

# Double Parton Scattering and Double PDFs

Jasper Amirante  
(In Collaboration with Jonathan Gaunt)

University of Manchester

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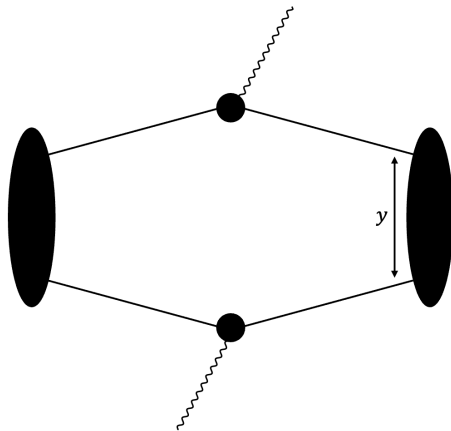


# Outline

- Double Parton Scattering (DPS) and Double Parton Distributions (DPDs)
- DPS in same-sign WW
- What same-sign WW tells us about correlations in the proton

# Double Parton Scattering

DPS: Two partons from each proton collide into two separate hard scatters.



$$\sigma_{DPS}^{(A,B)} = \int F_{ik}(x_1, x_2, \mathbf{y}) \otimes \hat{\sigma}_{ij}^A \hat{\sigma}_{kl}^B \otimes F_{jl}(x'_1, x'_2, \mathbf{y}) d^2\mathbf{y}$$

Factorisation now proven for double colour-singlet production in QCD.

(Diehl, Gaunt, Ostermeier, Plöchl, Schafer, JHEP 1601 (2016) 076, Diehl, Ostermeier, Schafer, JHEP 1203 (2012) 089, Diehl, Nagar, JHEP 1904 (2019) 124, Vladimirov, JHEP 1804 (2018) 045)

# Pocket Formula

This talk is about correlations, but for now let's ignore them

$$F^{ij}(x_1, x_2, \mathbf{y}) \rightarrow f^i(x_1) f^j(x_2) G(\mathbf{y}).$$

This gives the "DPS pocket formula":

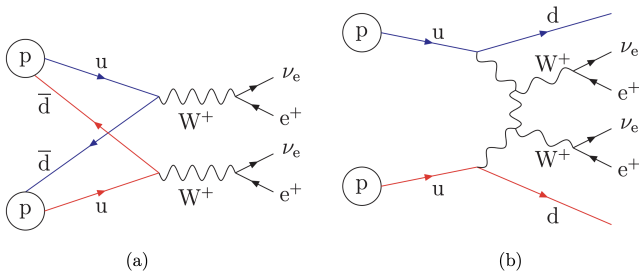
$$\sigma_D^{(A,B)} = \frac{\sigma_S^{(A)} \sigma_S^{(B)}}{\sigma_{\text{eff}}}$$

with

$$\sigma_{\text{eff}} = \frac{1}{\int G(\mathbf{y}) d^2 \mathbf{y}}$$

## DPS in Same-Sign WW

Same-sign WW in single scattering suppressed by multiple coupling constants and must involve jets to conserve charge. This allows DPS to compete.

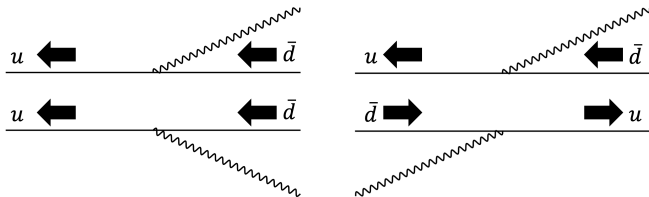


Formally DPS is power suppressed with respect to SPS

$$\sigma_{DPS} \sim \sigma_{SPS} \frac{\Lambda_{QCD}^2}{Q^2}$$

# Correlations in DPS

- A significant motivation to study DPS is to investigate correlations between partons in the proton.
- DPS offers insights into momentum, valence, spin and colour correlations.
- Same-sign WW offers a particularly nice insight into spin correlations due the W boson coupling only to LH particles/ RH antiparticles.



# Observing Correlations

Asymmetry can be used to probe correlations.

Gaunt, Kom, Kulesza, Stirling, Eur.Phys.J. C69 (2010) 53

$$\mathcal{A} = \frac{\text{Diagram 1} - \text{Diagram 2}}{\text{Diagram 3} + \text{Diagram 4}}$$

$$P \left[ \text{Diagram 1} \right] - P \left[ \text{Diagram 2} \right] = P \left[ \text{Diagram 1} \right] \left\{ P \left[ \text{Diagram 3} \right] - P \left[ \text{Diagram 4} \right] \right\} = 0$$

One must be careful with phase space cuts when using asymmetry to study correlations.

# Momentum and Valence Number Conservation

- The requirement that the momentum and valence numbers of the proton be conserved leads to correlations.
- These correlations induce positive asymmetry as they reduce the likelihood of obtaining two high- $x$  partons from one proton.
- dPDFs must obey the sum rules which are as follows:

Momentum sum rule encodes available momentum after “taking out”  $j_1$ :

$$\sum_{j_2} \int dx_2 x_2 D_{j_1 j_2}(x_1, x_2, Q) = (1 - x_1) f_{j_1}(x_1, Q)$$

Number sum rule encodes number of  $j_2$  quarks - number of  $\bar{j}_2$  quarks after “taking out”  $j_1$ :

$$\int dx_2 D_{j_1 j_2 \nu}(x_1, x_2, Q) = \left( N_{j_2 \nu} - \delta_{j_1 j_2} + \delta_{j_1 \bar{j}_2} \right) f_{j_1}(x_1, Q)$$

Gaunt, Stirling, 0910.4347

Blok, Dokshitzer, Frankfurt, Strikman, 1306.3763

Diehl, Plöchl, Schäfer, 1811.00289

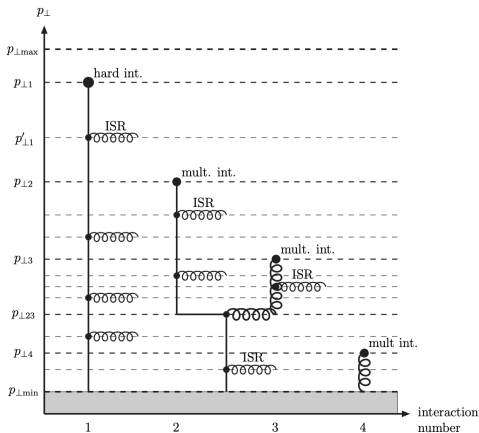


# Building dPDFs That Satisfy the Sum Rules

Unfortunately we do not have dPDFs from experimental fits, we must therefore build models which use information from single PDFs. Several different approaches:

- Pythia MPI model
- x-ordered modification to Pythia model
- GS09 dPDFs

# Pythia MPI model



- Pythia uses single PDFs for hardest interaction and evolves backwards from there, interleaving ISR and MPI
- Dynamically modifies raw sPDFs to account for number and valence effects:  $f_i^r(x, Q) \rightarrow f_i^{m \leftarrow j_1, x_1 \dots j_{n-1}, x_{n-1}}(x, Q)$

Sjostrand, Skands, JHEP 03 (2004) 053, Eur.Phys.J.C 39 (2005) 129-154

# Pythia MPI model

## Momentum “Squeezing”

$$f_i^{m \leftarrow j_1, x_1 \dots j_{n-1}, x_{n-1}}(x, Q) = \frac{1}{X} f_i^r\left(\frac{x}{X}, Q\right) \text{ where } X = 1 - \sum_{i=1}^{n-1} x_i$$

## Valence Number Subtraction

$$f_{i_v}^{m \leftarrow j_1, x_1 \dots j_{n-1}, x_{n-1}}(x, Q) = \frac{N_{i_v n}}{N_{i_v 0}} \frac{1}{X} f_{i_v}^r\left(\frac{x}{X}, Q\right)$$

## Companion Quark Addition

$$q_c(x, x_s) = C(x_s) P_{g \rightarrow q \bar{q}}\left(\frac{x_s}{x_s + x}\right) \frac{f_g(x_s + x)}{x_s + x}$$

## Sea Quark and Gluon Rescaling

Steps 2 and 3 break the fact that step 1 guarantees:

$$\sum_i \int dx x f_i^{m \leftarrow j_1, x_1 \dots j_{n-1}, x_{n-1}}(x, Q) = X$$

Sjostrand, Skands, hep-ph/0402078, hep-ph/0408302

# Pythia MPI model

The problem with this model is that it is based on hardness ordering and this is incompatible with equal scales processes such as same-sign WW. The dPDFs should be symmetric under  $x_1 \leftrightarrow x_2, j_1 \leftrightarrow j_2$ . This approach satisfies the sum rules to some extent, but can be improved

	Momentum sum rule ( $j_1 = u$ ). Should = 1.	$\bar{u}u$ number sum rule. Should = 3.
$x_1$		
$10^{-6}$	0.979	2.961
$10^{-3}$	0.980	3.351
$10^{-1}$	1.014	3.491
0.2	1.047	3.580
0.4	1.133	3.858
0.8	1.679	7.048

Naively symmetrised Pythia dPDFs

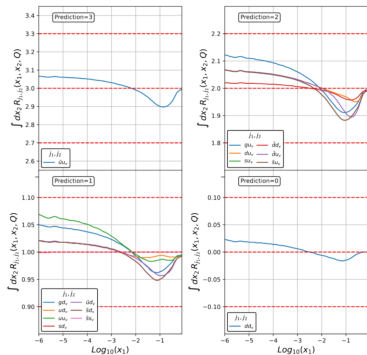
# Improved Model

An improved model has been developed which satisfies these symmetries and is therefore more suited to equal scales processes. The modifications are:

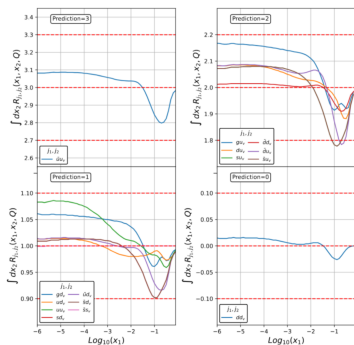
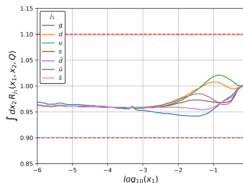
- Order scatters in  $x$  rather than  $Q$  and has smooth transitions
- Improve “companion quark mechanism” so that it is symmetric and conforms better to picture of quark and companion arising from  $g \rightarrow q\bar{q}$  splitting.

Fedkevych, Gaunt, Smith, in preparation

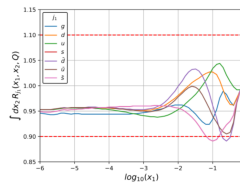
# Improved Model Sum Rules



dPDFs  
constructed at  
 $Q = M_Z$



dPDFs  
constructed at  
 $Q = 2 \text{ GeV}$



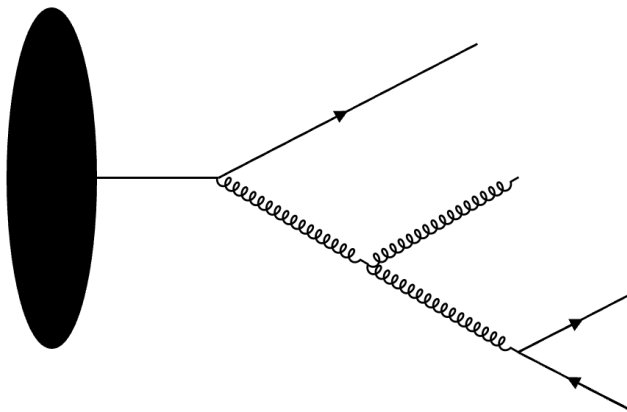
Sum Rules well satisfied.

- dPDFs created to satisfy the sum rules.
- Modelled as product of single PDFs and a phase space suppression factor plus valence and companion adjustments at low scale and evolved using dDGLAP:

$$\frac{d}{d\ln Q^2} \text{ (diagram with two horizontal lines entering a black oval)} = \text{diagram with two horizontal lines entering a black oval and one line exiting at an angle} + \text{diagram with two horizontal lines entering a black oval and one line exiting at an angle from the bottom} + \text{diagram with two horizontal lines entering a black oval and one line exiting to the right which then splits into two lines}$$

## Perturbative Splitting Contribution

Partons can originate from the perturbative splitting of a single mother parton. This correlation again induces positive asymmetry as it adds a contribution which strongly prefers the high- $x$  low- $x$  configuration.





# Spin Correlations

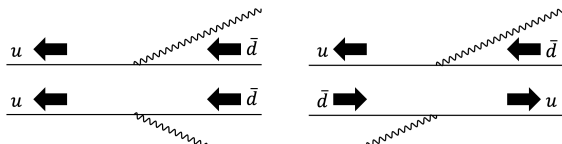
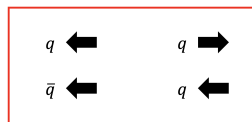
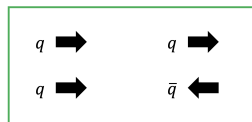
- Due to the coupling of the W boson only certain spin configurations will contribute to the cross section.
- Spin correlations encoded in polarised dPDFs which must satisfy the positivity bound:

$$f_{a,b} = \uparrow\uparrow + \downarrow\downarrow + \downarrow\uparrow + \uparrow\downarrow$$

$$f_{\Delta a, \Delta b} = \uparrow\uparrow + \downarrow\downarrow - \downarrow\uparrow - \uparrow\downarrow$$

$$|f_{\Delta a, \Delta b}| \leq f_{a,b}$$

Mekhfi, Phys. Rev. D32 (1985) 2380  
 Diehl, Ostermeier and Schafer (JHEP 1203 (2012))  
 Manohar, Waalewijn, Phys.Rev. D85 (2012) 114009



# Spin Correlations

Build models to understand the maximum possible effects of spin correlations. Construct at  $Q = 1\text{GeV}$  and evolve to  $W$  mass. 2 models proposed by Cotogno, Kasemets, Myska: *Phys.Rev.D*100 (2019) 1, 011503, *JHEP* 10 (2020) 214

## Positive Polarisation

Saturate the positivity bound with positive polarised dPDFs:

$$f_{\Delta a \Delta b}(x_1, x_2, \mathbf{y}; Q_0) = f_{ab}(x_1, x_2, \mathbf{y}; Q_0) = f(x_1; Q_0) f(x_2; Q_0) G(\mathbf{y})$$

## Mixed Polarisation

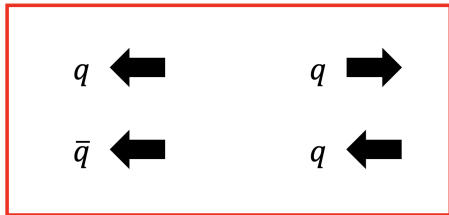
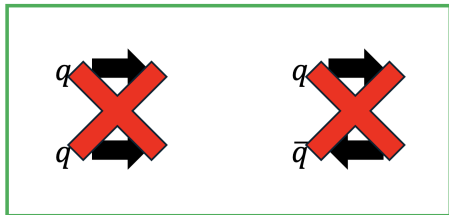
Saturate the positivity bound with flavour dependent polarised dPDFs:

$$f_{\Delta a \Delta b}(x_1, x_2, \mathbf{y}; Q_0) = (-1)^n f_{ab}(x_1, x_2, \mathbf{y}; Q_0)$$

Where  $n = 1$  if  $ab = qq, \bar{q}\bar{q}$  and  $n = 2$  otherwise.

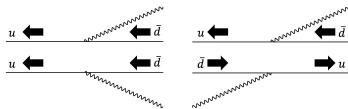
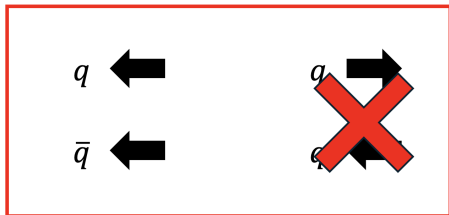
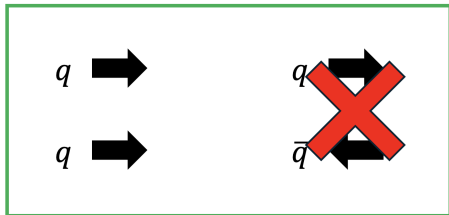
A third model, negative polarisation is useful to compare to mixpol and is just -pospol.

# Mixed Polarisation



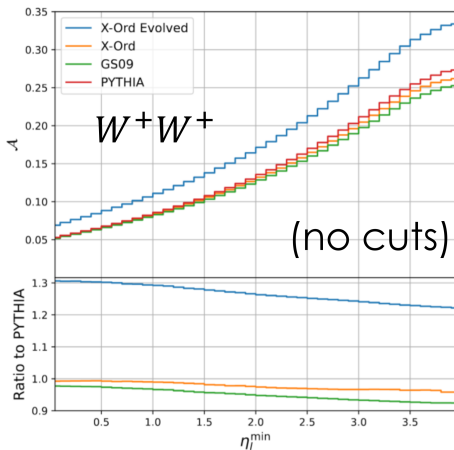
- Both configurations that contribute to  $W^+W^+$  are killed.
- Cross section fully generated by evolution
- Evolution should primarily generate positive asymmetry since two high- $x$  is unlikely.

# Positive Polarisation



- Only align contribution allowed to contribute at low scale.
- If we assume generally  $f_q > f_{\bar{q}}$ , two quarks will push the  $W$ s into the same hemisphere
- Can this effect withstand evolution?

# Effects of Moment and Valence Number Conservation

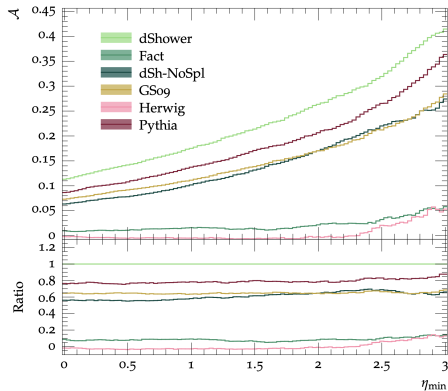


- Positive asymmetry, large increase with minimum rapidity cuts
- Very small differences between Pythia and x-ordered since ssWW doesn't probe companion quark
- X-order evolved and GS09 show large differences

Fedkevych, Gaunt, Smith, in preparation

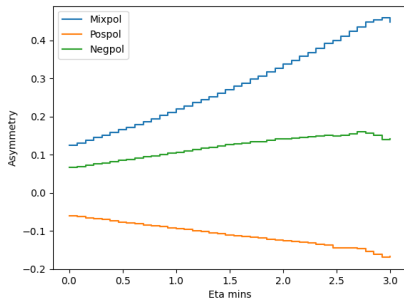
# Effects of Splitting Contribution

- dShower is an MC that properly accounts for  $y$  dependence in splitting contribution
- These results are for  $n_f = 3$ , effects are partially washed out for  $n_f = 5$
- GS09 captures about half the effect compared to dShower since it does not properly account for the  $y$  dependence



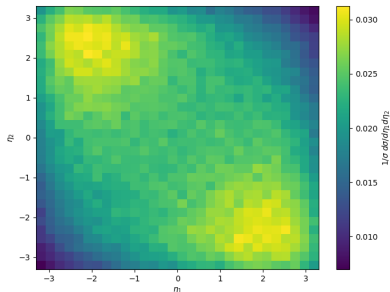
Cabouat, Gaunt, Ostrolenk, JHEP 11 (2019) 061

# Effects of Spin Correlations



- Mixpol and pospol generate positive and negative asymmetry as expected
- At  $|\eta| > 0$  approximately a 5-10% effect
- Mixpol especially good at generating high asymmetry when compared to negpol

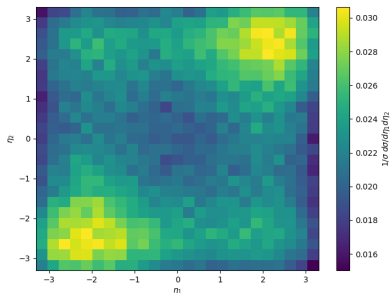
# Understanding Mixpol



- Mixpol behaves as expected producing strong bias towards opposite hemispheres
- Lower cross section since it must all be generated by evolution
- Very few events in same hemisphere corners, especially at high rapidities



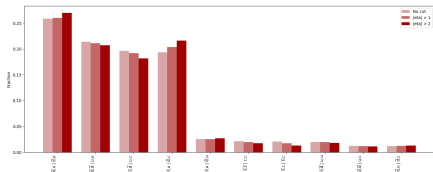
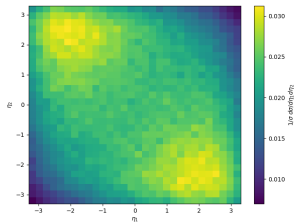
# Understanding Pospol



- Strong bias towards same hemisphere events
- Evolution still drives some opposite hemisphere events
- Could be problematic if it compensates the positive asymmetry from other sources

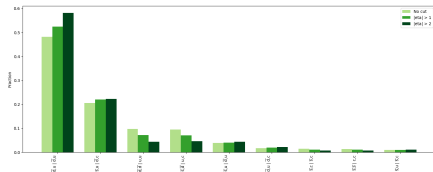
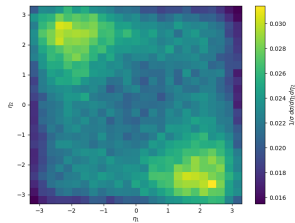
# Why is Mixpol So Good?

## Mixpol



Evolution gives strong preference towards large- $x$  small- $x$ .

## Negpol

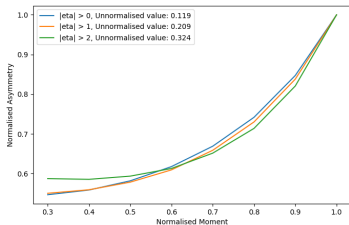
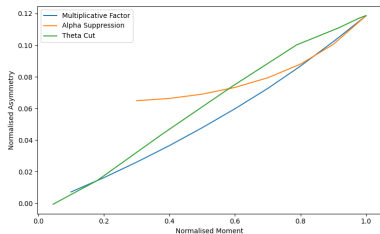


Valenceness of  $\bar{d}$  pushes some events into same hemisphere

# Placing Constraints on Spin Contributions

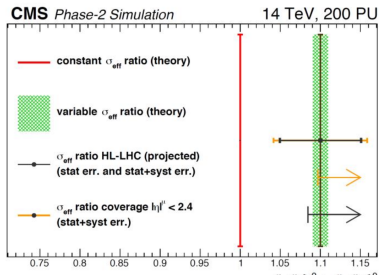
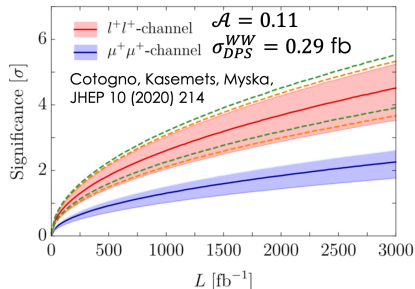
- Lattice-QCD should be able to constrain the moment of the polarised dPDFs
- We should be able to place a fairly strong constraint on the size of the spin contribution given this moment
- Different models of constraining asymmetry lead to similar predictions

$$M_{ij} = \int_{x_{\min}}^1 x_1 x_2 F_{ij}(x_1, x_2, Q^2) dx_1 dx_2.$$



# Prospects of Determining Correlations

- Should be able to distinguish correlations from the DPS pocket formula if the total effect of the correlations is  $\mathcal{A} > 0.05$
- Spin effects can have negative contribution to  $\mathcal{A}$  and potentially comparable to the effect from other correlations. Lattice info can help constrain spin effects and indicate if measurable  $\mathcal{A} \neq 0$  remains



CMS-TDR-016

# Summary

- DPS is an important contribution to same-sign WW production, and studying it can reveal information on correlations between partons in the proton.
- The lepton pseudorapidity asymmetry  $\mathcal{A}$  is an observable that probes such correlations.  $\mathcal{A} \neq 0$  implies correlations!
- The main correlations that affect  $\mathcal{A}$  are:
  - ▶ Momentum and number sum rule correlations
  - ▶ Correlations induced by perturbative splittings
  - ▶ Spin Correlations
- Each yields a few per cent impact on  $\mathcal{A}$ . First two effects give  $\mathcal{A} > 0$ , whilst spin correlations can give a positive or a negative contribution to  $\mathcal{A}$ .