# Ultrarelativistic Heavy Ion Physics: Theory

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## From elementary interactions to collective phenomena



How do collective phenomena and macroscopic properties of matter emerge from fundamental interactions?



QCD much richer than QED:

- non-abelian theory
- degrees of freedom change with  $Q^2$

# Question:Why do we need collider energies $\sqrt{s_{NN}} = 200 \, GeV \quad [RHIC]$ $\sqrt{s_{NN}} = 5500 \, GeV \quad [LHC]$ test properties of dense OCD matter

to test properties of dense QCD matter which arise on typical scales

 $T \approx 150 \, MeV, \quad Q_s \approx 1 - 2 \, GeV$ ?

# Answer 1: Large quantitative gains

Increasing the center of mass energy implies

- Denser initial system
- Longer lifetime
- Bigger spatial extension
- Stronger collective phenomena

A large body of experimental data from the CERN SPS and RHIC supports this argument.

# Elliptic Flow: Hallmark of a collective phenomenon





# Elliptic flow vs. hydrodynamic simulations

#### Assumptions:

- fluid with shear viscous term

$$T^{\mu\nu} = (\varepsilon + p) u^{\mu} u^{\nu} - p g^{\mu\nu} + shear$$

- Bjorken boost invariance
- 'realistic' equation of state
- 'realistic' initial conditions
- 'realistic' decoupling (freeze-out)

#### <u>Results</u>:

- initial transverse pressure gradient

 $\stackrel{\bullet}{\Longrightarrow} \phi \text{ - dependence of flow field } u_{\mu} \\ \stackrel{\bullet}{\Longrightarrow} \text{ elliptic flow } v_2(p_T)$ 

- size and pt-dependence of  $v_2$  data accounted for by fluid of  $\underline{\eta}_{<<1}$
- characteristic mass dependence, since all particle species emerge from common flow field  $u_{\mu}$



p<sub>r</sub>[GeV]

# <u>Viscosity: Bounds from theory</u>

- Viscosity  $\eta$  controls entropy s increase Hydrodynamics is valid, if  $\frac{d(\tau s)}{d\tau} = \frac{\frac{4}{3}\eta}{\tau T}$
- $\frac{\eta}{\tau T} \frac{1}{s} \ll 1$

Constraint from string theory



Same trend in exploratory **QCD** lattice calculation

H. Meyer, arXiv:0805.4567 [hep-lat]

# LHC 1<sup>st</sup> year running tests hallmark of collectivity



# Answer 2: Qualitatively novel access to properties of dense matter

To test properties of QCD matter, large-  $Q^2$  processes provide wellcontrolled tools (<u>example: DIS</u>).

Heavy Ion Collisions produce <u>auto-generated probes</u> at high  $\sqrt{s_{NN}}$ 



Q: How sensitive are such 'hard probes'?

# Bjorken's original estimate and its correction

Bjorken 1982: consider jet in p+p collision, hard parton interacts with underlying event <u>collisional energy loss</u>

 $dE_{coll}/dL \approx 10 \, GeV/fm$  (error in estimate!)

Bjorken conjectured monojet phenomenon in proton-proton



But: radiative energy loss expected to dominate

$$\Delta E_{rad} \approx \alpha_s \hat{q} L^2$$

Baier Dokshitzer Mueller Peigne Schiff 1995

- p+p:  $L \approx 0.5 \text{ fm}, \Delta E_{rad} \approx 100 \text{ MeV}$  Negligible !
- A+A:  $L \approx 5 \text{ fm}, \Delta E_{rad} \approx 10 \text{ GeV}$

Monojet phenomenon! Observed at RHIC

# Parton energy loss - a simple estimate



Medium characterized by transport coefficient:

$$\hat{q} \equiv \frac{\mu^2}{\lambda} \propto n_{density}$$

• How much energy is lost ?

Phase accumulated in medium:

$$\left< \frac{k_T^2 \Delta z}{2\omega} \right> \approx \frac{\hat{q}L^2}{2\omega} = \frac{\omega_c}{\omega}$$

Characteristic gluon energy

Number of coherent scatterings: 
$$N_{coh} \approx$$

 $\approx \frac{t_{coh}}{\lambda}, \text{ where } t_{coh} \approx \frac{2\omega}{k_T^2} \approx \sqrt{\omega/\hat{q}} \\ k_T^2 \approx \hat{q} t_{coh}$ 

Gluon energy distribution:

Average energy loss

$$\omega \frac{dI_{med}}{d\omega \, dz} \approx \frac{1}{N_{coh}} \omega \frac{dI_1}{d\omega \, dz} \approx \alpha_s \sqrt{\frac{\hat{q}}{\omega}}$$

$$\Delta E = \int_0^L dz \int_0^{\omega_c} d\omega \,\omega \frac{dI_{med}}{d\omega \, dz} \sim \alpha_s \omega_c \,\sqrt{\alpha_s q L^2}$$

# High p<sub>T</sub> Hadron Spectra



$$R_{AA}(p_T,\eta) = \frac{dN^{AA}/dp_T d\eta}{n_{coll} dN^{NN}/dp_T d\eta}$$

Centrality dependence:





# @ RHIC, radiative energy loss accounts for:

**H** 1.2

- Nuclear modification factor
- Centrality dependence
- Back-to-back correlations
- R<sub>AA</sub> = 0.2 is a natural limit due to <u>surface emission</u>



indicates very opaque medium.

• Particle species (in)dependence

 $\hat{q} = 0$ , no medium  $\hat{q} = 0$ , no medium  $\hat{q} = 1 \text{ GeV}^2/\text{fm}$   $\hat{q} = 1 \text{ GeV}^2/\text{fm}$   $\hat{q} = 5 \text{ GeV}^2/\text{fm}$   $\hat{q} = 5 \text{ GeV}^2/\text{fm}$   $\hat{q} = 10,15 \text{ GeV}^2/\text{fm}$   $\hat{q} = 10,15 \text{ GeV}^2/\text{fm}$   $\hat{q} = 10,15 \text{ GeV}^2/\text{fm}$  $\hat{q} = 10,15 \text{ GeV}^2/\text{fm}$ 

NPA747 (2005) 511

Eskola, Honkanen, Salgado, Wiedemann

PHENIX **π**<sup>0</sup>

PHOBOS h BRAHMS h

STAR h

Many aspects still under debate:

- role of other e-loss mechanisms?
- Suppression of heavy flavored hadrons?
- theoretical basis and numerical consistency of model parameters such as  $\hat{q}$

# "True" Jets only at the LHC





# Jet modifications in reach @ LHC ...

• <u>'Longitudinal Jet heating':</u> The entire longitudinal jet multiplicity distribution softens due to medium effects.

Borghini, Wiedemann, hep-ph/0506218



 Jets <u>'blown with the wind'</u> Hard partons are not produced in the rest frame comoving with the medium

Armesto, Salgado, Wiedemann, Phys. Rev. Lett. 93 (2004) 242301

# Jet Finding Algorithms



## Jet shapes, energy flows

• Thrust – baseline



Thrust – medium above baseline







#### Recent jet quenching MCs:

JEWEL (K.Zapp et al.) Q-Pythia (Santiago group) YaJEM (T.Renk)

## Monte Carlo including LPM-effect

<u>Needed</u>: probabilistic implementation of quantum interference

K. Zapp, J. Stachel, UAW arXiv:0812.3888

<u>Solution</u>: in vacuum => angular ordering only <sup>arXiv:0812.3888</sup> in medium => formation time constraint implements LPM



# Perturbative vs. non-perturbative description

 Can the properties of dense QCD matter produced in HICs be described in terms of a quasi-particle picture?
 QCD thermodynamics is qualitatively different for T > 2-3 T<sub>c</sub>



- What is the reason for the success of AdS/CFT inspired calculations in RHIC phenomenology?
  - reproducing shear viscosity, quenching parameter

# QCD vs. AdS/CFT

	<i>N</i> =4 SYM in vacuum	<i>QCD</i> in vacuum	<i>N</i> =4 SYM finite T	<i>QCD</i> at finite T
<ul> <li>chiral condensate</li> </ul>	NO	YES	NO	melted
<ul> <li>asymptotic freedom</li> </ul>	NO	YES	NO	less important
confinement	NO	YES	NO	deconfined
<ul> <li>superymmetric</li> </ul>	YES	NO	broken	NO
<ul> <li>conformal</li> </ul>	YES	NO	YES	approx at T >>T <sub>c</sub>
<ul> <li>degrees of freedom</li> </ul>			Very different, but may be taken care of by normalization	
	Physics near vac high energy is ve	uum and at ry different		
EXP: Conceptual questions in reach of LHC	• qu	uasi-particles	NO	Conceivable for T>2-3 Tc only
tools are best suited?	• 8	z-3p	Zero	Small for T>2-3 Tc only

@ RHIC, suppression of hadron spectra is strong (~ 5-fold) and unattenuated up to highest p<sub>T</sub> (~ 20 GeV), collectivity is strong and rising with cms energy

### => <u>Collectivity expected to be even larger a LHC</u>,

### suppression expected to persists in wide pT-range at LHC

## The probes:

- Jets
- identified hadron specta
- D-,B-mesons
- Quarkonia
- Photons
- Z-boson tags

Abundant yield

of hard probes

+ <u>robust</u> signal

(medium sensitivity >> uncertainties)

= <u>detailed understanding</u> of dense QCD matter

