$K \mbox{ and } B \mbox{ Physics in the } \mbox{Custodially Protected Randall-Sundrum Model}$

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Outline

1 RS Model with Custodial Protection

- (2) K and B Physics Observables
 - Meson-Antimeson Mixing
 - Rare Decays

3 Summary

based on: ALBRECHT, MB, BURAS, DULING, GEMMLER, 0903.2415
 MB, BURAS, DULING, GORI, WEILER, JHEP 03 (2009) 001 [0809.1073]
 MB, BURAS, DULING, GEMMLER, GORI, JHEP 03 (2009) 108 [0812.3803]

other related studies, see e.g.: GEDALIA, ISIDORI, PEREZ CASAGRANDE, GOERTZ, HAISCH, NEUBERT, PFOH BAUER, CASAGRANDE, GRUNDER, HAISCH, NEUBERT AGASHE, PEREZ, SONI CSAKI, FALKOWSKI, WEILER HUBER; HUBER, SHAFI

The Randall-Sundrum Framework

5D spacetime with warped metric:

Randall, Sundrum, hep-ph/9905221

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2 \,, \qquad 0 \le y \le L$$



• fermions and gauge bosons live in the bulk

• Higgs localised on IR brane

Chang et al., hep-ph/9912498 Grossman, Neubert, hep-ph/9912408 Gherghetta, Pomarol, hep-ph/0003129

➤ energy scales suppressed by warp factor e^{-ky} → natural explanation of gauge hierarchy problem

> Kaluza-Klein (KK) excitations live close to the IR brane

RS Model with Custodial Protection

Protection of T parameter and $Zb_L\bar{b}_L$ coupling

Agashe et al., hep-ph/0308036; Csaki et al., hep-ph/0308038 Agashe et al., hep-ph/0605341

Extended electroweak symmetry structure:



RS Model with Custodial Protection

Fermion Localisation and Yukawa Couplings

zero mode profile depends strongly on bulk mass parameter $c{:}$ $f^{(0)}(y,c) \propto e^{(\frac{1}{2}-c)ky}$



$$> \frac{1}{2}$$
: localisation around UV brane $< \frac{1}{2}$: localisation around IR brane

effective 4D Yukawa couplings:

$$(Y_{u,d})_{ij} = (\lambda_{u,d})_{ij} f_i^Q f_j^{u,d}$$

- > hierarchical structure can be naturally generated by exponential suppression of $f^{Q,u,d}$ (fermion profile on IR brane)
- light fermions live close to the UV brane third generation localised closest to the IR brane

GROSSMAN, NEUBERT, HEP-PH/9912408 Arkani-Hamed, Grossman, Hep-ph/9909411

Tree Level FCNCs & RS-GIM Mechanism

- different fermion localisation leads to flavour non-universal couplings of KK gauge bosons
- \succ transmitted to Z couplings via EWSB
- tree level FCNCs arise through rotation to mass eigenbasis

$$\begin{split} \bar{d}_{L}^{i}G_{\mu}d_{L}^{j} &\sim ig^{\text{4D}}\gamma_{\mu}\sqrt{kL} \ \boldsymbol{f_{i}^{Q}} \ \boldsymbol{f_{j}^{Q}} \\ \bar{d}_{R}^{i}G_{\mu}d_{R}^{j} &\sim ig^{\text{4D}}\gamma_{\mu}\sqrt{kL} \ \boldsymbol{f_{i}^{d}} \ \boldsymbol{f_{j}^{d}} \end{split}$$

> RS-GIM mechanism: suppression by flavour hierarchies

Agashe, Perez, Soni, hep-ph/0408134

new contributions to
 meson-antimeson mixing and rare K and B decays



 \boldsymbol{K} and \boldsymbol{B} Physics Observables

Detailed Study of K and B Physics Observables

Albrecht, MB, Buras, Duling, Gemmler, 0903.2415

theoretical framework & Feynman rules

MB, BURAS, DULING, GORI, WEILER, JHEP 03 (2009) 001 [0809.1073]

- calculation of new strong and EW contributions to $\Delta F=2$
- simultaneous analysis of all interesting $\Delta F = 2$ observables
- study of required fine-tuning

MB, BURAS, DULING, GEMMLER, GORI, JHEP 03 (2009) 108 [0812.3803]

- extension to include the most prominent rare K and B decays $(K \to \pi \nu \bar{\nu}, K_L \to \pi^0 \ell^+ \ell^-, K_L \to \mu^+ \mu^-, B \to K^{(*)} \nu \bar{\nu}, B \to X_s \nu \bar{\nu}, B_{s,d} \to \mu^+ \mu^-)$
- quantification of possible new physics effects
- study of correlations among various observables
- comparison to LHT predictions (C. TARANTINO'S TALK TUM TEAM, 06-09)

K and B Physics Observables

Messages from the $\Delta F = 2$ Sector

- large contributions to left-right operators by KK gluons
- stringent constraint from ε_K
 CSAKI, FALKOWSKI, WEILER, 0804.1954
- even for $M_{\rm KK} \simeq 2.5 \,{\rm TeV}$: all $\Delta F = 2$ data can be fulfilled without significant fine-tuning





★ distinction from LHT!

- large CP-violating effects in
 - $B_s \bar{B}_s$ mixing are possible
- $\circ -1 < S_{\psi \phi} < 1$ allowed
- ullet sizable NP effects also in $A^s_{\mathsf{SL}}, \Delta\Gamma_s$

Pattern of New Physics Effects in Rare Decays

• $Zd_L^i \bar{d}_L^j$ protected by enlarged custodial symmetry

see also Buras, Duling, Gori, 0905.2318

- Z_H, Z' contributions geometrically suppressed $(\propto 1/(kL))$
- \succ rare decays dominated by Z coupling to right-handed quarks

flavour hierarchy weaker in right-handed down sector



Rare Decays

The $K \rightarrow \pi \nu \bar{\nu}$ System





Rare Decays

The $K \rightarrow \pi \nu \bar{\nu}$ System





K and B Physics Observables

Rare Decays

Other Useful Correlations in the K System

$$K_L
ightarrow \mu^+ \mu^-$$
 vs. $K^+
ightarrow \pi^+
u ar{
u}$

(CP-conserving)



$$K_L
ightarrow \pi^0 \mu^+ \mu^-$$
 vs. $K_L
ightarrow \pi^0 e^+ e^-$

(CP-violating)



inverse correlation

> V+A structure of flavour violating coupling

★ distinction from LHT!

linear correlation

- ➤ no scalar operators
- ➤ universality of CP-phases

Isidori et al., hep-ph/0404127 Friot et al., hep-ph/0404136 Mescia et al., hep-ph/0606081

Rare Decays

K vs. B Physics



large effects possible in $S_{\psi\phi}$ and rare K decays, but not simultaneously

- much smaller effects
 (~10%) in rare *B* decays
- o difficult to measure
- MFV relations strongly violated



Summary

Main Messages from Flavour Phenomenology

Confronting $\Delta F = 2$ observables:

- (1) generic strong constraint from ε_{K} due to enhanced \mathcal{Q}_{LR} contribution
- 2 but also for $M_{\rm KK}\gtrsim (2-3)\,{\rm TeV}$ agreement with ε_K possible without relevant fine-tuning
- (3) possible tensions in the SM $(\varepsilon_K, S_{\psi K_S}, ...)$ can be solved
- ④ $B_s \bar{B}_s$ CP-violation can be large

Implications for rare K and $B_{d,s}$ decays:

- **(1)** dominant contribution from right-handed Z couplings
- ② large effects in K decays, smaller effects in $B_{s,d}$ decays
- 3 correlations allow for distinction from other NP models
- ④ sizable violations of MFV relations possible

Main Virtues of Warped Extra Dimensions

- warped extra dimensions with bulk fields provide a simultaneous explanation of gauge and flavour hierarchies
- ② custodially extended model is consistent with EW precision data for $M_{
 m KK}\simeq (2-3){
 m TeV}$
- 3 flavour exhibits a Froggatt-Nielsen-like structure:
 - effective 4D Yukawa couplings obtained from hierarchical bulk profiles and O(1) 5D couplings
 - natural explanation of mass and CKM hierarchy
 - new tree level FCNC effects strongly suppressed by RS-GIM mechanism

RS versus Froggatt-Nielsen

bulk fermions in RS

Froggatt-Nielsen symmetry

$$(Y_{u,d}^{\rm RS})_{ij} \propto (\lambda_{u,d})_{ij} \, e^{-kL(c_Q^i - c_{u,d}^j)}$$

$$(Y_{u,d}^{\sf FN})_{ij} \propto (\lambda_{u,d})_{ij} \, \epsilon^{a_i - b_j^{u,d}}$$

self-similarity along ybulk mass parameters $c^i_{Q,u,d}$ IR brane at y = Lwarp factor e^{-kL}

$$U(1)_F$$
 symmetry
 $U(1)_F$ charges $Q_F = a_i, b_i^{u,d}$
VEV of scalar Φ ($Q_F = 1$)
 $\epsilon = \langle \Phi \rangle / \Lambda \ll 1$

• geometric interpretation of flavour symmetry

FN formulae for masses and flavour mixings can be applied
 ➤ dependence on λ_{u,d} and CP phases made explicit

BBDGW; CASAGRANDE ET AL., 0807.4937

Explicit Expressions for Masses and Mixings

quark masses:

$$m_b = \frac{v}{\sqrt{2}} \lambda_{33}^d \mathbf{f_3^Q} \mathbf{f_3^d}$$

$$m_s = \frac{v}{\sqrt{2}} \frac{\lambda_{33}^d \lambda_{22}^d - \lambda_{23}^d \lambda_{32}^d}{\lambda_{33}^d} \mathbf{f_2^Q} \mathbf{f_2^d}$$

$$m_d = \frac{v}{\sqrt{2}} \frac{\det(\lambda^d)}{\lambda_{33}^d \lambda_{22}^d - \lambda_{23}^d \lambda_{32}^d} \mathbf{f_1^Q} \mathbf{f_1^d}$$

flavour mixing matrices (responsible for FCNCs):

$$\begin{split} (\mathcal{D}_L)_{ij} &= \omega_{ij}^d \frac{f_i^Q}{f_j^Q} \qquad (\mathcal{D}_R)_{ij} = \rho_{ij}^d \frac{f_i^d}{f_j^d} \qquad (i < j) \\ (\omega_{ij}^d, \rho_{ij}^d: \text{ functions of } \lambda_d) \end{split}$$

analogous formulae for the up-type quarks

Some Details on the Custodially Protected RS Model

Sources of Flavour Violation & Parameter Counting

Flavour is violated by:

- bulk mass terms c_Q, c_u, c_d: 3 × 3 hermitian matrices
- Yukawa couplings λ_u, λ_d :
 - 3×3 complex matrices

$U(3)^3$ flavour symmetry can be used to remove

physical flavour parameters:

Agashe, Perez, Soni, hep-ph/0408134

- $3\times 6~\mathrm{real}$ parameters
- $3\times 3~{\rm complex}$ phases
- $2\times9~\mathrm{real}$ parameters
- 2×9 complex phases

36 real parameters 27 complex phases

- 9 real parameters
- 17 complex phases

27 real parameters 10 complex phases

containing the SM 9+1 ones

Required Fermion Content (Quark Sector)

- Higgs transforms as $(\mathbf{2}, \mathbf{2})$
- due to P_{LR} .: $q_L \in (\mathbf{2}, \mathbf{2})$ with $T_L^3(b_L) = T_R^3(b_L) = -1/2$
- right-handed SM fields t_R, b_R are SU(2)_L singlets
 ➤ possible gauge invariant Yukawa structures:

$$\overline{(\mathbf{2},\mathbf{2})}\otimes(\mathbf{2},\mathbf{2})\otimes(\mathbf{1},\mathbf{1})$$
 $\overline{(\mathbf{2},\mathbf{2})}\otimes(\mathbf{1},\mathbf{3})\otimes(\mathbf{2},\mathbf{2})$

Contino et al., hep-ph/0612048, Carena et al., hep-ph/0607106

smoking gun signature: $SU(2)_R$ partner of t_L : Q = 5/3 quarks with mass ≤ 1 TeV

More on RS Flavour Violation

Generic Bound on KK Scale

BBDGW



average required tuning in ε_K , depending on $M_{\rm KK}$

▶ generic naturalness bound: $M_{\rm KK} \gtrsim 20 \,{\rm TeV}$

confirms CSAKI, FALKOWSKI, WEILER, 0804.1954

Operator Competition in $\Delta F = 2$



(no chiral LR enhancement in B system)

 Q_{RR} contribution generally small

BBDGW

Anatomy of Z, Z_H , Z' Contributions

• FCNC couplings to left-handed quarks:

 $\Delta_L(Z_H) : \Delta_L(Z') : \Delta_L(Z) \sim \mathcal{O}(1) : \mathcal{O}(10^{-1}) : \mathcal{O}(10^{-4})$

BBDGG

Z and Z^\prime coupling suppressed by custodial protection

• FCNC couplings to right-handed quarks:

 $\Delta_R(Z_H): \Delta_R(Z'): \Delta_R(Z) \sim \mathcal{O}(1): \mathcal{O}(1): \mathcal{O}(10^{-2})$

custodial protection not effective

• propagators $M_Z^2/M_{\mathsf{KK}}^2 \sim \mathcal{O}(10^{-3})$ and leptonic couplings $\Delta_{L,R}^{\ell}(Z_H) : \Delta_{L,R}^{\ell}(Z') : \Delta_{L,R}^{\ell}(Z) \sim \mathcal{O}(10^{-1}) : \mathcal{O}(10^{-1}) : 1$

 $\succ \Delta F = 1$ rare decays dominated by $\Delta_R(Z)$

LHT: CP-Violation in $B_s - \bar{B}_s$ Mixing

BBPTUW, hep-ph/0605214; BBRT, 0805.4393; BBDRT, 0906.5454

generally: LHT effects in B physics expected to be small but: CP-violation in B_s extremely suppressed in the SM due to $\beta_s\simeq -1^\circ$

Iarge LHT effects possible



• $S_{\psi\phi} \sim 0.5$ possible • naturally $S_{\psi\phi} \lesssim 0.2$ • strong correlation with A_{SL}^s LIGETI ET AL., HEP-PH/0604112 MB ET AL., HEP-PH/0604057 GROSSMAN ET AL., 0904.0305 BIGI ET AL., 0904.1545

Predictions of the Littlest Higgs with T-Parity

LHT: The $K \rightarrow \pi \nu \bar{\nu}$ System

BBDRT, 0906.5454



- factor 2–3 enhancements of $K \rightarrow \pi \nu \bar{\nu}$ possible
- strict correlation (two branches of possible points)
 - ▶ equal CP-phases in $K^0 \bar{K}^0$ and $K \to \pi \nu \bar{\nu}$
 - > no new operators in $K^0 \bar{K}^0$ mixing

MB, 0904.2528

Predictions of the Littlest Higgs with T-Parity

LHT: Correlations between Various Rare K Decays



strong linear correlation in both cases \succ V-A structure of flavour violating coupling $(K_L \rightarrow \mu^+ \mu^- \text{ vs. } K^+ \rightarrow \pi^+ \nu \bar{\nu})$ \downarrow

> universality of CP-phases $(K_L \to \pi^0 \mu^+ \mu^- \text{ vs. } K_L \to \pi^0 \nu \bar{\nu})$

Predictions of the Littlest Higgs with T-Parity

LHT: K versus B Physics

BBDRT, 0906.5454



simultaneous large effects in $S_{\psi\phi}$ and rare K decays unlikely, but not impossible