# Strongly coupled dynamics with a phase transition – some surprises

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RJ, M. Järvinen, H. Soltanpanahi, J. Sonnenschein, PRL '22 [2205.06274] RJ, M. Järvinen, J. Sonnenschein, to appear...

### Long term motivation

**Domain wall velocities** 

Conventional picture and a key question Results of holographic simulations A simple formula for domain wall velocity

**Boost-invariant expansion and hot remnants** Boost-invariant expansion and cooling Reheating of plasma remnants Entropy considerations

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Question: How to understand phase transitions in real time within holography?

- Phase transitions in equilibrium in holography correspond to switching between two spacetime geometries
- It is nontrivial what happens in real time...
- In some cases there is a classical gravity description, in others probably not...

#### Theoretically interesting even for its own sake!

But also interesting real-world applications:

What does holography tell us about hadronization??

- ▶ 1<sup>st</sup> order phase transition reappears at nonzero density...
- Some physics in early universe??

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caution for  $N_c = 3$ : crossover!

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# 1. Bubble wall velocities

- At a 1<sup>st</sup> order phase transition T = T<sub>c</sub>, we can have domains of coexisting phases separated by domain walls
- The pressures on both sides are balanced and the domain wall can be static...

Question: What happens when we move away from  $T = T_c$ ?

- This can occur for nucleated bubbles of a stable phase within an supercooled medium
- At an interface between phases away from  $T = T_c$
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Goal:

Understand bubble wall velocities at strong coupling...

...but this does not happen — the domain wall ultimately moves with a constant velocity...

Common lore: friction in the second phase balances the net force
 — challenging to calculate...



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$$\rho_{4} > \rho_{2}$$
 $r_{en}^{\nu(1)} \rho_{2}$ 



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$$\rho_{A} \supset \rho_{2}$$
  $\rho_{2}$   $\rho_{2$ 



- ► The net force across the domain wall implies that the pressure difference is localized close to the domain wall...
- It is not obvious *a-priori* if this is always the case...



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- Witten model in 3D confinement/deconfinement transition use simplified hydrodynamics+scalar field fitted to holography
- **2.** Holographic gravity+scalar model with a transition between two deconfined phases

full holographic simulation

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- ▶ We perform time evolution from the above static initial conditions...
- The pressure in each phase may be read off from the T<sup>ys</sup> component of the energy-momentum tensor



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- ► The large pressure difference appears **away** from the domain wall
- ► The pressure is essentially constant across the domain wall, and very close to p(T<sub>c</sub>)...



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#### Key features:

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- ▶ The change in pressures occurs in the high energy density phase

ightarrow hydrodynamic description

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## We can extract the hydrodynamic velocity from $\mathcal{T}^{\times t}$



- The hydrodynamic velocity is quite close to the domain wall velocity v<sub>dw</sub>
- This gets better with increasing ratio of entropies in the two phases...
- We can fomulate finding  $v_{dw}$  as a hydrodynamic problem:

This leads to the formula:

$$v_{dw} = anh \int_{
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Compare with holographic simulations:

for a nucleated bubble at rest one can correct for deviation of  $v_{dw}$  from  $v_{hydro}$ 

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## **Domain wall velocities**

- 1. Movement of a domain wall is accompanied by a hydrodynamic wave
- 2. The pressure difference is realized on the hydro wave and **not** on the domain wall
- No need to invoke friction domain wall velocity expressed in terms of EOS (perfect fluid hydro)
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# 2. Boost-invariant expansion and hot remnants



- The plasma undergoes boost-invariant expansion in the longitudinal direction
- The expansion induces generically cooling and thus naturally pushes the system across the phase transition
- A scale invariant plasma cools as  $T( au) \sim 1/ au^{rac{1}{3}}$
- Entropy per unity rapidity is preserved

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 $S \propto \tau V_{\perp} T^3 \sim const$ 

We would expect the expanding plasma to be more and more supercooled and transitioning to the other phase through nucleating more and more bubbles... ► We study transverse dynamics in a theory with a confinement deconfinement phase transition

(3D Witten model in a simplified hydro+scalar field description)

- If we start from a deconfined plasma, we expect to eventually arrive in the confining phase...
- There would be quantum tunnelling which we cannot describe, so we introduce a "seed" in the initial conditions

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- ▶ The plasma in-between does not seem to want to decay!
- $\blacktriangleright$  We would expect it to be more and more supercooled (  $T \sim 1/ au^{rac{1}{3}}$  ) X
- $\blacktriangleright$  But it effectively heats up and  $\mathit{T} \sim \mathit{T_c} \sim 1$   $\checkmark$
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Recall entropy per unit rapidity

$$S \propto au V_{\perp} T( au)^3 \sim \textit{const} \qquad ext{for} \quad T( au) \sim rac{1}{ au^{rac{1}{3}}}$$

But we can also have entropy conservation with

$$T( au)\sim {\it const}$$
 and  $V_\perp\sim rac{1}{ au}$ 

Indeed this holds until the size is of the order of the domain wall thickness...

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This may perhaps be seen as an analog of "hadronization" (or "glueballization")??

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21 / 22

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- The dynamics of domain walls at strong coupling is much simpler than one could expect
- ► The pressure difference between the phases is **not** localized in the vicinity of the **domain wall**
- but within a hydrodynamic wave in the high entropy phase
- This provides a very simple hydrodynamic formula for the domain wall velocity expressed purely in terms of the equation of state
- We studied phase transition in a boost-invariant setting
- We find evidence for hot remnants of plasma which shrink but do not become overcooled...

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