

Twist-2 and twist-3 gluon TMDs from back-to-back DIS dijets at N $\overline{\text{E}}$ ik accuracy

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- Introduction
 - Eikonal vs Next-to-eikonal
 - CGC vs TMD
- Quark propagator at NEik accuracy
- Scattering amplitude of DIS dijet production in the correlation limit
- Cross section: interplay of CGC and TMD

Eikonal approximation in the CGC

High-energy dense-dilute scattering in the CGC : Semiclassical and Eikonal approx.

Dense target represented by a **strong semiclassical gluon field** $\mathcal{A}^\mu(x) \propto 1/g$
 \Rightarrow Perturbative expansion in g needs improvement by all order resummation of $(g\mathcal{A}^\mu(x))^n$

Eikonal approx. : limit of **infinite boost** of $\mathcal{A}^\mu(x)$ along x^- :

- $\mathcal{A}^\mu(x)$ **independent on x^- (static limit)** due to Lorentz time dilation
 \Rightarrow No p^+ transfer from the target
- Lorentz contraction of $\mathcal{A}^\mu(x)$ (**shockwave limit**)
 \Rightarrow Partons from the projectile interact instantly in x^+ with the target, without transverse motion **within** the target
- Under a boost of parameter γ_t along the "–" direction, \mathcal{A}^- is enhanced and \mathcal{A}^+ is suppressed:
 $\mathcal{A}^- = O(\gamma_t) \gg \mathcal{A}_\perp = O(1) \gg \mathcal{A}^+ = O(1/\gamma_t)$

Background field in the eikonal limit: $\mathcal{A}^\mu(x^+, x^-, \mathbf{x}) \approx \delta^{\mu-} \mathcal{A}^-(x^+, \mathbf{x}) \propto \delta(x^+)$

\Rightarrow Only $(g\mathcal{A}^-(x^+, \mathbf{x}))^n$ needs all orders resummation \Rightarrow Wilson line along x^+

Next-to-Eikonal corrections to the CGC

Next-to-Eikonal (NEik) power corrections to the standard CGC formalism:

- Of order $1/\gamma_t$ at the level of the boosted background field
- Of order $1/s$ at the level of a cross section

→ They arise from relaxing either of the 3 approximations:

- 1 x^- dependence of $\mathcal{A}^\mu(x)$ beyond infinite Lorentz dilation
 → Treated as gradient expansion around a common x^- value:

$$\frac{\partial_- \mathcal{A}^-(x)}{\mathcal{A}^-(x)} = O(1/\gamma_t)$$
 ⇒ Possibility of (small) p^+ exchange with the target
- 2 Target with finite width
 ⇒ transverse motion of the projectile partons within the target
- 3 Interactions with \mathcal{A}_\perp field taken into account, not only \mathcal{A}^-

Note: Background quark field of the target also relevant at NEik.

CGC vs TMD approaches

connection between CGC and TMD \iff overlapping region between high-energy limit and correlation limit

For a process with a hard \mathbf{P} and a not so hard \mathbf{k} transverse momenta, and center-of-mass energy $\sqrt{s} \rightarrow$ **hierarchy of scales**

- CGC result: leading power (eikonal) in the limit $|\mathbf{k}| \sim |\mathbf{P}| \ll \sqrt{s}$
 - subleading power corrections suppressed as $\left(\frac{\mathbf{P}^2}{s}\right)^n$
- TMD factorization: leading power (twist 2) in the limit $|\mathbf{k}| \ll |\mathbf{P}| \sim \sqrt{s}$
 - subleading power corrections suppressed as $\left(\frac{|\mathbf{k}|}{|\mathbf{P}|}\right)^n$
- **CGC vs TMD: $|\mathbf{k}| \ll |\mathbf{P}| \ll \sqrt{s}$**
 - **what about power corrections in \mathbf{P}^2/s or $|\mathbf{P}||\mathbf{k}|/s$ beyond the eikonal limit?**

Consistency of both approaches shown in the double limit $|\mathbf{k}| \ll |\mathbf{P}| \ll \sqrt{s}$ at leading power ([Dominguez, Marquet, Xiao, Yuan, 2011](#))

Power corrections in $|\mathbf{k}|/|\mathbf{P}|$ in the regime $|\mathbf{k}| \ll |\mathbf{P}| \ll \sqrt{s}$ studied from the CGC approach ([Altinoluk, Boussarie, Kotko, 2019](#))

Full NEik quark propagator through a gluon background field

Propagator from y before the target to x after the target:

$$\begin{aligned}
 S_F(x, y) = & \int \frac{dq^+ d^2 \mathbf{q}}{(2\pi)^3} \int \frac{dk^+ d^2 \mathbf{k}}{(2\pi)^3} \theta(q^+) \theta(k^+) e^{-ix \cdot \tilde{q}} e^{iy \cdot \tilde{k}} \frac{(\not{\tilde{k}} + m)}{2q^+} \gamma^+ \\
 & \times \int d^2 \mathbf{z} e^{-iz \cdot (\mathbf{q} - \mathbf{k})} \left\{ \int dz^- e^{iz^- (q^+ - k^+)} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \right. \\
 & \left. + \frac{2\pi\delta(q^+ - k^+)}{(q^+ + k^+)} \left[-\frac{(\mathbf{q}^j + \mathbf{k}^j)}{2} \mathcal{U}_{F;j}^{(1)}(\mathbf{z}) - i \mathcal{U}_F^{(2)}(\mathbf{z}) + \frac{[\gamma^i, \gamma^j]}{4} \mathcal{U}_{F;ij}^{(3)}(\mathbf{z}) \right] \right\} \frac{(\not{\tilde{k}} + m)}{2k^+} \\
 & + \text{NNEik}
 \end{aligned}$$

Altinoluk, Beuf, Czajka, Tymowska (2021); Altinoluk, Beuf (2022).
 See also Kovchegov *et al.* (2016-...); Chirilli (2019,2021).

- Generalized Eikonal contribution: also includes the NEik non-static corrections: overall z^- dependence of the Wilson line.

$$\mathcal{U}_F(x^+, y^+; \mathbf{z}, z^-) \equiv \mathbf{1} + \sum_{N=1}^{+\infty} \frac{1}{N!} \mathcal{P}_+ \left[-ig \int_{y^+}^{x^+} dz^+ t \cdot \mathcal{A}^-(z) \right]^N$$

- NEik contributions beyond the shockwave approx or due to \mathcal{A}_\perp .

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 & \times \int d^2 \mathbf{z} e^{-iz \cdot (\mathbf{q} - \mathbf{k})} \left\{ \int dz^- e^{iz^- (q^+ - k^+)} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \right. \\
 & \left. + \frac{2\pi \delta(q^+ - k^+)}{(q^+ + k^+)} \left[-\frac{(\mathbf{q}^j + \mathbf{k}^j)}{2} \mathcal{U}_{F;j}^{(1)}(\mathbf{z}) - i \mathcal{U}_F^{(2)}(\mathbf{z}) + \frac{[\gamma^i, \gamma^j]}{4} \mathcal{U}_{F;ij}^{(3)}(\mathbf{z}) \right] \right\} \frac{(\not{\bar{\mathbf{k}}} + m)}{2k^+} \\
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 See also Kovchegov *et al.* (2016-...); Chirilli (2019,2021).

NEik correction due to the overall transverse drift of the quark during its interaction with the target:

$$\begin{aligned}
 \mathcal{U}_{F;j}^{(1)}(\mathbf{z}) = & \int dz^+ \mathcal{U}_F(+\infty, z^+; \mathbf{z}) \overleftrightarrow{D}_{z^+} \mathcal{U}_F(z^+, -\infty; \mathbf{z}) \\
 = & 2i \int dz^+ z^+ \mathcal{U}_F(+\infty, z^+; \mathbf{z}) g t \cdot \mathcal{F}_j^-(z^+, \mathbf{z}) \mathcal{U}_F(z^+, -\infty; \mathbf{z})
 \end{aligned}$$

Full NEik quark propagator through a gluon background field

Propagator from y before the target to x after the target:

$$\begin{aligned}
 S_F(x, y) = & \int \frac{dq^+ d^2 \mathbf{q}}{(2\pi)^3} \int \frac{dk^+ d^2 \mathbf{k}}{(2\pi)^3} \theta(q^+) \theta(k^+) e^{-ix \cdot \bar{q}} e^{iy \cdot \bar{k}} \frac{(\not{q} + m)}{2q^+} \gamma^+ \\
 & \times \int d^2 \mathbf{z} e^{-iz \cdot (\mathbf{q} - \mathbf{k})} \left\{ \int dz^- e^{iz^- (q^+ - k^+)} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \right. \\
 & \left. + \frac{2\pi\delta(q^+ - k^+)}{(q^+ + k^+)} \left[-\frac{(\mathbf{q}^j + \mathbf{k}^j)}{2} \mathcal{U}_{F;j}^{(1)}(\mathbf{z}) - i \mathcal{U}_{F;j}^{(2)}(\mathbf{z}) + \frac{[\gamma^i, \gamma^j]}{4} \mathcal{U}_{F;ij}^{(3)}(\mathbf{z}) \right] \right\} \frac{(\not{k} + m)}{2k^+} \\
 & + \text{NNEik}
 \end{aligned}$$

Altinoluk, Beuf, Czajka, Tymowska (2021); Altinoluk, Beuf (2022).
 See also Kovchegov *et al.* (2016-...); Chirilli (2019,2021).

NEik correction due to the transverse Brownian motion of the quark during its interaction with the target:

$$\begin{aligned}
 \mathcal{U}_F^{(2)}(\mathbf{z}) = & \int dz^+ \mathcal{U}_F(+\infty, z^+; \mathbf{z}) \overleftarrow{\mathcal{D}}_{\mathbf{z}^j} \overrightarrow{\mathcal{D}}_{\mathbf{z}^j} \mathcal{U}_F(z^+, -\infty; \mathbf{z}) \\
 = & - \int dz^+ \int^{z^+} dz'^+ (z^+ - z'^+) \mathcal{U}_F(+\infty, z^+, \mathbf{z}) g t \cdot \mathcal{F}_j^-(z^+, \mathbf{z}) \\
 & \times \mathcal{U}_F(z^+, z'^+, \mathbf{z}) g t \cdot \mathcal{F}_j^-(z'^+, \mathbf{z}) \mathcal{U}_F(z'^+, -\infty; \mathbf{z})
 \end{aligned}$$

Full NEik quark propagator through a gluon background field

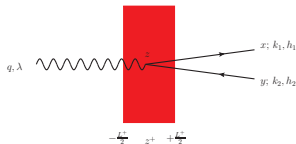
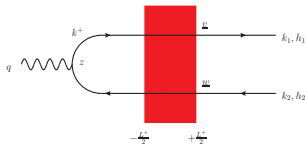
Propagator from y before the target to x after the target:

$$\begin{aligned}
 S_F(x, y) = & \int \frac{dq^+ d^2 \mathbf{q}}{(2\pi)^3} \int \frac{dk^+ d^2 \mathbf{k}}{(2\pi)^3} \theta(q^+) \theta(k^+) e^{-ix \cdot \tilde{q}} e^{iy \cdot \tilde{k}} \frac{(\not{\tilde{k}} + m)}{2q^+} \gamma^+ \\
 & \times \int d^2 \mathbf{z} e^{-iz \cdot (\mathbf{q} - \mathbf{k})} \left\{ \int dz^- e^{iz^- (q^+ - k^+)} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \right. \\
 & \left. + \frac{2\pi \delta(q^+ - k^+)}{(q^+ + k^+)} \left[-\frac{(\mathbf{q}^j + \mathbf{k}^j)}{2} \mathcal{U}_{F;j}^{(1)}(\mathbf{z}) - i \mathcal{U}_F^{(2)}(\mathbf{z}) + \frac{[\gamma^i, \gamma^j]}{4} \mathcal{U}_{F;ij}^{(3)}(\mathbf{z}) \right] \right\} \frac{(\not{\tilde{k}} + m)}{2k^+} \\
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 See also Kovchegov *et al.* (2016-...); Chirilli (2019,2021).

NEik coupling between the light-front helicity of the quark and the longitudinal chromomagnetic field of the target \mathcal{F}_{ij} :

$$\mathcal{U}_{F;ij}^{(3)}(\mathbf{z}) = \int dz^+ \mathcal{U}_F(+\infty, z^+; \mathbf{z}) g t \cdot \mathcal{F}_{ij}(z^+, \mathbf{z}) \mathcal{U}_F(z^+, -\infty; \mathbf{z})$$

DIS dijet at NEik accuracy: S-matrix for γ_L^* 

DIS dijet cross section calculated at NEik accuracy, at LO in α_s in the CGC.

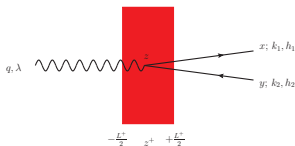
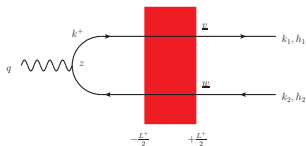
(Altinoluk, Beuf, Czajka, Tymowska, (2023))

- Second diagram vanishes in γ_L^* case, but matters in γ_T^* case.

S-matrix at NEik accuracy: $S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*} = S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{Gen. Eik}} + S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{dyn. target}} + S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{dec. on } q} + S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{dec. on } \bar{q}}$

$$S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{Gen. Eik}} = -2Q \frac{eef}{2\pi} \bar{u}(1) \gamma^+ v(2) \frac{(q^+ + k_1^+ - k_2^+)(q^+ + k_2^+ - k_1^+)}{4(q^+)^2} \int_{\mathbf{v}, \mathbf{w}} e^{-i\mathbf{v} \cdot \mathbf{k}_1} e^{-i\mathbf{w} \cdot \mathbf{k}_2} \\ \times K_0(\hat{Q} |\mathbf{v} - \mathbf{w}|) \int db^- e^{ib^-(k_1^+ + k_2^+ - q^+)} \left[\mathcal{U}_F(\mathbf{v}, b^-) \mathcal{U}_F^\dagger(\mathbf{w}, b^-) - 1 \right]$$

$$S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{dyn. target}} = 2\pi \delta(k_1^+ + k_2^+ - q^+) iQ \frac{eef}{2\pi} \bar{u}(1) \gamma^+ v(2) \frac{(k_1^+ - k_2^+)}{(q^+)^2} \int d^2\mathbf{v} e^{-i\mathbf{v} \cdot \mathbf{k}_1} \int d^2\mathbf{w} e^{-i\mathbf{w} \cdot \mathbf{k}_2} \\ \times \left[K_0(\bar{Q} |\mathbf{v} - \mathbf{w}|) - \frac{(\bar{Q}^2 - m^2)}{2\bar{Q}} |\mathbf{v} - \mathbf{w}| K_1(\bar{Q} |\mathbf{v} - \mathbf{w}|) \right] \left[\mathcal{U}_F(\mathbf{v}, b^-) \overleftrightarrow{\partial}_{b^-} \mathcal{U}_F^\dagger(\mathbf{w}, b^-) \right] \Big|_{b^- = 0}$$

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$$S_{q_1 \bar{q}_2 \leftarrow \gamma_L^*}^{\text{dec. on } q} = 2\pi\delta(k_1^+ + k_2^+ - q^+) \frac{eef}{2\pi} (-1)Q \frac{k_2^+}{(q^+)^2} \int d^2\mathbf{v} e^{-i\mathbf{v}\cdot\mathbf{k}_1} \int d^2\mathbf{w} e^{-i\mathbf{w}\cdot\mathbf{k}_2} K_0(\bar{Q}|\mathbf{v}-\mathbf{w}|) \\ \times \bar{u}(1)\gamma^+ \left[\frac{[\gamma^i, \gamma^j]}{4} \mathcal{U}_{F;ij}^{(3)}(\mathbf{v}) - i\mathcal{U}_F^{(2)}(\mathbf{v}) + \mathcal{U}_{F;j}^{(1)}(\mathbf{v}) \left(\frac{(\mathbf{k}_2^j - \mathbf{k}_1^j)}{2} + \frac{i}{2} \partial_{\mathbf{w}^j} \right) \right] \mathcal{U}_F^\dagger(\mathbf{w}) v(2)$$

Similar structure for the decoration on the anti-quark line.

Change of variables and back-to-back limit

Back-to-back limit of dijets is conveniently expressed in terms of:

(dijet momentum imbalance) $\mathbf{k} = \mathbf{k}_1 + \mathbf{k}_2$ and (relative momentum) $\mathbf{P} = (z_2 \mathbf{k}_1 - z_1 \mathbf{k}_2)$

$z_1 = k_1^+ / (k_1^+ + k_2^+)$ and $z_2 = k_2^+ / (k_1^+ + k_2^+) = 1 - z_1$ such that

$$\mathbf{k}_1 = \mathbf{P} + z_1 \mathbf{k}$$

$$\mathbf{k}_2 = -\mathbf{P} + z_2 \mathbf{k}$$

back-to-back correlation limit: $|\mathbf{k}| \ll |\mathbf{P}|$

In coordinate space:

(conjugate to \mathbf{k}) $\mathbf{b} = (z_1 \mathbf{v} + z_2 \mathbf{w})$ and (conjugate to \mathbf{P}) $\mathbf{r} = \mathbf{v} - \mathbf{w}$

such that

$$\mathbf{v} = \mathbf{b} + z_2 \mathbf{r}$$

$$\mathbf{w} = \mathbf{b} - z_1 \mathbf{r}$$

back-to-back correlation limit: $|\mathbf{r}| \ll |\mathbf{b}|$

Small \mathbf{r} expansion for the eikonal contribution

Open dipole from the Generalized Eikonal term for $\mathbf{r} = \mathbf{v} - \mathbf{w} \rightarrow 0$:

$$\begin{aligned}
 & \int_{\mathbf{b}} e^{-i\mathbf{b}\cdot\mathbf{k}} \left[\mathcal{U}_F(\mathbf{b} + z_2 \mathbf{r}, b^-) \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}, b^-) - 1 \right] \\
 = & \int_{\mathbf{b}} e^{-i\mathbf{b}\cdot\mathbf{k}} \left[-i \left(1 + \frac{i(z_2 - z_1)}{2} \mathbf{r}\cdot\mathbf{k} \right) \mathbf{r}^j t^{\alpha'} \int dv^+ \mathcal{U}_A(+\infty, v^+; \mathbf{b}, b^-)_{a'a} g\mathcal{F}_{j-a}^-(v^+, \mathbf{b}, b^-) \right. \\
 & - \frac{1}{2} \mathbf{r}^i \mathbf{r}^j t^{\alpha'} t^{b'} \int dv^+ \int dw^+ \mathcal{U}_A(+\infty, v^+; \mathbf{b}, b^-)_{a'a} g\mathcal{F}_{i-a}^-(v^+, \mathbf{b}, b^-) \\
 & \left. \times \mathcal{U}_A(+\infty, w^+; \mathbf{b}, b^-)_{b'b} g\mathcal{F}_{j-b}^-(w^+, \mathbf{b}, b^-) + \mathcal{O}(|\mathbf{r}|^3) \right]
 \end{aligned}$$

- \Rightarrow 0th order in the \mathbf{r} expansion trivial \rightarrow first order in \mathbf{r} is the leading power
- \Rightarrow interplay between subleading power corrections \rightarrow terms of order \mathbf{r}^2 needed as well!
- \Rightarrow order $|\mathbf{r}|^2$ correction: contributions with either one or two field strength \mathcal{F}_\perp^-

Small r limit for the non-static and helicity-dependent NEik correction

For the open decorated dipole due to the dynamics of the target:

$$\left[\mathcal{U}_F(\mathbf{b} + z_2 \mathbf{r}, b^-) \overleftrightarrow{\partial}_{b^-} \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}, b^-) \right] \Big|_{b^- = 0} = 2it^{a'} \int_{z^+} \mathcal{U}_A(+\infty, z^+; \mathbf{b})_{a'a} g\mathcal{F}_a^{+-}(z^+, \mathbf{b})$$

\Rightarrow involves the longitudinal chromoelectric field \mathcal{F}^{+-} instead of the transverse field \mathcal{F}_j^-

For the term with $\mathcal{U}_{F;ij}^{(3)}$ decorating the quark line:

$$\mathcal{U}_{F;ij}^{(3)}(\mathbf{b} + z_2 \mathbf{r}) \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}) = t^{a'} \int_{z^+} \mathcal{U}_A(+\infty, z^+; \mathbf{b})_{a'a} g\mathcal{F}_{ij}^a(z^+, \mathbf{b})$$

\Rightarrow involves the longitudinal chromomagnetic field \mathcal{F}_{ij} instead of \mathcal{F}_j^-

Emergence of \mathcal{F}^{+-} and \mathcal{F}_{ij} \rightarrow NEik suppression with respect to \mathcal{F}_j^-

Small r limit for the NEik corrections in $\mathcal{U}_{F;j}^{(1)}$ and $\mathcal{U}_F^{(2)}$

Terms with $\mathcal{U}_{F;j}^{(1)}$ and $\mathcal{U}_F^{(2)}$ decorating the quark line (remembering that $|\mathbf{r}| \sim 1/|\mathbf{P}|$):

$$\int_{\mathbf{b}} e^{-i\mathbf{b}\cdot\mathbf{k}} \left[- \left(\mathbf{P}^j + \frac{(z_1 - z_2)}{2} \mathbf{k}^j \right) \mathcal{U}_{F;j}^{(1)}(\mathbf{b} + z_2 \mathbf{r}) \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}) \right. \\ \left. + \frac{i}{2} \mathcal{U}_{F;j}^{(1)}(\mathbf{b} + z_2 \mathbf{r}) \partial_j \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}) - i \mathcal{U}_F^{(2)}(\mathbf{b} + z_2 \mathbf{r}) \mathcal{U}_F^\dagger(\mathbf{b} - z_1 \mathbf{r}) \right]$$

$$\mathcal{U}_{F;j}^{(1)}(\mathbf{b}) \mathcal{U}_F^\dagger(\mathbf{b}) = 2it^{a'} \int_{z^+} z^+ \mathcal{U}_A(+\infty, z^+; \mathbf{b})_{a'a} g\mathcal{F}_j^{a-}(z^+, \mathbf{b})$$

$$\mathcal{U}_F^{(2)}(\mathbf{b}) \mathcal{U}_F^\dagger(\mathbf{b}) = -t^{a'} t^{b'} \int_{z^+, z'^+} (z^+ - z'^+) \theta(z^+ - z'^+) \mathcal{U}_A(+\infty, z^+; \mathbf{b})_{a'a} g\mathcal{F}_j^{a-}(z^+, \mathbf{b}) \\ \times \mathcal{U}_A(+\infty, z'^+; \mathbf{b})_{b'b} g\mathcal{F}_j^{b-}(z'^+, \mathbf{b})$$

$$\mathcal{U}_{F;j}^{(1)}(\mathbf{b}) \partial_i \mathcal{U}_F^\dagger(\mathbf{b}) = -2t^{a'} t^{b'} \int dz^+ \int dz'^+ z^+ \mathcal{U}_A(+\infty, z^+; \mathbf{b})_{a'a} g\mathcal{F}_j^{a-}(z^+, \mathbf{b}) \\ \times \mathcal{U}_A(+\infty, z'^+; \mathbf{b})_{b'b} g\mathcal{F}_i^{b-}(z'^+, \mathbf{b})$$

⇒ like in the GEik term → contributions with either 1 or 2 \mathcal{F}_\perp^-

⇒ now with an extra factor z^+ or $(z^+ - z'^+)$ → NEik suppression with the target width

⇒ similar results for decorations on the antiquark line

Back-to-back cross section: (Generalized) Eikonal piece

Squaring the single \mathcal{F}_\perp^- part of the Generalized Eikonal contribution in the back-to-back limit:

$$\begin{aligned} \frac{d\sigma_{\gamma_L^* \rightarrow q_1 \bar{q}_2}}{d\text{P.S.}} \Bigg|_{\text{Gen.Eik}}^{\mathcal{F}_\perp^- \mathcal{F}_\perp^-} &= g^2 (ee_f)^2 Q^2 (q^+ + k_1^+ - k_2^+)^2 (q^+ - k_1^+ + k_2^+)^2 \frac{k_1^+ k_2^+}{4(q^+)^6} \left[\frac{4\mathbf{P}^i \mathbf{P}^j}{(\mathbf{P}^2 + \hat{Q}^2)^4} - 2(z_2 - z_1) \frac{(\mathbf{P}^i \mathbf{k}^j + \mathbf{k}^i \mathbf{P}^j)}{[\mathbf{P}^2 + \hat{Q}^2]^4} \right. \\ &\quad \left. + 16(z_2 - z_1) \frac{(\mathbf{k} \cdot \mathbf{P}) \mathbf{P}^i \mathbf{P}^j}{[\mathbf{P}^2 + \hat{Q}^2]^5} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^8}\right) \right] \int d(\Delta b^-) e^{i\Delta b^- (k_1^+ + k_2^+ - q^+)} \int_{\mathbf{b}, \mathbf{b}'} e^{-i\mathbf{k} \cdot (\mathbf{b} - \mathbf{b}')} \int_{z^+, z'^+} \\ &\quad \times \left\langle \mathcal{F}_i^a \left(z'^+, \mathbf{b}', -\frac{\Delta b^-}{2} \right) \left[\mathcal{U}_A^\dagger \left(+\infty, z'^+, \mathbf{b}', -\frac{\Delta b^-}{2} \right) \mathcal{U}_A \left(+\infty, z^+, \mathbf{b}, \frac{\Delta b^-}{2} \right) \right]_{ab} \mathcal{F}_j^b \left(z^+, \mathbf{b}, \frac{\Delta b^-}{2} \right) \right\rangle \end{aligned}$$

Strict Eikonal result found by neglecting Δb^- in the fields:

$$\begin{aligned} \frac{d\sigma_{\gamma_L^* \rightarrow q_1 \bar{q}_2}}{d\text{P.S.}} \Bigg|_{\text{Strict.Eik}}^{\mathcal{F}_\perp^- \mathcal{F}_\perp^-} &= (2q^+) 2\pi \delta(k_1^+ + k_2^+ - q^+) (ee_f)^2 g^2 4z_1^3 z_2^3 Q^2 \\ &\quad \times \left[\frac{4\mathbf{P}^i \mathbf{P}^j}{(\mathbf{P}^2 + \hat{Q}^2)^4} - 2(z_2 - z_1) \frac{(\mathbf{P}^i \mathbf{k}^j + \mathbf{k}^i \mathbf{P}^j)}{[\mathbf{P}^2 + \hat{Q}^2]^4} + 16(z_2 - z_1) \frac{(\mathbf{k} \cdot \mathbf{P}) \mathbf{P}^i \mathbf{P}^j}{[\mathbf{P}^2 + \hat{Q}^2]^5} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^8}\right) \right] \\ &\quad \times \int_{\mathbf{b}, \mathbf{b}'} e^{-i\mathbf{k} \cdot (\mathbf{b} - \mathbf{b}')} \int_{z^+, z'^+} \left\langle \mathcal{F}_i^a \left(z'^+, \mathbf{b}' \right) \left[\mathcal{U}_A^\dagger \left(+\infty, z'^+, \mathbf{b}' \right) \mathcal{U}_A \left(+\infty, z^+, \mathbf{b} \right) \right]_{ab} \mathcal{F}_j^b \left(z^+, \mathbf{b} \right) \right\rangle \end{aligned}$$

- **Correlator related to twist-2 gluon TMDs in the target**, with momentum fraction $x = 0$ and transverse momentum \mathbf{k} , with a *future staple* gauge link.
- **Kinematical twist 3 corrections**, suppressed by an extra $|\mathbf{k}|/|\mathbf{P}|$ in the back-to-back dijet limit $|\mathbf{k}| \ll |\mathbf{P}|$
- Not shown here: **Genuine twist 3 corrections**, involving a correlator of the type $\langle \mathcal{F}_\perp^- \mathcal{F}_\perp^- \mathcal{F}_\perp^- \rangle$
- Difference between Gen. Eik and strict Eik. : involves correlator $\langle \mathcal{F}_\perp^- \mathcal{F}_\perp^- \mathcal{F}^{+-} \rangle$ or $\mathbf{k} \langle \mathcal{F}_\perp^- \mathcal{F}^{+-} \rangle$
 \Rightarrow NEik correction but twist 4: beyond our accuracy here!

Back-to-back cross section: twist 2 TMDs contribution

Including all contributions of the form $\langle \mathcal{F}_{\perp}^- \mathcal{F}_{\perp}^- \rangle$, of order Eik or NEik, and twist 2 or twist 3:

$$\begin{aligned}
 \frac{d\sigma_{\gamma_L^+ \rightarrow q_1 \bar{q}_2}}{d\text{P.S.}} \Big|_{\mathcal{F}_{\perp}^- \mathcal{F}_{\perp}^-} &= (2q^+) 2\pi \delta(k_1^+ + k_2^+ - q^+) (e e_f)^2 g^2 4z_1^3 z_2^3 Q^2 \\
 &\times \left[\frac{4\mathbf{P}^i \mathbf{P}^j}{(\mathbf{P}^2 + \bar{Q}^2)^4} - 2(z_2 - z_1) \frac{(\mathbf{P}^i \mathbf{k}^j + \mathbf{k}^i \mathbf{P}^j)}{[\mathbf{P}^2 + \bar{Q}^2]^4} + 16(z_2 - z_1) \frac{(\mathbf{k} \cdot \mathbf{P}) \mathbf{P}^i \mathbf{P}^j}{[\mathbf{P}^2 + \bar{Q}^2]^5} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^8}\right) \right] \\
 &\times \int_{\mathbf{b}, \mathbf{b}'} e^{-i\mathbf{k} \cdot (\mathbf{b} - \mathbf{b}')} \int_{z^+, z'^+} \left[1 + i(z^+ - z'^+) \frac{(\mathbf{P}^2 + \bar{Q}^2)}{2q^+ z_1 z_2} + \text{NNEik} \right] \\
 &\times \left\langle \mathcal{F}_i^a(z'^+, \mathbf{b}') \left[\mathcal{U}_A^\dagger(+\infty, z'^+; \mathbf{b}') \mathcal{U}_A(+\infty, z^+; \mathbf{b}) \right]_{ab} \mathcal{F}_j^b(z^+, \mathbf{b}) \right\rangle
 \end{aligned}$$

- NEik corrections and kinematic twist 3 corrections to the $\langle \mathcal{F}_{\perp}^- \mathcal{F}_{\perp}^- \rangle$ contribution factorize from each other!
- Result of the same form can be obtained for transverse photon case.

Back-to-back cross section: twist 2 TMDs contribution

Introducing notation: TMD $\langle \mathcal{F}^{\mu\nu} \mathcal{F}^{\rho\sigma} \rangle$ correlator in unpolarized target:

$$\begin{aligned} \Phi^{\mu\nu;\rho\sigma}(x, \mathbf{k}) &\equiv \frac{1}{x P_{tar}^-} \frac{1}{(2\pi)^3} \int d^2\mathbf{z} e^{-i\mathbf{k}\cdot\mathbf{z}} \int dz^+ e^{ixP_{tar}^- z^+} \left\langle P_{tar} \left| \mathcal{F}_a^{\mu\nu}(0) \left[\mathcal{U}_A^\dagger(+\infty, 0; 0) \mathcal{U}_A(+\infty, z^+; \mathbf{z}) \right]_{ab} \mathcal{F}_b^{\rho\sigma}(z^+, \mathbf{z}) \right| P_{tar} \right\rangle \\ &= \frac{1}{(2\pi)^3} \int_{\mathbf{b}, \mathbf{b}'} e^{-i\mathbf{k}\cdot(\mathbf{b}-\mathbf{b}')} \int_{z^+, z'^+} e^{ixP_{tar}^- (z^+ - z'^+)} \left\langle \mathcal{F}_a^{\mu\nu}(z'^+, \mathbf{b}') \left[\mathcal{U}_A^\dagger(+\infty, z'^+; \mathbf{b}') \mathcal{U}_A(+\infty, z^+; \mathbf{b}) \right]_{ab} \mathcal{F}_b^{\rho\sigma}(z^+, \mathbf{b}) \right\rangle, \end{aligned}$$

$$\begin{aligned} \frac{d\sigma_{\gamma_{L,T}^* \rightarrow q_1 \bar{q}_2}}{dz_1 d^2\mathbf{P} d^2\mathbf{k}} \Big|_{\mathcal{F}_\perp^- \mathcal{F}_\perp^-} &= \alpha_{em} \alpha_s e_f^2 C_{T,L}^{ij}(z_1, \mathbf{P}, \mathbf{k}) \left\{ \left[1 + \frac{(\mathbf{P}^2 + \bar{Q}^2)}{z_1 z_2 W^2} \partial_x + \text{NNEik} \right] \left[x \Phi^{i-j-}(x, \mathbf{k}) \right] \right\} \Big|_{x=0} \\ &= \alpha_{em} \alpha_s e_f^2 C_{T,L}^{ij}(z_1, \mathbf{P}, \mathbf{k}) \left\{ \left[x \Phi^{i-j-}(x, \mathbf{k}) \right] \Big|_{x=\frac{(\mathbf{P}^2 + \bar{Q}^2)}{z_1 z_2 W^2}} + \text{NNEik} \right\} \end{aligned}$$

Non-zero value of momentum fraction x in the twist 2 gluon TMDs recovered from NEik corrections!

Standard parametrization of $\Phi^{i-j-}(x, \mathbf{k})$ (with the target mass M):

$$\Phi^{i-j-}(x, \mathbf{k}) = \frac{\delta^{ij}}{2} f_1^g(x, \mathbf{k}) + \left[\mathbf{k}^i \mathbf{k}^j - \frac{\mathbf{k}^2}{2} \delta^{ij} \right] \frac{1}{2M^2} h_1^{\perp g}(x, \mathbf{k})$$

Back-to-back cross section: twist 3 TMDs from NEik

From the interference between the non-static NEik correction and the strict Eikonal amplitudes: **terms in $\langle \mathcal{F}^{+-} \mathcal{F}_{\perp}^{-} \rangle$**

$$\frac{d\sigma_{\gamma_{L,T}^* \rightarrow q_1 \bar{q}_2}}{dz_1 d^2\mathbf{P} d^2\mathbf{k}} \Big|_{NEik}^{\mathcal{F}_{\perp}^{-} \mathcal{F}^{+-}} = \alpha_{em} \alpha_s e_f^2 \frac{1}{2q^+} C_{T,L}^j(z_1, \mathbf{P}) \left[x \Phi^{j-;+-}(x, \mathbf{k}) + x \Phi^{+-;j-}(x, \mathbf{k}) \right] \Big|_{x=0}$$

\Rightarrow NEik. correction beyond the static approximation for the target involves a **twist-3 gluon TMD**, ([Mulders, Rodrigues \(2001\)](#)) with momentum fraction $x = 0$.

From the interference between the NEik correction with $U_{F;ij}^{(3)}$ and the strict Eikonal amplitude: **terms in $\langle \mathcal{F}^{ij} \mathcal{F}_{\perp}^{-} \rangle$**

$$\frac{d\sigma_{\gamma_{L,T}^* \rightarrow q_1 \bar{q}_2}}{dz_1 d^2\mathbf{P} d^2\mathbf{k}} \Big|_{NEik}^{\mathcal{F}_{\perp}^{-} \mathcal{F}^{ij}} = \alpha_{em} \alpha_s e_f^2 \frac{1}{2q^+} C_{T,L}^{ijl}(z_1, \mathbf{P}) \left[x \Phi^{l-;ij}(x, \mathbf{k}) + x \Phi^{ij;l-}(x, \mathbf{k}) \right] \Big|_{x=0}$$

- Contribution from another type of **twist-3 gluon TMD**, ([Mulders, Rodrigues \(2001\)](#)), still with $x = 0$.
- Note: term absent in the γ_{L}^* : $C_{L}^{ijl}(z_1, \mathbf{P}) = 0$ due to Dirac algebra.

Parametrization (from [Lorcé, Pasquini \(2013\)](#)):

$$\begin{aligned} \Phi^{j-;+-}(x, \mathbf{k}) + \Phi^{+-;j-}(x, \mathbf{k}) &= \frac{2\mathbf{k}^j}{P_{tar}^-} f^{\perp g}(x, \mathbf{k}) \\ \Phi^{l-;ij}(x, \mathbf{k}) + \Phi^{ij;l-}(x, \mathbf{k}) &= \epsilon^{ij} \epsilon^{ln} \frac{2\mathbf{k}^n}{P_{tar}^-} \bar{g}^{\perp g}(x, \mathbf{k}) \end{aligned}$$

Final result: $\langle \mathcal{FF} \rangle$ contributions to the cross sections

$$\frac{d\sigma_{\gamma_{T,L}^* \rightarrow q_1 \bar{q}_2}}{dz_1 d^2\mathbf{P} d^2\mathbf{k}} \Big|_{\text{Eik+NEik}}^{\mathcal{FF}} = \alpha_{\text{em}} e_f^2 \alpha_s \left\{ C_{T,L}^{f_1^g}(z_1, \mathbf{P}, \mathbf{k}) \times f_1^g(x, \mathbf{k}) + C_{T,L}^{h_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times h_1^{+g}(x, \mathbf{k}) \right. \\ \left. + \frac{1}{W^2} C_{T,L}^{f_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times f_1^{+g}(x, \mathbf{k}) + \frac{1}{W^2} C_{T,L}^{\bar{g}^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times \bar{g}^{+g}(x, \mathbf{k}) \right\} \Big|_{x=\frac{|\mathbf{P}^2+\bar{Q}^2|}{z_1 z_2 W^2}}$$

Contribution from longitudinal photon exchange:

$$C_L^{f_1^g}(z_1, \mathbf{P}, \mathbf{k}) = \frac{8Q^2 z_1^2 z_2^2}{[\mathbf{P}^2 + \bar{Q}^2]^4} \left\{ \mathbf{P}^2 + (z_2 - z_1)(\mathbf{k} \cdot \mathbf{P}) \left[-1 + \frac{4\mathbf{P}^2}{[\mathbf{P}^2 + \bar{Q}^2]} \right] \right\} + O\left(\frac{Q^2 \mathbf{k}^2}{\mathbf{P}^8}\right)$$

$$C_L^{h_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = \frac{4Q^2 z_1^2 z_2^2}{[\mathbf{P}^2 + \bar{Q}^2]^4} \frac{\mathbf{k}^2}{M^2} \left\{ \left(\frac{2(\mathbf{k} \cdot \mathbf{P})^2}{\mathbf{k}^2 \mathbf{P}^2} - 1 \right) \mathbf{P}^2 \right. \\ \left. + (z_2 - z_1)(\mathbf{k} \cdot \mathbf{P}) \left[-1 + \frac{4\mathbf{P}^2}{[\mathbf{P}^2 + \bar{Q}^2]} \left(\frac{2(\mathbf{k} \cdot \mathbf{P})^2}{\mathbf{k}^2 \mathbf{P}^2} - 1 \right) \right] \right\} + O\left(\frac{Q^2 \mathbf{k}^2}{\mathbf{P}^8}\right)$$

$$C_L^{f_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = -32Q^2 z_1 z_2 \frac{(z_2 - z_1)(\mathbf{k} \cdot \mathbf{P})[\mathbf{P}^2 + m^2]}{[\mathbf{P}^2 + \bar{Q}^2]^4} + O\left(\frac{Q^2 \mathbf{k}^2}{\mathbf{P}^6}\right)$$

$$C_L^{\bar{g}^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = 0$$

Final result: $\langle \mathcal{FF} \rangle$ contributions to the cross sections

$$\frac{d\sigma_{\gamma_{T,L}^* \rightarrow q_1 \bar{q}_2}}{dz_1 d^2\mathbf{P} d^2\mathbf{k}} \Big|_{\text{Eik+NEik}}^{\mathcal{FF}} = \alpha_{\text{em}} e_f^2 \alpha_s \left\{ C_{T,L}^{f_1^g}(z_1, \mathbf{P}, \mathbf{k}) \times f_1^g(x, \mathbf{k}) + C_{T,L}^{h_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times h_1^{+g}(x, \mathbf{k}) \right. \\ \left. + \frac{1}{W^2} C_{T,L}^{f_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times f_1^{+g}(x, \mathbf{k}) + \frac{1}{W^2} C_{T,L}^{\bar{g}^{+g}}(z_1, \mathbf{P}, \mathbf{k}) \times \bar{g}^{+g}(x, \mathbf{k}) \right\} \Big|_{x=\frac{[\mathbf{P}^2+\bar{Q}^2]}{z_1 z_2 W^2}}$$

Contribution from transverse photon exchange:

$$C_T^{f_1^g}(z_1, \mathbf{P}, \mathbf{k}) = -\frac{2[(z_1^2+z_2^2)\bar{Q}^2-m^2]}{[\mathbf{P}^2+\bar{Q}^2]^4} \left\{ \mathbf{P}^2 + (z_2-z_1)(\mathbf{k}\cdot\mathbf{P}) \left[-1 + \frac{4\mathbf{P}^2}{[\mathbf{P}^2+\bar{Q}^2]} \right] \right\} \\ + \frac{(z_1^2+z_2^2)}{[\mathbf{P}^2+\bar{Q}^2]^2} \left[1 + \frac{2(z_2-z_1)(\mathbf{k}\cdot\mathbf{P})}{[\mathbf{P}^2+\bar{Q}^2]} \right] + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^6}\right)$$

$$C_T^{h_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = -\frac{[(z_1^2+z_2^2)\bar{Q}^2-m^2]}{[\mathbf{P}^2+\bar{Q}^2]^4} \frac{\mathbf{k}^2}{M^2} \left\{ \left(\frac{2(\mathbf{k}\cdot\mathbf{P})^2}{\mathbf{k}^2\mathbf{P}^2} - 1 \right) \mathbf{P}^2 \right. \\ \left. + (z_2-z_1)(\mathbf{k}\cdot\mathbf{P}) \left[-1 + \frac{4\mathbf{P}^2}{[\mathbf{P}^2+\bar{Q}^2]} \left(\frac{2(\mathbf{k}\cdot\mathbf{P})^2}{\mathbf{k}^2\mathbf{P}^2} - 1 \right) \right] \right\} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^6}\right)$$

$$C_T^{f_1^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = \frac{4(z_2-z_1)(\mathbf{k}\cdot\mathbf{P})}{[\mathbf{P}^2+\bar{Q}^2]^4} \left\{ [\mathbf{P}^2+\bar{Q}^2 + (z_1^2+z_2^2)Q^2][\mathbf{P}^2-\bar{Q}^2] + 2m^2Q^2 \right\} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^4}\right)$$

$$C_T^{\bar{g}^{+g}}(z_1, \mathbf{P}, \mathbf{k}) = \frac{4(z_2-z_1)(\mathbf{k}\cdot\mathbf{P})}{[\mathbf{P}^2+\bar{Q}^2]^2} + O\left(\frac{\mathbf{k}^2}{\mathbf{P}^4}\right)$$

Back-to-back cross section: 3 \mathcal{F} correlators

Leftover contributions in $\langle \mathcal{F}_\perp^- \mathcal{F}_\perp^- \mathcal{F}_\perp^- \rangle$, of order Eik or NEik, starting at twist 3.

Various terms in the cross section, with a correlator of the form:

$$2\text{Re tr}_F(t^{a'} t^{b'} t^{c'}) \int_{\mathbf{b}, \mathbf{b}'} e^{-i\mathbf{k} \cdot (\mathbf{b} - \mathbf{b}')} \int_{v^+, w^+, v'^+} \left[1 + i \frac{(\mathbf{P}^2 + \bar{Q}^2)}{2q^+ z_1 z_2} \phi(v^+ - v'^+, w^+ - v'^+) \right] \\ \times \left\langle \mathcal{U}_A(+\infty, v'^+; \mathbf{b}')_{c'c} g \mathcal{F}_l^c(v'^+, \mathbf{b}') \mathcal{U}_A(+\infty, v^+; \mathbf{b})_{a'a} g \mathcal{F}_i^a(v^+, \mathbf{b}) \mathcal{U}_A(+\infty, w^+; \mathbf{b})_{b'b} g \mathcal{F}_j^b(w^+, \mathbf{b}) \right\rangle$$

with several possible functions ϕ in the NEik correction.

In the TMD formalism, such 3 fields correlator should in principle depend on two different momentum fractions x_1 and x_2 via phase factors. However: not yet studied in the gluon case in the TMD literature.

Not entirely clear at the moment how to resum the NEik corrections unambiguously in order to recover this dependence on x_1 and x_2 .

Note: $\text{tr}_F(t^{a'} t^{b'} t^{c'}) = (d^{a'b'c'} + i f^{a'b'c'})/4$

- Antisymmetric piece in $f^{a'b'c'}$: Eikonal contribution cancel \Rightarrow starts at NEik only
- Symmetric piece in $d^{a'b'c'}$: both Eik and NEik terms. But still unclear how to interpret in terms of dependence on x_1 and x_2 .

Summary

To further understand the interplay between CGC and TMD, we studied the NEik DIS dijet cross-section in the back-to-back jets limit, including twist 3 power corrections. Various types of contributions are obtained:

- $\langle \mathcal{F}_i^- \mathcal{F}_j^- \rangle$: twist 2 gluon TMDs f_1^g and $h_1^{\perp g}$
 - Factorization of kinematic twist 3 and of NEik corrections
 - NEik correction is the first order correction in the Taylor expansion of the TMDs around $x = 0$
 \Rightarrow x dependence of the TMDs recovered by resumming terms of all powers beyond the eikonal approximation
- Twist 3 gluon TMDs found as contributions to the cross section at NEik and twist 3 order:
 - $f^{\perp g}$, of the type $\langle \mathcal{F}_i^- \mathcal{F}^{+-} \rangle$
 - $\bar{g}^{\perp g}$, of the type $\langle \mathcal{F}_l^- \mathcal{F}_{ij} \rangle$, for γ_T^* case
- 3 fields twist 3 correlators $\langle \mathcal{F}_i^- \mathcal{F}_j^- \mathcal{F}_l^- \rangle$: beyond TMD partonic distributions
 - Symmetric part in $d^{a'b'c'}$ already appear in Eikonal contributions
 - NEik corrections should be related to the dependence on x_1, x_2
 - Antisymmetric part in $f^{a'b'c'}$ starts only at NEik order

Recap: Corrections beyond the static approximation

- Relative z^- dependence along the same propagator : NNEik, higher twist ?
- Relative z^- dependence between Wilson lines in the amplitude: NEik, twist 3
→ $f^{\perp g}$ twist 3 gluon TMD
- Relative z^- dependence between amplitude and cc. amplitude: NEik, twist 4
→ Associated with the overall k^+ exchange with the target
- Overall z^- dependence at cross section level:
→ Disappear when performing target average, in (semi-)inclusive processes.

All these non-static corrections can be written as insertions of the longitudinal chromoelectric field \mathcal{F}^{+-} of the target.

More about NEik corrections beyond the static approx

Effect of relative z^- dependence of \mathcal{A}^- insertions along **one** propagator:

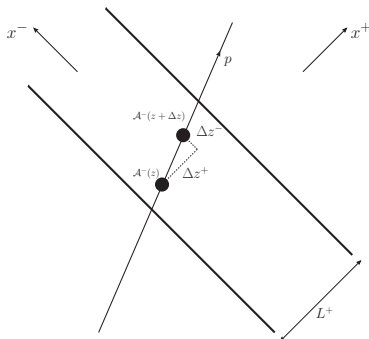
$$\mathcal{A}^-(z^- + \Delta z^-) - \mathcal{A}^-(z^-) \simeq \Delta z^- \partial_- \mathcal{A}^-(z^-)$$

- Slow z^- dependence from time dilation:

$$\partial_- \mathcal{A}^- \propto \frac{1}{\gamma t} \mathcal{A}^-$$

- Small Δz^- displacement of the trajectory within the target width L^+ :

$$\Delta z^- \sim \frac{p^-}{p^+} \Delta z^+ \leq \frac{p^-}{p^+} L^+ = O\left(\frac{1}{\gamma t}\right)$$



Double power suppression, beyond static approx and beyond shockwave approx:

⇒ NNEik effect within a single propagator!

More about NEik corrections beyond the static approx

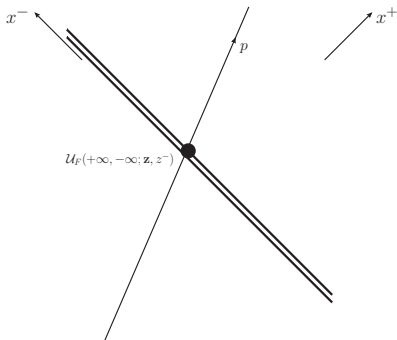
Effect of relative z^- dependence of \mathcal{A}^- insertions along **one** propagator is NNEik.

However, dependence on average z^- is suppressed only once.

⇒ Use Wilson lines with overall z^- dependence

$$\partial_- \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \propto \frac{1}{\gamma t} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-)$$

→ Accounts for NEik effects beyond static approx



More about NEik corrections beyond the static approx

Effect of relative z^- dependence of \mathcal{A}^- insertions along **one** propagator is NNEik.

However, dependence on average z^- is suppressed only once.

⇒ Use Wilson lines with overall z^- dependence

$$\partial_- \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-) \propto \frac{1}{\gamma_t} \mathcal{U}_F(+\infty, -\infty; \mathbf{z}, z^-)$$

→ Accounts for NEik effects beyond static approx

In particular: NEik corrections induced by the difference in z^- between different Wilson lines.

