Particle Physics for Specialists

Forward Physics

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Introduction



David d'Enterria, arXiv:0708.0551

Physics of low x

Diffraction

Photon-induced processes

Gluons at low x



Saturation



Underlying event

- · Complex structure of hadron-hadron interaction
- Underlying event activity in addition to the hard interaction:
 - initial state radiation
 - final state radiation
 - multiple parton interaction
 - colour reconnections with beam remnants
- Non-perturbative effects
- No clear soft/hard separation
- Phenomenological model in MC generators
- A need for tuning to experimental data



Figure from [arXiv:1411.4085]

Principle of the measurements

- Regions in ϕ defined w.r.t. the direction of the hard object
- Transverse region sensitive to UE
- Two transverse regions \rightarrow trans-min and trans-max (distinguished on the event-by-event basis according to $\sum p_T$)
- UE observables:
 - $\cdot \,\, N_{\rm ch}/\delta\eta\,\delta\phi$
 - $\sum p_T / \delta \eta \, \delta \phi$
 - Mean p_T



Results from inclusive pp interactions





- $\Delta \phi$ distribution initially not well described by MC
- Flattening of p_T dependence: entering dense region in proton



FIG. 6. Schematic illustration of the expected dependence of the transverse multiplicity, $N(p_T)$, on the p_T of the trigger.

arXiv:1009.2559 [hep-ph]

Physics of low x

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Diffraction in particle physics



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Physics of elastic scattering



Scattering angle vs impact parameters



Figure 1: Schematic (qualitative) view of the "break", followed by the diffraction minimum ("dip"), shown both as function in t and its Fourier transform (impact parameter representation), in b. While the "break" reflects the presence of the pion "atmosphere" (clouding) around the nucleon at peripheral values of b, the dip results from absorption corrections, suppressing the impact parameter amplitude at small b.

István Szanyi, arXiv:1711.04743

Proton hollowness



FIG. 1: Normalised inelasticity density, $G_{\rm in}$, for LHC and ISR energies as a function of the impact parameter. Sub-pannel: fits to $d\sigma_{\rm el}/dt$ data.

S matrix and the Optical Theorem $\sum P(i \to n) = 1 = \sum |\langle n|S|i \rangle|^2 = \sum \langle i|S^{\dagger}|n \rangle \langle n|S|i \rangle = \langle i|S^{\dagger}S|i \rangle = 1$ true for any $|i\rangle$, so $S^{\dagger}S = I$. Introduce trans matrix T: S = I + iT $(I - iT^{\dagger})(I + iT) = I$ $i(T^{\dagger} - T) = T^{\dagger}T$ $i\langle f|T^{\dagger} - T|i\rangle = \sum \langle f|T^{\dagger}|n\rangle \langle n|T|i\rangle$ $2 \operatorname{Im} T(i \to f) = \sum \langle n | T^* | f \rangle \langle n | T | i \rangle$ forward elastic scatt. \rightarrow Optical theorem put f = i, $2 \text{ Im}T_{\rm el}(t=0) = \sum |T(i \to n)|^2 = \sigma_{\rm tot}$

from prof. Alan Marin¹⁴

Total cross section



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High- β optics



ALFA Detectors



Diffractive processes



Modeling soft diffractive dissociation



Regge theory Triple pomeron vertex



$$d\sigma_{\rm diff}/d^2b = \langle T^2 \rangle$$

Good-Walker

 Ψ_k – mass eigenstates Φ_n – diffractive eigenstates

$$\Psi_{k} = \sum c_{kn} \Phi_{n}$$
$$\langle \Psi_{1} | T | \Psi_{1} \rangle = \langle T \rangle$$
$$d\sigma_{el} / d^{2}b = \langle T \rangle^{2}$$
$$d\sigma_{diss} / d^{2}b = \langle T^{2} \rangle - \langle T \rangle^{2}$$

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Hard diffraction

Resolved pomeron

- Ingelman-Schlein model
- pomeron has partonic structure

Soft colour interactions

- QCD-inspired model
- additional gluon exchanges screen the color flow





HERA: Diffractive PDFs





- QCD fits
- dominated by gluons

Factorisation breaking





- Hard diffractive events rarer than naive extrapolations from HERA
- Suppression factor: gap survival probability
- Origin: additional interactions
- Confirmed in many processes

Physics of low x

Diffraction

Photon-induced processes

Consider a charged nucleus at rest. The associated electromagnetic field can be represented by:



As a charged nucleus moves with nearly the speed of light, the electromagnetic field becomes transverse to its velocity.



Since the electric and magnetic field associated to the nucleus take on the same absolute value, this transverse electromagnetic field can be simulated by an equivalent swarm of photons^a.



^aE. Fermi (1924), E. J. Williams (1933), C. F. Von Weizacker (1934)

Thus the collision of two charged nuclei at large impact parameter can be described as the collision of two equivalent swarms of photons.



Two-photon processes

- Two-photon processes can be computed within QED
- + Exclusive $\gamma\gamma \rightarrow ll$
 - Standard candle for photon-induced physics
 - Non-negligible background to Drell-Yan like reactions
- + Test of SM γ WW and $\gamma\gamma$ WW couplings

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Forward proton spectrometer



(drawing from Jesse Liu)

Trajectories of forward protons





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Roman pots









Feedthrough flange



ATLAS Forward Proton detectors - one arm



Kinematic matching



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Signal evidence



Forward physics

- \cdot physics between perturbative and non-perturbative QCD
- wide range of different topics
- standard and dedicated experimental methods