



EVOLUTION OF RADIATION INDUCED POROSITY IN METASTABLE MATERIALS APPLIED AT CRYOGENIC TEMPERATURES

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Outline:

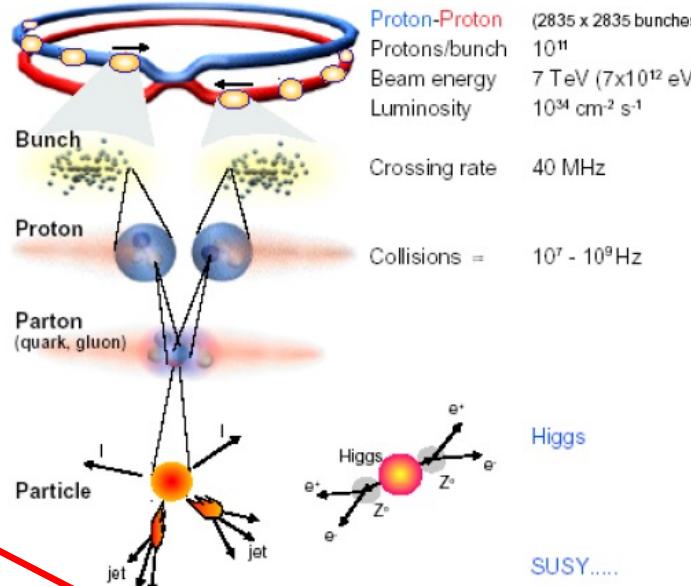
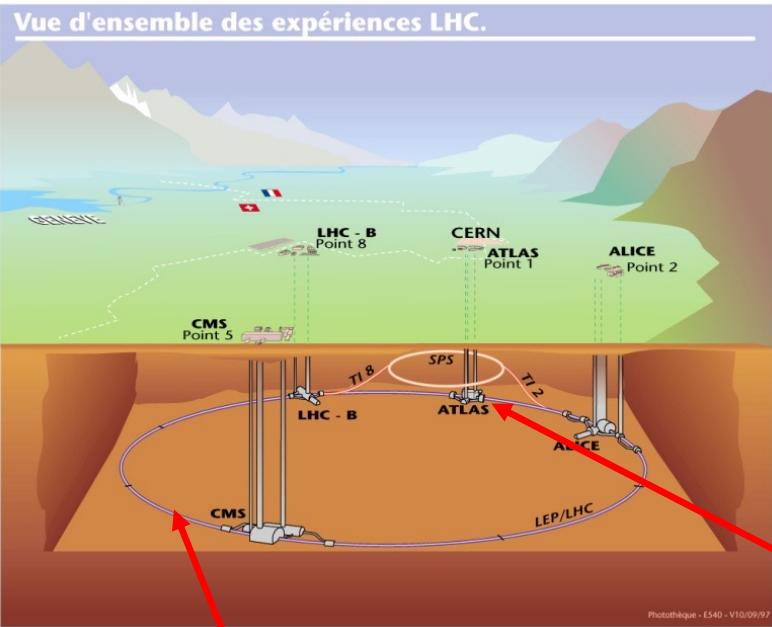
1. Motivation: radiation sources at cryogenic temperatures
2. Strain induced fcc-bcc phase transformation
3. Radiation induced damage
4. Radiation induced hardening
5. Conclusions



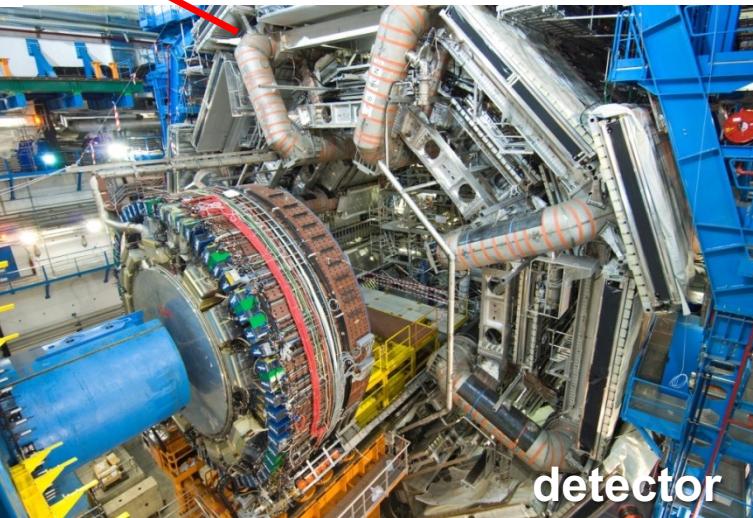
Motivation: CERN Large Hadron Collider

...ative phenomena at cryogenic temperatures

Vue d'ensemble des expériences LHC.



accelerator



detector

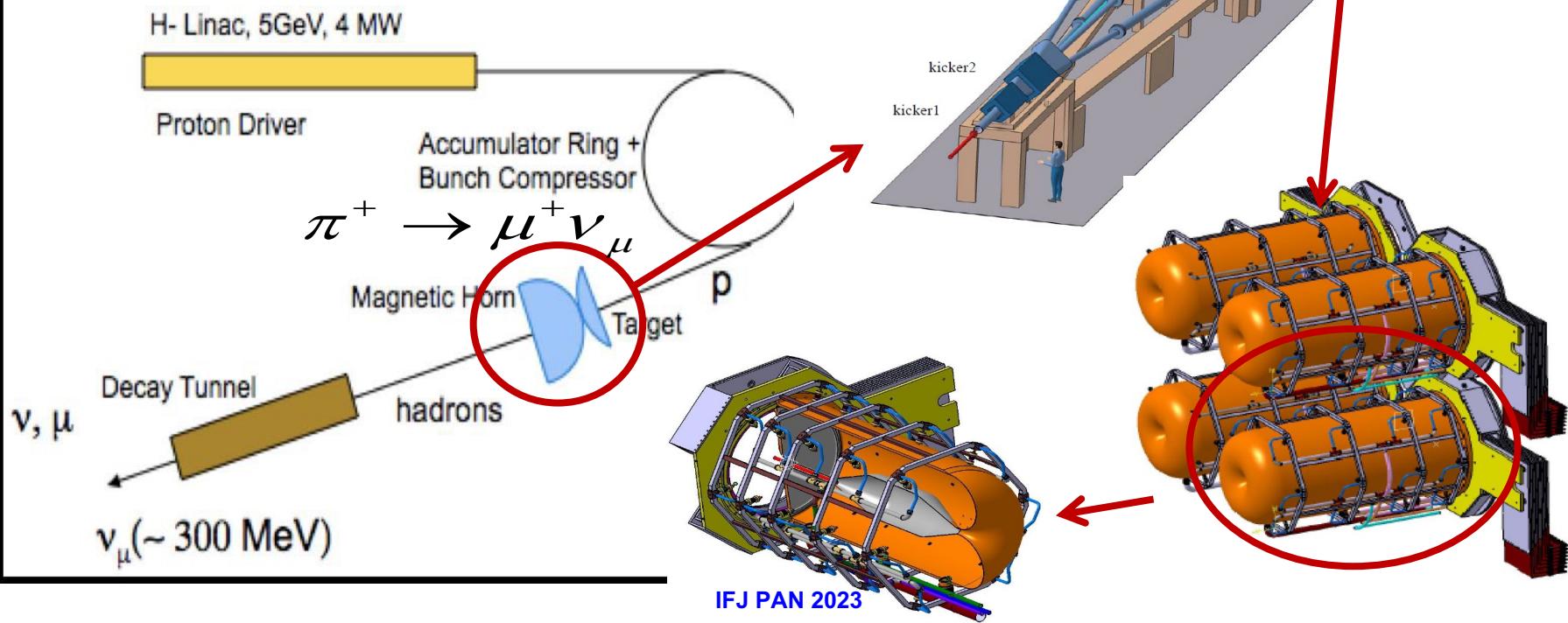
LHC
is the largest
scientific
instrument
in the world
based on the
principle of
super-
conductivity!

LHC
operates in
super-fluid
helium
at 1.9 K



EUROnu: High Intensity Neutrino Oscillation Facility in EU

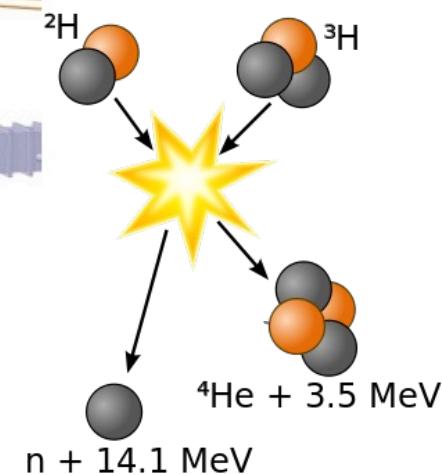
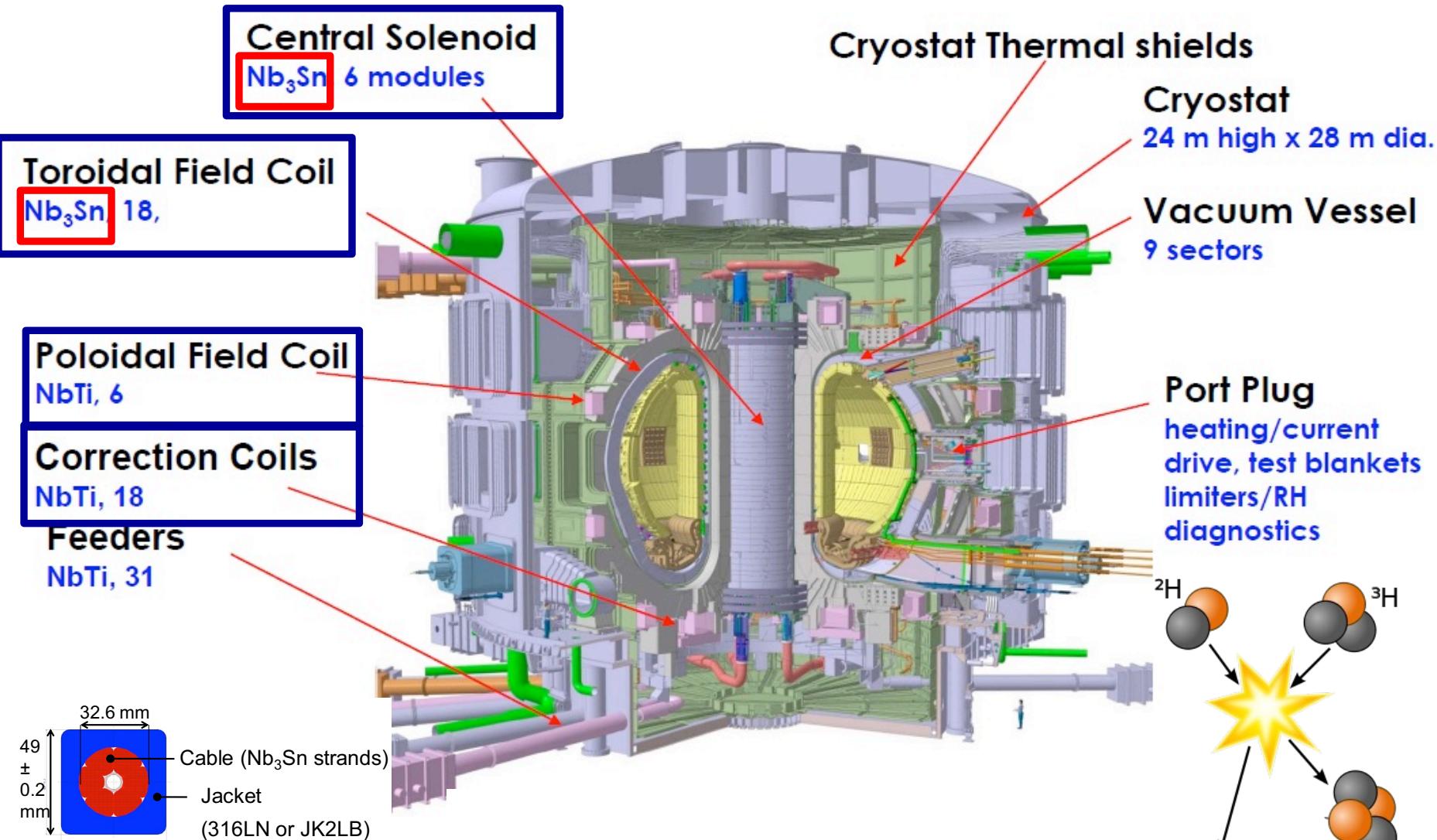
operatures





International Thermonuclear Experimental Reactor ITER

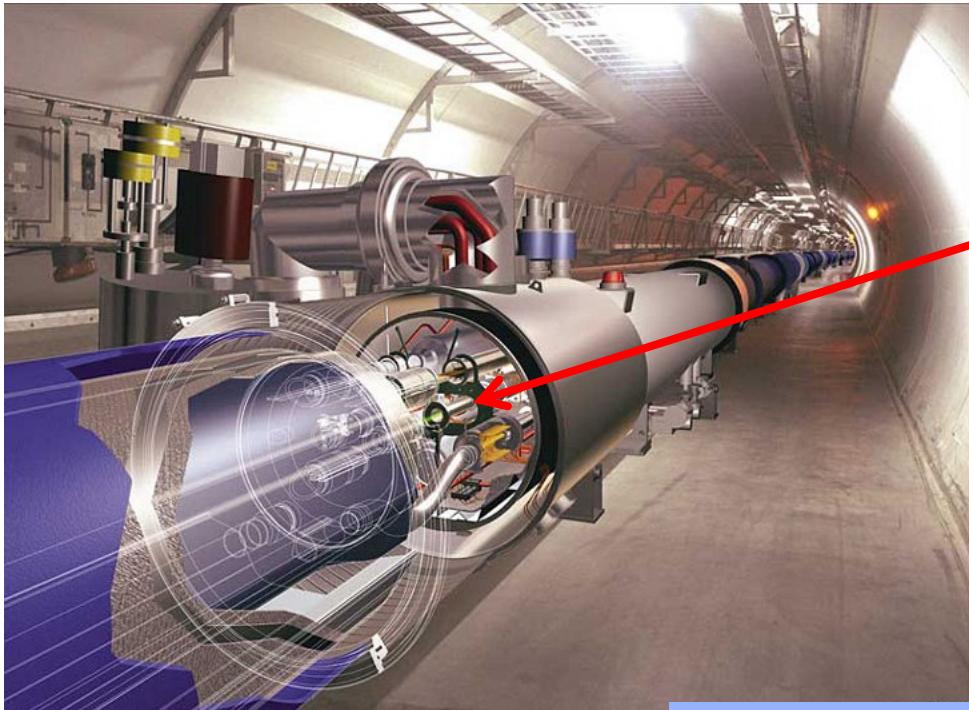
temperatures



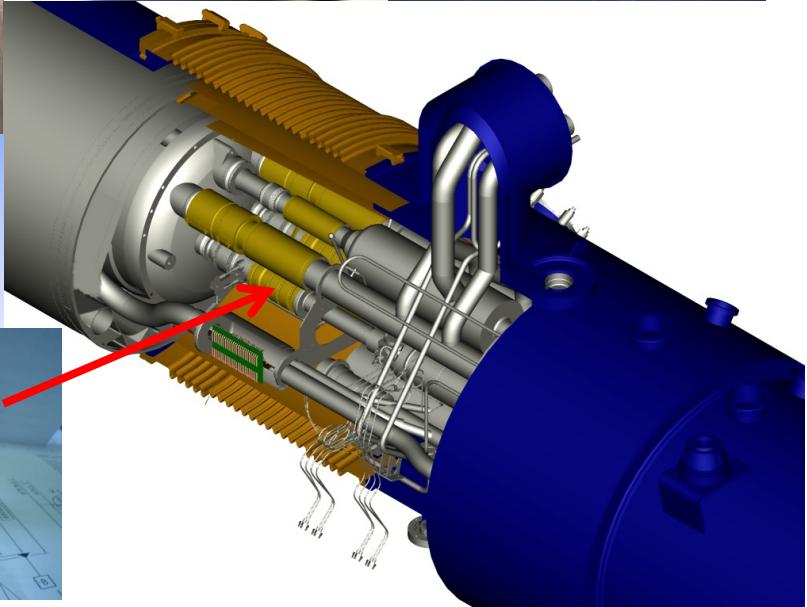


Motivation: CERN Large Hadron Collider

Superconductive phenomena at cryogenic temperatures

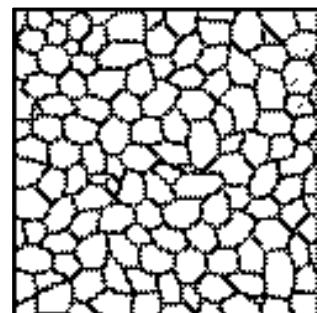
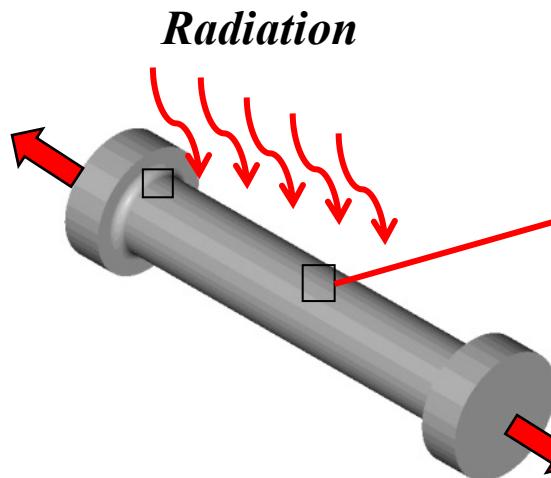


20000
expansion
bellows





Coupled field problems: radiation versus phase transformation

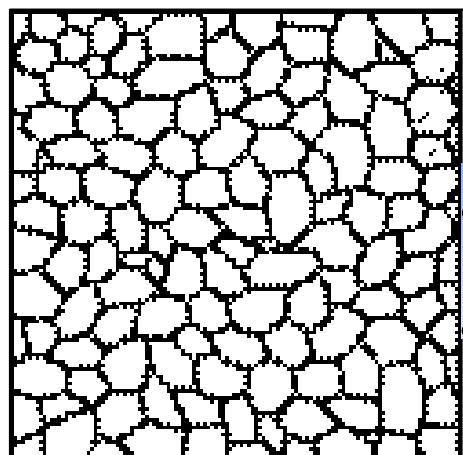


Scale: mezo

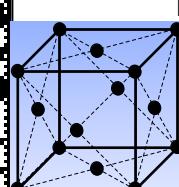


Scale: micro

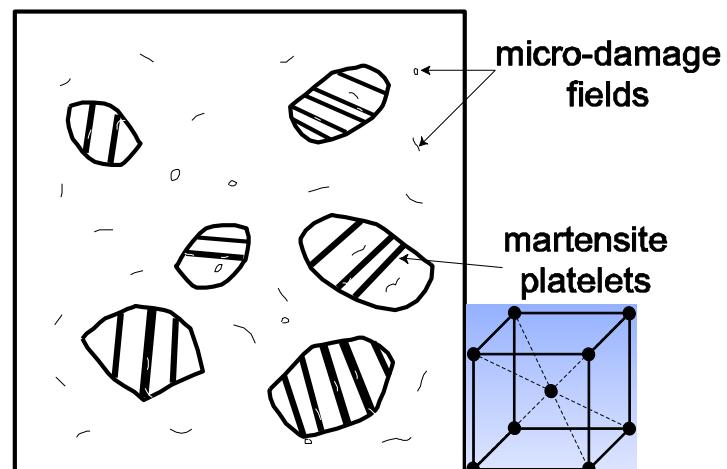
Single-phase continuum γ



Plastic strain

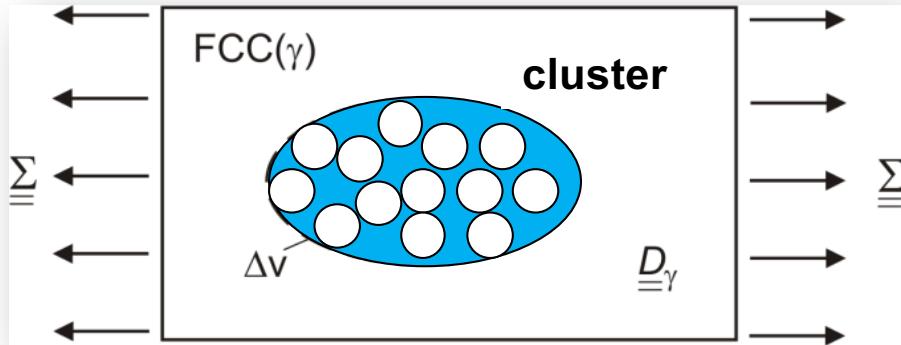


Two-phase continuum $\gamma + \alpha'$

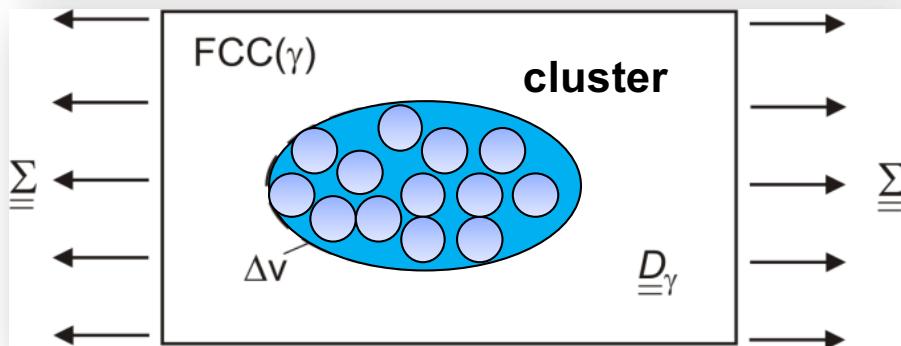




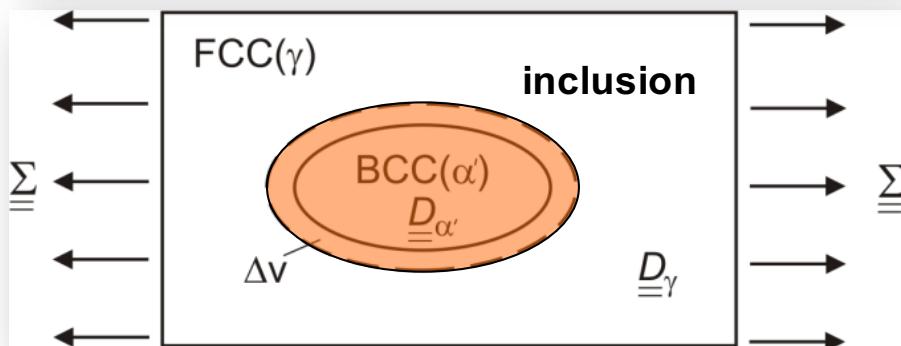
Coupled field problems: radiation versus phase transformation res



3D vacancy clusters: ξ_c



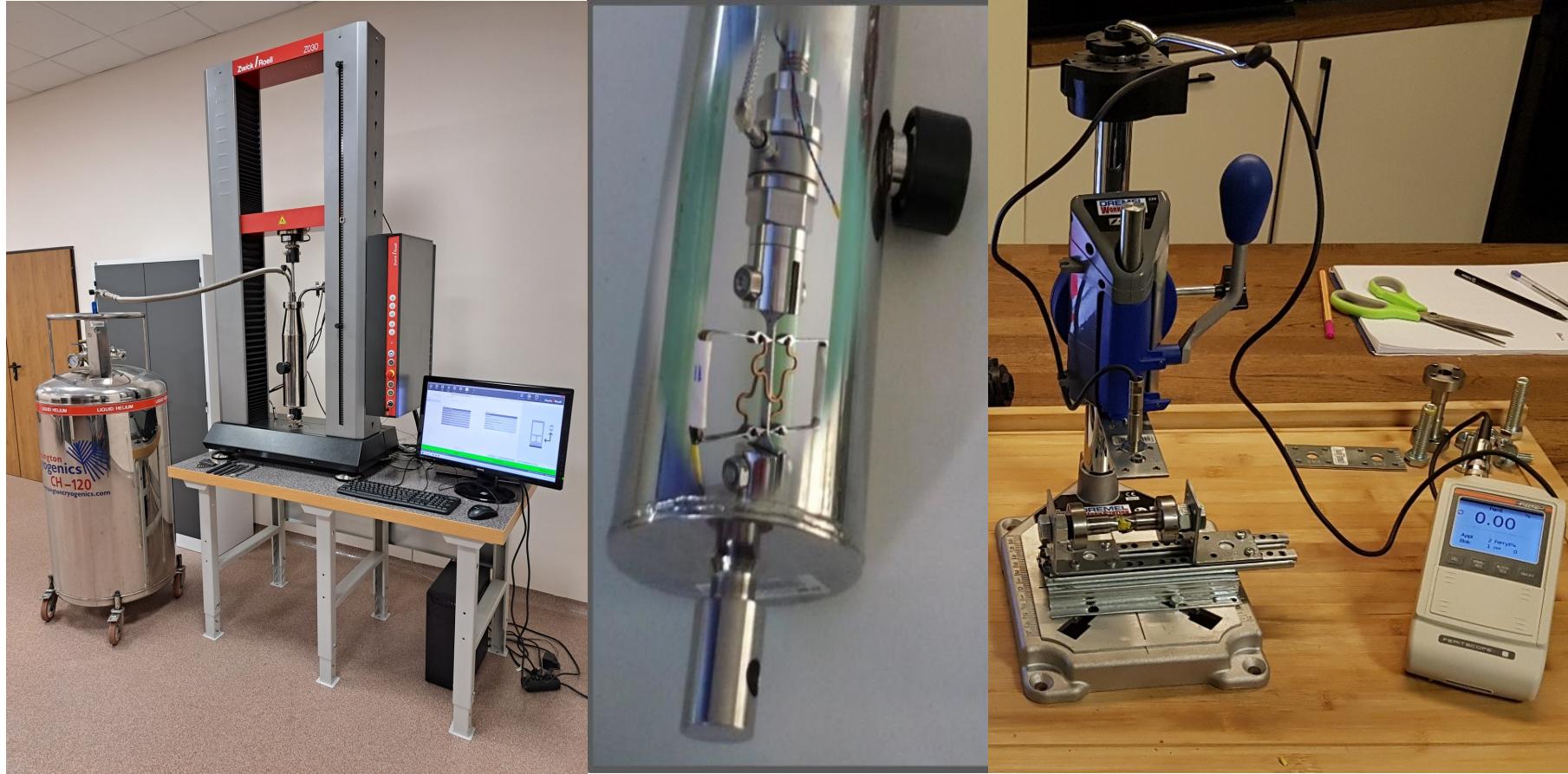
3D vacancy cluster with
impurities (He): ξ_c



Inclusions of secondary
phase: ξ_α



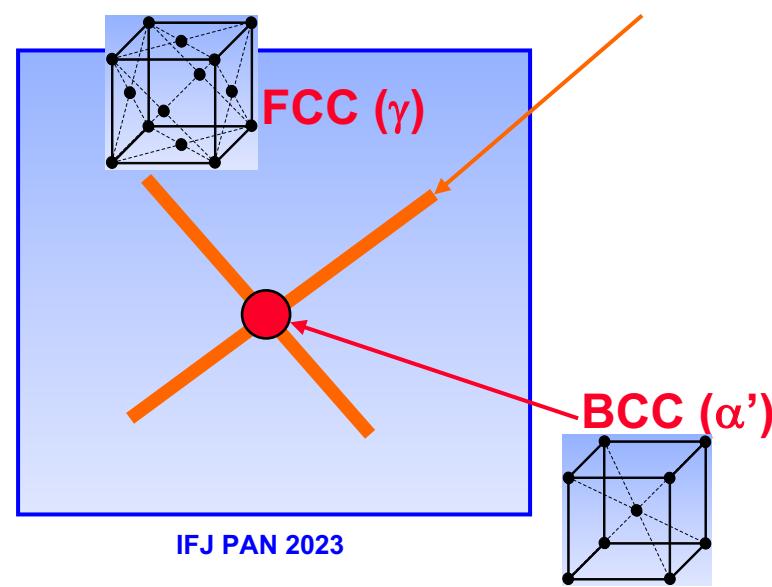
Experiments at extremely low temperatures (77 K, 4.2 K)



Dedicated cryogenic set-up for materials testing at extremely low temperatures (liquid nitrogen 77 K, liquid helium 4.2 K)



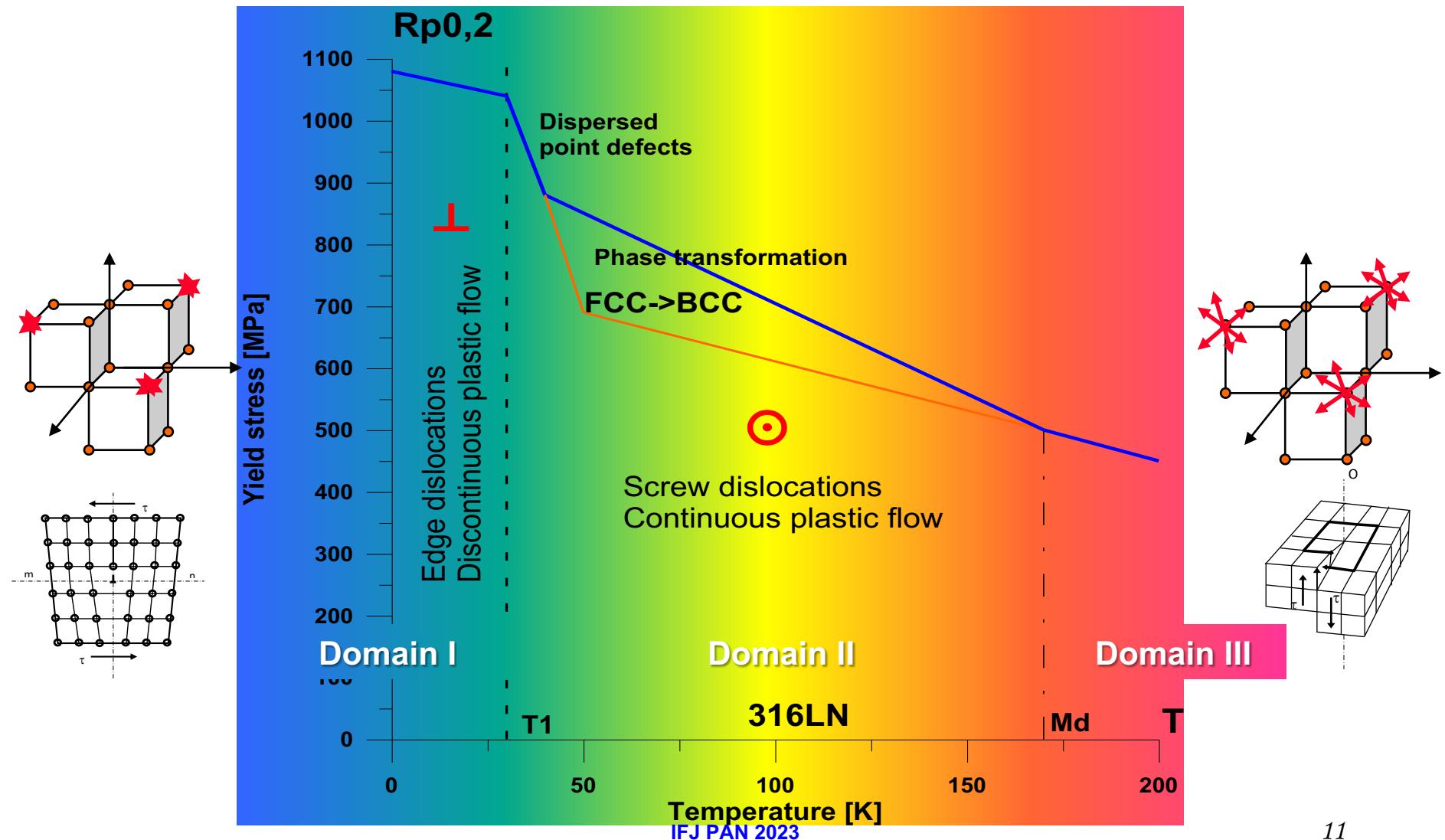
Plastic strain induced fcc-bcc phase transformation





Mechanisms of plastic flow at cryogenic temperatures

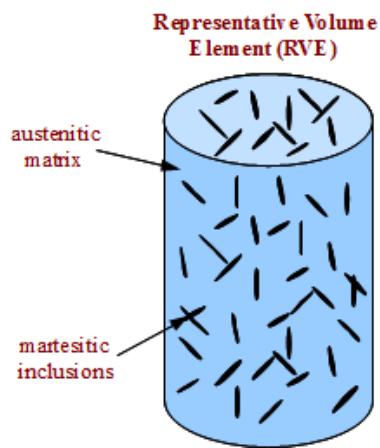
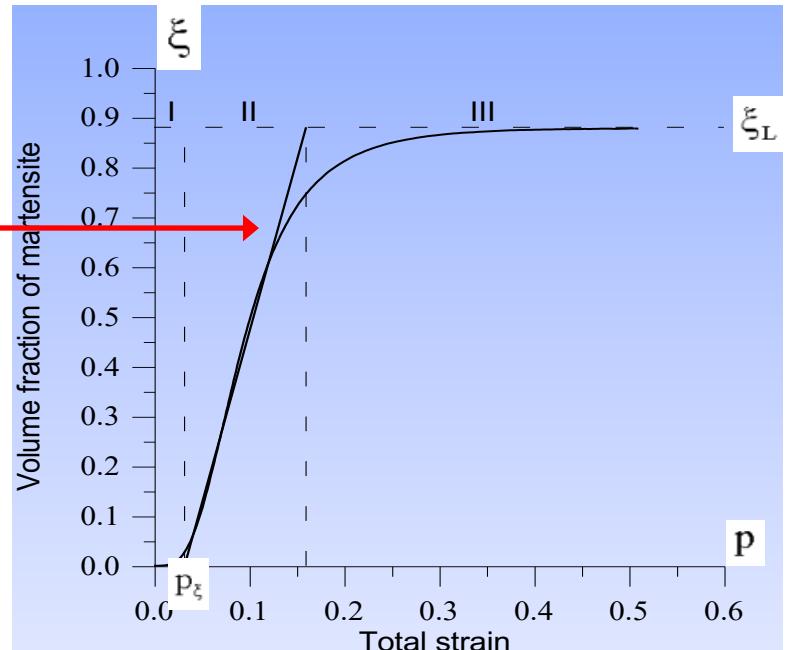
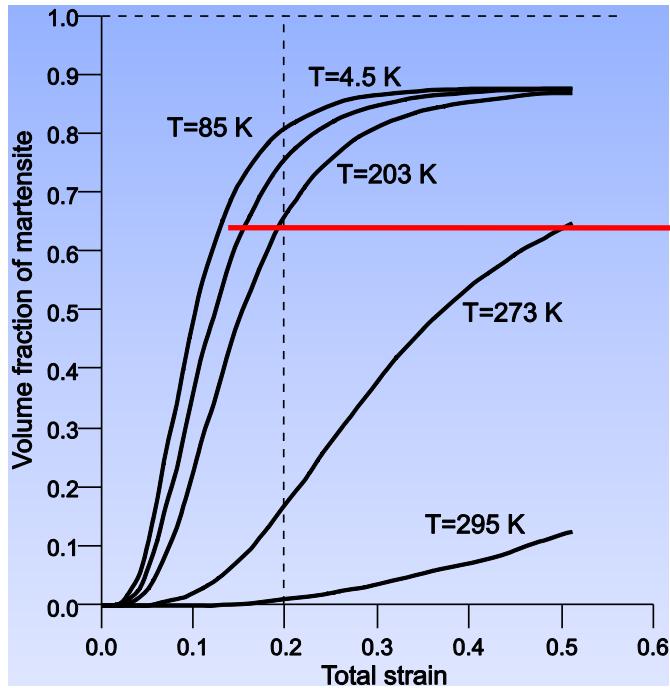
anic temperatures





Kinetics of Fcc-Bcc phase transformation

phenomena at cryogenic temperatures



$$\xi = \frac{dV_\xi}{dV} ; \quad 0 \leq \xi \leq 1$$

$$\dot{\xi} = A(T, \dot{\varepsilon}^p, \underline{\underline{\sigma}}) \dot{p} H((p - p_\xi)(\xi_L - \xi))$$

ξ – volume fraction of α' phase



Micromechanics: transformation strain

utive phenomena at cryogenic temperatures

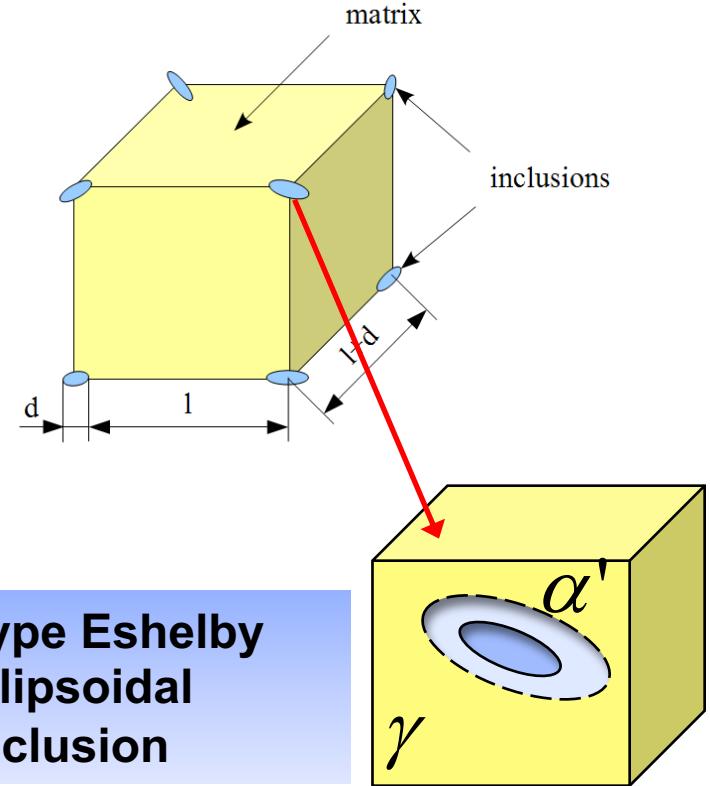
$$\underline{\underline{\varepsilon}}^{bs} = \frac{1}{V} \int_V \underline{\underline{\varepsilon}}_{\mu}^{bs} dV$$

$$\underline{\underline{\varepsilon}}^{bs} = \frac{V_\gamma}{V} \frac{1}{V_\gamma} \int_{V_\gamma} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV + \frac{V_\alpha}{V} \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV$$

$$\underline{\underline{\varepsilon}}^{bs} = \frac{V_\alpha}{V} \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV = \xi \frac{1}{V_\alpha} \int_{V_\alpha} \underline{\underline{\varepsilon}}_{\mu}^{bs} dV = \xi \left\langle \underline{\underline{\varepsilon}}_{\mu}^{bs} \right\rangle$$

$$\underline{\underline{\varepsilon}}_{\mu}^{bs} = \begin{pmatrix} 0 & 0 & \frac{\gamma}{2} \\ 0 & 0 & 0 \\ \frac{\gamma}{2} & 0 & \Delta\nu \end{pmatrix}_{(\vec{x}, \vec{y}, \vec{z})}$$

$$\left\langle \underline{\underline{\varepsilon}}_{\mu}^{bs} \right\rangle = \frac{1}{3} \Delta\nu I_3$$



Type Eshelby
ellipsoidal
inclusion

$$\underline{\underline{\sigma}} = E : \left(\underline{\underline{\varepsilon}} - \underline{\underline{\varepsilon}}^p - \underline{\underline{\varepsilon}}^{th} - \xi \left\langle \underline{\underline{\varepsilon}}_{\mu}^{bs} \right\rangle \right)$$

$$\underline{\underline{\varepsilon}}^{bs} = \xi \frac{1}{3} \Delta\nu I_3$$



Constitutive description of two-phase continuum

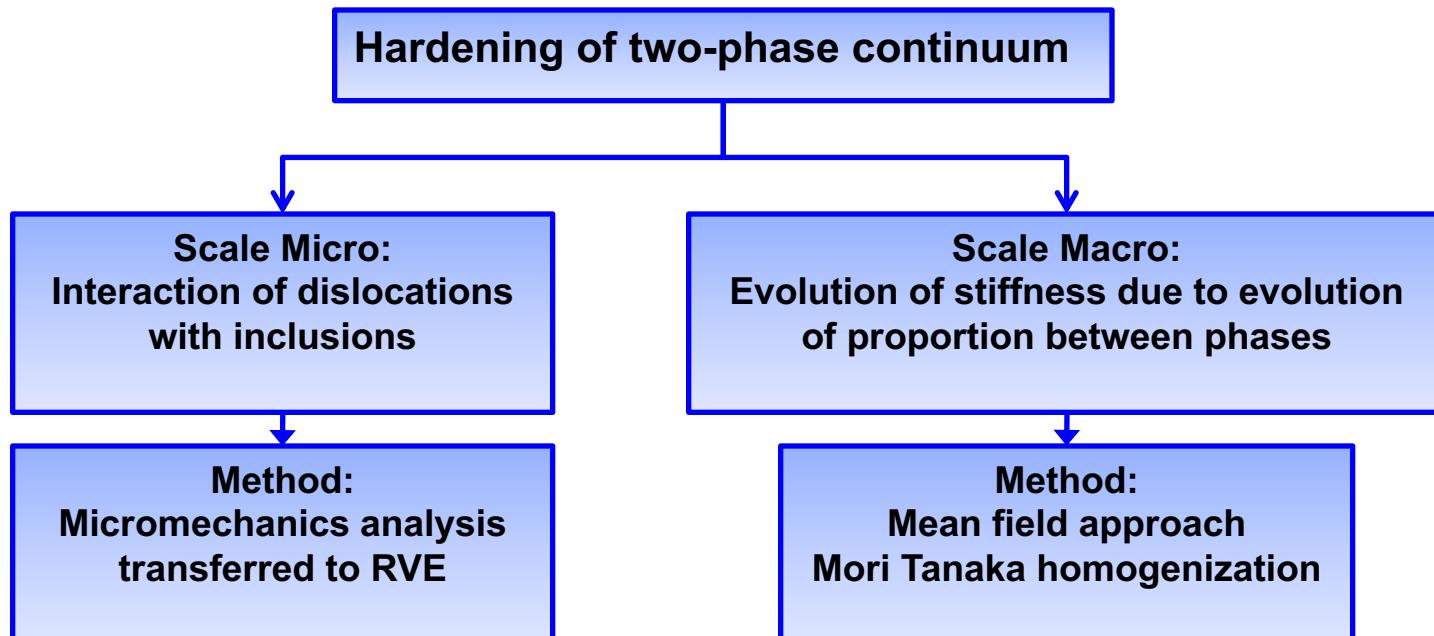
a at cryogenic temperatures

Yield condition:

$$f_c(\underline{\underline{\sigma}}, \underline{\underline{X}}, R) = J_2(\underline{\underline{\sigma}} - \underline{\underline{X}}) - \sigma_y - R = 0 \quad J_2(\underline{\underline{\sigma}} - \underline{\underline{X}}) = \sqrt{\frac{3}{2}(\underline{s} - \underline{\underline{X}}) : (\underline{s} - \underline{\underline{X}})}$$

Mixed hardening depending on the phase transformation parameter:

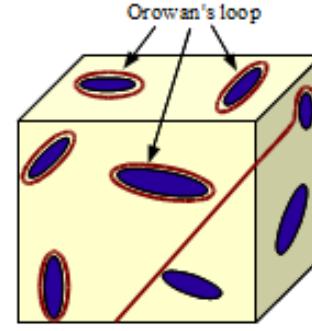
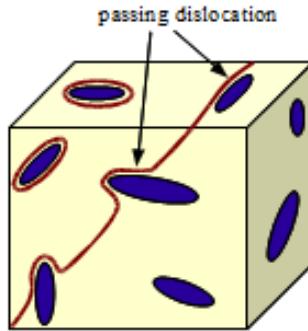
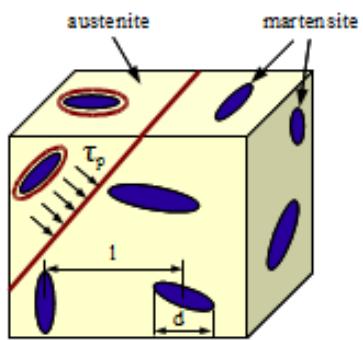
$$\begin{aligned} d\underline{\underline{X}} &= d\underline{\underline{X}}_a + d\underline{\underline{X}}_{a+m} = \frac{2}{3}C_X(\xi)d\underline{\underline{\varepsilon}}^p \\ dR &= C_R(\xi)dp \end{aligned}$$





Scale Micro: interaction of dislocations with inclusions

kinetic temperatures

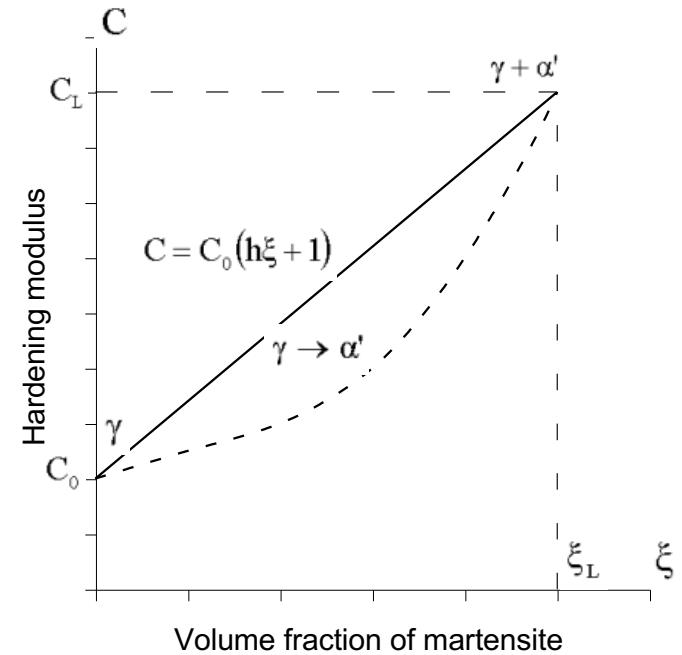


$$dX_a = \frac{2}{3} C_0 d\varepsilon^p$$

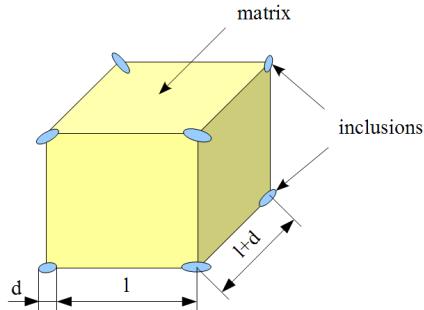
**Initial hardening of matrix
(austenite)**

$$dX_a = \frac{2}{3} C_0 \phi(\xi) d\varepsilon^p$$

**Hardening of matrix
containing inclusions**



Micromechanics analysis



$$\tau_p = \frac{Gb}{d} \left(\frac{6\xi_0}{\pi} \right)^{\frac{1}{3}} \left(1 + \frac{\xi - \xi_0}{3\xi_0} \right)$$



$$\phi(\xi) = 1 + h\xi ; \quad 0 \leq \xi \leq 1$$

$$C = C_0 \phi(\xi)$$



Scale Macro: evolution of proportion between phases

ogenic temperatures

Elastic-plastic matrix:

$$\Delta \underline{\underline{\sigma}}_a = \underline{\underline{E}}_{ta} : \Delta \underline{\underline{\varepsilon}}$$

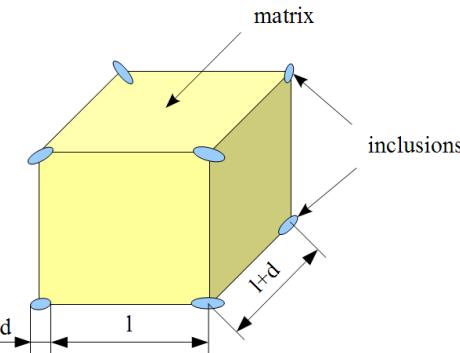
$$\underline{\underline{E}}_{ta} = 3k_a \underline{\underline{J}} + 2\mu_a \underline{\underline{K}} - 2\mu_a \frac{\underline{\underline{n}} \otimes \underline{\underline{n}}}{1 + \frac{C(\xi)}{3\mu_a}}$$

„Linearization”: extraction of isotropic part of tangent stiffness operator

$$\underline{\underline{E}}_{ta} = 3k_{ta} \underline{\underline{J}} + 2\mu_{ta} \underline{\underline{K}}$$

$$\mu_{ta} = \frac{E_t}{2(1+\nu)} \quad k_{ta} = \frac{E_t}{3(1-2\nu)}$$

$$E_t = \frac{EC}{E+C}$$



Elastic inclusions:

$$\Delta \underline{\underline{\sigma}}_m = \underline{\underline{E}}_m : \Delta \underline{\underline{\varepsilon}}$$

$$\underline{\underline{E}}_m = 3k_m \underline{\underline{J}} + 2\mu_m \underline{\underline{K}}$$

$$\mu_m = \frac{E}{2(1+\nu)} \quad k_m = \frac{E}{3(1-2\nu)}$$



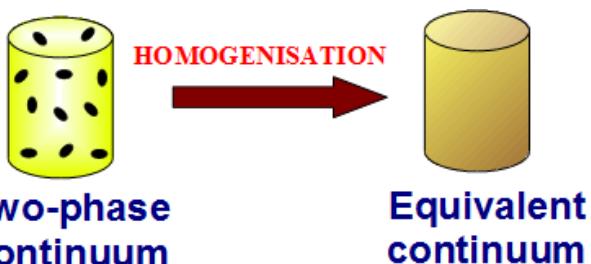
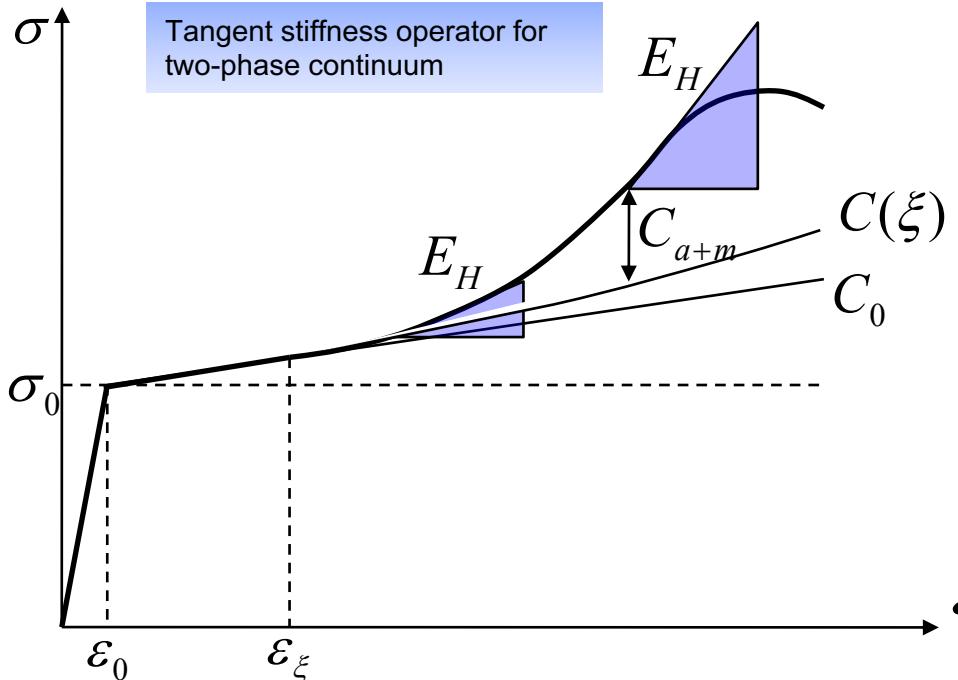
Homogenization:

$$\Delta \underline{\underline{\sigma}} = \underline{\underline{E}}_H : \Delta \underline{\underline{\varepsilon}}$$



Scale Macro: evolution of proportion between phases

at homogenetic temperatures



$$\underline{\Delta \sigma} = \underline{\underline{E}}_{ta} : \underline{\underline{\Delta \varepsilon}}$$

$$\underline{\Delta \sigma} = \underline{\underline{E}}_m : \underline{\underline{\Delta \varepsilon}}$$

$$\underline{\Delta \sigma} = \underline{\underline{E}}_H : \underline{\underline{\Delta \varepsilon}}$$

$$\underline{\underline{E}}_H = \underline{\underline{E}}_{MT} = 3\underline{k}_{MT} \underline{\underline{J}} + 2\underline{\mu}_{MT} \underline{\underline{K}}$$

$$\left[\underline{\underline{E}}_{MT} + \underline{\underline{E}}^* \right]^{-1} = \sum_{i=a,m} f_i \left[\underline{\underline{E}}_i + \underline{\underline{E}}^* \right]^{-1}$$

$$\underline{\underline{E}}_{MT} = \left[(1-\xi) \left(\underline{\underline{E}}_a + \underline{\underline{E}}^* \right)^{-1} + \xi \left(\underline{\underline{E}}_m + \underline{\underline{E}}^* \right)^{-1} \right]^{-1} - \underline{\underline{E}}^*$$

$$3k_{MT} + 3k^* = \left[\frac{1-\xi}{3(k_{ta} + k^*)} + \frac{\xi}{3(k_m + k^*)} \right]^{-1} \quad k^* = \frac{4}{3} \mu_{ta}$$

$$2\mu_{MT} + 2\mu^* = \left[\frac{1-\xi}{2(\mu_{ta} + \mu^*)} + \frac{\xi}{2(\mu_m + \mu^*)} \right]^{-1} \quad 2\mu^* = \frac{\mu_{ta}(9k_{ta} + 8\mu_{ta})}{3(k_{ta} + 2\mu_{ta})}$$



Constitutive description of two-phase continuum

a at cryogenic temperatures

Kinematic hardening

$$\beta = \frac{\sigma' + \sigma'^-}{2(\sigma' - \sigma_0)}$$

Isotropic hardening

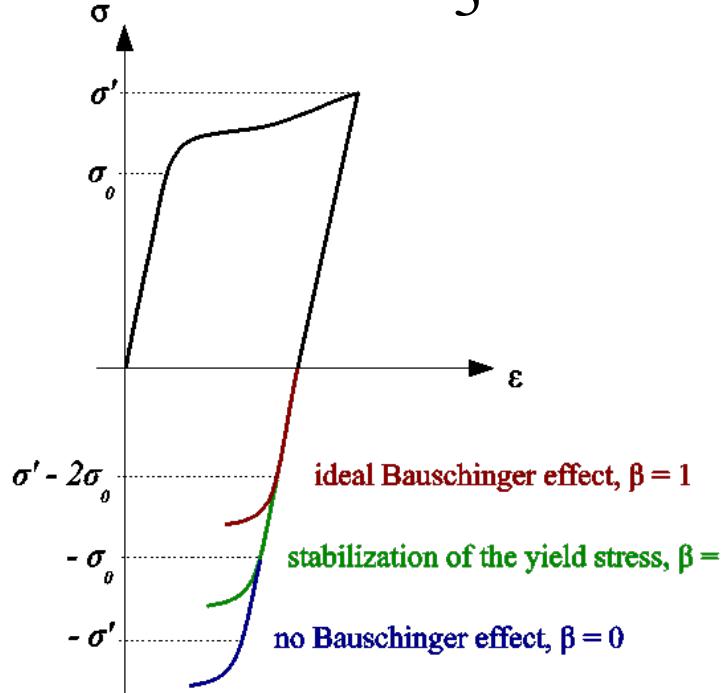
$$\Delta X_{\underline{\underline{a+m}}} = \Delta \sigma_{\underline{\underline{a+m}}}$$

Parametization: *Życzkowski, 1981*

$$\Delta R = \|\Delta \sigma_{a+m}\|$$

$$dX_{\underline{\underline{a+m}}} = \frac{2}{3} \beta C_{a+m}(\xi) d\varepsilon^p$$

$$dR = (1 - \beta) C_{a+m}(\xi) dp$$



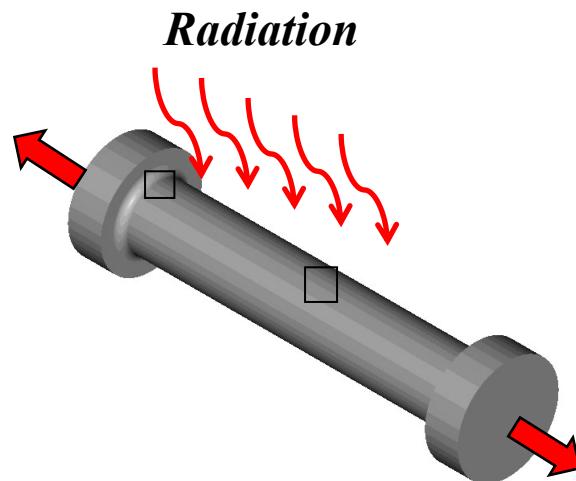
Evolution laws of hardening parameters

$$dR = C_R(\xi) dp = (1 - \beta) C_{a+m}(\xi) dp$$

$$dX_{\underline{\underline{a+m}}} = \frac{2}{3} C_X(\xi) d\varepsilon^p = \frac{2}{3} \beta C_{a+m}(\xi) d\varepsilon^p$$



Radiation induced damage





Research program

Coupled dissipative phenomena at cryogenic temperatures

Experiments including proton and neutron irradiated samples subjected to loading/unloading technique



Building well calibrated multi-scale 3D constitutive models of damage/porosity evolution in the framework of CDM



Combining CDM with fracture mechanics in order to predict transition from critical damage/porosity to fracture



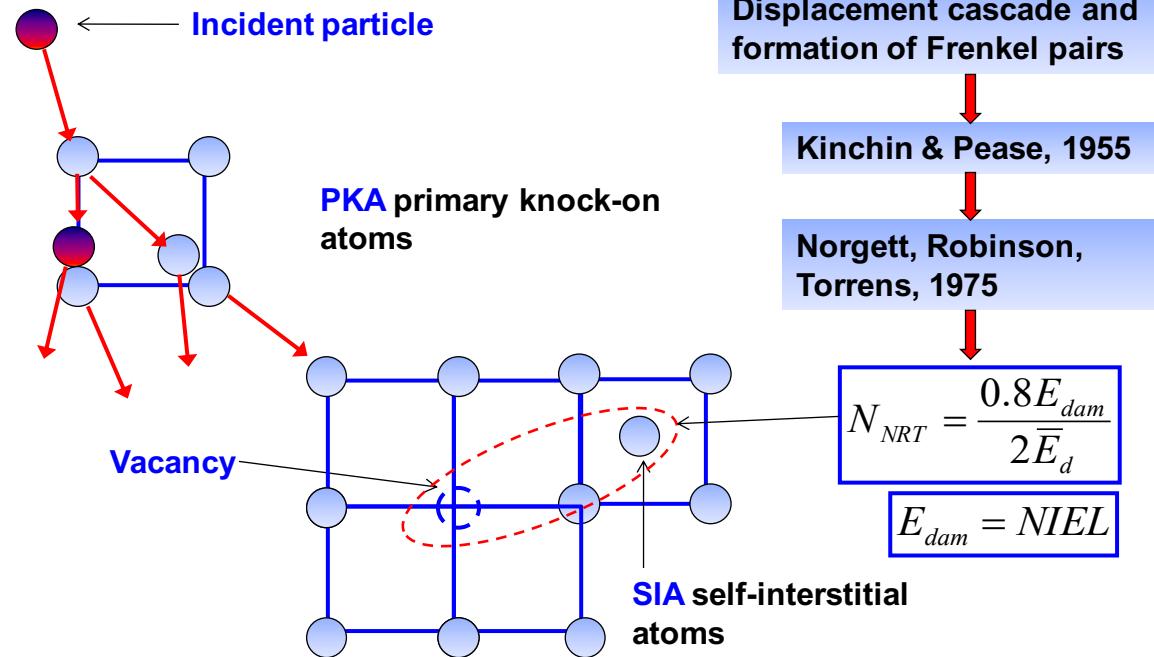
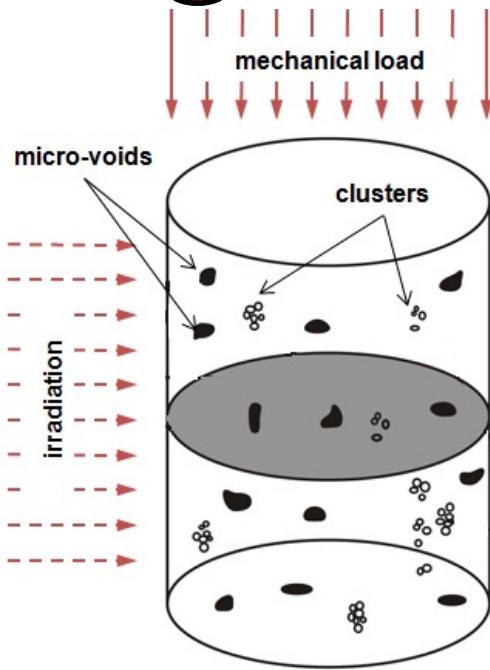
Computing evolution of nano/micro damage fields and macro-crack propagation in the irradiated components

Lifetime prediction



Radiation induced micro-damage

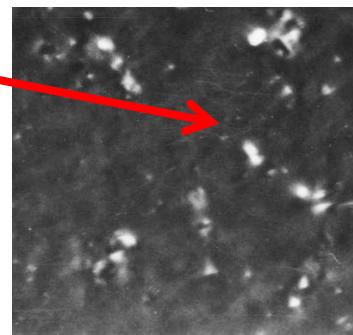
coupled dissipative phenomena at cryogenic temperatures



Transverse View

Cluster
of micro-voids

SRIM
(Monte-Carlo
method)



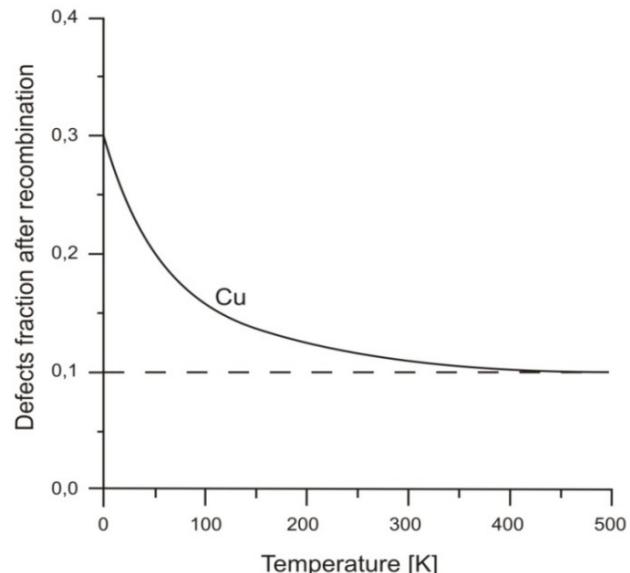
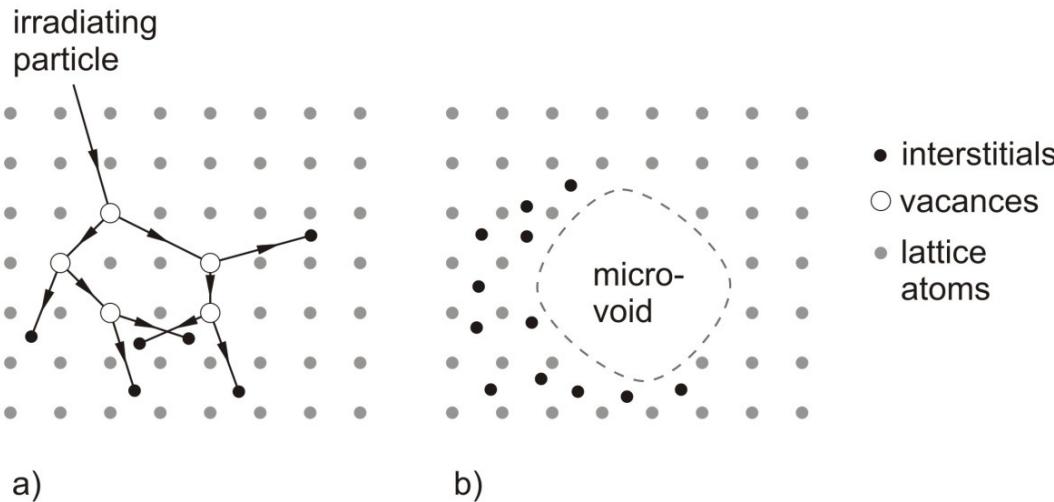
- Defects due to irradiation:
1. SFT – stacking fault tetrahedron
 2. Faulted or perfect dislocation loops
 3. Voids – 3D vacancy clusters
 4. Cavities – 3D vacancy clusters with impurities (He)



Measure of irradiation induced damage

Five phenomena at cryogenic temperatures

1 displacement per atom (dpa):
corresponds to stable displacement from their lattice site of all atoms in the material during irradiation near absolute zero (no thermally-activated point defect diffusion)

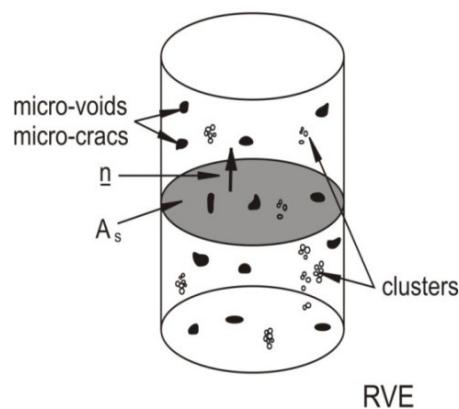
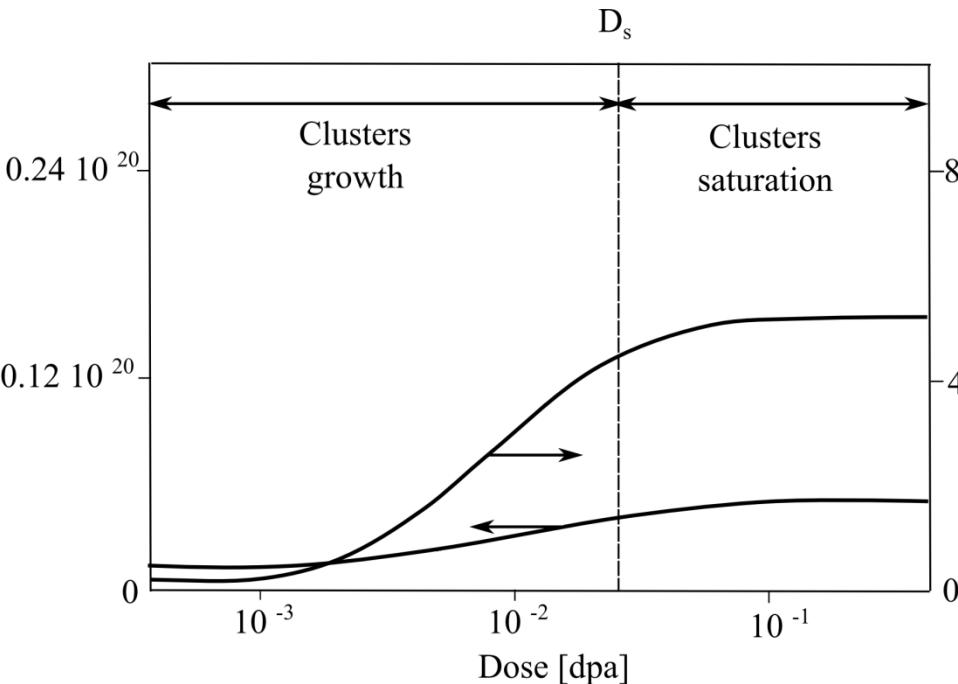




Lattice defects after irradiation

Coupled dissipative phenomena at cryogenic temperatures

Clusters density [cm⁻³]



q_c – number of clusters per unit volume

r_c – average radius of clusters

$$D = \frac{dS_D}{dS} ; \quad 0 \leq D \leq 1$$

$$\xi = \frac{dV_D}{dV} ; \quad 0 \leq \xi \leq 1$$

Clusters size [atoms number/cluster]

Physics

$$q_c = \begin{cases} C_{qI}(dpa)^{n_{qI}} & \text{for } dpa < D_S \\ C_{qII}(dpa)^{n_{qII}} & \text{for } dpa \geq D_S \end{cases}$$

$$r_c = \begin{cases} C_r(dpa)^{n_r} & \text{for } dpa < D_S \\ r_{cr} & \text{for } dpa \geq D_S \end{cases}$$

$$q_A = \left(\sqrt[3]{q_V} \right)^2 = q_c^{2/3}$$

$$D_{r0} = q_A \pi r_{c0}^2$$

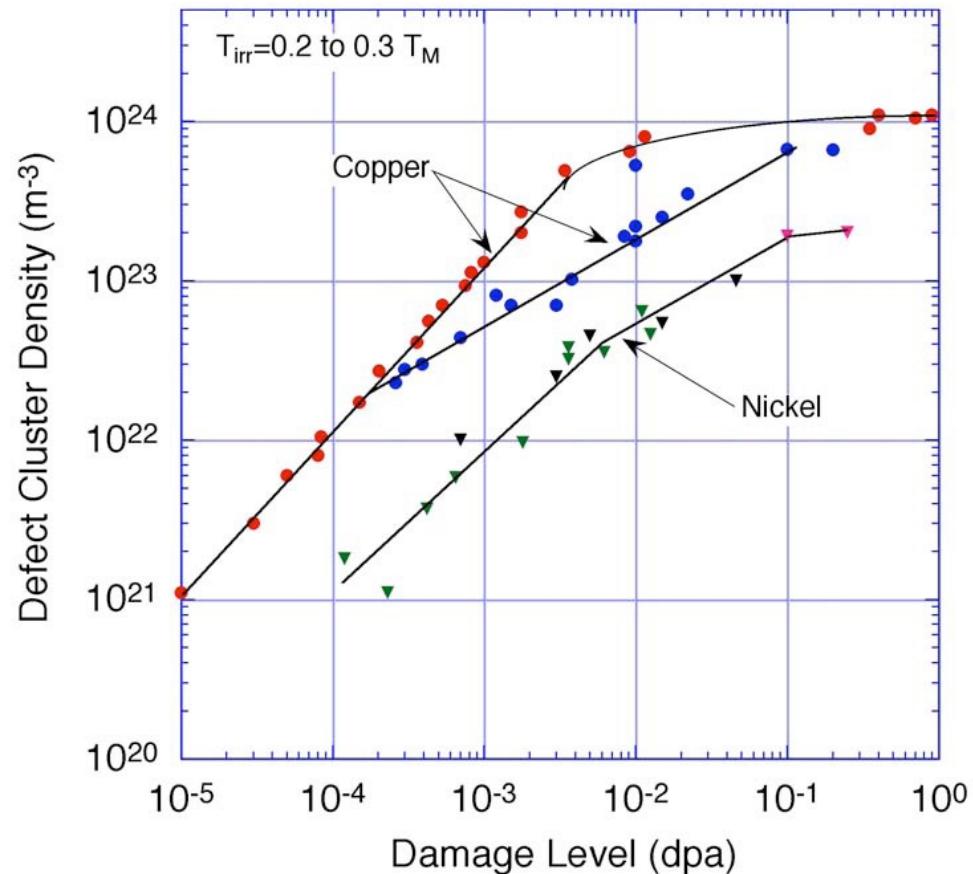
Mechanics



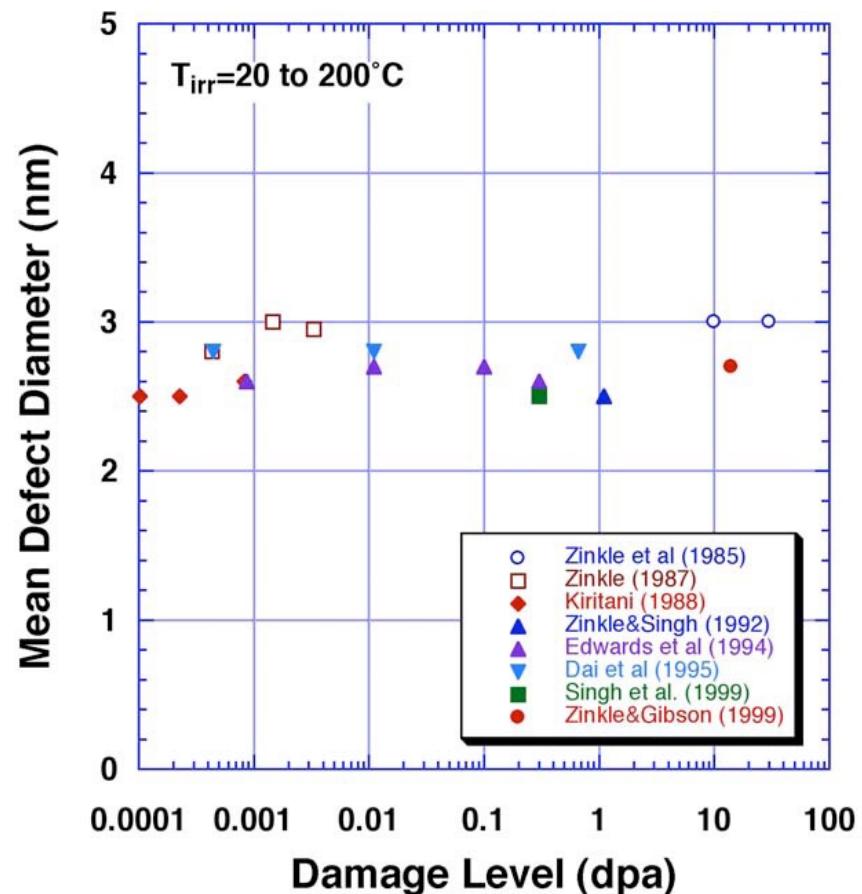
Irradiated metals and alloys: Nickel and Copper

phenomena at cryogenic temperatures

COMPARISON OF DEFECT CLUSTER ACCUMULATION
IN NEUTRON-IRRADIATED NICKEL AND COPPER



Measured Average Image Width of Defect Clusters in Neutron and Ion-Irradiated Copper



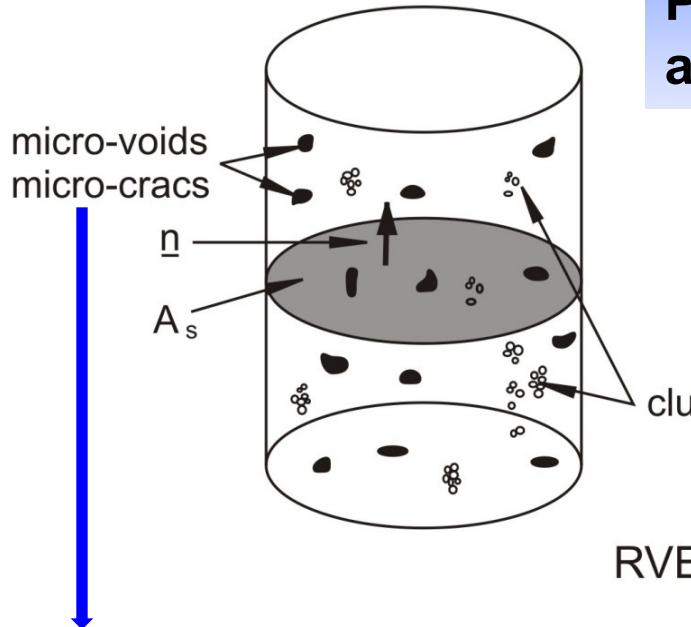
Source: S.J. Zinkle „Microstructure evolution in irradiated metals and alloys: fundamental aspects”, Italy, 2004.

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Radiation and mechanical damage: additive formulation

at temperatures



Postulate: both micro-damage components are treated in additive way

$$D_r = D_{r0} + \int_0^{\hat{p}} dD_{rm}$$

radiation induced damage

$$\underline{\underline{\underline{D}}}_r = \frac{D_r}{3} \underline{\underline{\underline{I}}} =$$

$\underline{\underline{\underline{I}}}$ - identity tensor

isotropic

$$d\bar{D}_m = \underline{\underline{\underline{C}}}\underline{\underline{Y}}\underline{\underline{C}}^T dp$$

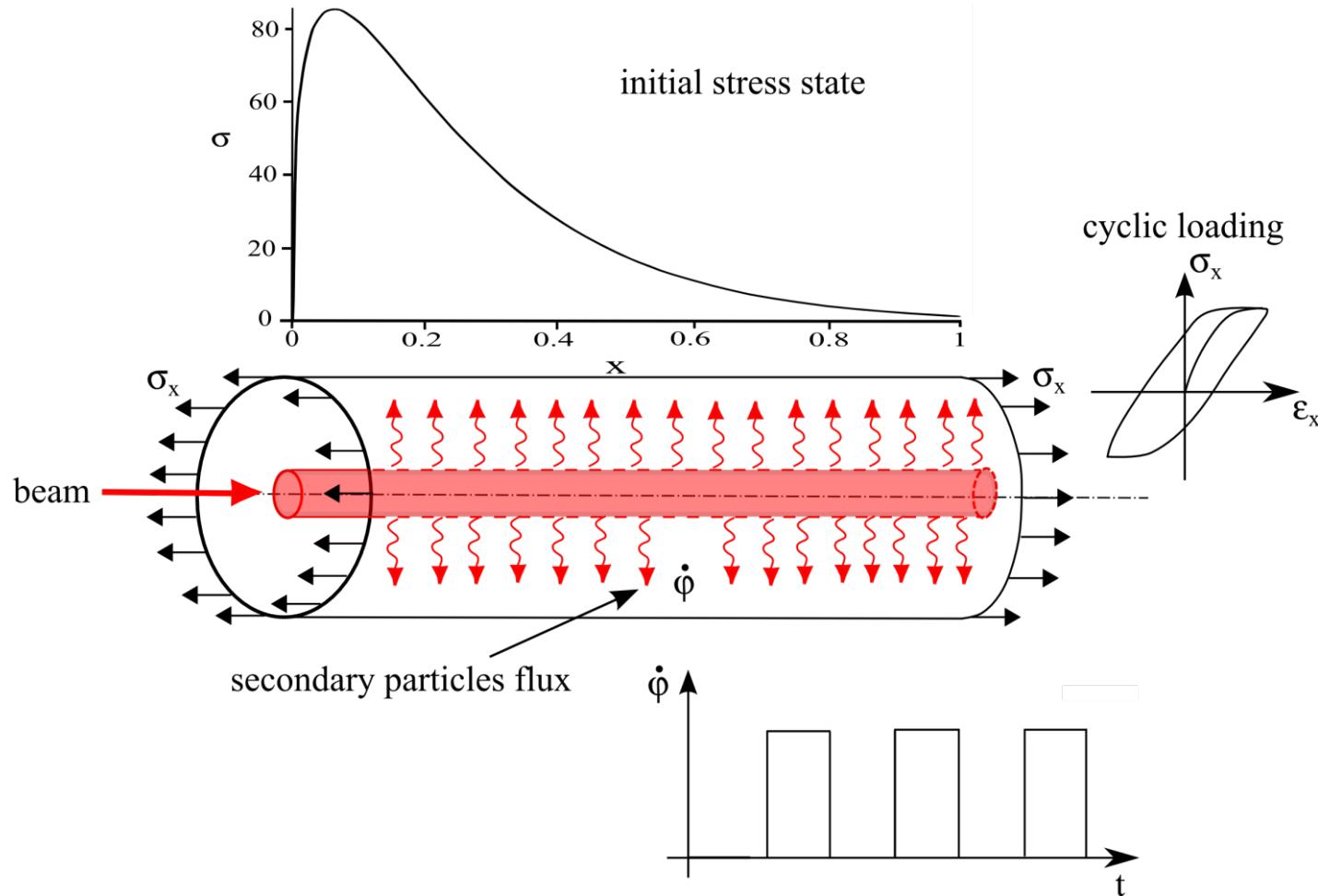
anisotropic

$$\underline{\underline{\underline{D}}} = \underline{\underline{\underline{D}}}_m + \underline{\underline{\underline{D}}}_r = \underline{\underline{\underline{D}}}_m + \frac{1}{3} D_r \underline{\underline{\underline{I}}}$$



Lifetime estimation for irradiated components

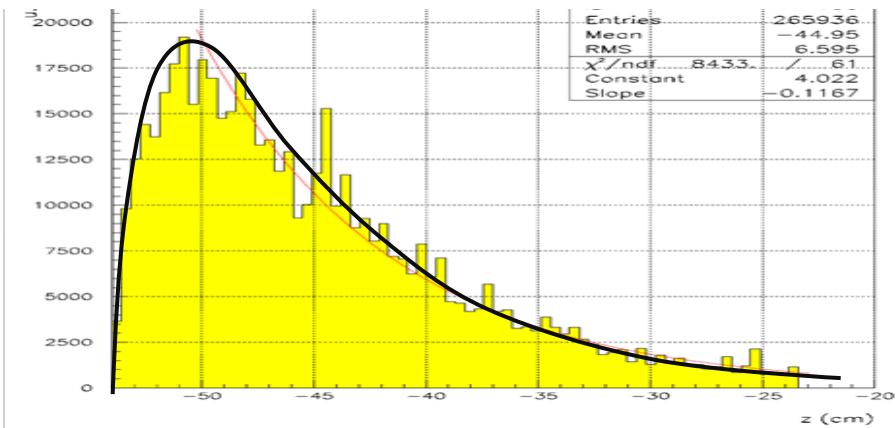
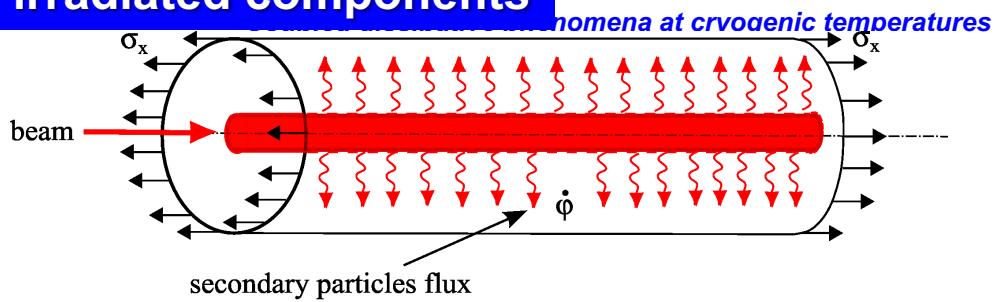
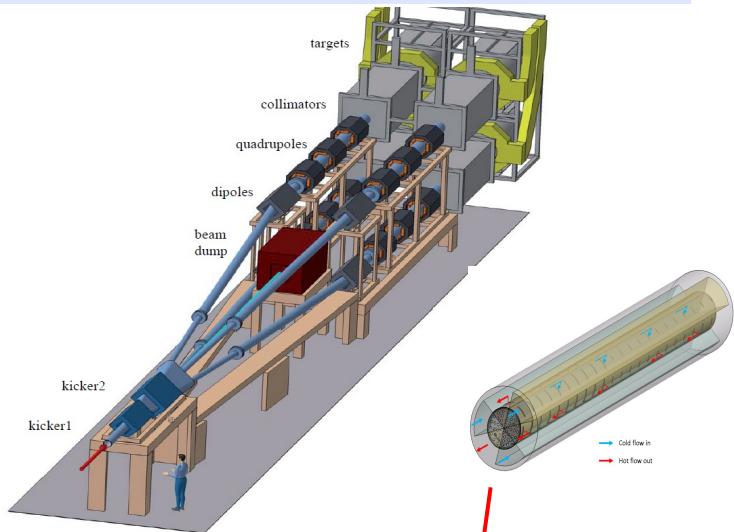
phenomena at cryogenic temperatures



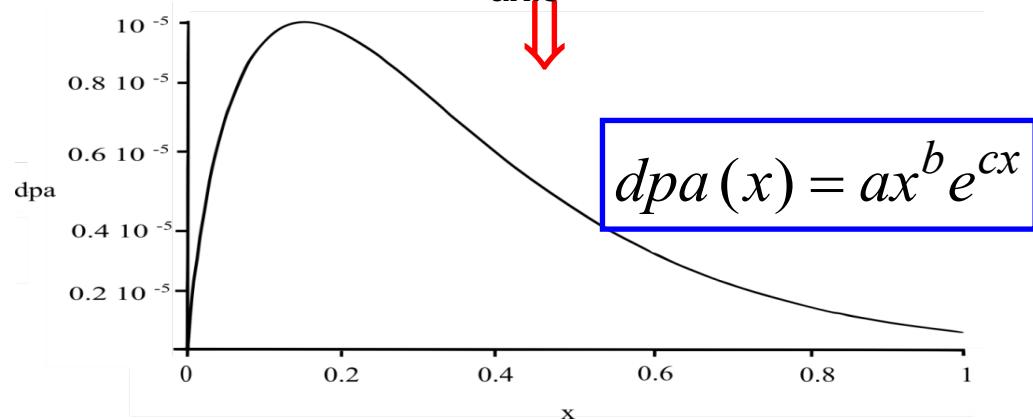
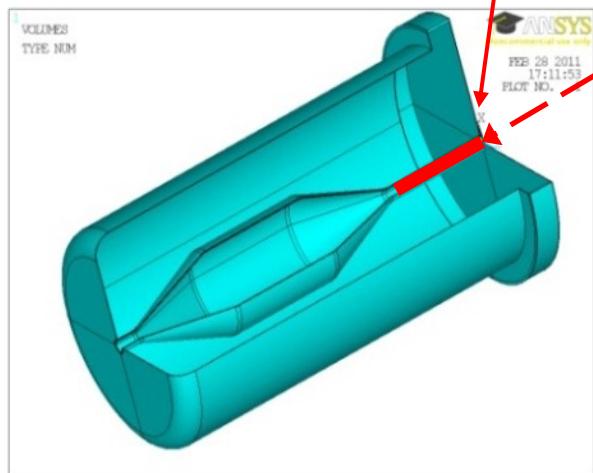


Lifetime estimation for irradiated components

Secondary particles flux: γ ,
 n , p^+ , π^\pm and e^\pm



Typical distribution of particle flux along the target axis





Kinetics of evolution of radiation induced damage (clusters of voids) under mechanical loads

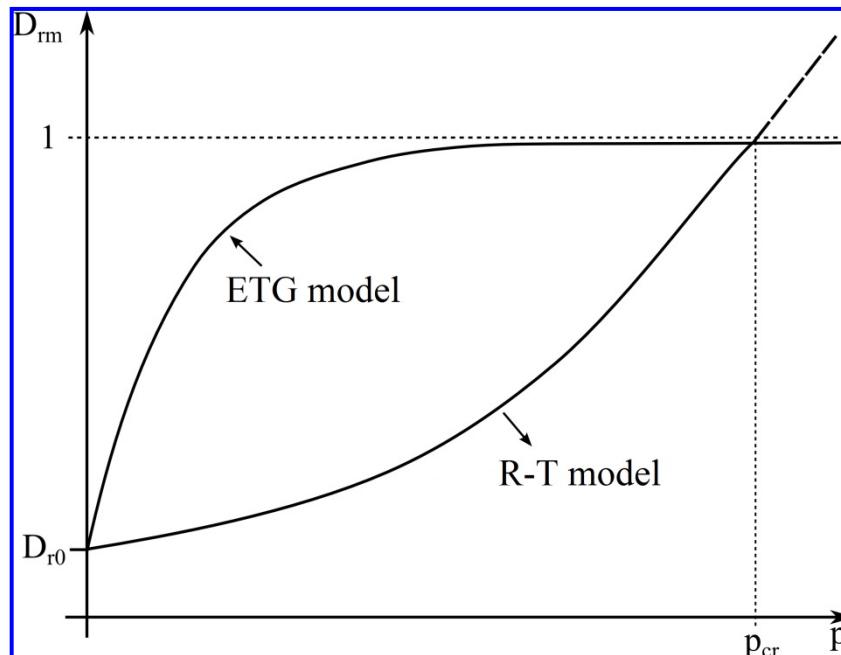
Rice&Tracey (R-T) model:

$$dr_c = r_c \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right) dp$$

Gurson (ETG) model:

$$d\xi = (1 - \xi)dp$$

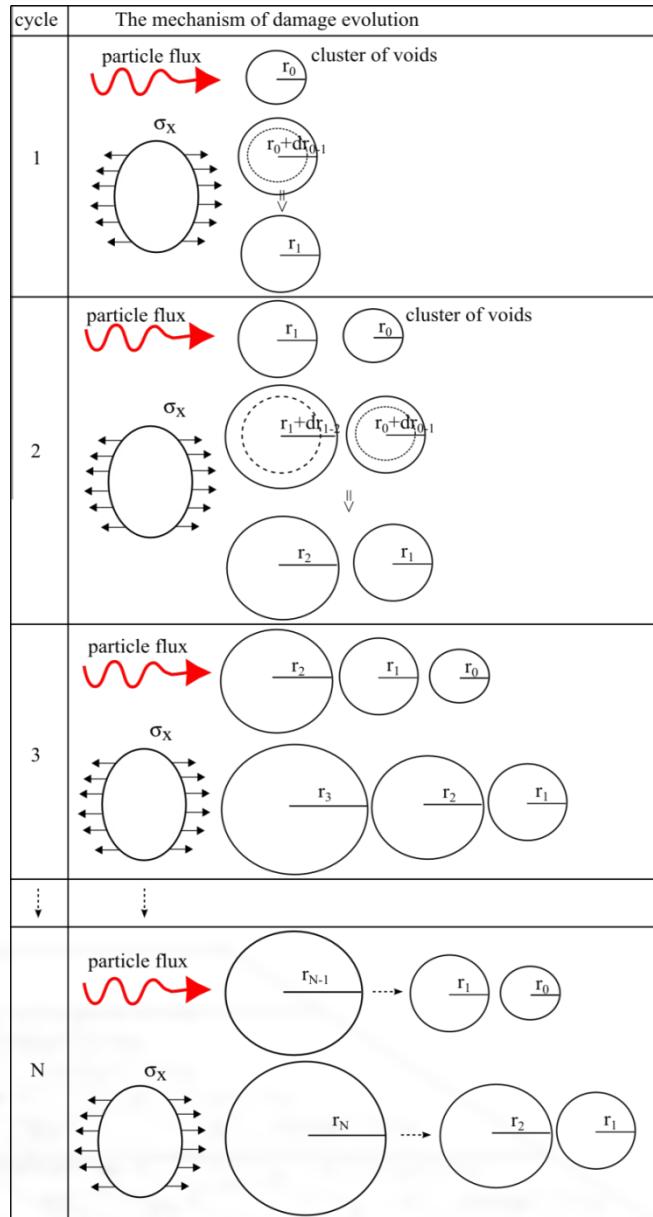
$$\dot{p} = \sqrt{\frac{2}{3} \dot{\varepsilon}_{ij}^p \dot{\varepsilon}_{ij}^p}$$





Lifetime estimation for irradiated components

Coupled dissipative phenomena at cryogenic temperatures



Rice & Tracey law

$$\int_{D_i}^{D_{i+1}} dD = q_A 2\pi \int_{r_i}^{r_{i+1}} r dr$$



$$\Delta D_{i \rightarrow i+1} = q_A \pi (r_{i+1}^2 - r_i^2)$$

$$\int_{r_i}^{r_{i+1}} \frac{dr_c}{r_c} = \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right) \int_0^{\tilde{p}} dp$$



$$r_{i+1} = r_i e^{A \tilde{p}} \quad A := \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_{eq}}\right)$$



$$\Delta D_{i \rightarrow i+1} = q_A \pi r_i^2 (e^{2A \tilde{p}} - 1)$$



Lifetime estimation for irradiated components

Phenomena at cryogenic temperatures

$$D_{r0} = q_A \pi r_{c0}^2$$

$$D_{rm1} = D_{r0} + \Delta D_{rm(0 \rightarrow 1)} = D_{r0} + q_A \pi r_{c0}^2 (e^{2A\tilde{p}} - 1)$$

$$D_{rm2} = D_{rm1} + \Delta D_{rm(1 \rightarrow 2)} + D_{r0} + \Delta D_{rm(0 \rightarrow 1)} = 2D_{r0} + q_A \pi r_{c0}^2 e^{2A\tilde{p}} + q_A \pi r_{c0}^2 e^{4A\tilde{p}} - 2q_A \pi r_{c0}^2$$

⋮

$$D_{rm_{i+1}} = D_{rm_i} + D_{r0} + \Delta D_{rm(i \rightarrow i+1)} + \Delta D_{rm(i-1 \rightarrow i)} + \dots + \Delta D_{rm(0 \rightarrow 1)}$$

$$D_{rmN} = \underbrace{q_A \pi r_{c0}^2 e^{2A\tilde{p}} + q_A \pi r_{c0}^2 e^{4A\tilde{p}} + q_A \pi r_{c0}^2 e^{6A\tilde{p}} + \dots + q_A \pi r_{c0}^2 e^{2NA\tilde{p}}}_{\text{Geometric series}}$$

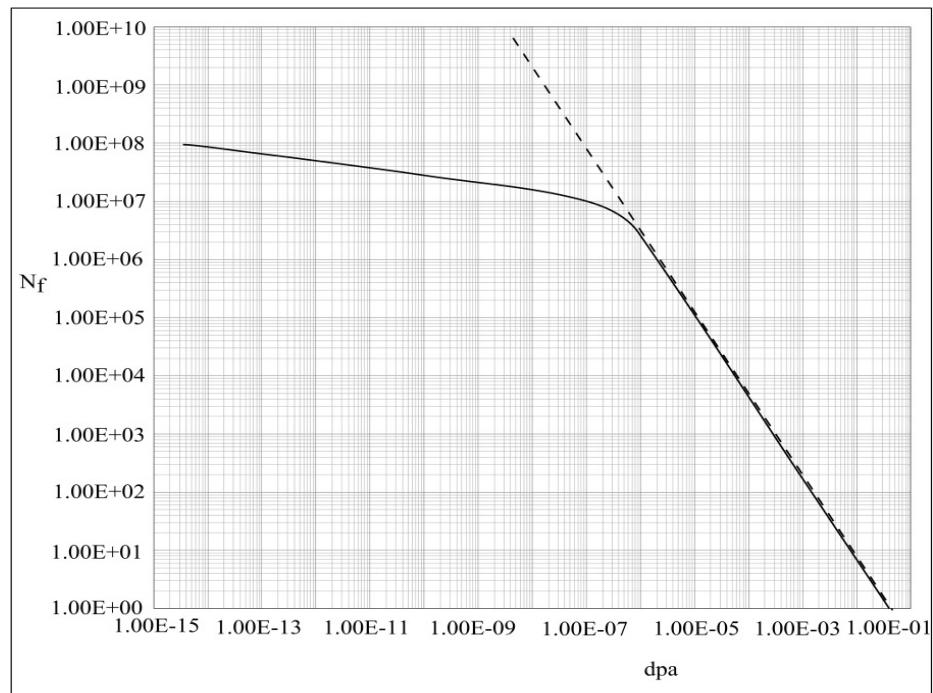
$$D_{rmN} = q_A \pi r_{c0}^2 \sum_{n=1}^N e^{2nA\tilde{p}}$$

$$S_N = a_1 \frac{1-q^N}{1-q} \quad S_N = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1-e^{2ApN}}{1-e^{2Ap}}$$

$$D_{rmN} = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1-e^{2ApN}}{1-e^{2Ap}}$$

$$D_{rmN_f} = q_A \pi r_{c0}^2 e^{2A\tilde{p}} \frac{1-e^{2A\tilde{p}N_f}}{1-e^{2A\tilde{p}}} = D_{cr}$$

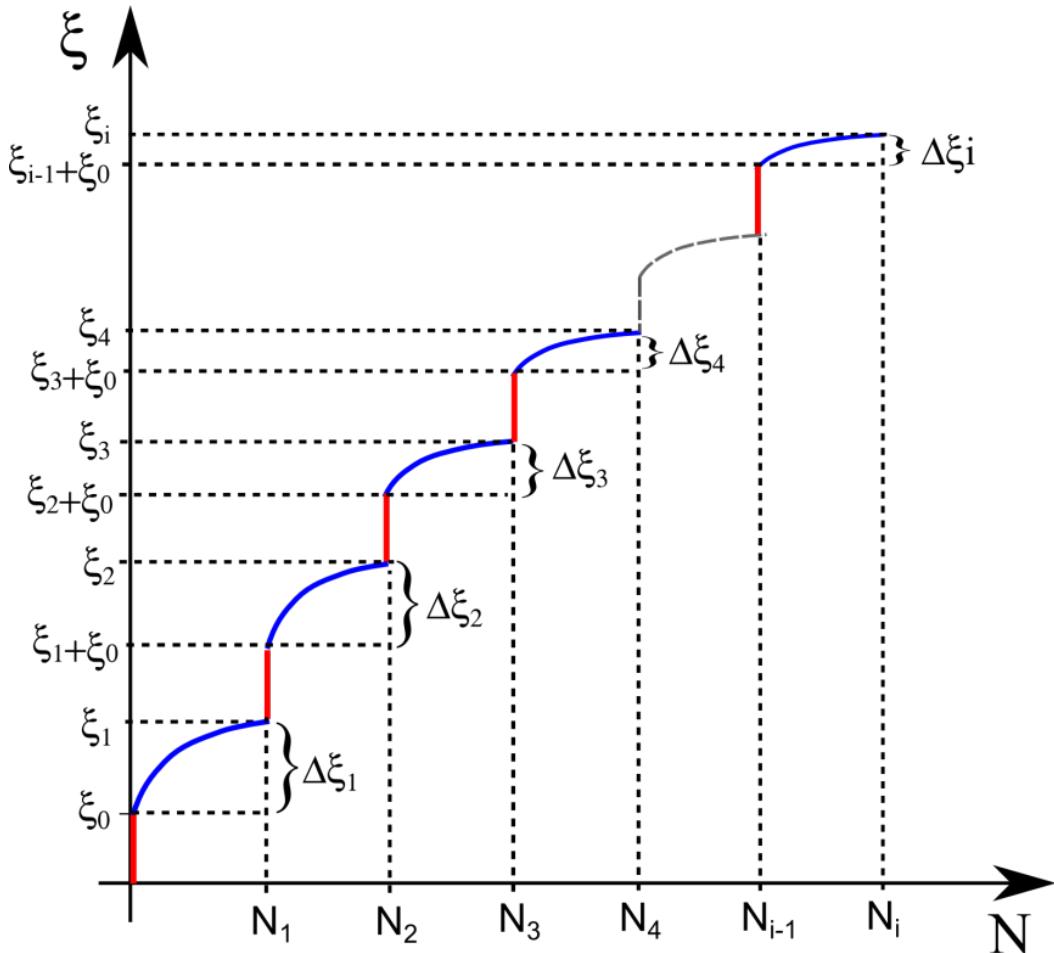
Number of cycles to failure N_f based on the critical damage criterion: $D_{rm}=D_{cr}$





Lifetime estimation for irradiated components

Phenomena at cryogenic temperatures



Gurson law

$$d\xi = (1 - \xi) dp$$

$$\int_{\xi_i + \xi_0}^{\xi_{i+1}} \frac{d\xi}{1 - \xi} = \int_0^{\tilde{p}} dp \quad K := e^{-\tilde{p}}$$

$$\boxed{\xi_{i+1} = 1 - (1 - \xi_0 - \xi_i) K}$$

Porosity parameter ξ_i increases from cycle to cycle by ξ_0 due to emission of secondary particles flux



Lifetime estimation for irradiated components

nomena at cryogenic temperatures

$$\xi_1 = 1 - (1 - \xi_0)K = 1 + \xi_0 K - K$$

$$\xi_2 = 1 - (1 - \xi_0 - \xi_1)K = 1 + \xi_0 K + \xi_0 K^2 - K^2$$

$$\xi_3 = 1 - (1 - \xi_0 - \xi_1 - \xi_2)K = 1 + \xi_0 K + \xi_0 K^2 + \xi_0 K^3 - K^3$$

$$\xi_4 = 1 + \xi_0 K + \xi_0 K^2 + \xi_0 K^3 + \xi_0 K^4 - K^4$$

⋮

$$\xi_i = (1 - K^i) + \xi_0 \sum_{n=1}^i K^n$$

$$\xi_N = 1 + \underbrace{\xi_0 K + \xi_0 K^2 + \xi_0 K^3 + \xi_0 K^4 + \dots + \xi_0 K^N}_{\text{Geometric series}} - K^N$$

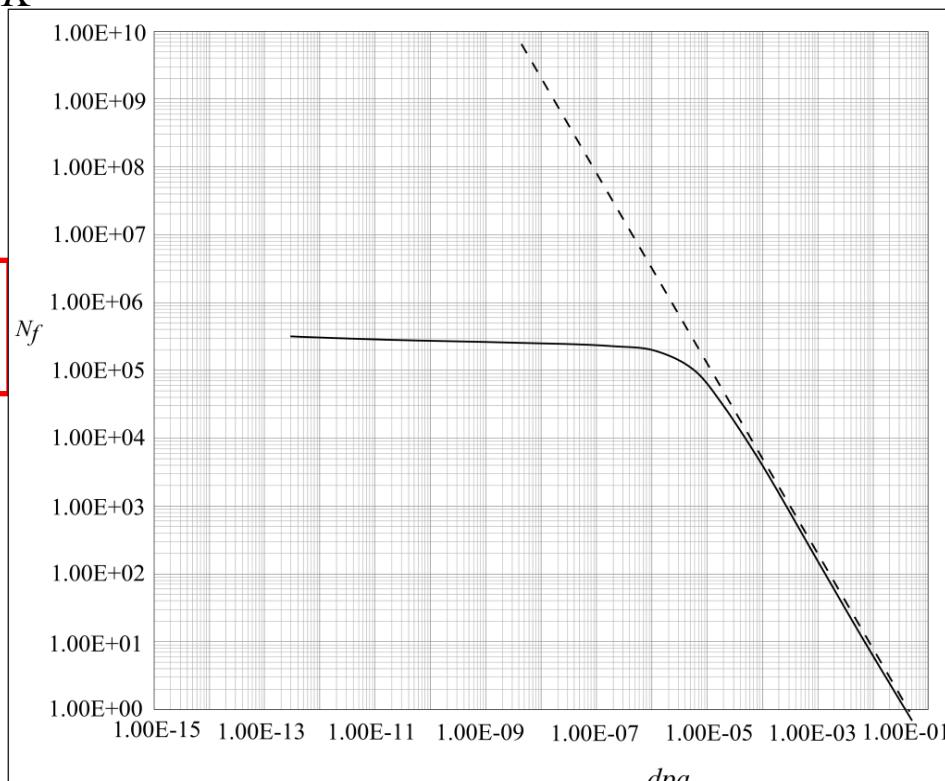
$$\xi_N = (1 - K^N) + \xi_0 \sum_{n=1}^N K^n$$

$$S_N = \xi_0 K \frac{1 - K^N}{1 - K}$$

$$\xi_N = 1 + \xi_0 K \frac{1 - K^N}{1 - K} - K^N$$

$$\xi_N = 1 + \xi_0 K \frac{1 - K^N}{1 - K} - K^N = \xi_{cr}$$

Number of cycles to failure N_f is based on the criterion $\xi_N = \xi_{cr}$



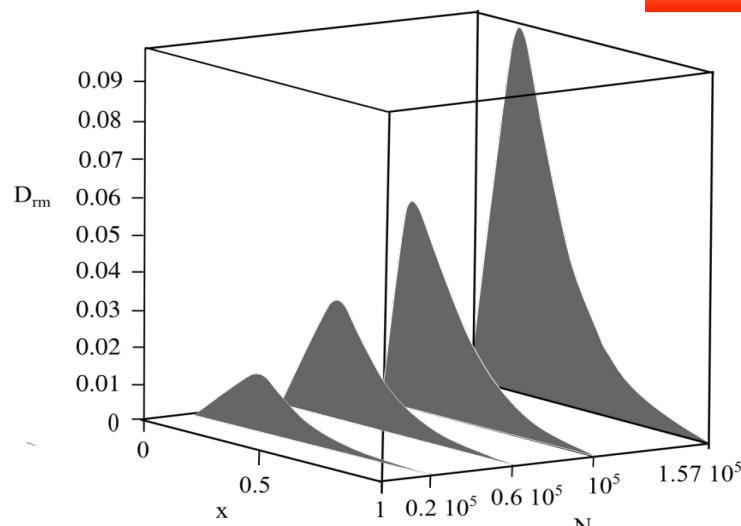
Mechanical loads have significant impact below $dpa \approx 10^{-5}$



Evolution of damage parameter in the horn (magnetic lens)

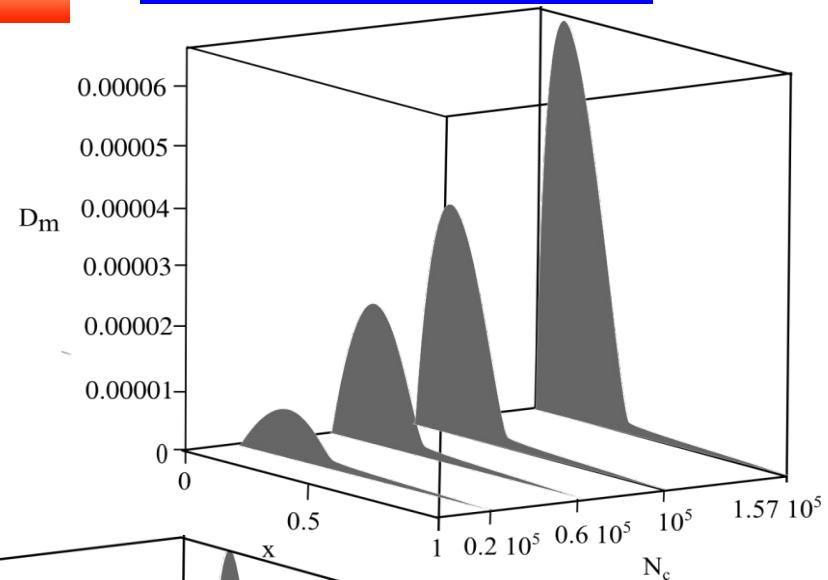
Coupled dissipative phenomena at cryogenic temperatures

radiation induced damage

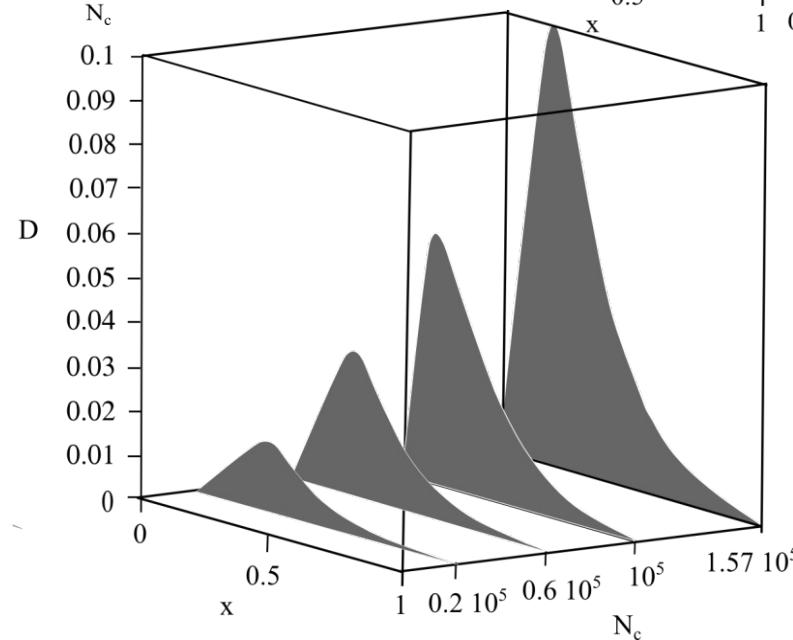


$D_{cr}=0.1$

mechanical damage

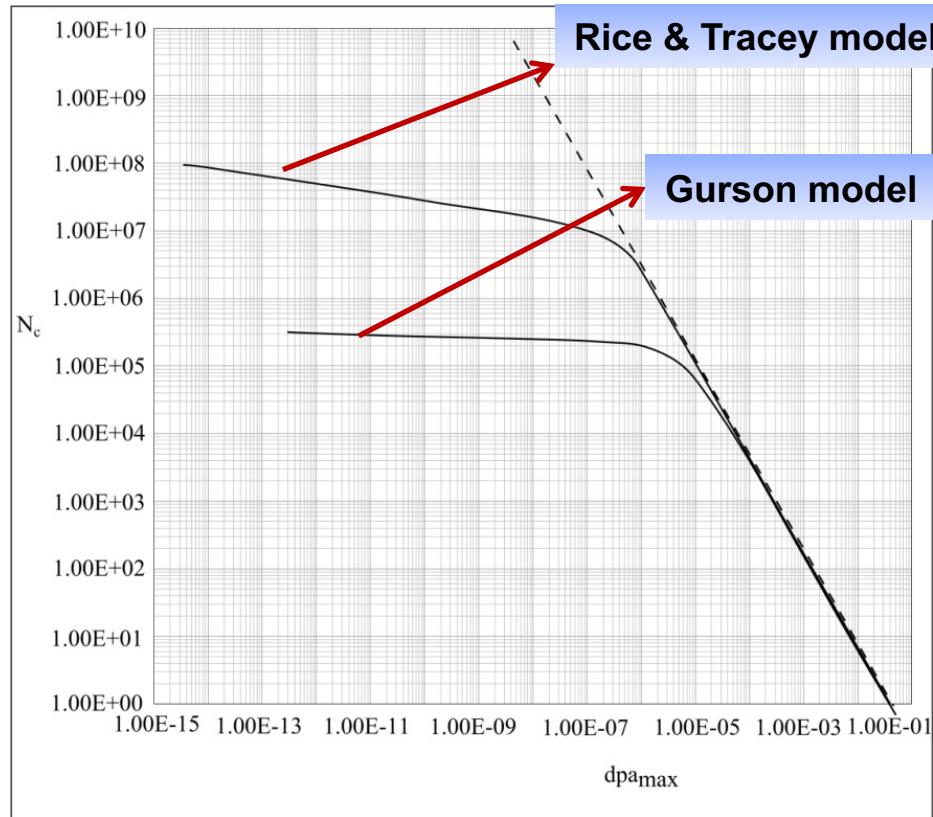


total damage parameter

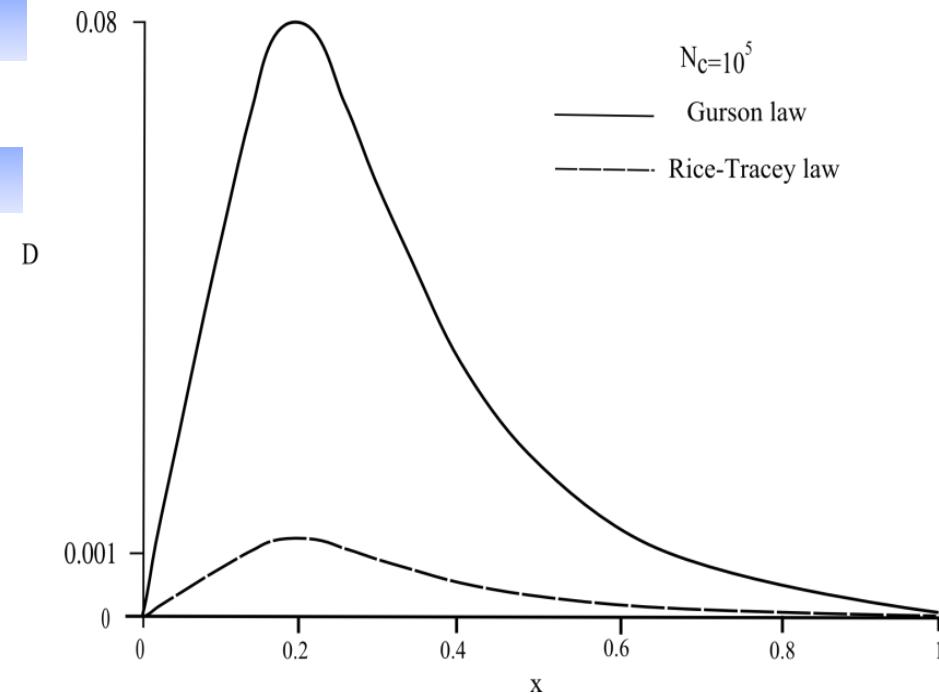




Performance of Rice-Tracey and Gurson models (log-log)



different sensitivity of both models



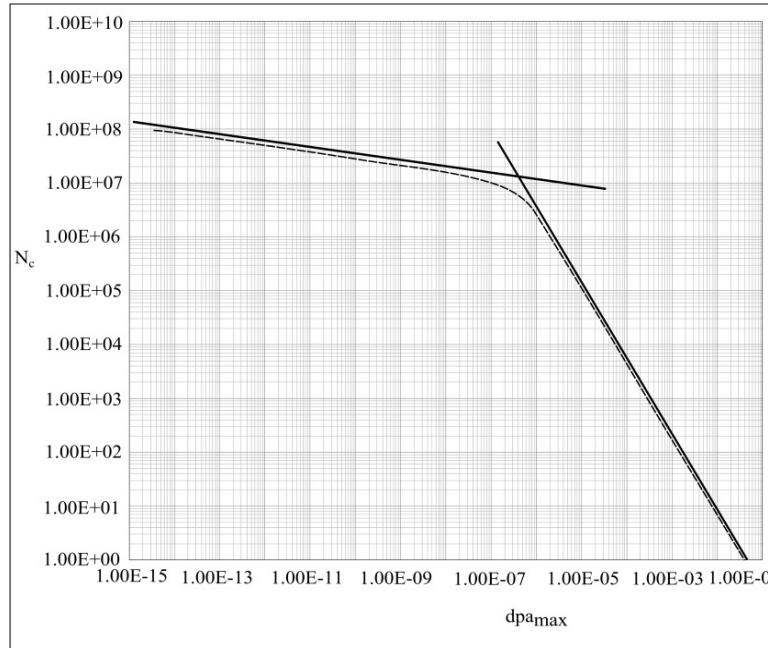
Rice & Tracey model predicts lower values of damage



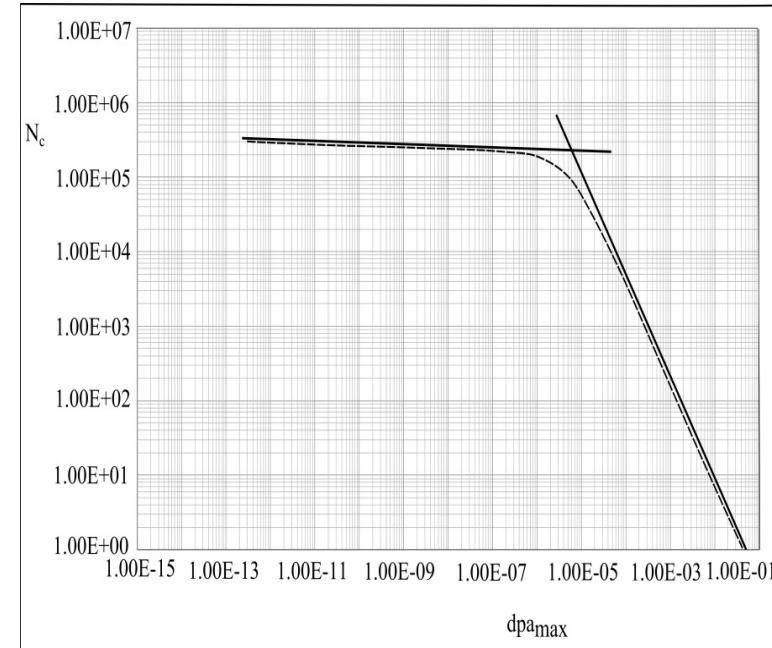
Bilinear approximations for R-T and Gurson models

Coupled dissipative phenomena at cryogenic temperatures

Rice & Tracey model



Gurson model



$$\log(N_c) = a + b \log(dpa_{\max})$$

Analytical formula - useful tool for estimation of number of cycles to failure



$$N_c = 10^a dpa_{\max}^b$$

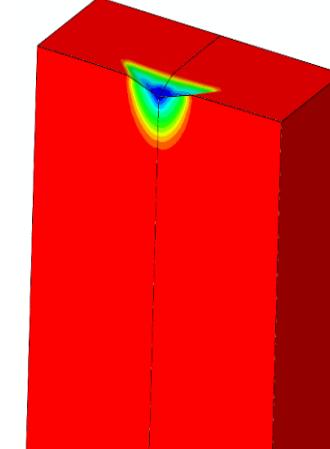
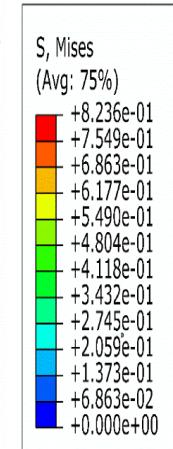
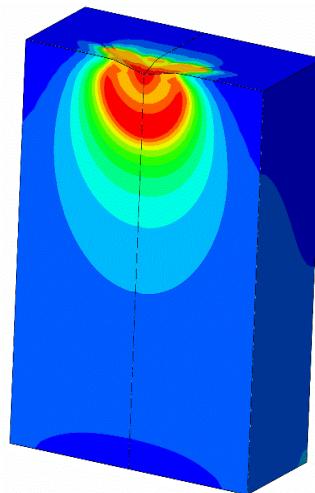
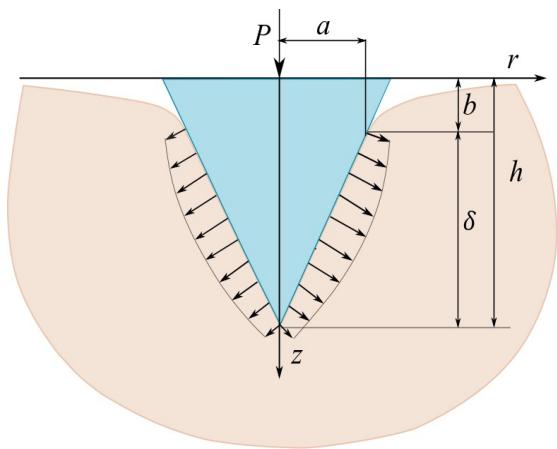
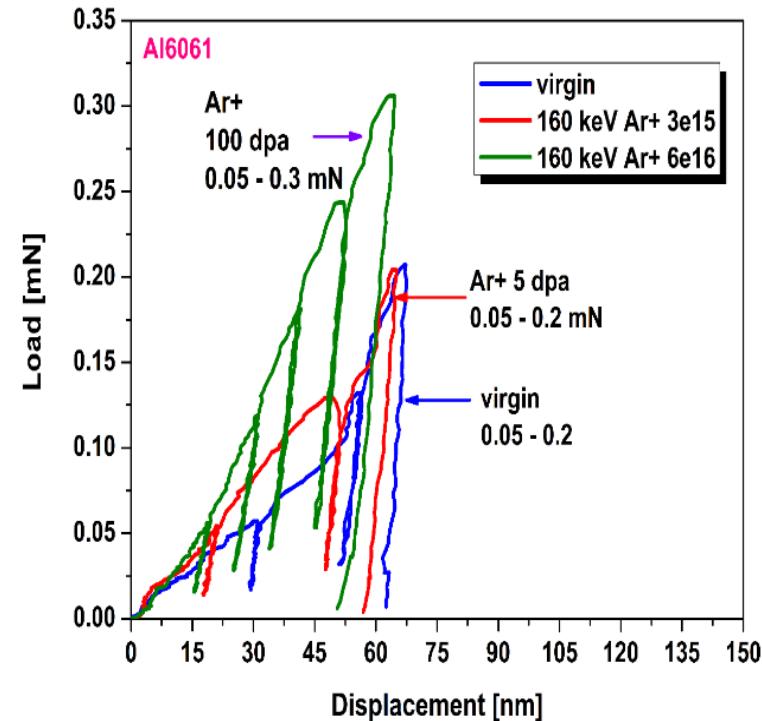
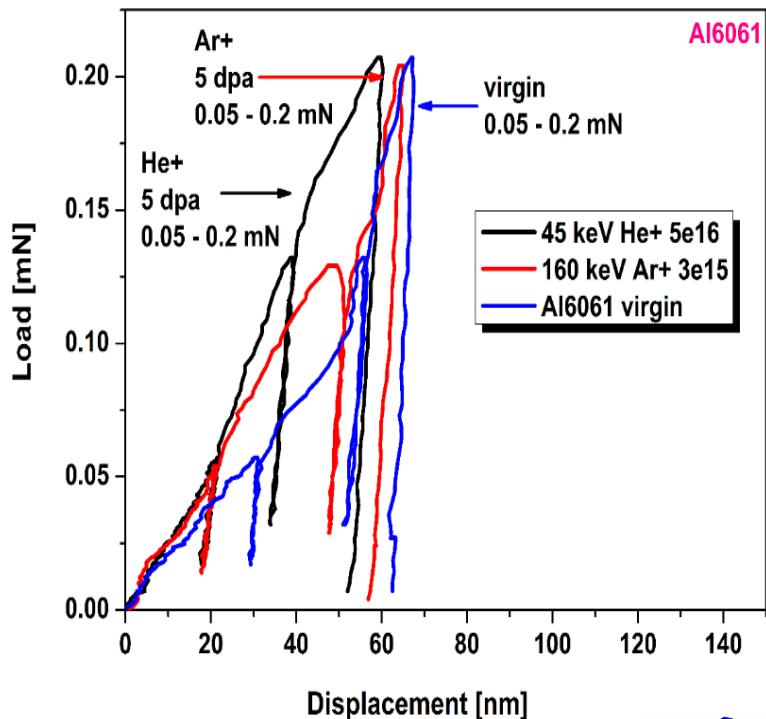
$$N_c = \begin{cases} 10^{-1.9} dpa_{\max}^{-1.4} & \text{for } dpa_{\max} \geq 10^{-6} \\ 10^{6.1} dpa_{\max}^{-0.13} & \text{for } dpa_{\max} < 10^{-6} \end{cases}$$

$$N_c = \begin{cases} 10^{-1.9} dpa_{\max}^{-1.4} & \text{for } dpa_{\max} \geq 10^{-5} \\ 10^{5.43} dpa_{\max}^{-0.016} & \text{for } dpa_{\max} < 10^{-5} \end{cases}$$



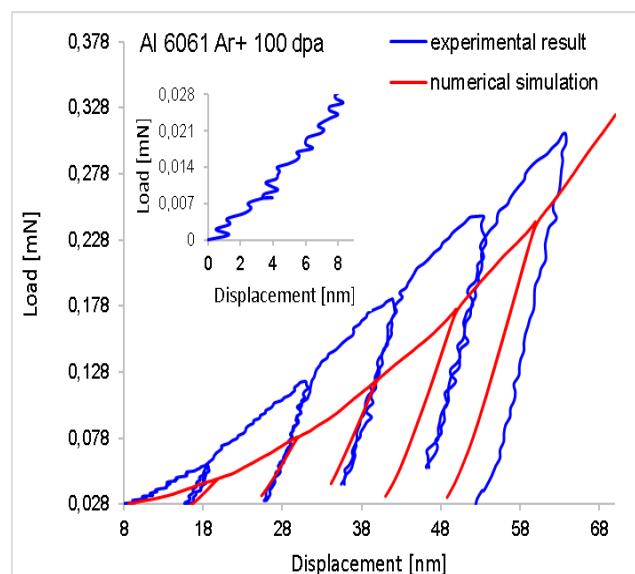
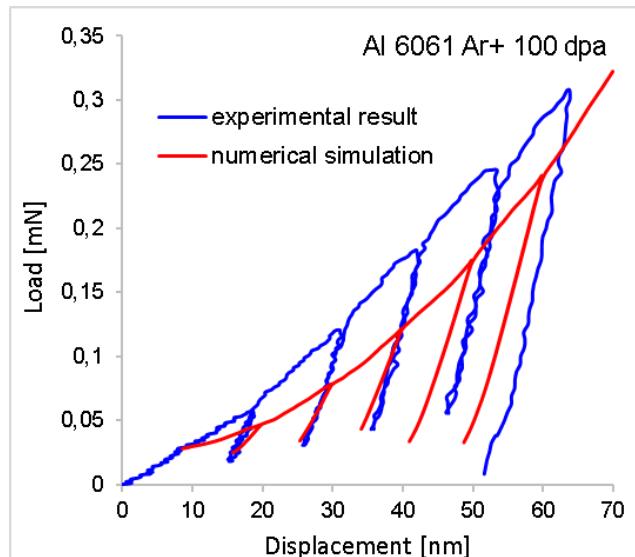
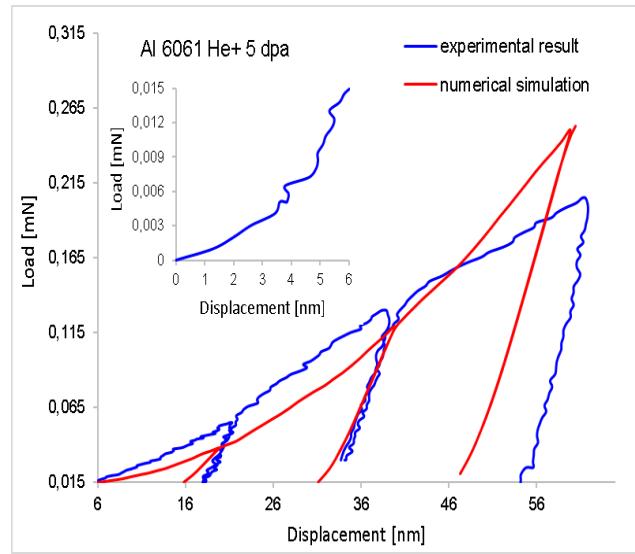
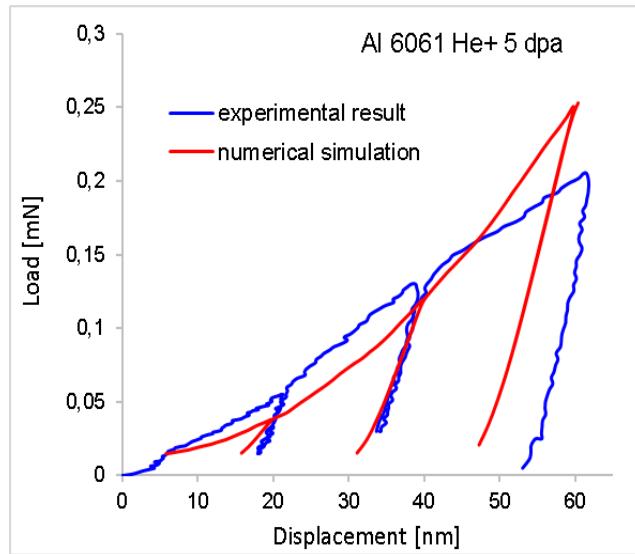
Nanoindentation of irradiated Al6061

Failure phenomena at cryogenic temperatures



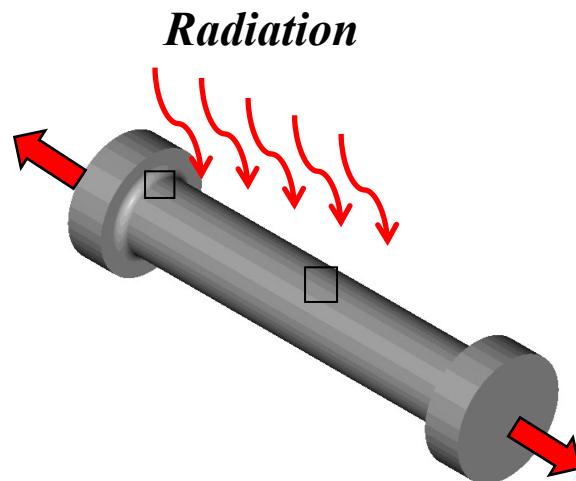


Nanoindentation of irr. Al6061: experiment vs. Gurson model





Radiation induced hardening



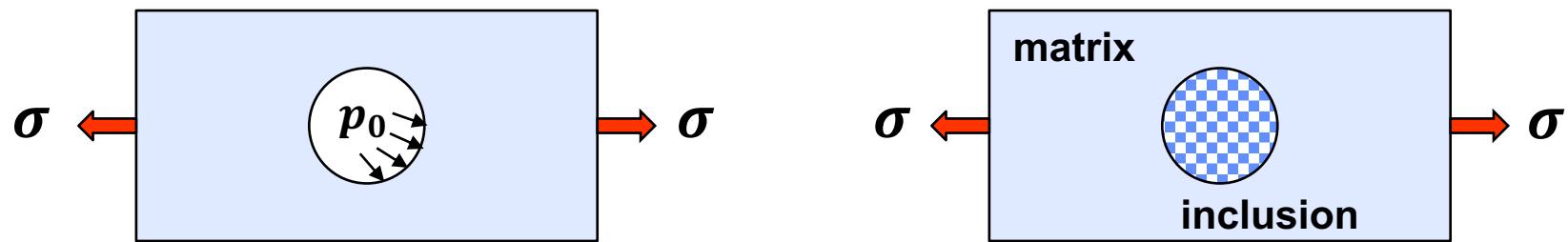


Type Eshelby entities: the equivalent inclusion

phenomena at cryogenic temperatures

Assumptions:

- small strains approach
- perfect gas inside the void at a constant temperature T
- pressurized void is equivalent to inclusion subjected to hydrostatic stress



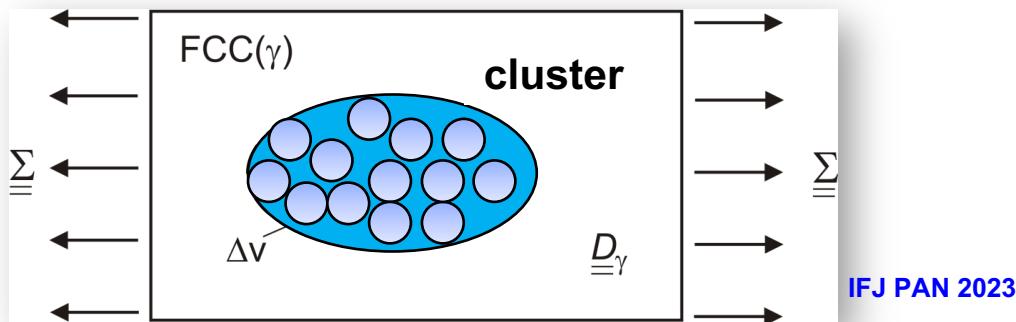
$$\Delta p = -3p_0\Delta\varepsilon$$

$$E_{pijkl} = 3k_p J_{ijkl} + 2\mu_p K_{ijkl}$$

$$\mu_p = 0 \ ; \ \nu_p = 0 \ ; \ k_p \neq 0$$

$$\Delta\sigma = E_p \Delta\varepsilon \ ; \ E_p = -3p_0$$

$$\Delta\sigma = -3p_0 \Delta\varepsilon$$





Evolution of voids – Rice & Tracey kinetics

ic temperatures

Evolution of voids corresponding to the far stress field: R&T kinetics

$$\frac{\dot{r}}{r} = \alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_i}\right) \dot{p}$$

$$\int_{r_0}^{\tilde{r}} \frac{dr}{r} = \alpha_r \int_0^{\tilde{p}} \exp\left(\frac{3\sigma_m}{2\sigma_i}\right) dp$$

For the uniaxial stress state:

$$\begin{aligned} \tilde{r} &= r_0 \left[\alpha_r \exp\left(\frac{3\sigma_m}{2\sigma_i}\right) \tilde{p} \right] \\ \tilde{r} &= r_0 [C_1 \tilde{p}] \end{aligned} \quad \left. \right\} \quad \begin{aligned} r &= Re^{C_1 p} \\ \Delta\xi &= \frac{n(V_d - V_D)}{V} \end{aligned}$$

$$\Delta\xi = \xi_0 (e^{3C_1 \Delta p} - 1)$$

Linearization:

$$\xi = \xi_0 + \Delta\xi = \xi_0 (1 + 3C_1 \Delta p)$$



Interaction of dislocations with voids

Dissipative phenomena at cryogenic temperatures

The Orowan mechanism:

$$\tau_p = \frac{\mu b}{d} \sqrt[3]{\frac{6\xi_0}{\pi}} \left(1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right) \quad \Delta\xi = 3C_1 \xi_0 \Delta p$$

$$\tau_p = \frac{\mu b}{d} \sqrt[3]{\frac{6\xi_0}{\pi}} (1 + C_1 \Delta p)$$

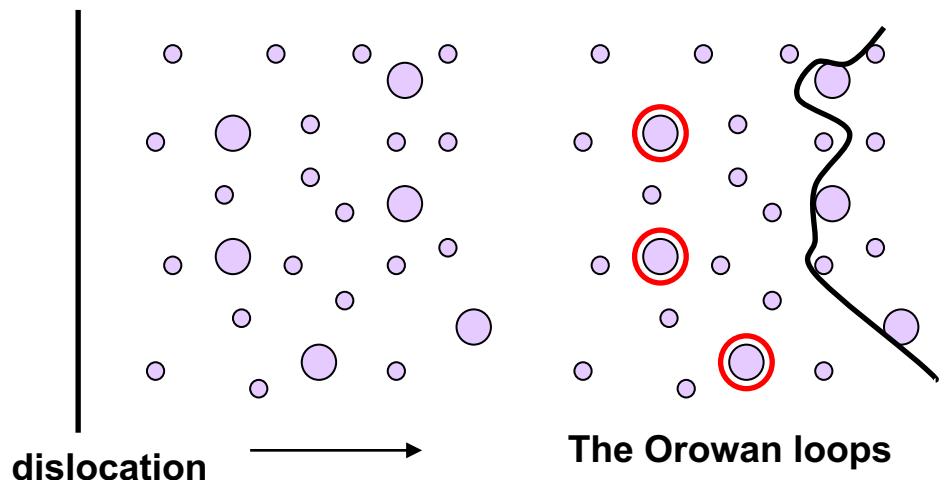
Using the Taylor factor:

$$\sigma_p = M\tau_p = MA_0 \sqrt[3]{\xi_0} \left(1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right)$$

Hardening modulus:

$$C = C_0 \left(1 + \frac{1}{3} \frac{\Delta\xi}{\xi_0} \right)$$

$$C = C_0 (1 + h\Delta\xi)$$



The Orowan loops



Mean field methods: homogenization

Dissipative phenomena at cryogenic temperatures

Stiffness of the matrix:

$$\Delta\sigma_{a_{ij}} = E_{ta_{ijkl}} \Delta\varepsilon_{kl}$$

$$E_{ta_{ijkl}} = 3k_{ta}J_{ijkl} + 2\mu_{ta}K_{ijkl}$$

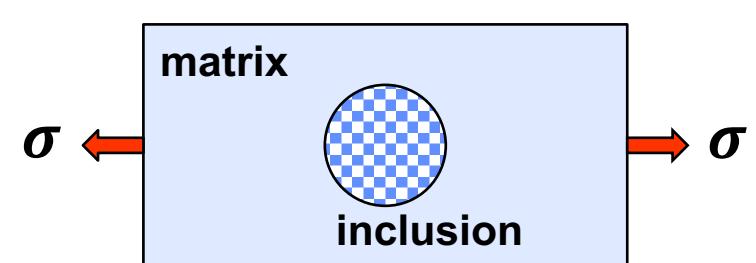
$$\mu_{ta} = \frac{E_t}{2(1+\nu)} \quad k_{ta} = \frac{E_t}{3(1-2\nu)} \quad E_t = \frac{EC}{E+C}$$

Stiffness of the inclusion:

$$\Delta\sigma_{p_{ij}} = E_{p_{ijkl}} \Delta\varepsilon_{kl}$$

$$E_{p_{ijkl}} = 3k_p J_{ijkl}$$

$$\mu_p = 0 \quad ; \quad \nu_p = 0 \quad ; \quad k_p = \frac{E_p}{3(1-2\nu_p)} = \frac{E_p}{3}$$



Mori-Tanaka homogenization:

$$\Delta\sigma_{ij} = E_{H_{ijkl}} \Delta\varepsilon_{kl}$$



Uniaxial case – tension/compression

dissipative phenomena at cryogenic temperatures

$$d\sigma = d\sigma_i + d\sigma_{MT}$$

$$d\sigma_i = C_0(1 + h\Delta\xi)d\varepsilon^p \quad ; \quad d\sigma_{MT} = E_H d\varepsilon^p = C_{MT} d\varepsilon^p$$

$$\Delta\sigma_i = C_0(\xi_0) \left(\varepsilon^p + \frac{1}{2} C_1 (\varepsilon^p)^2 \right) \quad \text{Interaction}$$

$$\Delta\sigma_{MT} = -\frac{5}{2} \mu \eta_0 \frac{\xi_0}{C_1} \left[\frac{1}{4} (\chi^2 - 1) - \frac{4}{81} \xi_0 (\chi^3 - 1) + -\frac{2}{81} \xi_0^2 (\chi^4 - 1) \right]$$

$$\chi = 1 + 3C_1 \varepsilon^p \quad \eta_0 = \frac{C_i}{E} = \frac{M \frac{\mu b}{d} \sqrt[3]{\frac{6}{\pi}} \sqrt[3]{\xi_0}}{E} \quad \text{MT homogenization}$$

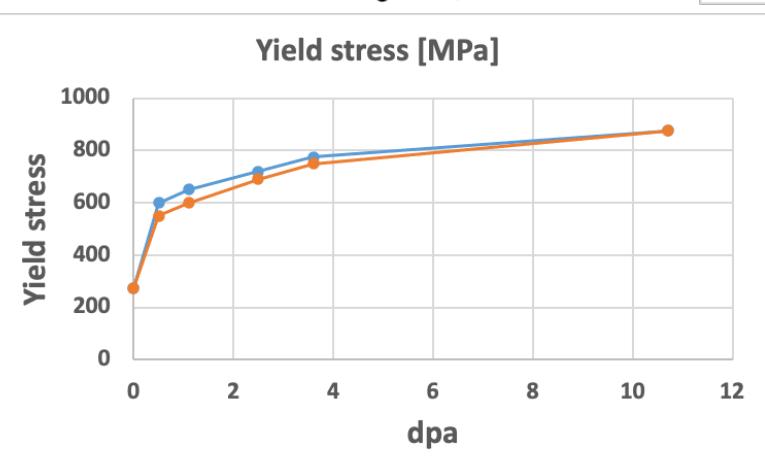
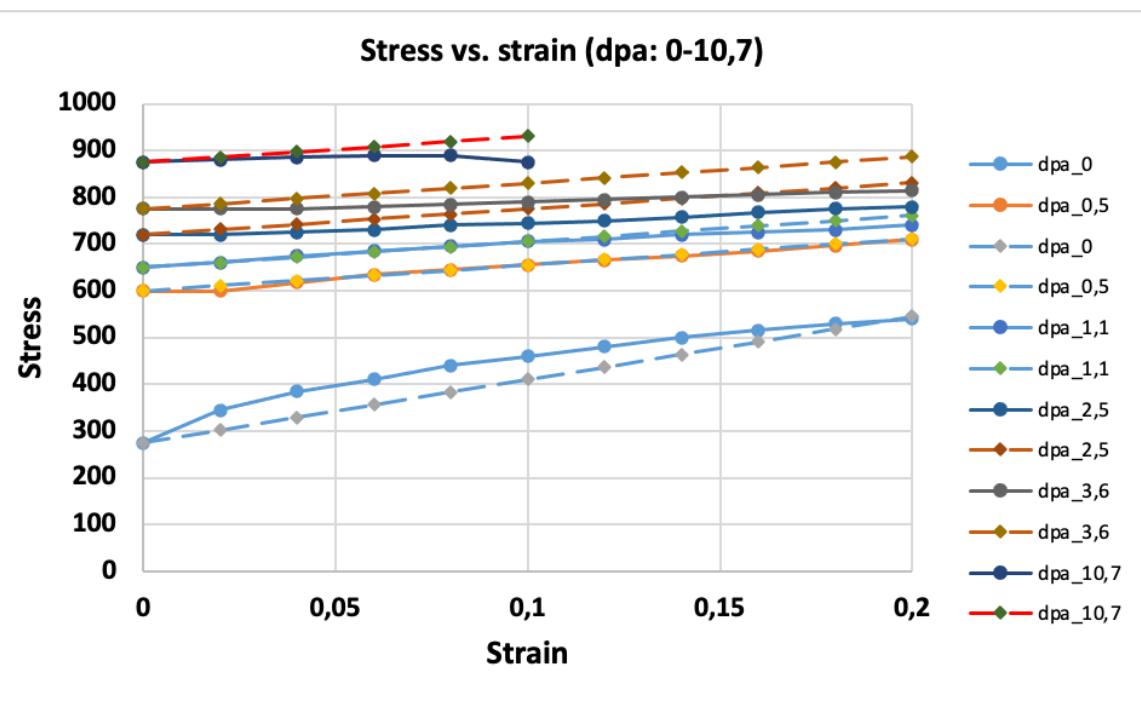
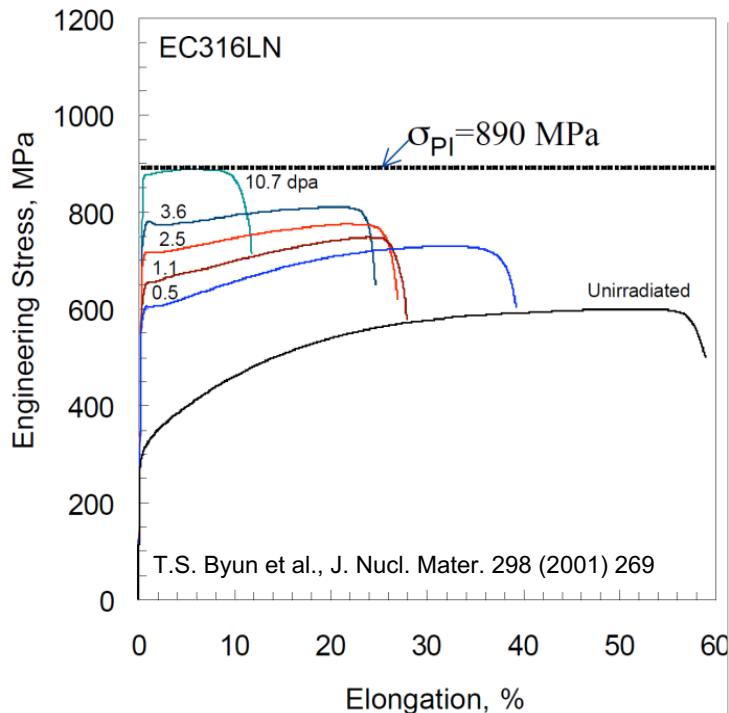
$$\begin{aligned} \sigma &= \sigma_0 + C_0(\xi_0) \left(\varepsilon^p + \frac{1}{2} C_1 (\varepsilon^p)^2 \right) \\ &- \frac{5}{2} \mu \frac{C_i(\xi_0)}{E} \frac{\xi_0}{C_1} \left[\frac{1}{4} (\chi^2 - 1) - \frac{4}{81} \xi_0 (\chi^3 - 1) + -\frac{2}{81} \xi_0^2 (\chi^4 - 1) \right] \end{aligned}$$





Uniaxial case: 316LN stainless steel

Dissipative phenomena at cryogenic temperatures

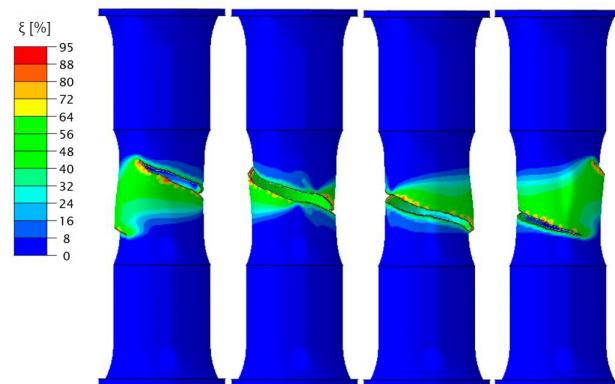


Radiation induced hardening comprising:

- massive interaction of dislocations with the pressurized voids,
- evolution of tangent stiffness expressed by the Mori-Tanaka homogenization.



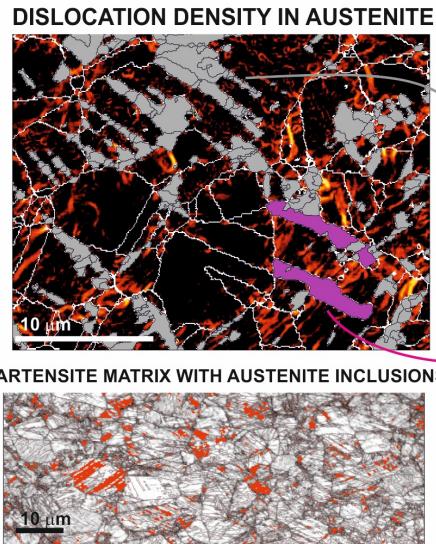
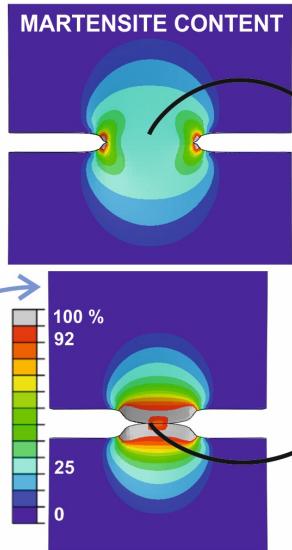
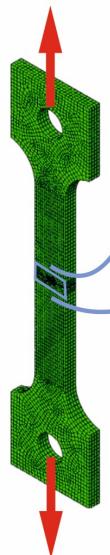
Fracture at extremely low temperatures



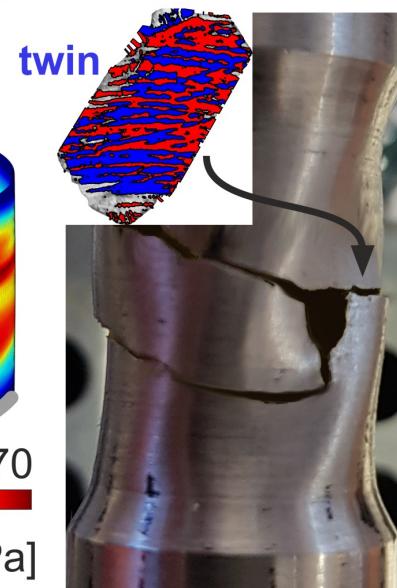
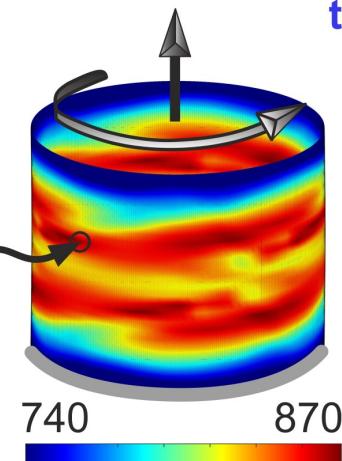
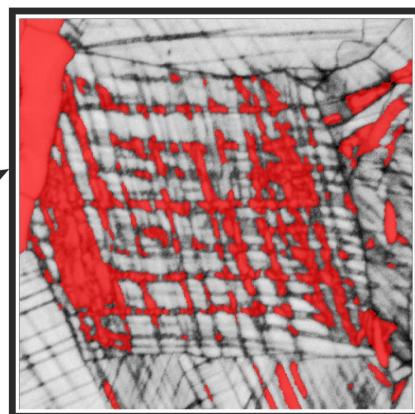
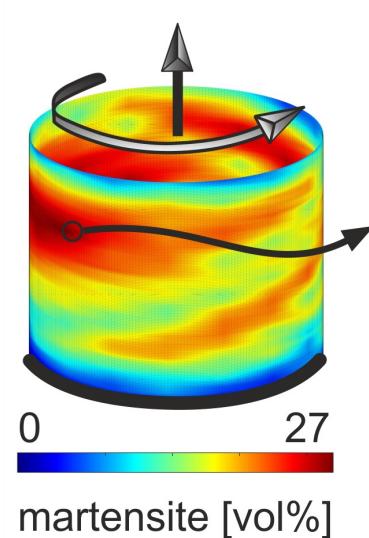
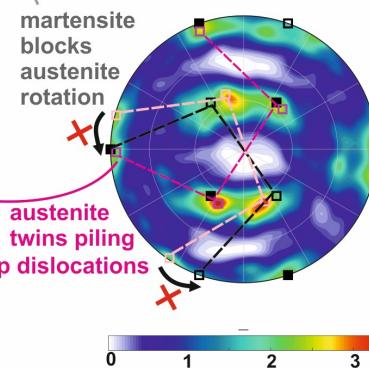


Fracture at extremely low temperatures

phenomena at cryogenic temperatures



POLE DISTRIBUTION FUNCTION
(1 1 1) FOR AUSTENITE
stable Brass orientation
□ I variant ■ II variant





Conclusions:

1. Radiation induced defects in the lattice constitute obstacles for the motion of dislocations.
2. Microvoids filled with impurities (gas) induce two physical effects: hardening and swelling.
3. Hardening is related to the interaction of dislocations with the defects, in particular with voids filled with impurities.
4. Tangent stiffness corresponds to the proportion between the volume fraction of matrix, and the volume fraction of voids with impurities.
5. Good correlation between the experiment and the numerical results was obtained.



**Thank you for your
attention!**