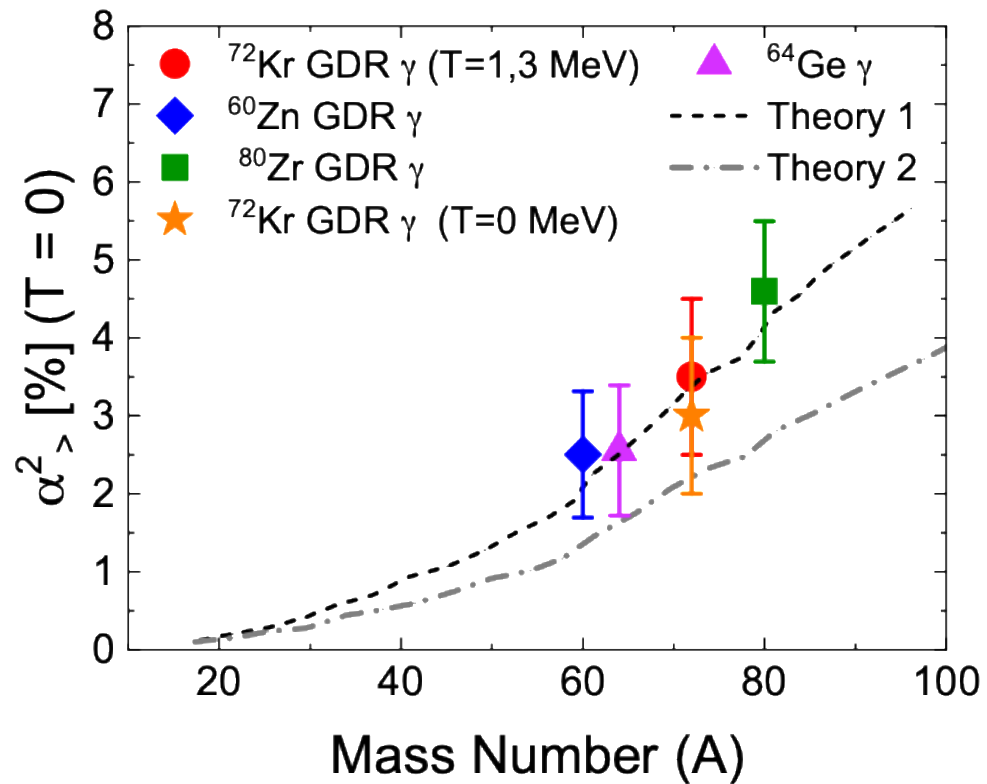
[1] G. Gosta et al., Phys. Rev. C **103**, (2021) L041302

[2] S. Ceruti et al., Phys. Rev. Lett. **115**, (2015) 222502

[3] A. Giaz et al., Phys. Lett. B, **868** (2025) 13965

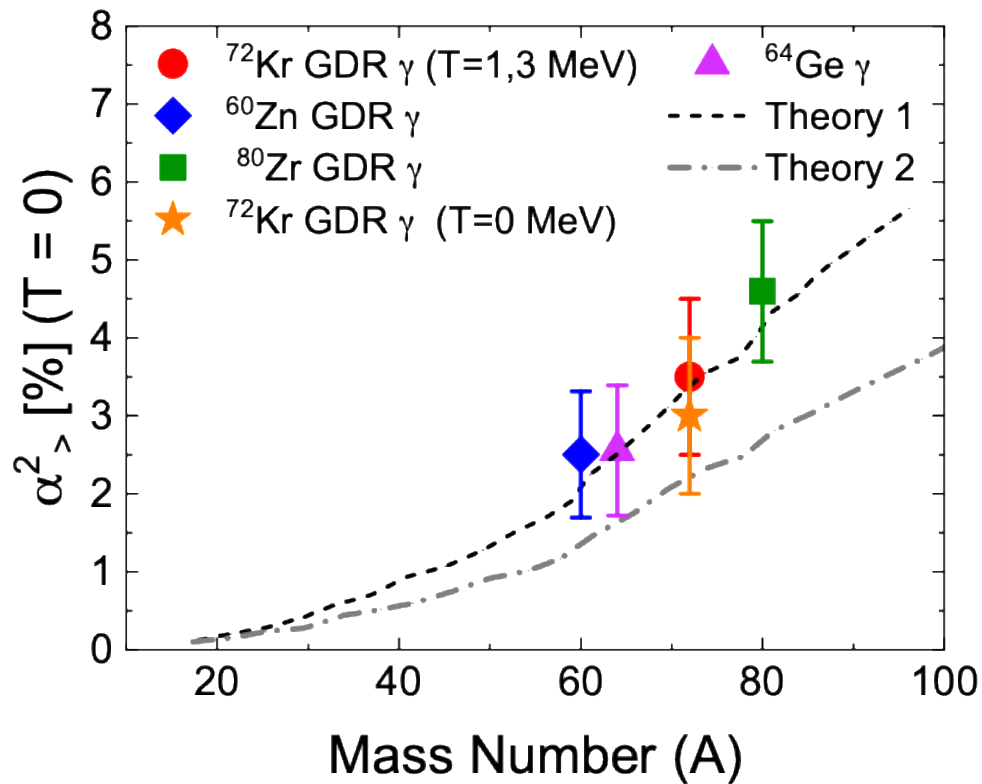
RESULTS: CORRECTION TERM δ_c

From α^2



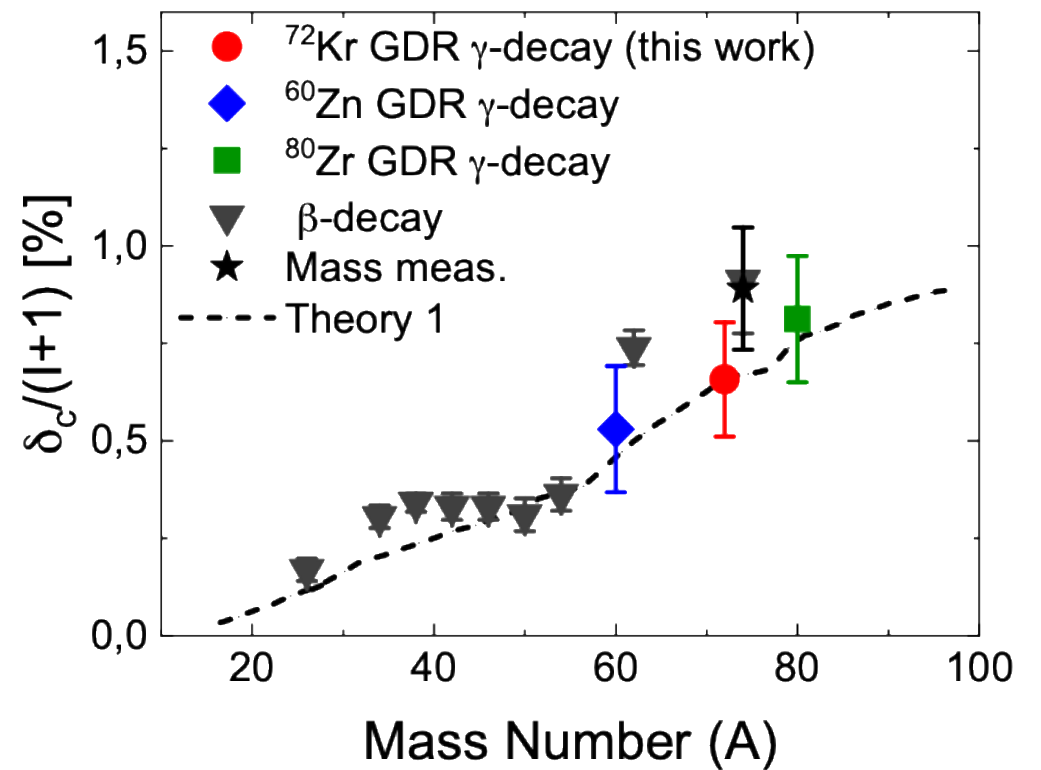
RESULTS: CORRECTION TERM δ_c

From α^2



$$\text{To } \delta_c = 4(I + 1) \frac{V_1}{41\xi A^{2/3}} \alpha^2 [1]$$

$$\delta_c = (0.66 \pm 0.14)\%$$



[1] N. Auerbach, Phys. Rev. **C 79**, 035502 (2009)

SUMMARY AND PRESPECTIVES

SUMMARY

- We measure the isospin mixing via GDR in ^{72}Kr at ELIGANT setup at the IFIN-HH Tandem (Bucharest)
- We found $\alpha^2 = (3.5 \pm 0.8)\%$ at $T=1.3$ MeV and $\alpha^2 = (3.0 \pm 0.8)\%$ at $T=0$
- We obtained the correction term $\delta_c = (0.66 \pm 0.14)\%$

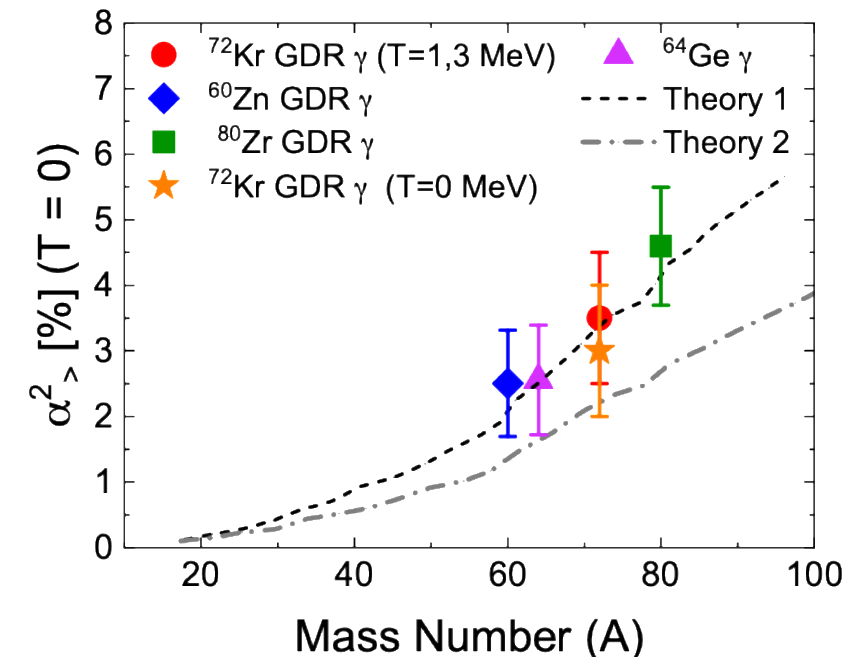
SUMMARY AND PRESPECTIVES

SUMMARY

- We measure the isospin mixing via GDR in ^{72}Kr at ELIGANT setup at the IFIN-HH Tandem (Bucharest)
- We found $\alpha^2 = (3.5 \pm 0.8)\%$ at $T=1.3$ MeV and $\alpha^2 = (3.0 \pm 0.8)\%$ at $T=0$
- We obtained the correction term $\delta_c = (0.66 \pm 0.14)\%$

PRESPECTIVES

- **Exotic beam** to reach $N=Z$ nuclei with $A > 80$
- Measurement of **isospin mixing at $T=0$** via GDR through inelastic scattering



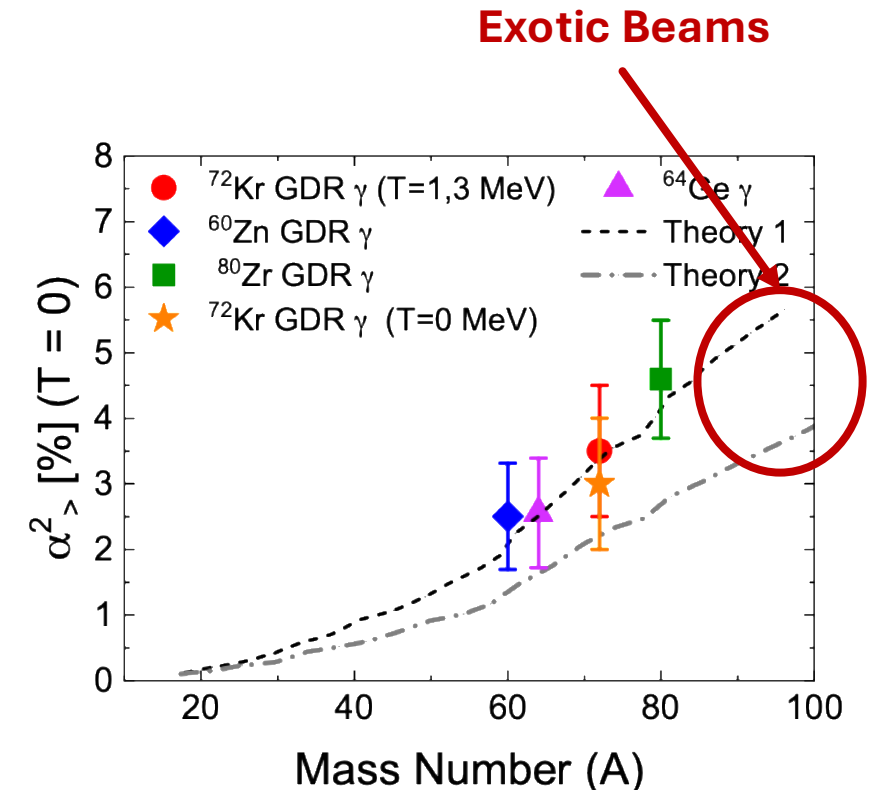
SUMMARY AND PRESPECTIVES

SUMMARY

- We measure the isospin mixing via GDR in ^{72}Kr at ELIGANT setup at the IFIN-HH Tandem (Bucharest)
- We found $\alpha^2 = (3.5 \pm 0.8)\%$ at $T=1.3$ MeV and $\alpha^2 = (3.0 \pm 0.8)\%$ at $T=0$
- We obtained the correction term $\delta_c = (0.66 \pm 0.14)\%$

PRESPECTIVES

- **Exotic beam** to reach $N=Z$ nuclei with $A > 80$
- Measurement of **isospin mixing at $T=0$** via GDR through inelastic scattering



COLLABORATORS

Agnese Giaz^a, Oliver Wieland^a, Angela Bracco^{b,a}, Franco Camera^{b,a}, Soichiro Aogaki^c, Dimiter L. Balabanskic, Maria Brezeanu^c, Ruxandra Borcea^d, Marian Boromiza^d, Ion Burducea^e, Stefana Calinescu^d, Adina Coman^d, Paul Constantin^c, Cristian Costache^d, Michał Ciemata^f, Gheorghe Ciocan^d, Cristina Clisu^d, Fabio C.L. Crespi^{b,a}, Mihai Cuciuc^c, Anukul Dhal^c, Irina Dinescu^d, Nikolay Djourellov^c, Nicoleta M. Florea^d, Ioana Gheorghe^d, Decebal Iancu^e, Violeta Iancu^c, David M. Kahl^c, Maria Kmiecik^f, Asli Kuşoğlu^{c,g}, Razvan Lica^d, Adam Maj^f, Nicolae M. Mărginean^d, Raluca Mărginean^d, Constantin Mihai^d, Radu E. Mihai^{d,h}, Bénédicte Million^a, Catalin Neacsu^d, Dragos Nichita^c, Cristina R. Nită^d, Haridas Pai^c, Alfio D. Pappalardo^c, Teodora Petruse^c, Adrian Rotaru^c, Andreea B. Serban^c, Pär-Anders Söderström^c, Cristophe O. Sotty^d, Anamaria Spătaru^c, Lucian Stan^d, Alexandru N. State^c, Alexandru Stoica^d, Dmitry A. Testov^c, Sebastian Toma^d, Tatiana Tozar^c, Andrei Turturică^d, Gabriel V. Turturică^c, Sorin Ujeniuc^d, Vlad Vasilca^c, Yi Xu^c, Fan Zhu^c

^a*INFN, sezione di Milano, Via Celoria, 16, Milano, 20133, MI, Italy*

^b*Università degli studi di Milano and INFN, sezione di Milano, Via Celoria, 16, Milano, 20133, MI, Italy*

^c*Extreme Light Infrastructure-Nuclear Physics, Str. Reactorului 30, Bucharest-Măgurele 077125, Romania*

^d*Department of Nuclear Physics, Str. Reactorului 30, Bucharest-Măgurele, 077125, Romania*

^e*Applied Nuclear Physics Department, Str. Reactorului 30, Bucharest-Măgurele, 077125, Romania*

^f*The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, ul. Radzikowskiego 152, Krakow, 31-342, Poland*

^g*Department of Physics, Faculty of Science, Istanbul University, Vezneciler/Fatih, Istanbul, 34134, Turkey*

^h*Institute of Experimental and Applied Physics, Czech Technical University in Prague, Husova 240/5, Prague, 110 00, Czech Republic*



THANK YOU FOR THE ATTENTION!