

# The TAUOLA generator for $\tau$ lepton decays

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- **(1)** The  $\tau$  lepton decays: fascinating laboratory for intermediate energy QCD; that may explain, why I am a bit biased against New Physics in  $\tau$  decays, but ...
- **(2)** I will address also TAUOLA in context of  $\tau$  as precision tool to measure properties of SM physics (and beyond). This is mostly because of  $\tau$  lepton production processes.
- **(3)** These two regimes separate perfectly, because  $\tau$  leptons are extremely narrow resonances and production processes are relatively free of strong interactions.
- **(4)** How to optimize work of inhomogeneous community. From model builders to people managing large experimental data files. From F77 to C++ and Python.
- **(5)** My concern is also on how to handle different component of systematic errors: experiment, theor. Choice of quantities for measurements and tests.

## Target points: what people may need

- (1) Simulate detector response
- (2) Provide distributions of  $\tau$  decay products and of the  $\tau$  itself: starting from lagrangian of Old and New physics
- (3) Environment to study prototypes for matrix elements and **prototypes for  $\tau$  decay observables.**  
Technical detail: narrow width limit for intermediate resonances is often needed.
- (4) For studies where  $\tau$  leptons are used to constrain else, like Higgs CP or B physics.
- (5) **New challenges: multidimensional distributions? ML? Experimental systematic errors for that?**

- **First some theory and software organization**

## Formalism for $\tau^+\tau^-$

- Because narrow  $\tau$  width approximation can be obviously used for phase space, cross-section for the process  $f\bar{f} \rightarrow \tau^+\tau^-Y; \tau^+ \rightarrow X^+\bar{\nu}; \tau^- \rightarrow \nu\nu$  reads:

$$d\sigma = \sum_{spin} |\mathcal{M}|^2 d\Omega = \sum_{spin} |\mathcal{M}|^2 d\Omega_{prod} d\Omega_{\tau^+} d\Omega_{\tau^-}$$

- This formalism is fine, but, e.g. for 20  $\tau$  decay channels we would have 400 distinct processes. Also picture of production and decay are mixed.
- Below only  $\tau$  spin indices are explicitly written:

$$\mathcal{M} = \sum_{\lambda_1 \lambda_2 = 1}^2 \mathcal{M}_{\lambda_1 \lambda_2}^{prod} \mathcal{M}_{\lambda_1}^{\tau^+} \mathcal{M}_{\lambda_2}^{\tau^-}$$

- Cross section can be re-written into **core formula of spin algorithms**

$$d\sigma = \left( \sum_{spin} |\mathcal{M}^{prod}|^2 \right) \left( \sum_{spin} |\mathcal{M}^{\tau^+}|^2 \right) \left( \sum_{spin} |\mathcal{M}^{\tau^-}|^2 \right) wt d\Omega_{prod} d\Omega_{\tau^+} d\Omega_{\tau^-}$$

- where

$$wt = \left( \sum_{i,j=0,3} R_{ij} h^i h^j \right)$$

$$R_{00} = 1, \quad \langle wt \rangle = 1, \quad 0 \leq wt \leq 4.$$

$R_{ij}$  can be calculated from  $\mathcal{M}_{\lambda_1 \lambda_2}$  by contraction with Pauli  $\sigma^i$  matrices and similarly  $h^i, h^j$  respectively from  $\mathcal{M}^{\tau^+}$  and  $\mathcal{M}^{\tau^-}$ .

- Bell inequalities tell us that it is impossible to re-write  $wt$  in the following form

$$wt \neq \left( \sum_{i,j=0,3} R_i^A h^i \right) \left( \sum_{i,j=0,3} R_j^B h^j \right)$$

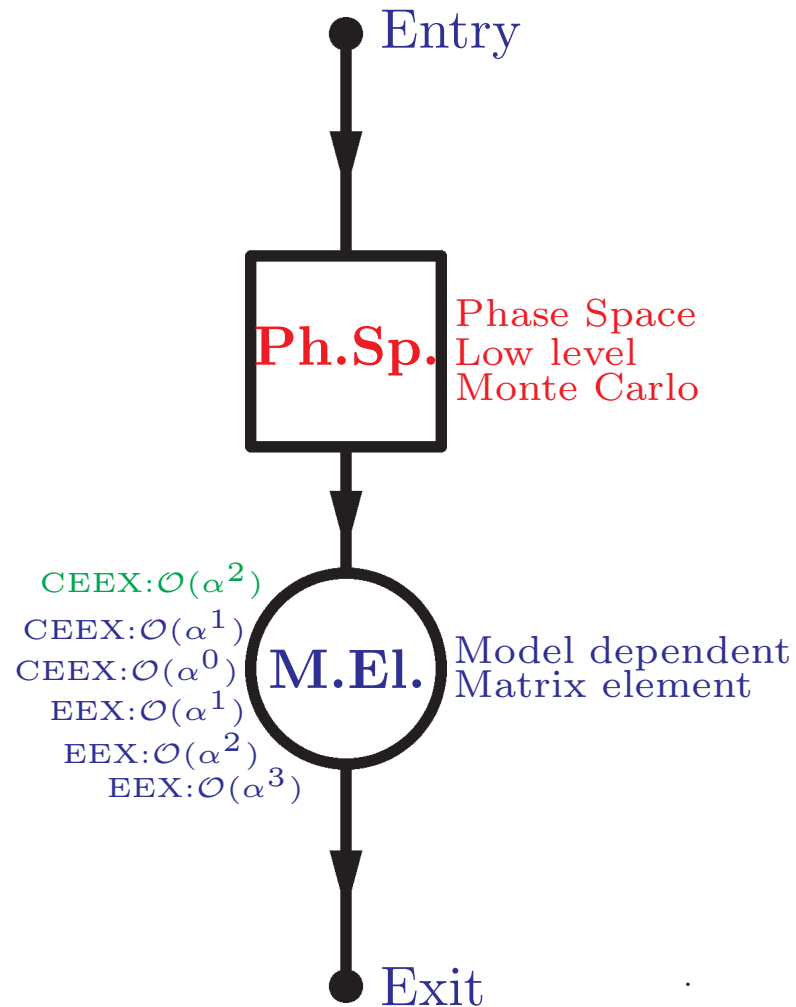
that means it is impossible to generate first  $\tau^+$  and  $\tau^-$  first in some given 'quantum state' and later perform separately decays of  $\tau^+$  and  $\tau^-$

- It can be done only if approximations are used !!!
- May be reasonable in e.g. ultrarelativistic regime, but nonetheless approximation.

## TAUOLA universal interface

- To run, generator for tau decays must be combined with production.
- In cases of packages for  $e^+e^-$  colliders, such as ours KORALB, KORALZ, KKMC, they provide environment for TAUOLA use.
- often information from event stored in production files can be used.
- I will adress only fraction of technicalities, there is a lot to that!
- In KKMC quantization frames, boosts to/from Lab to  $\tau$  rest-frames are carefully prepared **HOWEVER:**  $\tau$  lepton(s) spin states can be calculated from kinematical configurations of hard processes then it is up to user to control.
- Like in f77 version of EvtGen-interface to TAUOLA and PHOTOS. Also in our TAUOLA universal interface. Information from event record.
- It can work on stored production events. flexibility but require user attention.
- Algrithmic-wise, the same sophistication level, like solution in KKMC (e.g. quantum entanglement included).

Textbook principle “matrix element  $\times$  full phase space” useful



- Phase-space Monte Carlo module producing “raw events”.
- Library of models for provides input for “model weight”
- **Useful for any application, not only  $\tau$  production/decay.**
- Ratios of matrix elements squared define probability that event could be of model B if generated with mode A.
- Convenient for Machine Learning too.
- No compromises on precision are required.

## Formalism for semileptonic decays at 0.2% precision level

- Matrix element used in TAUOLA for semileptonic decay of  $\tau$  with  $P$  momentum and spin  $s$

$$\tau(P, s) \rightarrow \nu_\tau(N) X$$

$$\mathcal{M} = \frac{G}{\sqrt{2}} \bar{u}(N) \gamma^\mu (v + a\gamma_5) u(P) J_\mu$$

- $J_\mu$  – the current, depends on the momenta of all hadrons ( $h_\mu = H_\mu/H_t$ )

$$|\mathcal{M}|^2 = G^2 \frac{v^2 + a^2}{2} (\omega + H_\mu s^\mu)$$

$$\omega = P^\mu (\Pi_\mu - \gamma_{va} \Pi_\mu^5)$$

$$H_\mu = \frac{1}{M} (M^2 \delta_\mu^\nu - P_\mu P^\nu) (\Pi_\nu^5 - \gamma_{va} \Pi_\nu)$$

$$\Pi_\mu = 2[(J^* \cdot N) J_\mu + (J \cdot N) J_\mu^* - (J^* \cdot J) N_\mu]$$

$$\Pi^{5\mu} = 2 \operatorname{Im} \epsilon^{\mu\nu\rho\sigma} J_\nu^* J_\rho N_\sigma$$

$$\gamma_{va} = -\frac{2va}{v^2 + a^2}$$

$$\hat{\omega} = 2 \frac{v^2 - a^2}{v^2 + a^2} m_\nu M (J^* \cdot J)$$

$$\hat{H}^\mu = -2 \frac{v^2 - a^2}{v^2 + a^2} m_\nu \operatorname{Im} \epsilon^{\mu\nu\rho\sigma} J_\nu^* J_\rho P_\sigma$$

- Hadronic currents have to fulfil Lorentz invariance.
- In  $\tau^\pm \rightarrow \rho^\pm \rightarrow \pi^\pm \pi^0 \nu$  channel fits are technically straightforward: single distribution to be fitted with single real function to fit:

$$J^\mu = (p_{\pi^\pm} - p_{\pi^0})^\mu F_V(Q^2) + (p_{\pi^\pm} + p_{\pi^0})^\mu F_S(Q^2) \quad (F_S \simeq 0).$$

- For 3-scalar channels: 4 complex functions each of 3 variables to fit. Role of theoretical assumptions (oversimplifications?) is essential. Agreement on 1-dim distribution is just a consistency check.  $J^\mu = J_1^\mu + J_2^\mu + J_4^\mu + J_5^\mu$ . **CP breaking contribution installed so far in TAUOLA in some decays with  $K_L, K_S$  only.** Directly or through event weights.

- **No go for model independent measurements? Not necessarily.** Use of all dimensions for data distributions: invariant masses  $Q^2, s_1, s_2$  as arguments of form-factors. Angular asymmetries help to separate currents: scalar

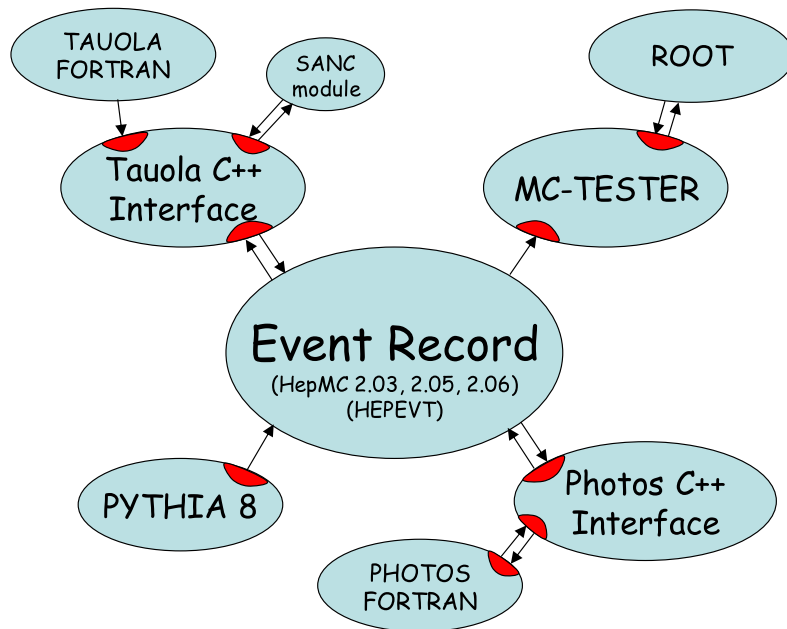
$$J_4^\mu \sim Q^\mu = (p_1 + p_2 + p_3)^\mu, \text{ vector } J_1^\mu \sim (p_1 - p_3)^\mu |_{\perp Q} \text{ and } J_2^\mu \sim (p_2 - p_3)^\mu |_{\perp Q} \text{ and finally } \underline{\text{pseudovector}} J_5^\mu \sim \epsilon(\mu, p_1, p_2, p_3).$$

- Model independent methods, if: (i) enough data, (ii) absolute precision, (iii) no background, (iv) full detector coverage of decay phase-space helpful. We need that for orthogonality of fitted functions. ML techniques instead?
- It is a challenge but worth a try. **I am ready to talk it over, any time any place.**



# An example when $\tau$ decays are modified.

Solutions: rigorous exact phase space  $\times$  approximate  $|M|^2$ .



## Parts:

- hard process: (Born, weak, new physics),
- parton shower,
- $\tau$  decays
- QED bremsstrahlung
- High precision achieved
- Detector studies: acceptance, resolution lepton with or without photon.

## Such organization requires:

- Good control of factorization (theory)
- Good understanding of tools on user side.

## Note:

- (1)  $\tau$  leptons are long lived
- (2) QED low energy limit - solvable, analyticity helps semi-factorization too.

# An example when $\tau$ decays are modified.

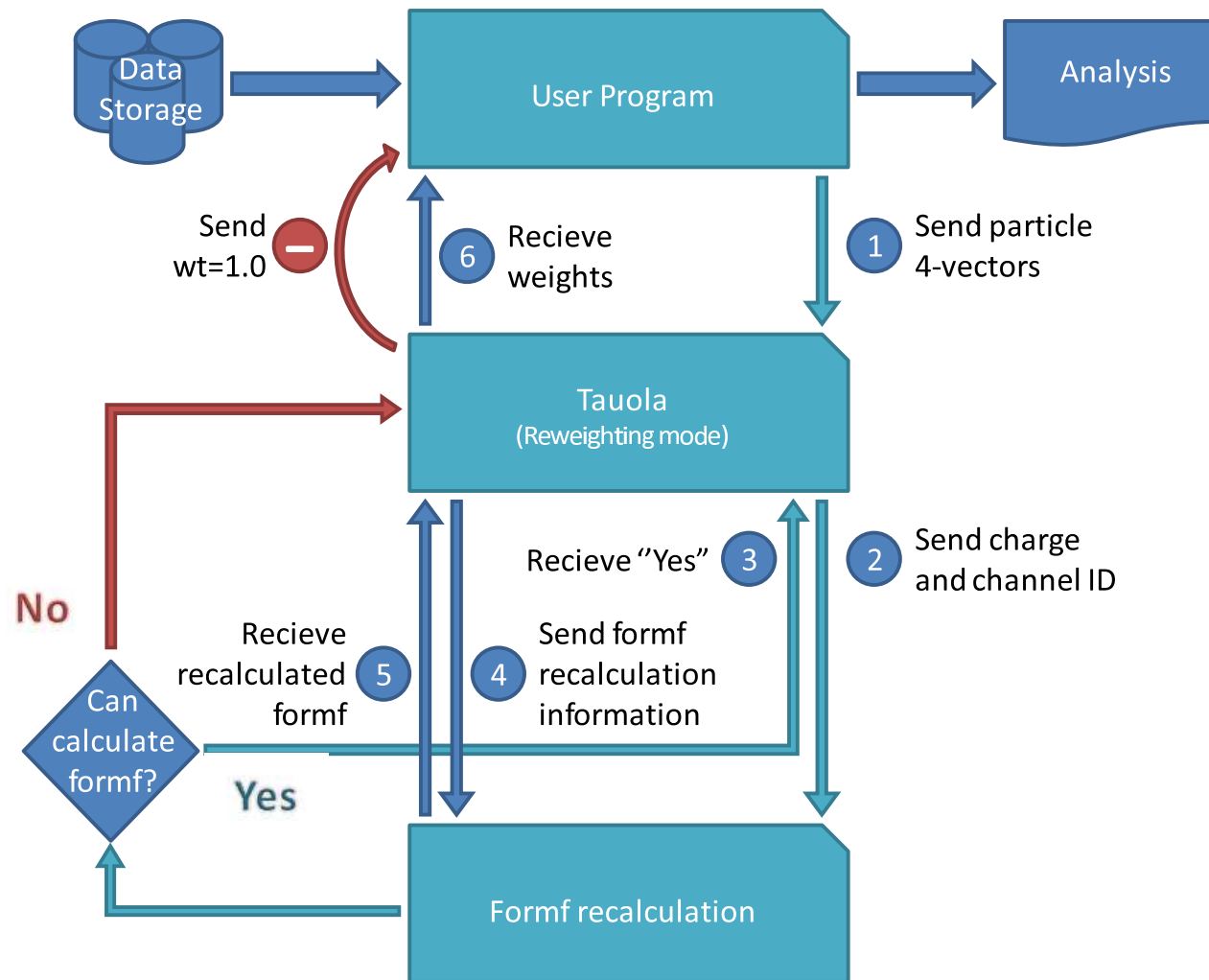


Figure 1: Flow chart for communication when already stored events are modified with the weights. Useful at LHC and at low energy applications as well.

1. **The following contexts: fit strategy, experimental, theoretical syst. errors., cooperation between sub-communities.**
2. Does it open a path to understand intermediate energy QCD?  
**I am ready to talk it over, any time any place.**
3. “Recently” I have introduced into TAUOLA some changes.
4. Quality stamps from the side of theory, experiment, technical precision.
5. TAUOLA of new hadronic currents, 200+ decay channels, which can be manipulated by user with c++ coded currents, ME and with any decay products:  
[Comput.Phys.Commun. 232 \(2018\) 220](#)
6. What should be included (acceptable in collaborations software) in standard initialization(s).
7. Constraints of software organization in Belle forced a step back: no user provided currents, migration to c++ is not completed.

- New hadronic currents (more than 88 % of hadronic  $\tau$  decay width) version installed with the 0.05 % technical tag:
- **But** physics precision was definitely **NOT** as good as 0.05 %.
- Over two years we worked on preparing confrontation env. with the data keeping precision in mind, but ...
- Despite efforts , we are left as far from the complete solution as many years ago.
- We have investigated technical aspects for fitting using weights.  
Seemed interesting when experimental cuts are present, multidimensional distributions are used and no semi-analytical results can be easily prepared for fits.
- We have returned to the semi-analytical 1-dim distributions for fits. Because multi-dimensional data were not public, no systematic ambiguities evaluated.
- We had encountered difficulties with fits as well.
- **WARNING:** results of distinct fits gave at 10 % the same 1-dim distributions, but in more than 10 % of the contributions to total widths differences were at a level of factor 2.
- I am ready to talk it over, any time any place.
- challenge for intermediate energy QCD?

To progress in case of  $\tau \rightarrow 3\pi\nu_\tau$  we had to:

- Modify the model (contribution of  $\sigma$  or CP sensitive terms)
- Work simultaneously with fits using weights. Note difficulties with stability if strong fitted parameters correlations. Template method I have learned in Orsay (at ALEPH time) requires confirmation that model parameters dependencies are not correlated. Necessity to linearize dependencies because of CPU-time constraints in case when model was not giving perfect predictions complicated things further.
- Finally we relied on fitting semi-analytical formulas:
  - We had to assure that derivatives of results are continuous.
  - We had to speed up calculations using different methods of pre-tabulation/interpolation of results for Q-dependent  $a_1$  width (unitarity constraint).
  - Only 1-dimensional invariant mass distributions were available.

- That was not enough constraining.
- What is the best, **and still safe**, input from experimental side  
Multidimensionality? Ambiguities controlled?
- I am ready to talk it over, any time any place.

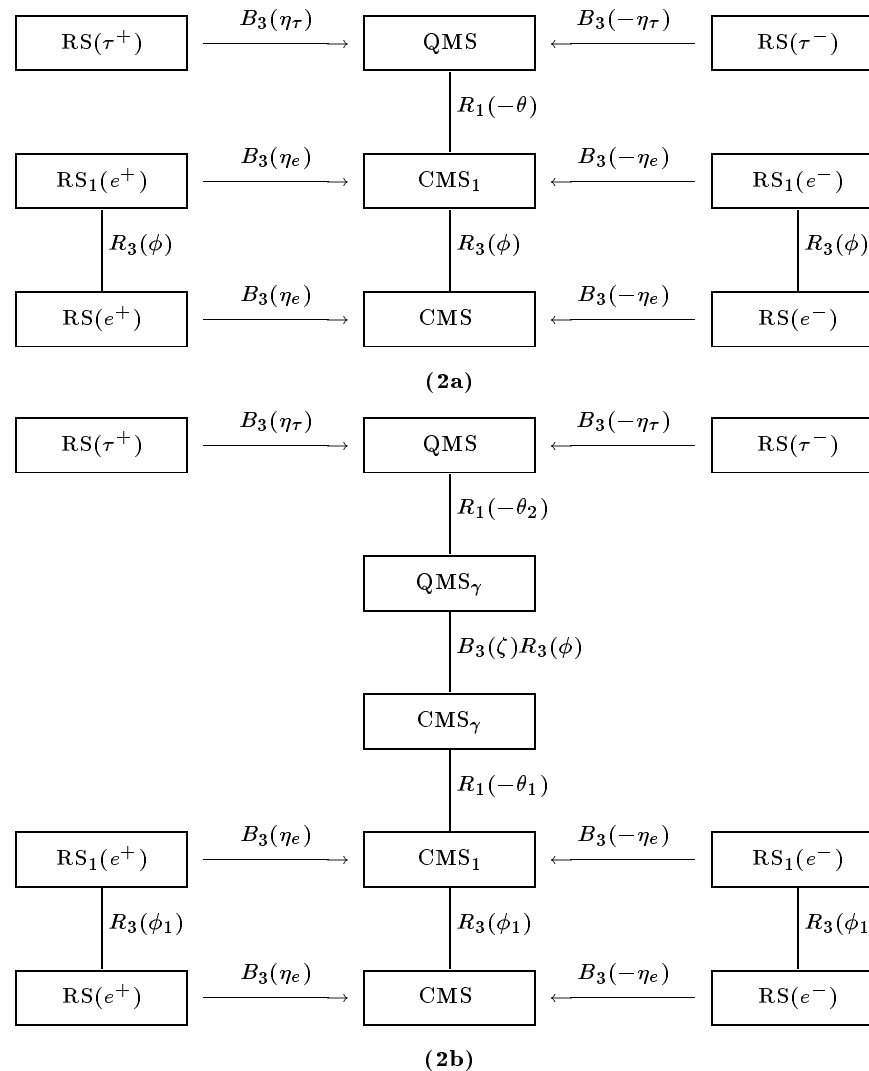


- Already for 3-scalar final states theoretical predictions and experimental data: distributions over 8-dimensional space. We fit 1- ( 2-) dim. histos. Result depend on model assumptions. Models inspired with results ... **Fitting setup → biases.**
- Our pattern recognition algorithms are far less elaborate than human eye/brain.
- How to facilitate dialog, role of MC.

- Biases in art, Giuseppe Arcimboldo (1572 - 1593).

Ref. frames for spin; production, decay. Geometry of QED amplitudes optimized

Figure 2





- *Does it make sense to have so many frames.*
- *Automatizations make it obsolete.*
- *but it is helpful for intuition build and for construction of optimal variables*
- *Who needs them at a time of ML?*
- *But then, what about ambiguities?*

## *Phenomenology Of Mixed Parity: also from M.E.*

- Higgs boson Yukawa coupling expressed with the help of the scalar–pseudo-scalar mixing angle  $\phi$

$$\bar{\tau} N (\cos \phi + i \sin \phi \gamma_5) \tau$$

- *Decay probability for the mixed scalar–pseudo-scalar case*

$$\Gamma(h_{mix} \rightarrow \tau^+ \tau^-) \sim 1 - s_{\parallel}^{\tau^+} s_{\parallel}^{\tau^-} + s_{\perp}^{\tau^+} R(2\phi) s_{\perp}^{\tau^-}$$

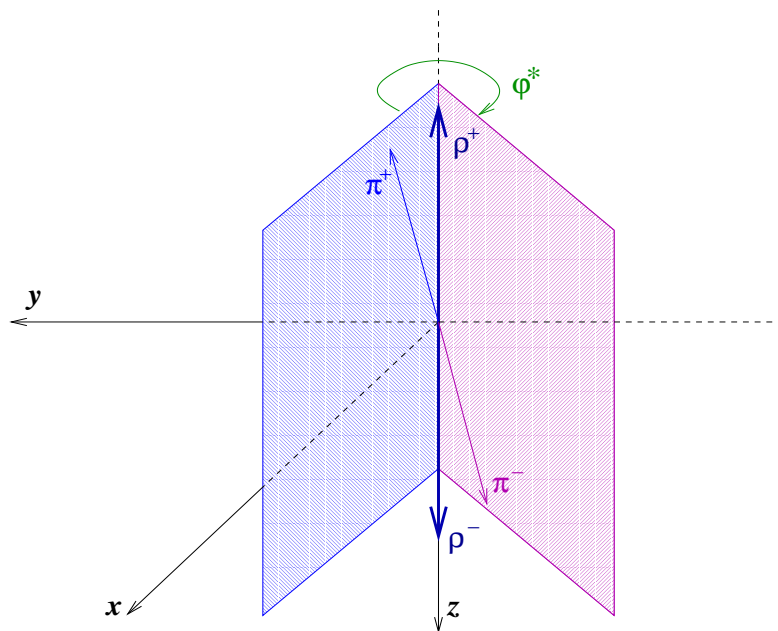
- *$R(2\phi)$  – operator for the rotation by angle  $2\phi$  around the  $\parallel$  direction.*

$$R_{11} = R_{22} = \cos 2\phi \quad R_{12} = -R_{21} = \sin 2\phi$$

- *Pure scalar case is reproduced for  $\phi = 0$ .*
- *For  $\phi = \pi/2$  we reproduce the pure pseudo-scalar case.*

### Optimal Observable Mixed Scalar–Pseudoscalar Case

- For mixing angle  $\phi$ , transverse component of  $\tau^+$  spin polarization vector is correlated with the one of  $\tau^-$  rotated by angle  $2\phi$ .
- Acoplanarity  $0 < \varphi^* < 2\pi$  is of physical interest, not just  $\arccos \mathbf{n}_- \cdot \mathbf{n}_+$ .
- Distinguish between the two cases  $0 < \varphi^* < \pi$  and  $2\pi - \varphi^*$
- If no separation made the parity effect would wash itself out.



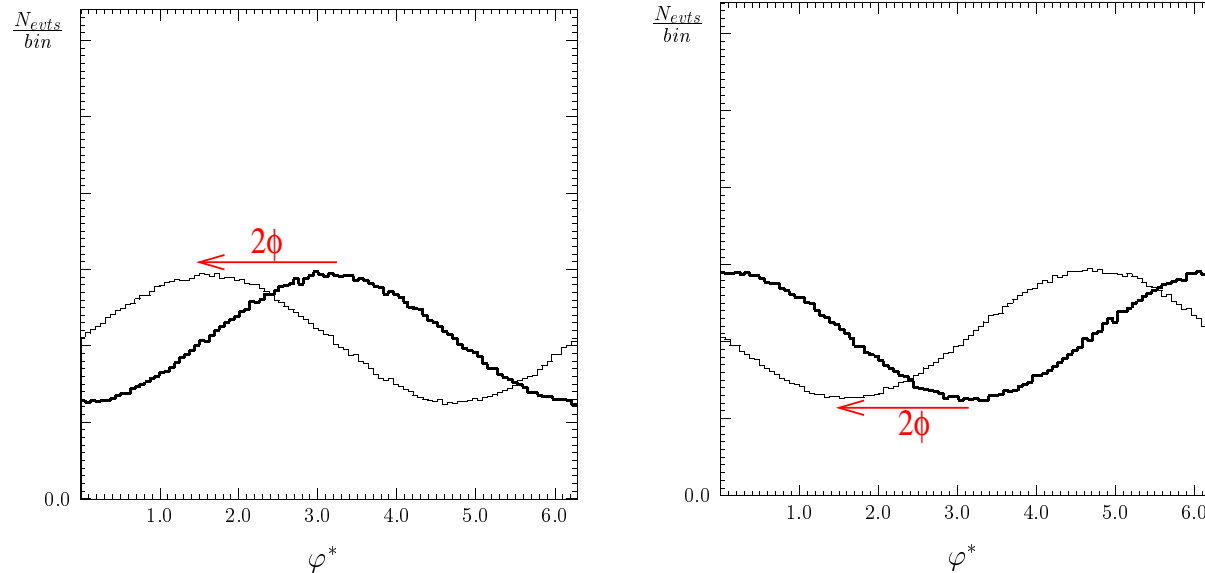
$$\text{Normal to planes: } \mathbf{n}_{\pm} = \mathbf{p}_{\pi^{\pm}} \times \mathbf{p}_{\pi^0}$$

$$\text{Find the sign of } \mathbf{p}_{\pi^-} \cdot \mathbf{n}_+$$

$$\text{Negative } 0 < \varphi^* < \pi$$

$$\text{Otherwise } 2\pi - \varphi^*$$

# From visible products in $H \rightarrow \tau^+ \tau^- \rightarrow \pi^+ \pi^0 \pi^- \pi^0$ 20

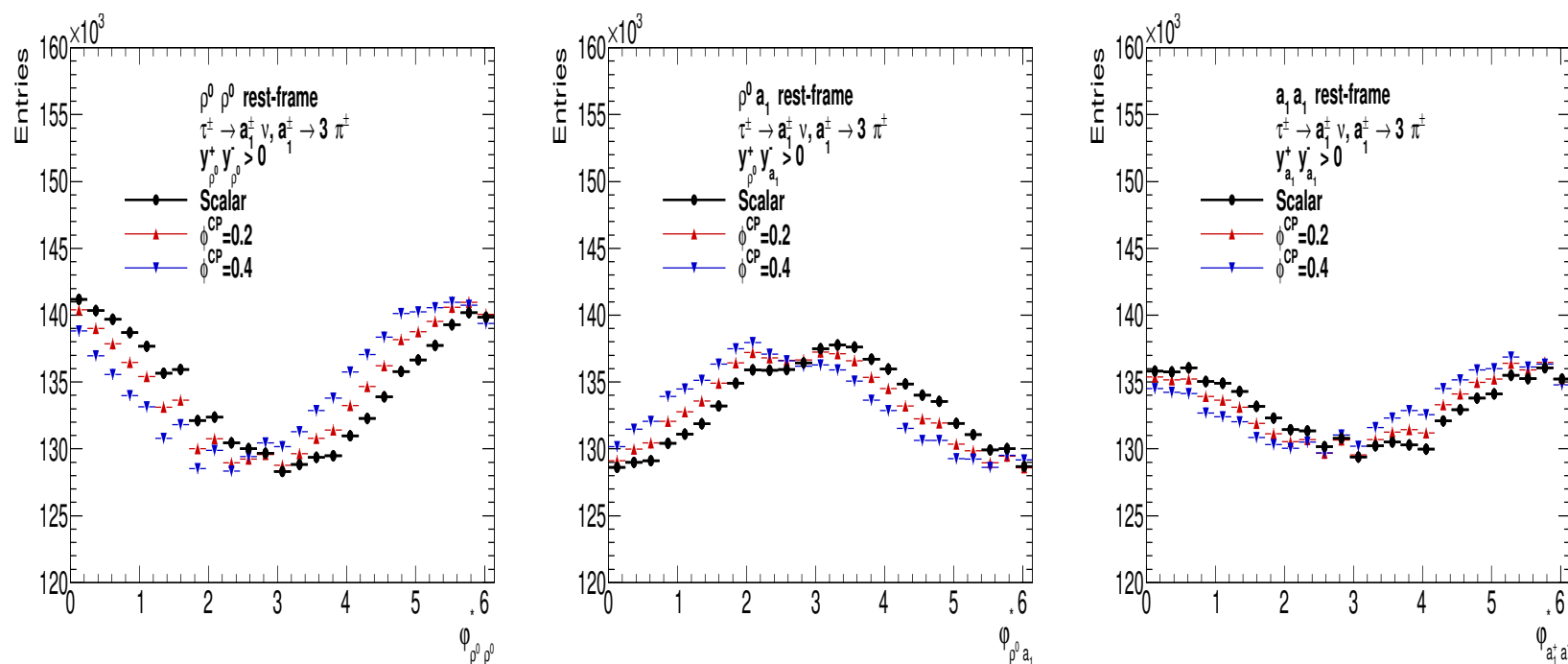


- Only events where the signs of  $y_1$  and  $y_2$  are the same whether calculated using the method without or with the help of the  $\tau$  impact parameter.
- The thick line corresponds to a scalar Higgs boson, the thin line to a mixed one.
- This Tesla-time observable , *Phys.Lett. B* 543 (2002) 227, *B*579 (2004) 157, survived into LHC times.

Old time experience with frames was helpful

# LHC, enrich small samples with $\tau \rightarrow 3\pi\nu$

Acoplanarity angles of oriented half decay planes:  $\varphi_{\rho^0\rho^0}^*$  (left),  $\varphi_{a_1\rho^0}^*$  (middle) and  $\varphi_{a_1a_1}^*$  (right), for events grouped by the sign of  $y_{\rho^0}^+ y_{\rho^0}^-$ ,  $y_{a_1}^+ y_{\rho^0}^-$  and  $y_{a_1}^+ y_{a_1}^-$  respectively. Three CP mixing angles  $\phi^{CP} = 0.0$  (scalar), 0.2 and 0.4. Note scale, effect on individual plot is so much smaller now. But up to **16 plots like that** have to be measured, correlations understood. Physics model depends on 1 parameter only  $\phi^{CP}$  mixing scalar pseudo-scalar angle, which brings linear shift. **I remained frustrated for 15 years, how to digest...**



- The purpose of my talk was to push some ideas forward and what is needed for that.
- It was not disciplined talk. Sorry for that.
- (i) Experimental systematic errors (ii) Theoretical systematic errors
- What are the constraints on organization of Monte Carlo and fitting environments?
- We have collected some experience on requirements for building fitting environments.
- Systematic ambiguities, in case of fits to multi-dimensional representation of data, is a challenge.
- ML bring benefits but a price too ...
- Question of manpower and training as well as motivation of involved people.
- $\tau$  leptons as tools for other high and medium energy physics points.
- Easy to ignore: e.g. narrow width approximation availability in MC is useful for: (i) tests (ii) model development and tuning (iii) observable construction interpretation (iv) intuition buildup.

- The purpose of my talk was to push some ideas forward and what is needed for that.
- It was not disciplined talk. Sorry for that.
- My aim was to underline directions, which in my opinion, are worth to follow.
- I have sketched also what are the main aspects of the activities till now.
- I hope it may be useful for somebody here.