

String Cosmology

Marco Zagermann
(MPI for Physics, Munich)



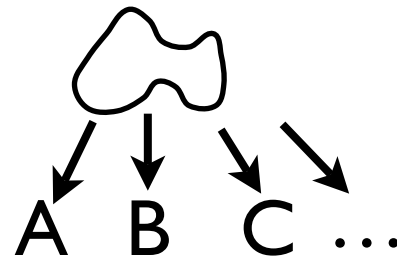
Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

String Theory

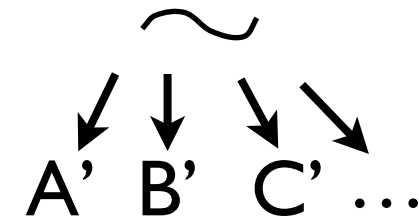
- “Unified” theory of all particles and interactions

– Particles:

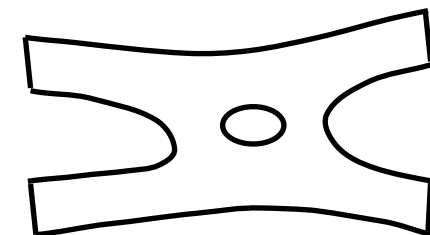
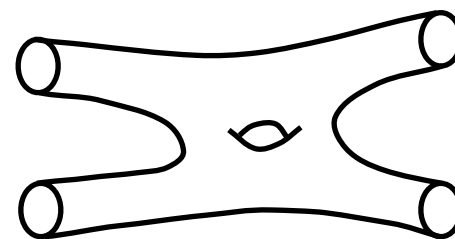
Closed



Open



– Interactions:

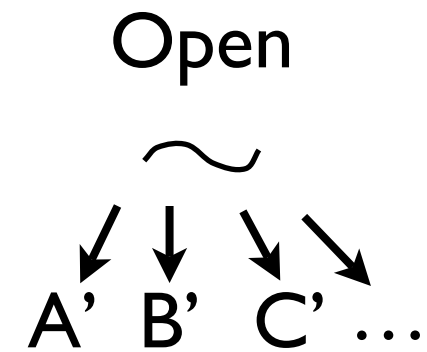
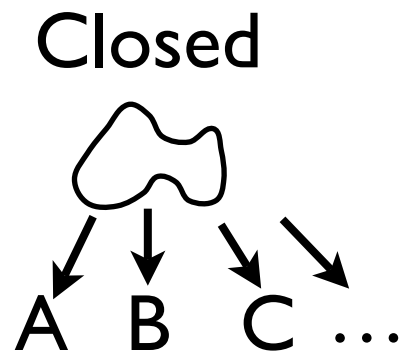


etc.

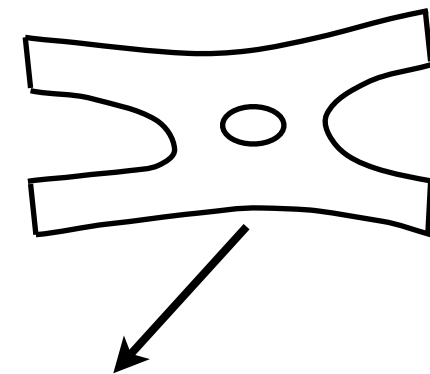
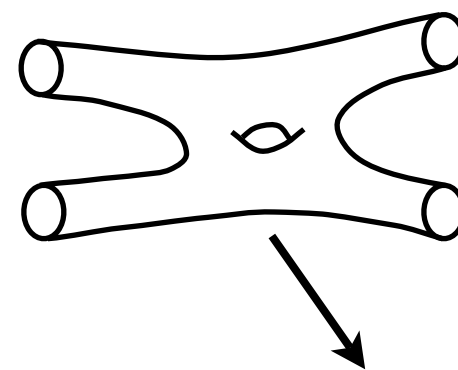
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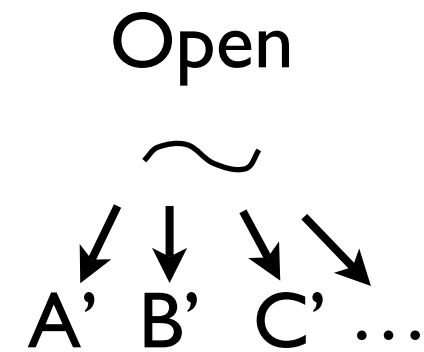
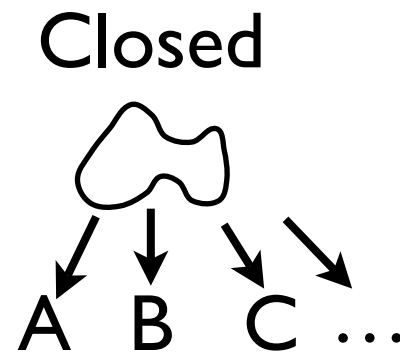
- UV-Completeness:

Finite

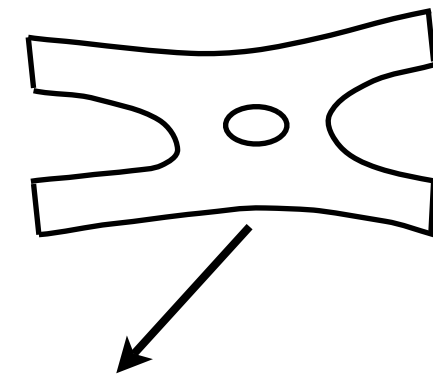
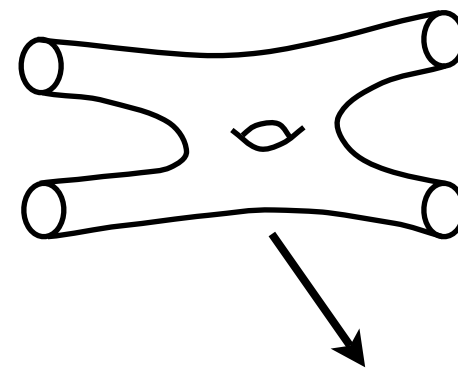
String Theory

- “Unified” theory of all particles and interactions

– Particles:



– Interactions:



etc.

- UV-Completeness:

Finite

⇒ Perturbatively consistent theory of quantum gravity
(+ Gauge interactions + Supersymmetry)

Cf. talks by Conlon, Gray, Camara, Uranga

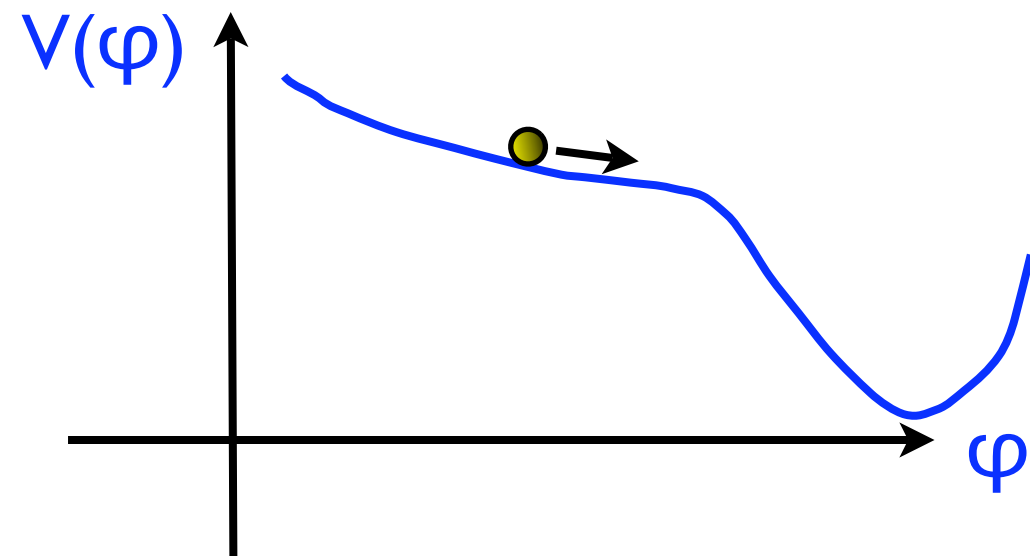
Why cosmology and string theory?

- Completeness of string theory:
String theory, if truly fundamental, should be able to
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- Completeness of string theory:
String theory, if truly fundamental, should be able to reproduce early Universe cosmology
- Cosmology may probe very high energy scales
⇒ Detailed dynamics may be highly UV-sensitive.
Prime example: The UV-sensitivity of inflation

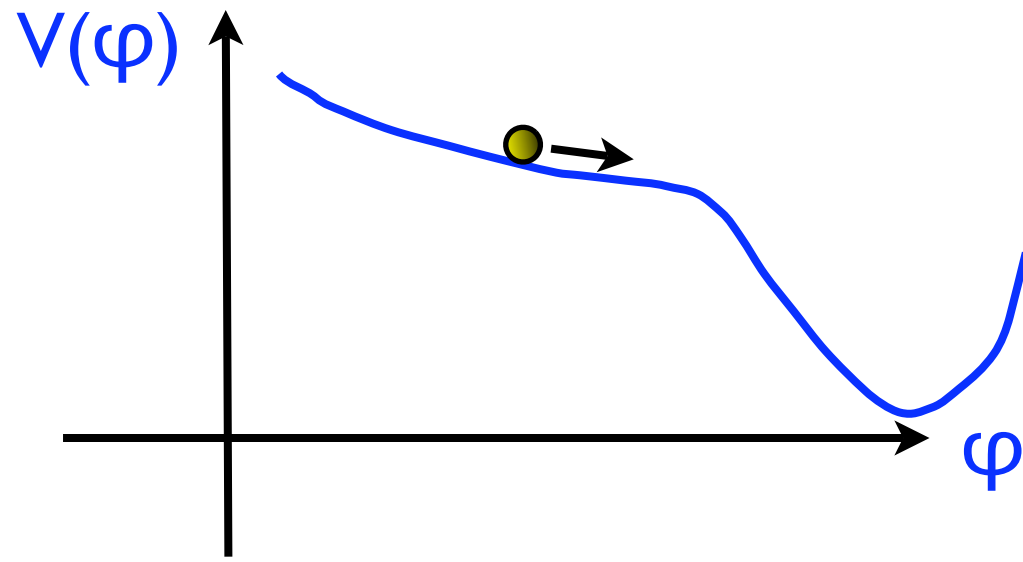
Slow-roll inflation:



$$\epsilon \equiv \frac{1}{2} \left(\frac{M_{\text{P}} V'}{V} \right)^2 \ll 1$$

$$|\eta| \equiv \left| \frac{M_{\text{P}}^2 V''}{V} \right| \ll 1$$

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- UV-sensitivity of inflation:

- η is sensitive to Planck suppressed operators:

$$\Delta V \sim \frac{\langle V \rangle \varphi^2}{M_{\text{P}}^2} \implies \Delta \eta = \mathcal{O}(1)$$

- Possibly very high energy scales

$$V_{\text{inf}}^{1/4} \sim M_{\text{GUT}} \epsilon^{1/4}$$

- Detectable tensor modes require large field excursions:

$$\Delta \varphi \sim M_{\text{P}}$$

- Certain “stringy” features may be relevant for cosmology

- Extra dimensions: $\mathcal{M}^{(10)} = \mathcal{M}^{(6)} \times \mathcal{M}^{(4)}$

Small & compact

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Small & compact

- Moduli:

4D scalars describing deformations of
compactification data

Cf. Joe Conlon’s talk

⇒ May cause phenomenological problems

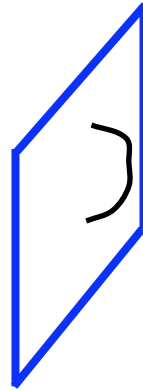
(E.g. 5th force, overclosure, BBN, ...)

⇒ Possibly good inflaton candidates

– Extended objects:



Fundamental string



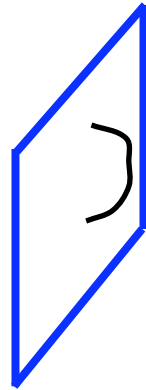
Dp-brane

⇒ Cosmological defects (e.g. cosmic (super)strings)?

- Extended objects:



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Dp-brane

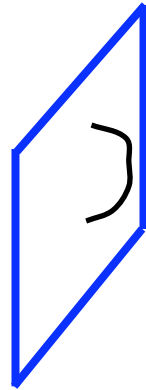
⇒ Cosmological defects (e.g. cosmic (super)strings)?

- Quantum gravity → Big Bang singularity?

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Dp-brane

⇒ Cosmological defects (e.g. cosmic (super)strings)?

- Quantum gravity → Big Bang singularity?

In any case: Constraints on string theory from cosmology are quite complementary to constraints from particle physics → Test string theory?

Rest of the talk:

1. Moduli stabilization
2. de Sitter vacua
3. Inflation
4. Cosmic (super)strings?
5. Tensor modes?
6. de Sitter and Inflation beyond IIB?
7. Conclusions

I. Moduli stabilization

$10D \rightarrow 4D \Rightarrow$ Moduli fields in $\mathcal{L}_{\text{eff}}^{4D}$

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From 6D components of 10D fields

$(g_{mn}, \phi, B_{mn}, C_{m\dots p})$

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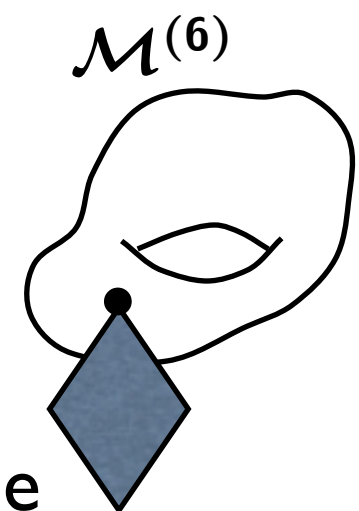
(i) Closed string moduli:

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(ii) Open string moduli:

E.g. D-brane positions/orientations

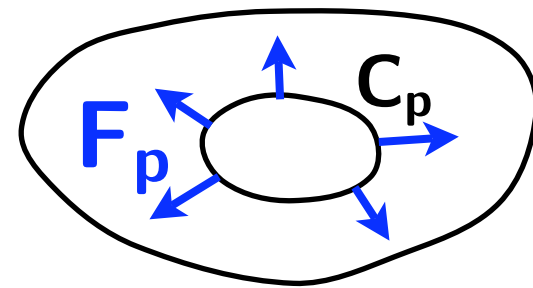


A simple way to avoid phenomenological problems from moduli: Make them sufficiently heavy

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Two important mechanisms:

(i) Fluxes of p-form field strengths F_p

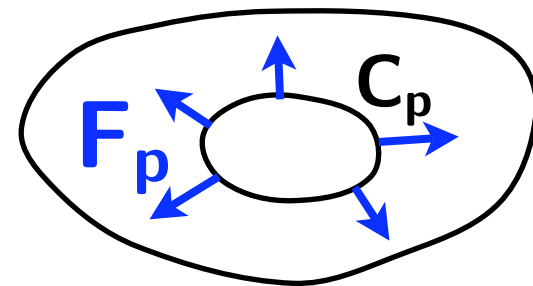


\Rightarrow Potential for C_p -deformation modulus

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Two important mechanisms:

(i) Fluxes of p-form field strengths F_p



\Rightarrow Potential for C_p -deformation modulus

(ii) Quantum corrections

\Rightarrow Most relevant for moduli **not** stabilized by fluxes

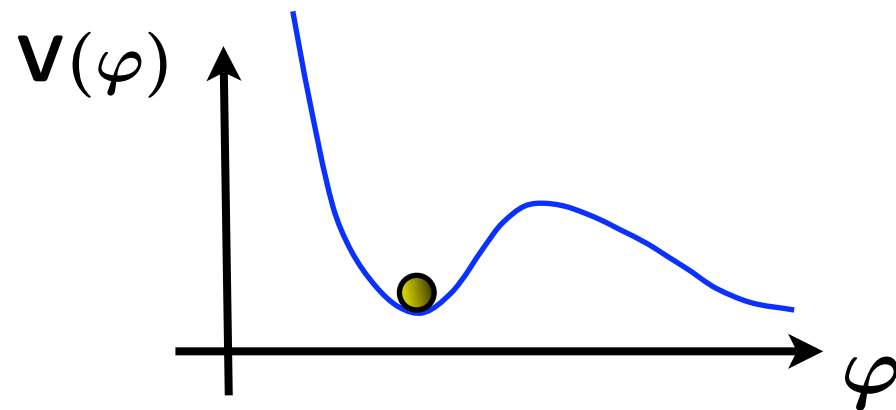
\Rightarrow Popular inflaton candidates

The interplay of fluxes and quantum corrections is presently best understood in type IIB string theory

⇒ So far, most work on string cosmology in type IIB

2. de Sitter vacua

Idea:



($\Lambda_{\text{eff}} > 0 \Rightarrow$ Today's accelerated expansion)

A popular scenario:

Giddings, Kachru, Polchinski (2001)
Kachru, Kallosh, Linde, Trivedi (2003)

IIB string theory on Calabi-Yau orientifolds
with 3-form fluxes & quantum corrections

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 - Complex structure moduli (“shape moduli”)
 - Dilaton (\rightarrow string coupling)

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IIB string theory on Calabi-Yau orientifolds
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- Stabilized by fluxes:
 - Complex structure moduli (“shape moduli”)
 - Dilaton (\rightarrow string coupling)
- Stabilized by subleading quantum effects:
 - Kähler moduli (“size moduli” \ni volume modulus)
(– D3-brane positions)

Volume stabilization:

$$W = W_{\text{flux}} + Ae^{-a\rho}$$

Gaugino cond. on D7-branes
or D3-brane instantons

$$K = -3 \ln[(\rho + \bar{\rho})]$$

$$\Rightarrow V_F^{\text{stab}} = e^K \left[\mathcal{D}_i W \overline{\mathcal{D}^i W} - 3|W|^2 \right] \Rightarrow \text{AdS}$$

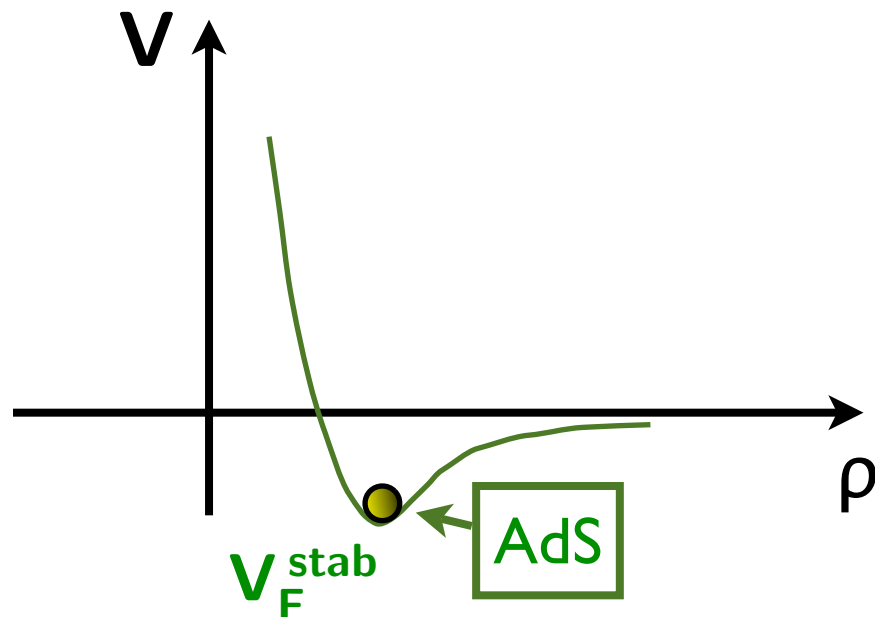
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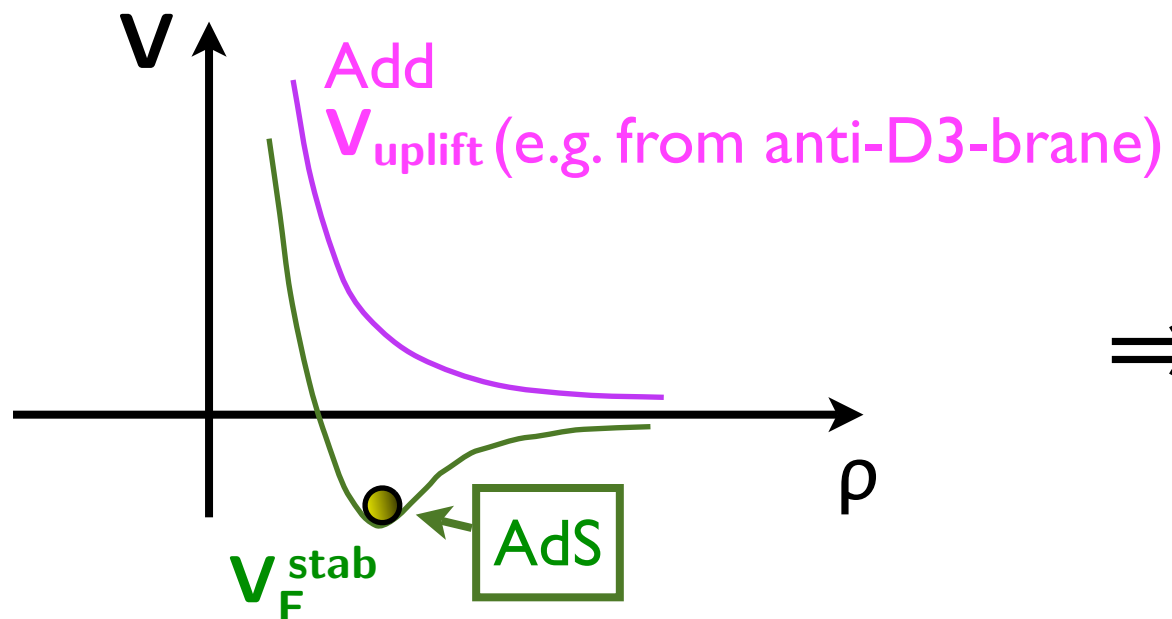
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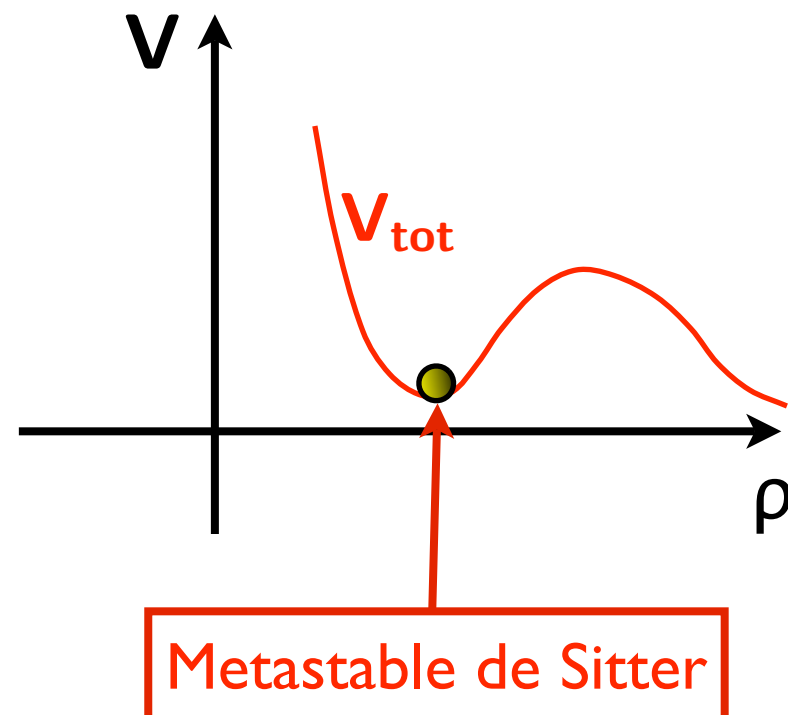
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\Rightarrow



A related interesting setup: Large volume scenario

Balasubramanian, Berglund (2004)

Balasubramanian, Berglund, Conlon, Quevedo (2005)

$$K = \text{tree level} + \alpha'$$

Becker, Becker, Haack, Louis (2002)

$$W = W_{\text{flux}} + \sum_i A_i e^{-a_i T_i}$$

Cf. Joe Conlon's talk

3. Inflation

Most scenarios:

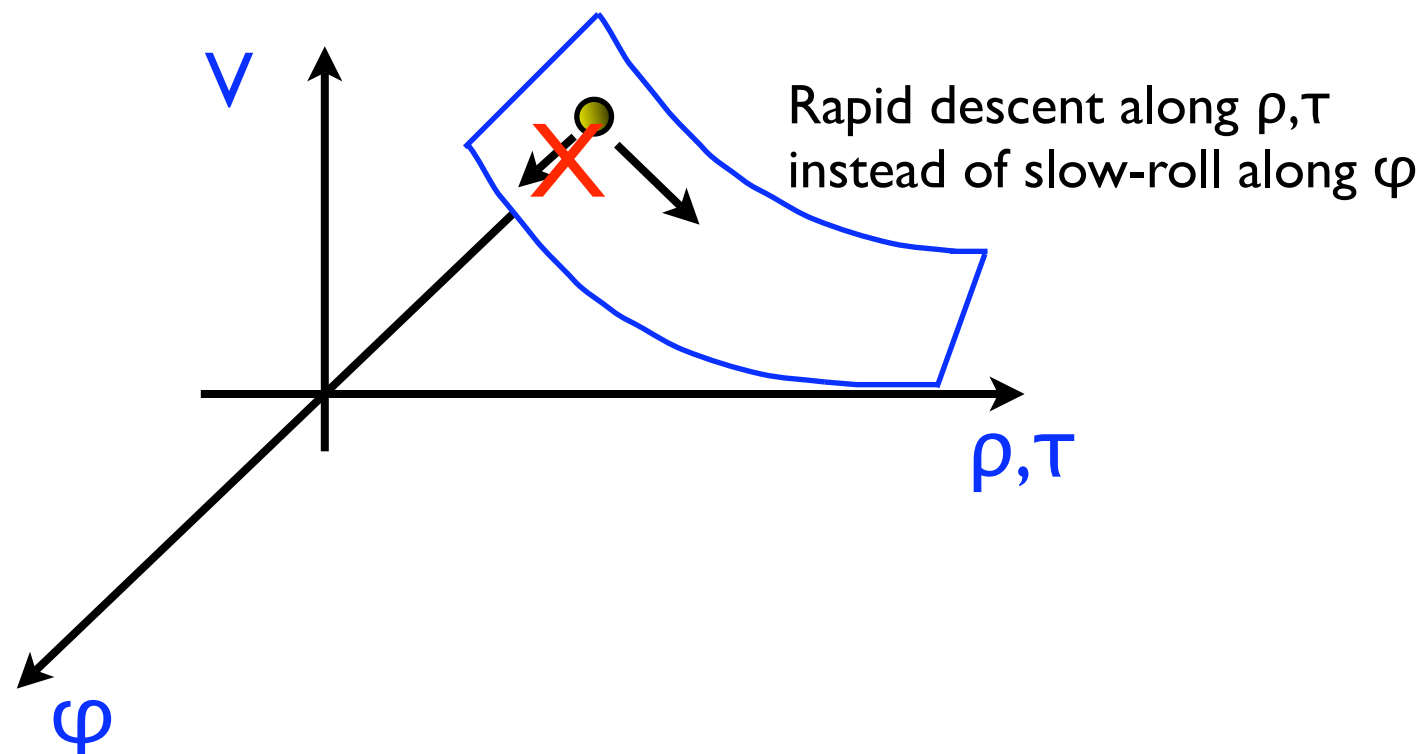
Inflaton = a modulus that receives a potential only at subleading order (quantum corrections)

Problem:

Some moduli tend to **interfere severely with inflation**

E.g. - Volume modulus ρ
- Dilaton τ

⇒ Tend to be **steep runaway directions**
of simplest potentials:





Inflation and moduli stabilization can not
be discussed separately

Most popular framework:

Type IIB string theory on CY-orientifolds with fluxes
(because moduli stabilization is so well understood)

Fixed by fluxes at tree-level:

- Complex structure moduli
- Axion-dilaton

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Remaining light moduli:

- D3-brane moduli (e.g. D3-positions on CY)
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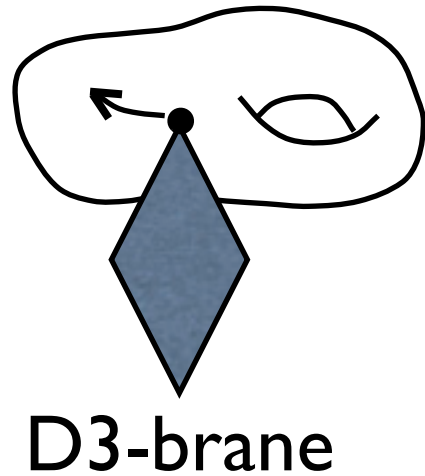
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Good inflaton candidates!

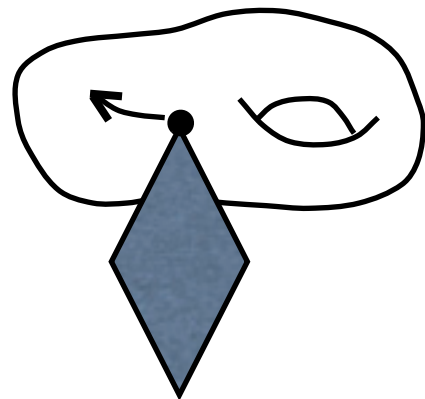
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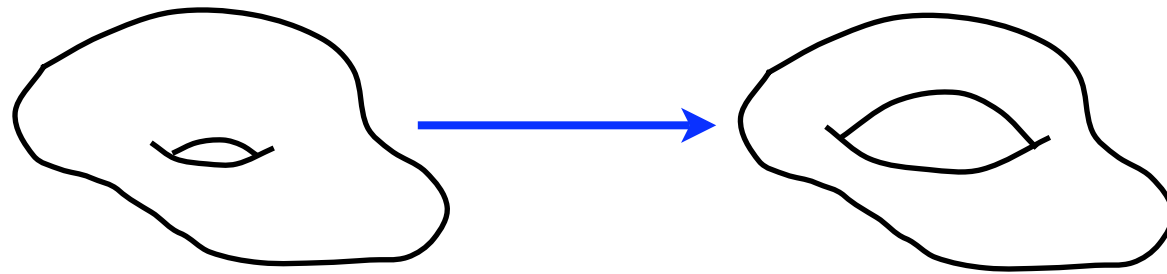


D3-brane

- (i) Warped $D3/\overline{D3}$ -inflation
- (ii) $D3/D7$ -inflation
- (iii) DBI-inflation
-
-
-

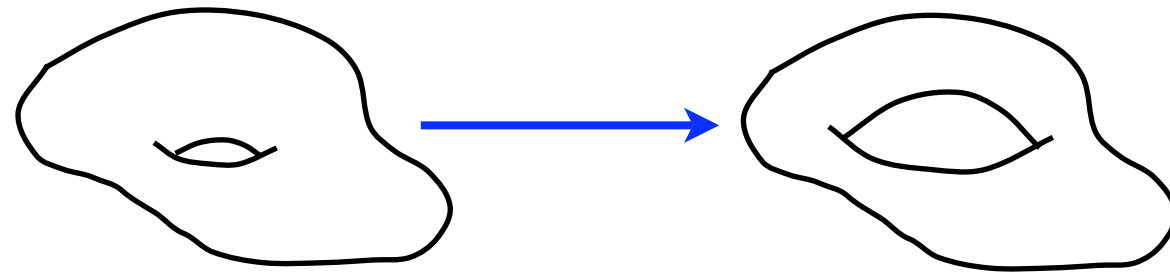
- Kähler moduli inflation models

Inflaton = a Kähler modulus
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(i) Racetrack inflation

(ii) Blow-up modulus inflation

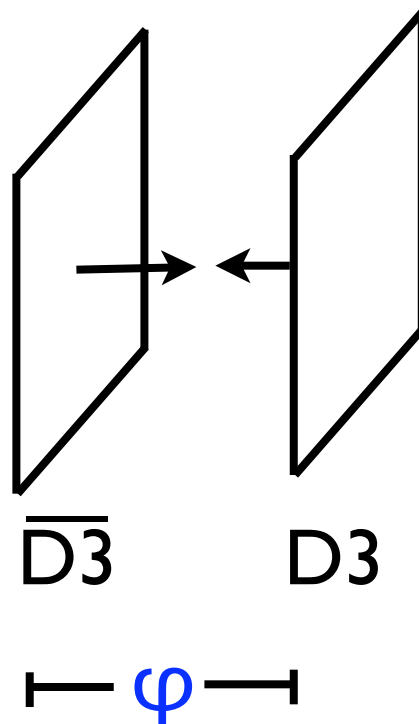
(iii) Fiber inflation

Also: Volume modulus inflation, axion inflation,...

Warped D3/ $\overline{\text{D3}}$ -brane inflation

Original idea:

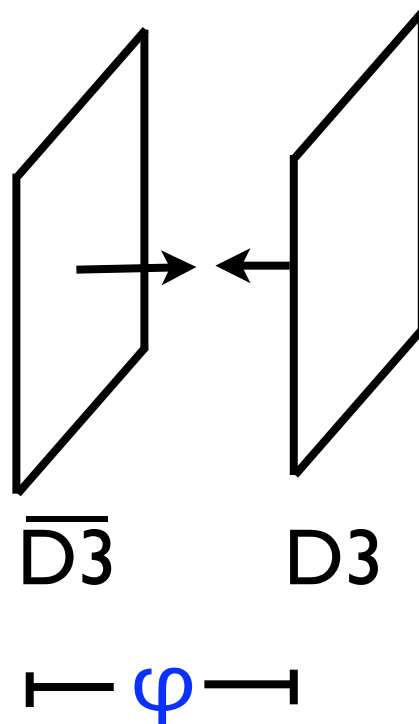
Dvali, Tye (1999),...



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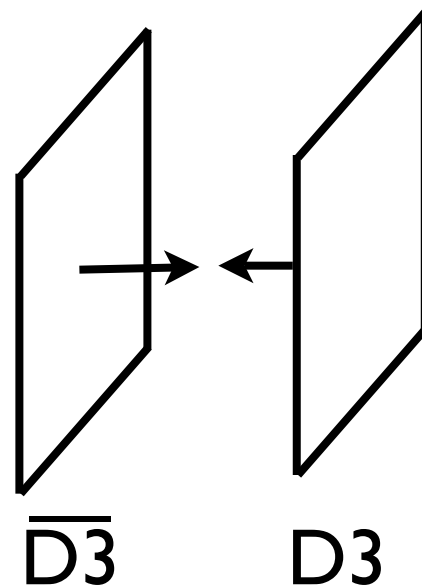


$$V_{\text{Coulomb}} \propto \left(1 - \frac{c}{\varphi^4}\right)$$

Warped D3/D $\overline{3}$ -brane inflation

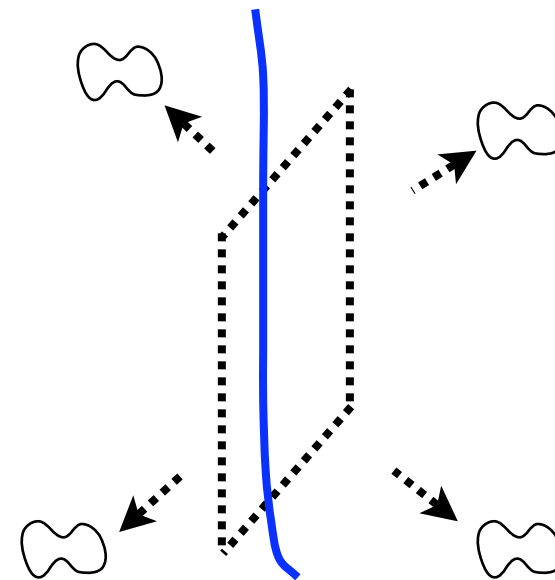
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φ

$$V_{\text{Coulomb}} \propto \left(1 - \frac{c}{\varphi^4}\right)$$



Annihilation into
closed string radiation
+ cosmic strings
(D1-branes)

Sen (1999)

Jones, Stoica, Tye (2002)

Sarangi, Tye (2002)

Dvali, Vilenkin (2002)

Problem: Finite separation φ in compact space

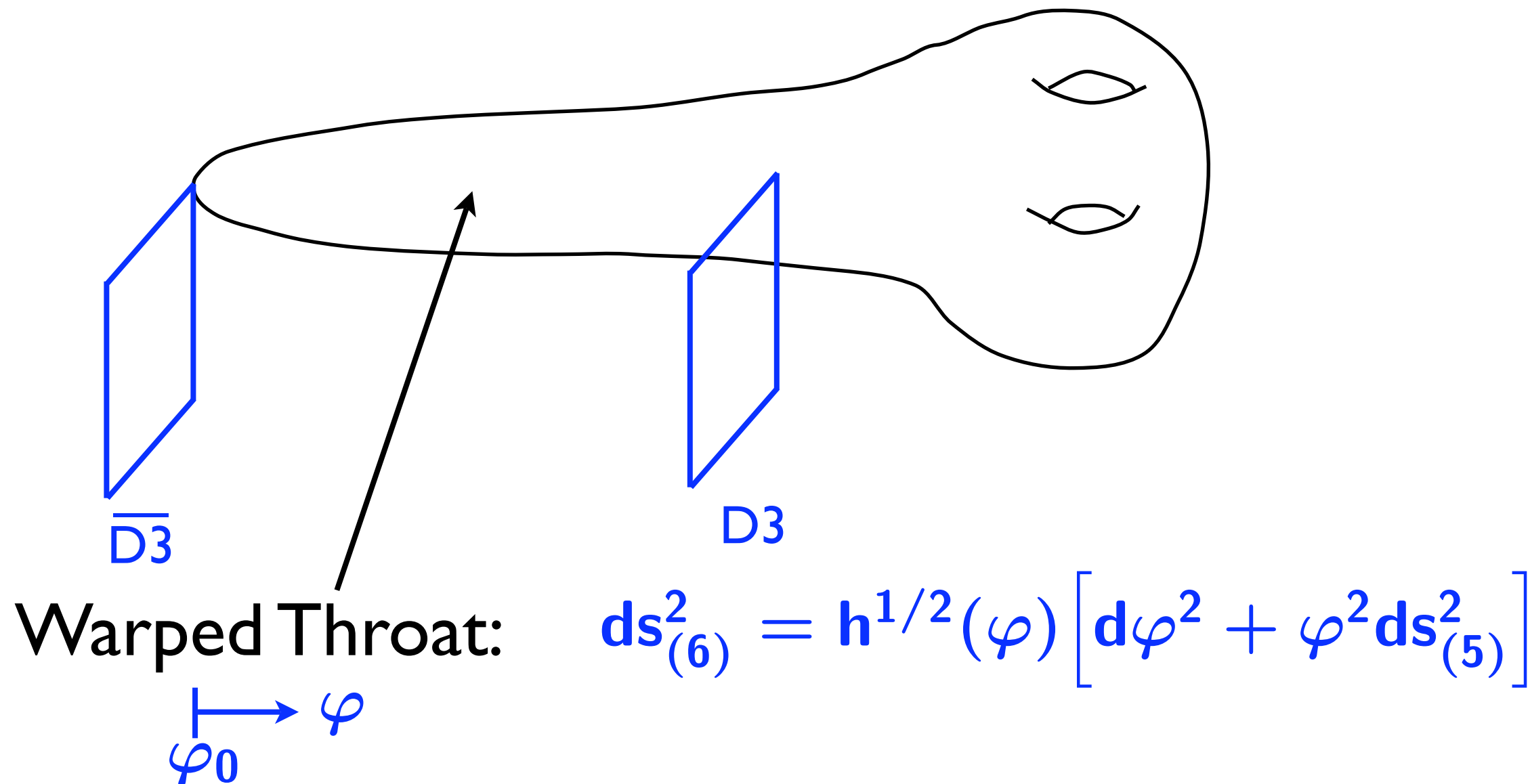
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\Rightarrow Warped D3/D3-brane inflation

KKLMMT (2003)



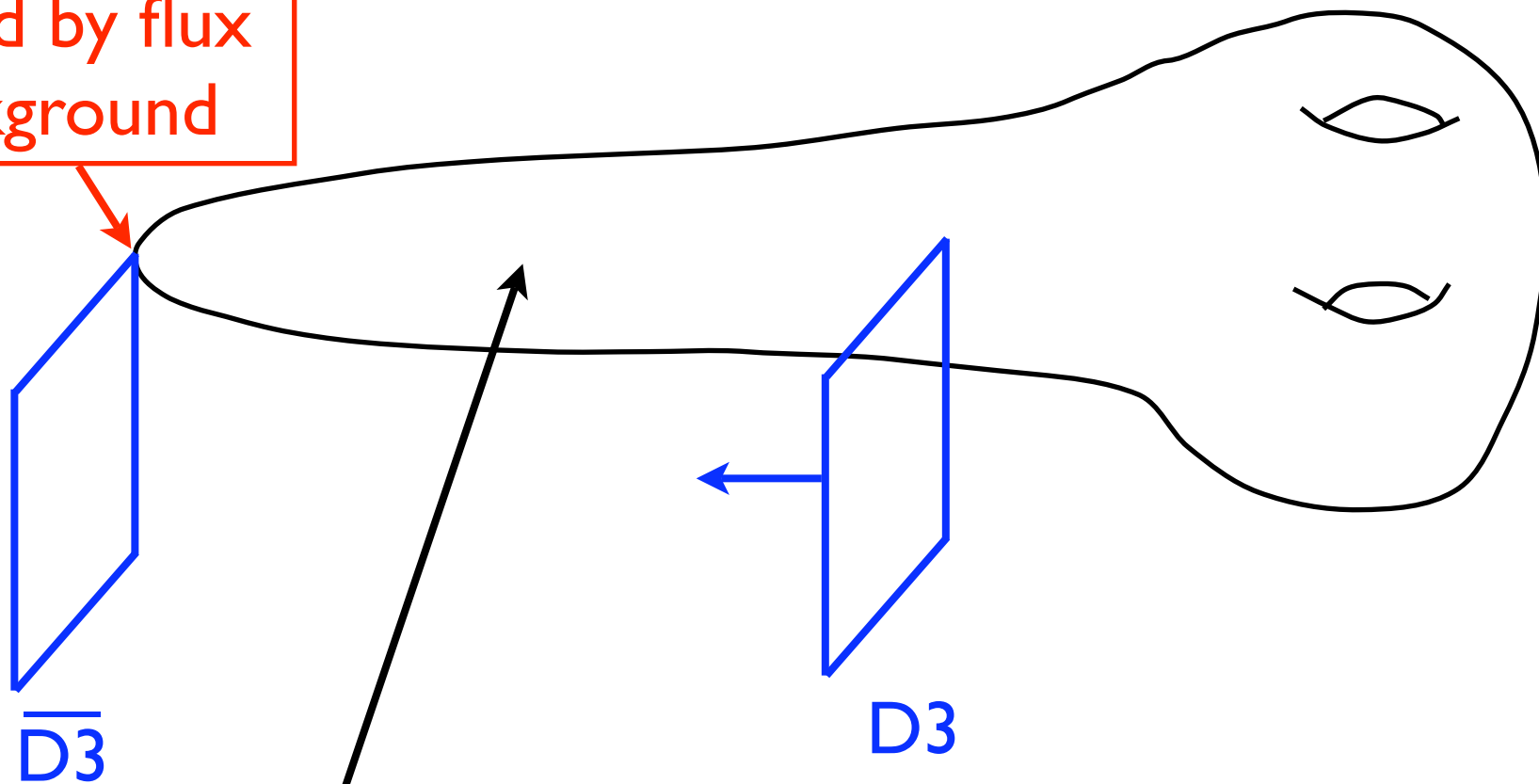
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KKLMMT (2003)

Fixed by flux
background



Warped Throat:

$$ds_{(6)}^2 = h^{1/2}(\varphi) \left[d\varphi^2 + \varphi^2 ds_{(5)}^2 \right]$$

$$\Rightarrow \mathbf{V}(\varphi)_{\text{Coulomb}} \propto \mathbf{h}^{-1}(\varphi_0) \left(\mathbf{1} - \frac{\mathbf{h}^{-1}(\varphi_0) \mathbf{C}}{\varphi^4} \right)$$

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$\eta \propto \mathbf{h}^{-1}(\varphi_0) \ll 1$ for strong warping $\mathbf{h}(\varphi_0) \gg 1$

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Additional benefit:

KKLMMT (2003)

Strong warping also redshifts the cosmic string tension

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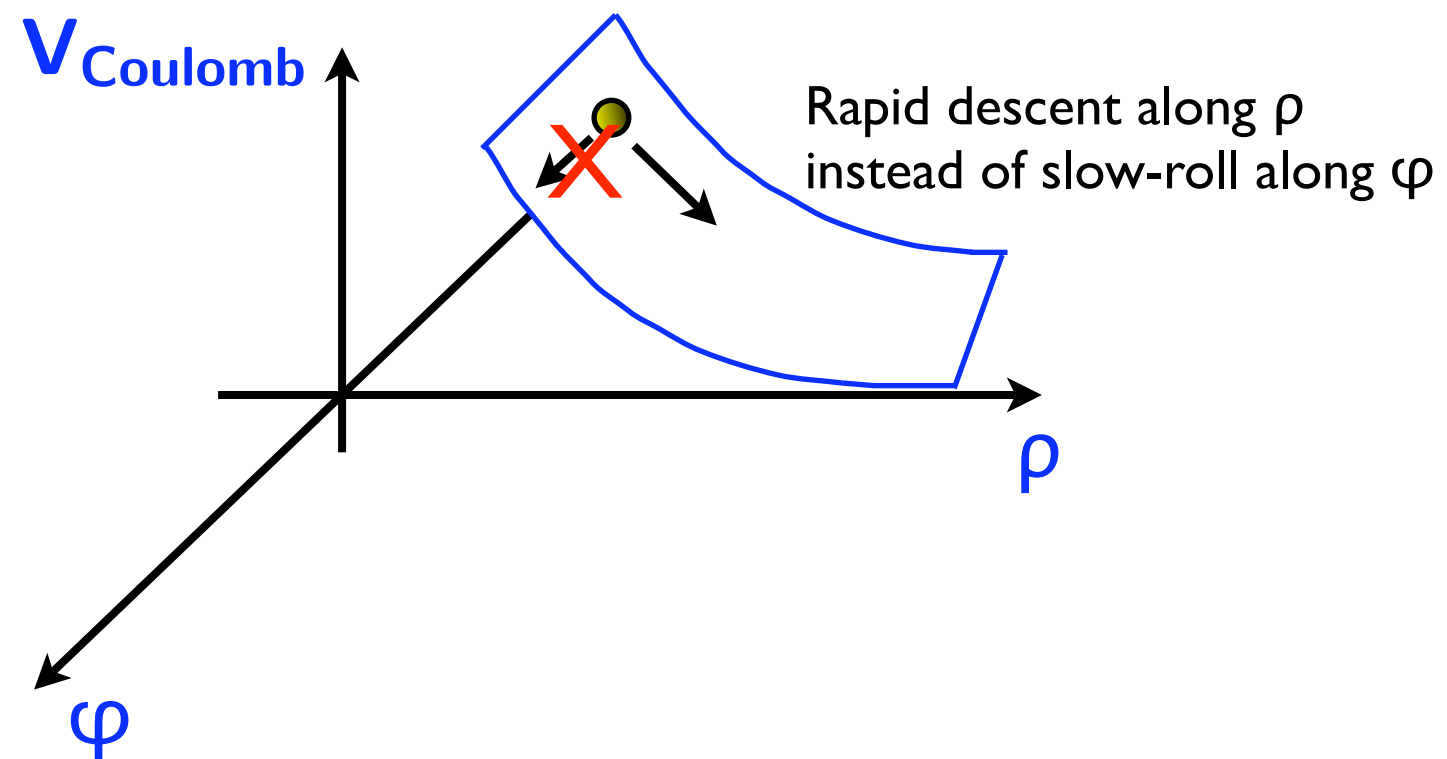
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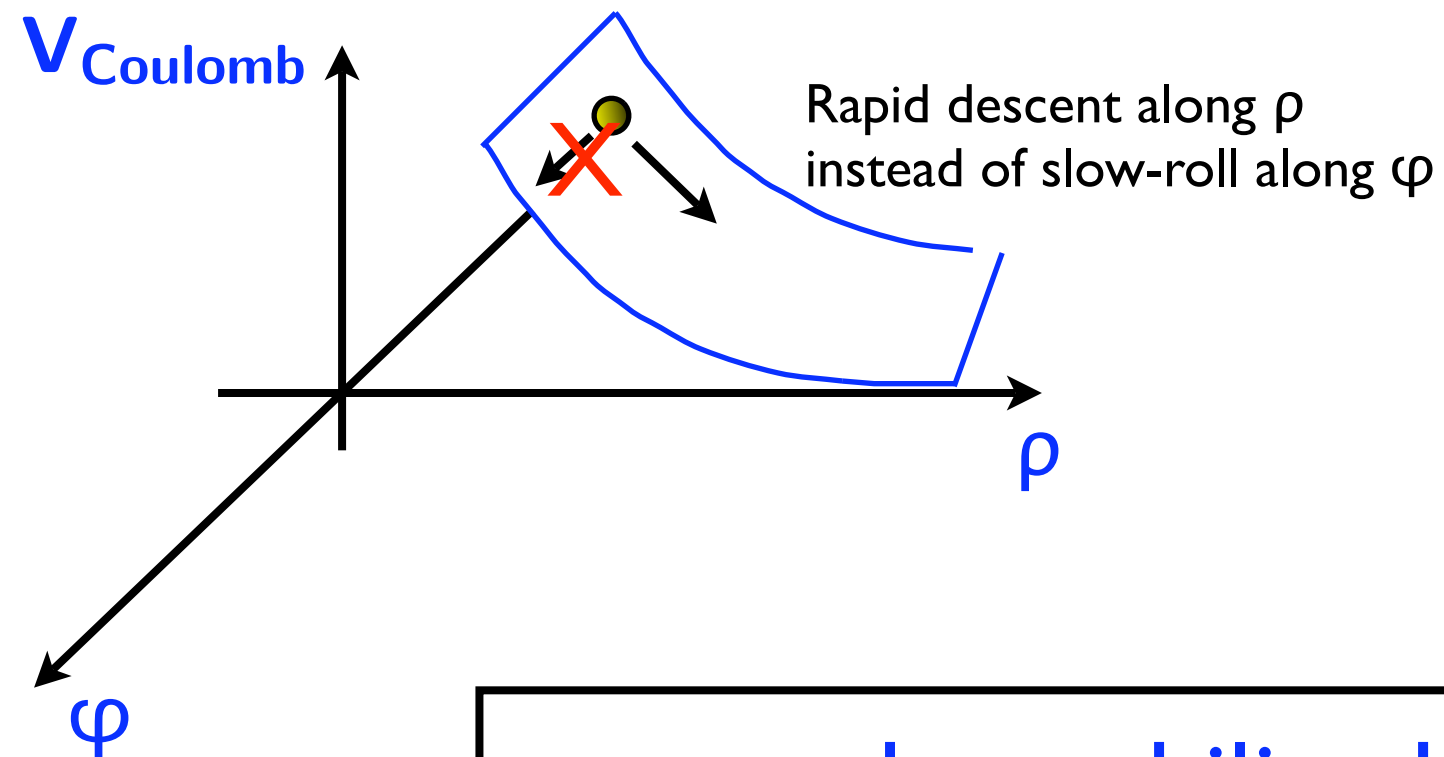
KKLMMT (2003)

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A remaining problem:

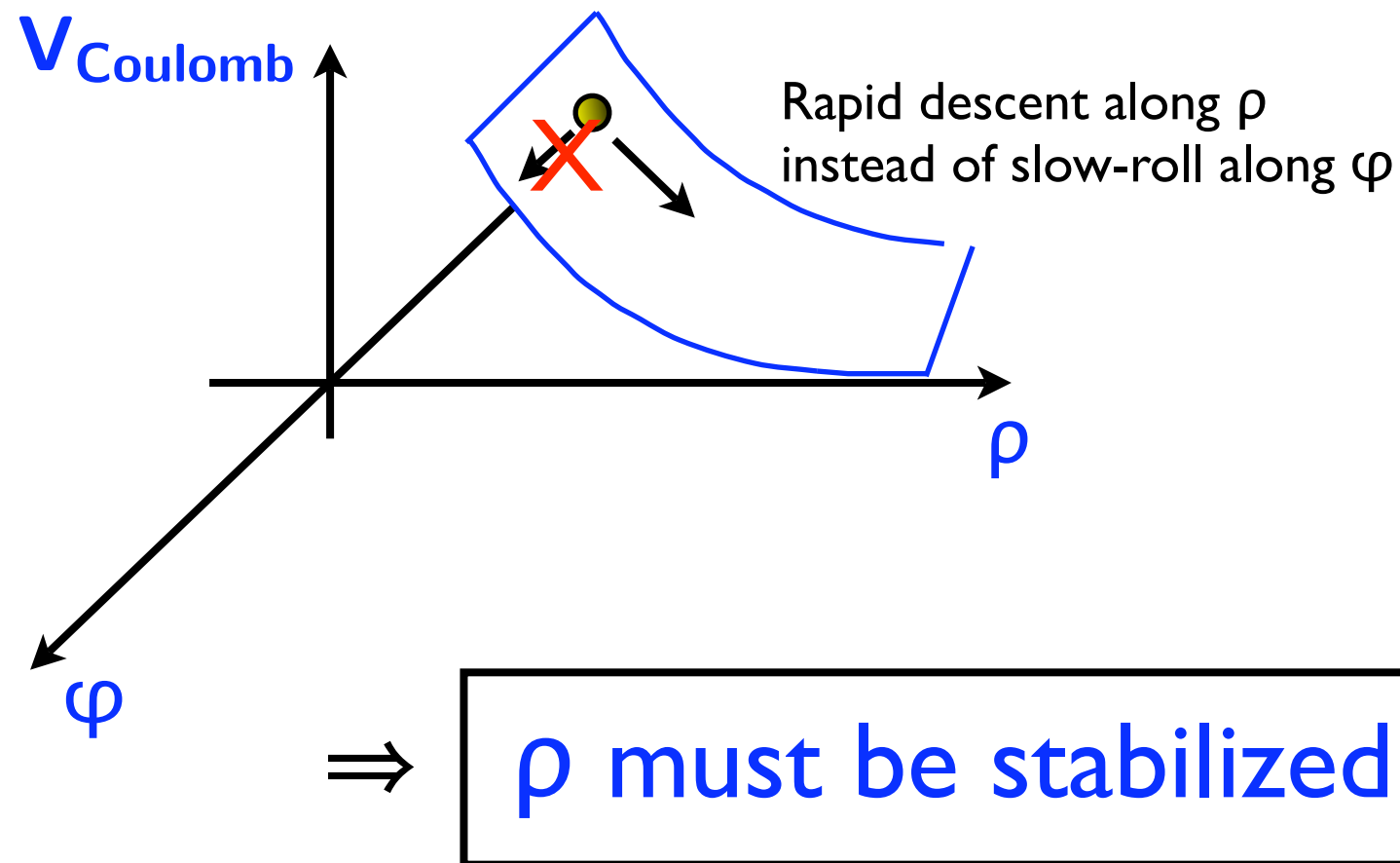
V_{Coulomb} is rapidly decreasing function
of volume modulus ρ





\Rightarrow

ρ must be stabilized !

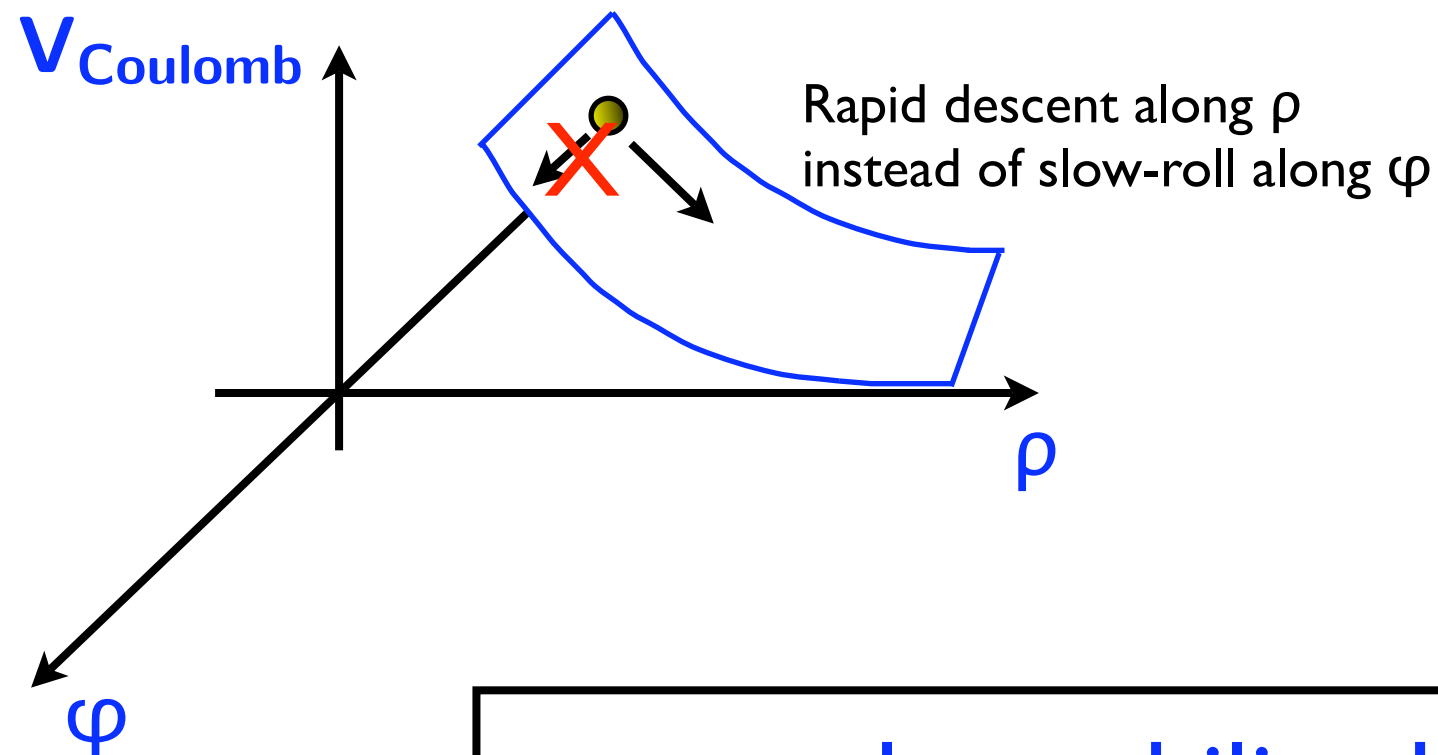


E.g. via nonperturbative superpotential: **KKLT (2003)**

$$W = W_{\text{flux}} + Ae^{-a\rho}$$

Gaugino cond. on D7-branes
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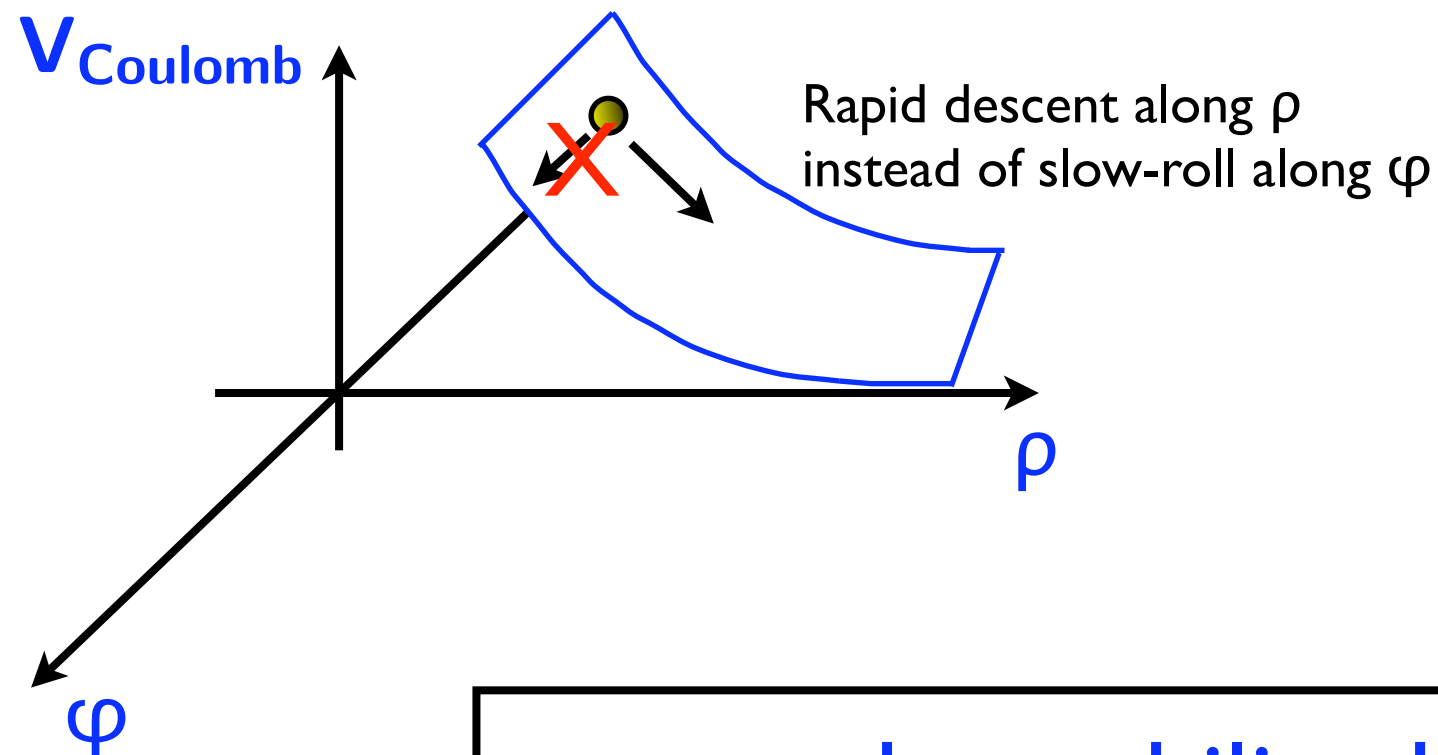
E.g. via nonperturbative superpotential: KKLT (2003)

In the presence
of D3-branes:

$$W = W_{\text{flux}} + A(\varphi, \dots) e^{-a\rho}$$

$$K = -3 \ln[(\rho + \bar{\rho} + |\varphi|^2)]$$

De Wolfe, Giddings (2002)



\Rightarrow

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E.g. via nonperturbative superpotential: KKLT (2003)

In the presence
of D3-branes:

$$W = W_{\text{flux}} + A(\varphi, \dots) e^{-a\rho}$$

$$K = -3 \ln[(\rho + \bar{\rho} + |\varphi|^2)]$$

\Rightarrow New inflaton dependences !

KKLMMT (2003)

\Rightarrow Compute $A(\varphi)$:

Backreaction of D3 on warp factor (Tree-level supergravity calculation)

Giddings, Maharana (2005)

Baumann, Dymarsky, Klebanov, Maldacena, McAllister, Murugan (2006)

Baumann, Dymarsky, Klebanov, McAllister (+ Steinhardt) (2007)

Krause, Pajer (2007)

Baumann, Dymarsky, Kachru, Klebanov, McAllister (2008)

One finds:

Generically $\varphi^{3/2}$ -terms

\Rightarrow Inflection point inflation scenarios

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Baumann et al. (2008)

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One lesson:

Detailed computational control is important

4. Cosmic (super)strings

Brane inflation
models



Typically cosmic strings
at the end of inflation

Closed string
inflation models



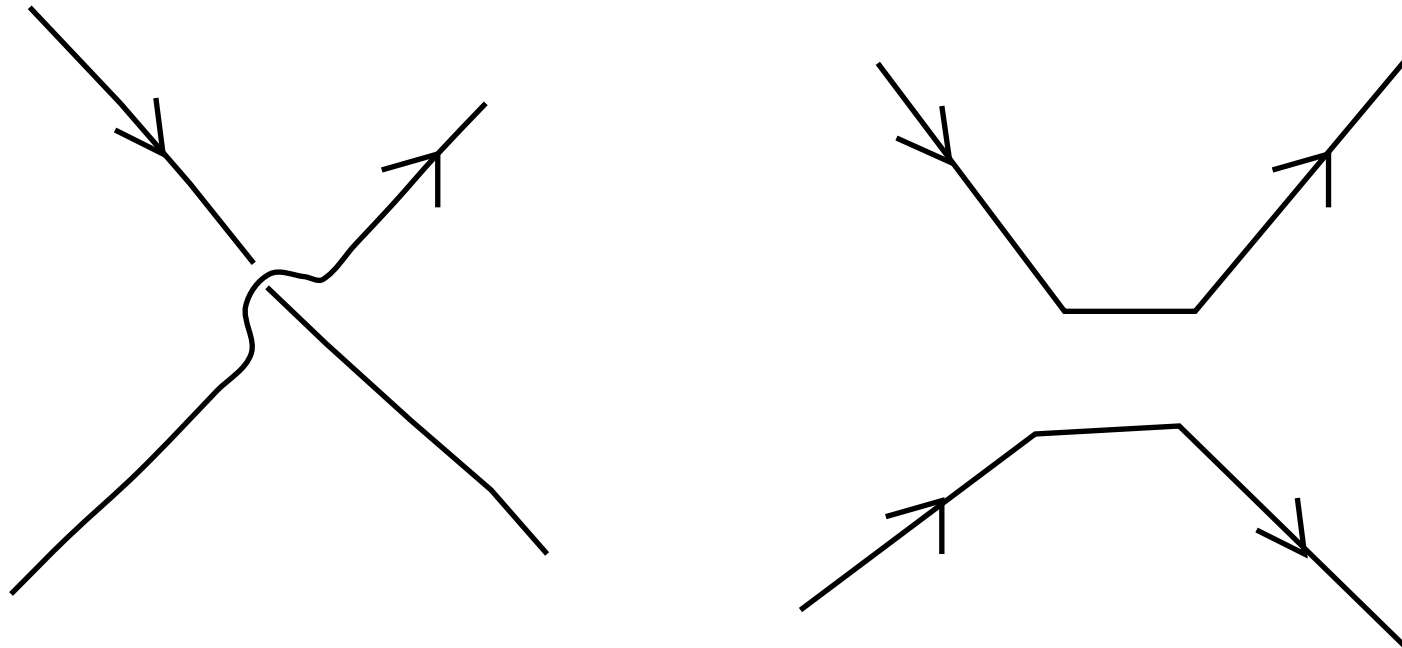
Typically no
cosmic strings

- D1-branes
- Fundamental strings
- (p,q)-strings
- wrapped Dp-branes ($p > 1$)

May have **different properties** than ordinary
gauge theory solitons

E.g. Copeland, Myers, Polchinski (2003)

E.g.:



Reconnection probability:

$P=1$ for standard gauge theory vortices

$P \ll 1$ possible for cosmic superstrings

\Rightarrow

Cosmic strings could be very interesting
window into string theory **if realized in nature**

5. Observable tensor modes?

$$\left(\frac{\Delta\varphi}{M_{\text{P}}} \right) \sim \left(\frac{r}{0.01} \right)^{\frac{1}{2}}$$

Lyth bound

Lyth (1996)

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Lyth bound

Lyth (1996)

Strong geometrical constraints for D3-inflation models:

- D3 on symmetric torus
- D3 in warped throat

E.g. Baumann, McAllister (2006)

⇒

r unobservably small

Some recent suggestions:

- Wrapped Dp-branes

Kobayashi, Mukohyama, Kinoshita (2007)

Becker, Leblond, Shandera (2007)

- Large complex structure

Haack, Kallosh, Krause, Linde, Lüst, M.Z. (2008)

- Monodromy inflation

Silverstein, Westphal (2008)

McAllister, Silverstein, Westphal (2008)

- Fiber inflation

Cf. Cicoli's talk

- Wilson line DBI inflation

Cf. Zavala's talk

- Multi axion inflation models

Dimopoulos, Kachru, McGreevy, Wacker (2005)

Easter, McAllister (2005)

Kallosh, Sivanandam, Soroush (2007)

Grimm (2007)

Tye, Xu, Zhang (2008)

- Multibrane models

E.g. Becker, Becker, Krause (2005)

Krause (2007)

6. Inflation beyond type IIB?

Not much explored

E.g. Becker, Becker, Krause (2005)

Moduli stabilization best understood in IIB

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Moduli stabilization best understood in IIB

But:

In a sense, moduli stabilization is **even simpler in IIA**

Grimm, Louis (2004)

Kachru, Kashani-Poor (2004)

Derendinger, Kounnas, Petropoulos, Zwirner (2004, 2005)

Villadoro, Zwirner (2005)

de Wolfe, Giryavets, Kachru, Taylor (2005)

Type IIA on Calabi-Yau spaces with

- p-form fluxes
- D6-branes/O6-planes

All geometric moduli may be stabilized
in parameterically controlled classical regime
 \Rightarrow Quantum corrections negligible

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\Rightarrow Very controlled inflation models?

Unfortunately:

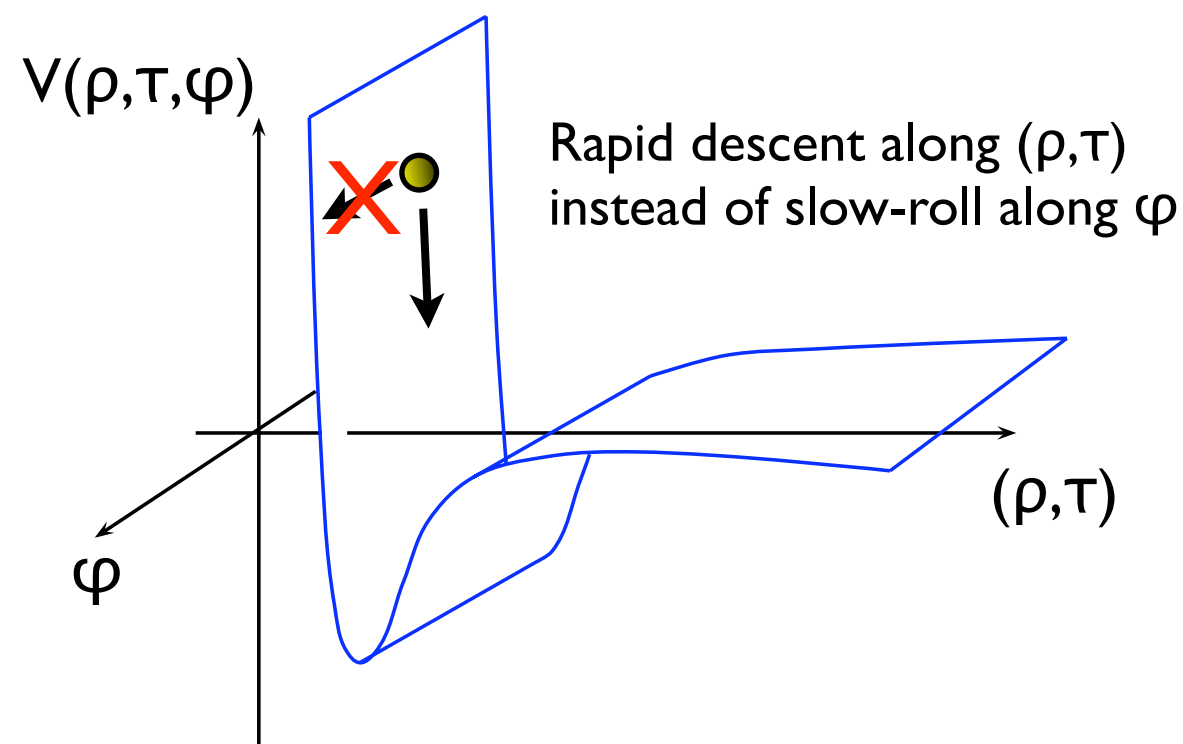
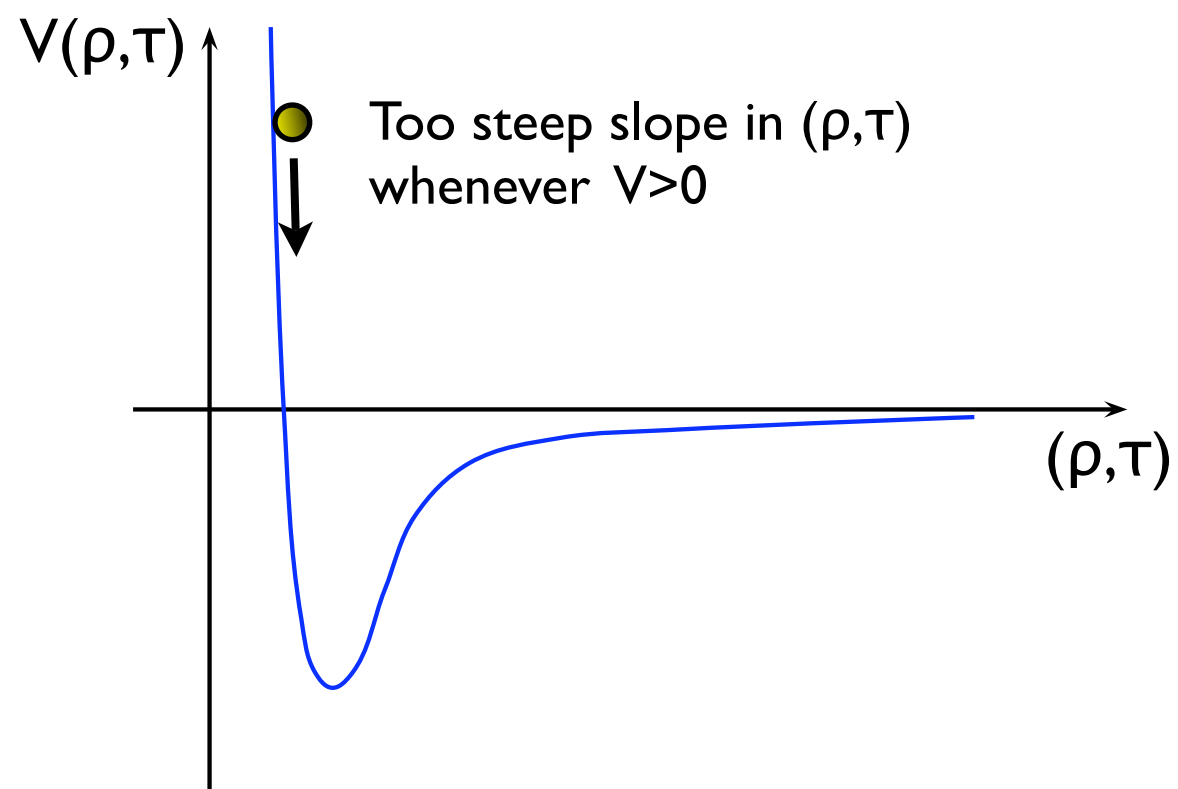
No-go theorem:

Classical IIA compactifications with

- $\mathcal{M}^{(6)} =$ Calabi-Yau (\rightarrow Ricci-flatness)
- O6/D6 sources
- p-form fluxes (incl. Romans' mass)

\Rightarrow No de Sitter vacua and no slow-roll inflation !

Hertzberg, Kachru, Taylor, Tegmark (2007)



$$\epsilon \geq 27/13$$

A simple way to circumvent the no-go:

Abandon Ricci-flatness of compact space

Several recent attempts:

- (i) $\mathcal{M}^{(6)} = (\text{Nil}_3 \times \text{Nil}'_3) / \mathcal{O}$
Haque, Shiu, Underwood, Van Riet (2008)
- (ii) $\mathcal{M}^{(6)} = \text{Cosets with SU(3)-structure}$
Caviezel, Koerber, Körs, Lüst, Wrase, M.Z. (2008)
- (iii) $\mathcal{M}^{(6)} = \text{More general twisted tori}$
Flauger, Paban, Robbins, Wrase (2008)

But:

New No-go theorems along different
moduli directions or tachyons

7. Conclusions

- Early Universe cosmology is UV-sensitive (E.g. inflation) \Rightarrow String theory
- Moduli stabilization cannot be ignored
- Best studied models in IIB (de Sitter vacua & inflation)
- Cosmic superstrings?
- Detectable tensor modes? Many recent suggestions
- Inflation elsewhere in the landscape? So far unsuccessful in IIA