# Diffractive charm production from the dipole model analysis

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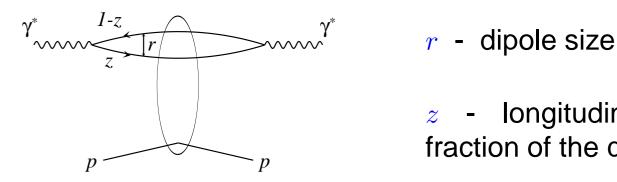
based on Phys.Rev.D79:114010,2009

### **Outline**

- Dipole model approach GBW and CGC parametrizations
- Diffractive structure functions in dipole models
- Diffractive charm quark production
- Comparision with HERA data
- Conclusions

## **Dipole model of DIS**

Dipole picture of DIS at small x in the proton rest frame



- z longitudinal momentum fraction of the quark/antiquark

Factorization: dipole formation + dipole interaction

$$\sigma^{\gamma p} = \frac{4\pi^2 \alpha_{em}}{Q^2} F_2 = \sum_f \int d^2 r \int_0^1 dz \, |\Psi^{\gamma}(r, z, Q^2, m_f)|^2 \, \hat{\sigma}(r, x)$$

Dipole-proton interaction  $\hat{\sigma}(r,x)$  is parameterized.

### **GBW** parametrization

(Golec-Biernat Wusthoff,99)

GBW parametrization

$$\hat{\sigma}(r,x) = \sigma_0 \left( 1 - \exp(-r^2/R_s^2) \right), \qquad R_s^2 = 4 \cdot (x/x_0)^{\lambda} \text{ GeV}^2$$

The dipole scattering amplitude in such a case reads

$$\hat{N}(\mathbf{r}, \mathbf{b}, x) = \theta(b_0 - b) \left( 1 - \exp(-r^2/R_s^2) \right)$$

where

$$\hat{\sigma}(r,x) = 2 \int d^2b \, \hat{N}(\mathbf{r}, \mathbf{b}, x)$$

Parameters  $b_0$ ,  $x_0$  and  $\lambda$  from fits of scattering amplitude to  $F_2$  data

$$\lambda = 0.288$$
  $x_0 = 4 \cdot 10^{-5}$   $\sigma_0 = 2\pi b_0^2 = 29 \text{ mb}$ 

### **CGC** parametrization

(lancu, Itakura, Munier, Soyez 04-07)

The dipole scattering amplitude in this parametrization

$$\hat{N}(\mathbf{r}, \mathbf{b}, x) = S(\mathbf{b}) N(\mathbf{r}, x)$$
  $S(\mathbf{b}) = \exp(-b^2/(2B_d))$ 

• Dipole cross section is given by  $(Q_s = 1/R_s)$ 

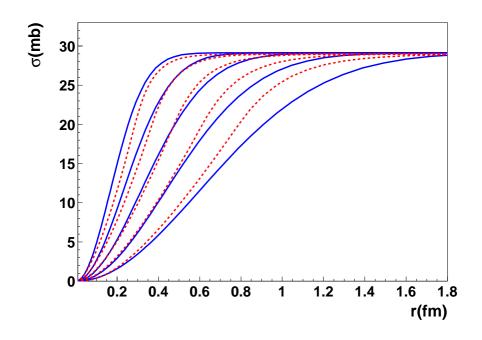
$$\hat{\boldsymbol{\sigma}}(\mathbf{r},x) = 4\pi B_d N(\mathbf{r},x)$$
.

$$N(\mathbf{r},x) = \begin{cases} N_0 \left(\frac{rQ_s}{2}\right)^{2\gamma_s} e^{\frac{2\ln^2(rQ_s/2)}{\kappa\lambda \ln(x)}} & \text{for} \quad rQ_s \le 2\\ 1 - e^{-4\alpha \ln^2(\beta rQ_s)} & \text{for} \quad rQ_s > 2 \end{cases}$$

lacksquare Parameters  $\lambda$ ,  $x_0$ ,  $N_0$ ,  $B_d$  are equal to

$$\lambda = 0.22, \quad x_0 = 1.63 \cdot 10^{-5}, \quad N_0 = 0.7, \quad B_d = 6 \text{ GeV}^{-2}$$

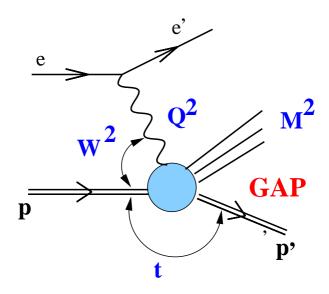
### **Comparison of GBW with CGC**



- The same normalization.
- In GBW from a fit to  $F_2$  data:  $\sigma_0 = 29 \text{ mb}$ .
- In CGC from measured diffractive slope at HERA:  $B_d = 6 \; GeV^{-2}$

$$\sigma_0 = 4\pi B_d = 29 \text{ mb}$$

### **DIS** diffraction



Kinematic variables

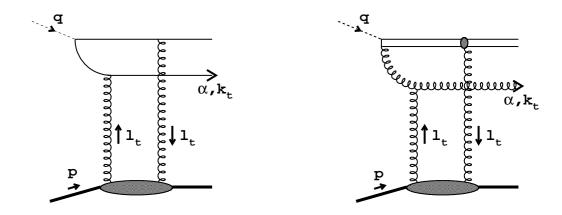
$$x_{I\!\!P} = rac{M^2 + Q^2}{W^2 + Q^2}, \qquad \beta = rac{Q^2}{Q^2 + M^2}$$

Diffractive structure functions

$$F_2^D(x_{I\!\!P}, \beta, Q^2)$$
  $F_L^D(x_{I\!\!P}, \beta, Q^2)$ 

### **DIS** diffraction in dipole models

Successful description with two component diffractive state:  $q\overline{q}$  and  $q\overline{q}g$  from photon with transverse and longitudinal polarization T,L



Diffractive structure functions are given by

$$F_2^D = F_T^{(q\overline{q})} + F_L^{(q\overline{q})} + F_T^{(q\overline{q}g)}$$

$$F_L^D = F_L^{(q\overline{q})}$$

### Diffractive structure functions: $q\overline{q}$ component

**●** The  $q\bar{q}$  components from T and L polarised photons are given by

$$\begin{split} x_{I\!\!P} F_T^{(q\overline{q})} &= \frac{3Q^4}{64\pi^4\beta B_d} \sum_f e_f^2 \int_{z_f}^{1/2} dz \, z (1-z) \\ &\times \left\{ \left[ z^2 + (1-z)^2 \right] Q_f^2 \, \phi_1^2 + m_f^2 \, \phi_0^2 \right\} \\ x_{I\!\!P} F_L^{(q\overline{q})} &= \frac{3Q^6}{16\pi^4\beta B_d} \sum_f e_f^2 \int_{z_f}^{1/2} dz \, z^3 (1-z)^3 \, \phi_0^2 \end{split}$$

• The functions  $\phi_i$  take the following form for i=0,1

$$\phi_i = \int_0^\infty dr r K_i(Q_f r) J_i(k_f r) \,\hat{\boldsymbol{\sigma}}(\boldsymbol{r}, \boldsymbol{x}_{I\!\!P})$$

• In diffractive DIS substitution in  $\hat{\sigma}(r, x_{I\!\!P})$ 

$$x = \frac{Q^2}{Q^2 + W^2} \qquad \to \qquad x_{I\!\!P} = \frac{Q^2 + M^2}{Q^2 + W^2}$$

## Diffractive structure functions: $q\overline{q}g$ component

 $\blacksquare$  The  $q\overline{q}g$  component from transverse photons with massless quarks

$$x_{\mathbb{I}\!P} F_{T}^{(q\overline{q}g)} = \frac{81\beta\alpha_{s}}{512\pi^{5}B_{d}} \sum_{f} e_{f}^{2} \int_{\beta}^{1} \frac{dz}{(1-z)^{3}} \left[ \left(1 - \frac{\beta}{z}\right)^{2} + \left(\frac{\beta}{z}\right)^{2} \right]$$

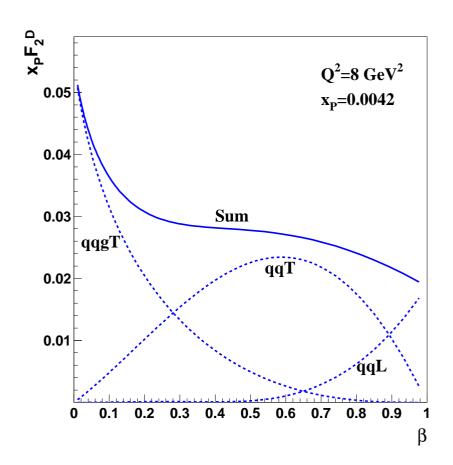
$$\times \int_{0}^{(1-z)Q^{2}} dk^{2} \log\left(\frac{(1-z)Q^{2}}{k^{2}}\right) \phi_{2}^{2}(x_{\mathbb{I}\!P}, z, k) ,$$

• Since  $k_{Tq} \approx k_{T\overline{q}} \gg k_{Tg}$  we have a gluon dipole and the GBW dipole cross section is given by

$$\hat{\sigma}(r, x_{\mathbb{I}P}) \equiv \hat{\sigma}_{gg} = \sigma_0 \left( 1 - e^{-(C_A/C_F)r^2 Q_s^2(x_{\mathbb{I}P})/4} \right)$$

• Color factor modyfication:  $C_A/C_F = 9/4$  for  $N_c = 3$ 

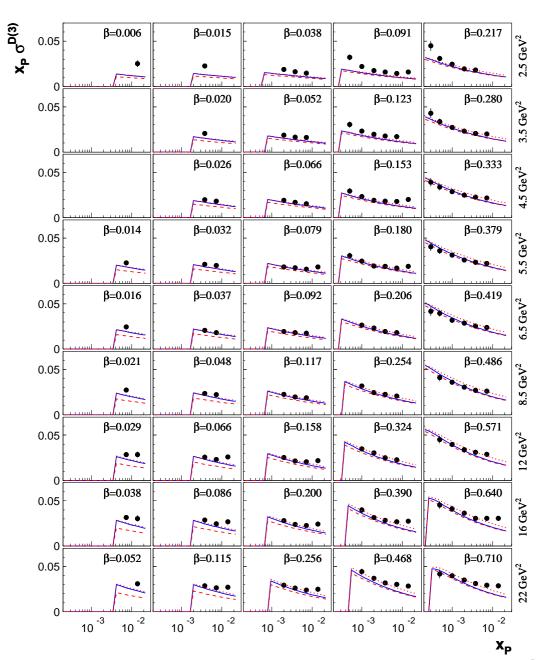
# Summary of the three contributions to ${\cal F}_2^D$



$$\beta = \frac{Q^2}{Q^2 + M^2}$$

- $q\overline{q}$  from transverse photons for etapprox 1/2
- $q\overline{q}$  from longitudinal photons for  $\beta\approx 1$  (for small diffractive mass)
- $q\overline{q}g$  for  $\beta \ll 1$  (for large diffractive mass)

# The comparison for $\sigma_r^D=F_2^D-\frac{y^2}{1+(1-y)^2}F_L^D$ : ZEUS 2009



### Diffractive charm production: only $c\overline{c}$

Standard dipole model formula with  $m_c = 1.4 \text{ GeV}$  and  $e_c = 2/3$ 

$$x_{I\!\!P} F_T^{(c\overline{c})} = \frac{3Q^4 e_c^2}{64\pi^4 \beta B_d} \int_{z_c}^{1/2} dz \, z (1-z)$$

$$\times \left\{ \left[ z^2 + (1-z)^2 \right] Q_c^2 \, \phi_1^2 + m_c^2 \, \phi_0^2 \right\}$$

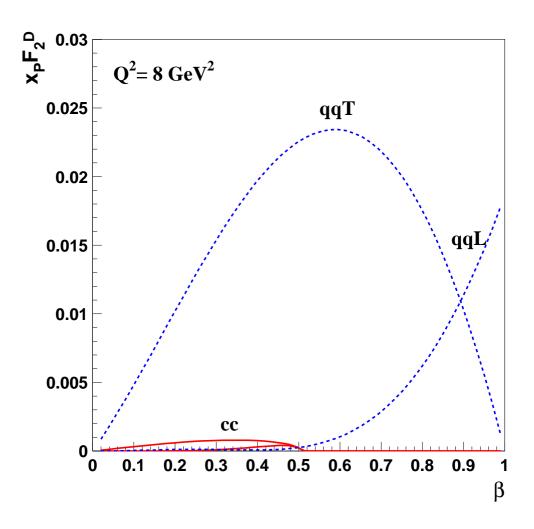
with 
$$z_c = (1 - \sqrt{1 - 4m_c^2/M^2})/2$$

Minimal diffractive mass  $M_{min}^2 = 4m_c^2$  gives maximal value of  $\beta$ 

$$\beta_{max} = \frac{Q^2}{Q^2 + 4m_c^2} < 1$$

Pure  $c\bar{c}$  contributions is very small.

### **Exclusive** $c\bar{c}$ **contribution**



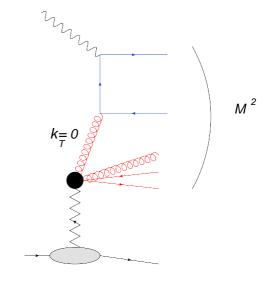
**•** Exclusive  $c\overline{c}$  production is negligible.

### Diffractive charm production: $c\overline{c}X$ state

We use collinear factorisation formula for  $c\bar{c}X$  diffractive production

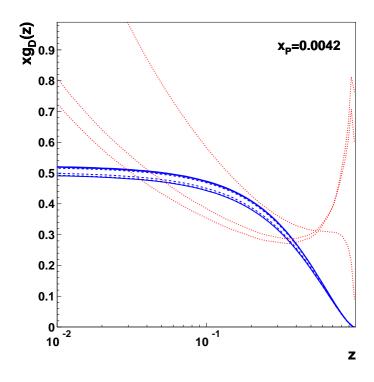
$$F_{2,L}^{D(c\overline{c}X)} = 2\beta \, e_c^2 \, \frac{\alpha_s(\mu_c^2)}{2\pi} \int_{a\beta}^1 \frac{d\beta'}{\beta'} \, C_{2,L}\left(\frac{\beta}{\beta'}, \frac{m_c^2}{Q^2}\right) \, g^D(x_{I\!\!P}, \beta', \mu_c^2)$$

where  $a=1+4m_c^2/Q^2$  and the factorization scale  $\mu_c^2=4m_c^2$ .



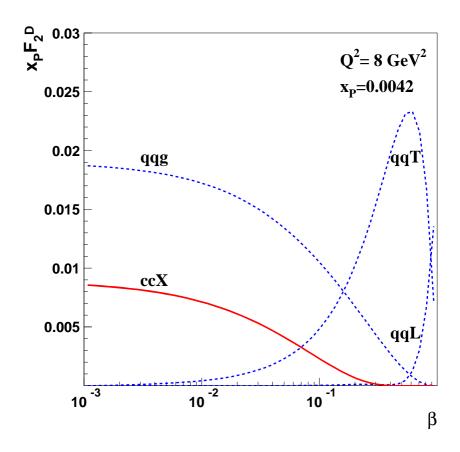
The standard gluon distribution is replaced by the diffractive gluon distribution  $g^D(x_{\mathbb{I}\!P}, \beta', \mu_c^2)$ .

## **Diffractive gluon distribution**



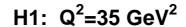
- Gluon from the  $q\overline{q}g$  component of the dipole model (blue lines)
- Gluon from DGLAP fits to HERA diffractive data (red lines)

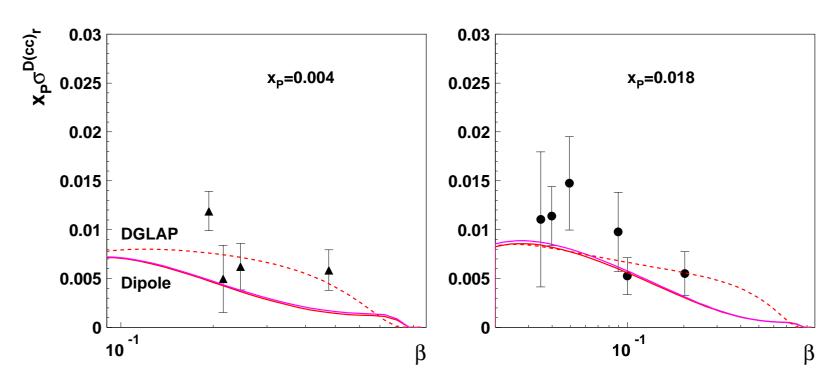
# **Diffractive charm production:** $c\overline{c}X$



•  $c\overline{c}X$  contributes up to 30% to  $F_2^D$  for large diffractive mass (small  $\beta$ ).

# **Comparison with HERA data**



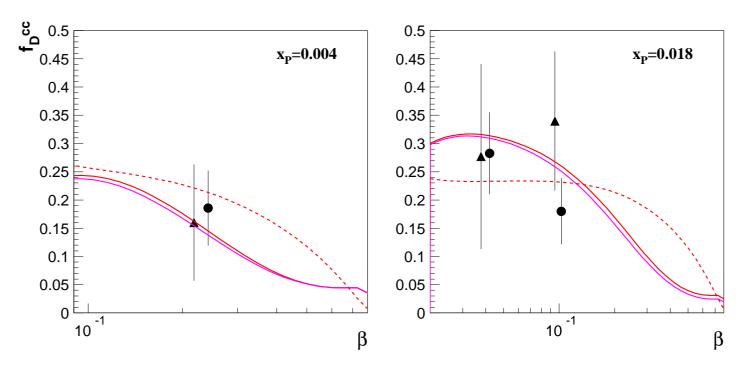


### Reduced cross section

$$\sigma_r^{D(charm)} = F_2^{D(charm)} - \frac{y^2}{1 + (1 - y)^2} F_L^{D(charm)}$$

### **Fractional charm contribution**

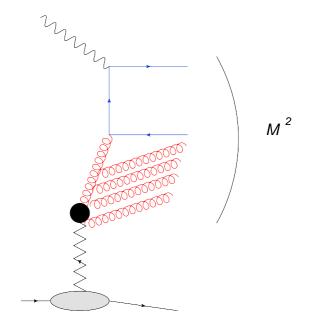
H1: 35 GeV<sup>2</sup>



- Fractional contribution  $f^{D(charm)} = \sigma_r^{D(charm)}/\sigma_r^D$  of the diffractive charm to total diffractive cross section.
- Up to 20 30%.

### **Conclusions**

- ullet Good overall agreement with the diffractive HERA data on  $F_2^D$ .
- For small  $\beta$  curves below the data more complicated diffractive state than  $q\overline{q}g$  is necessary.



• Charm contribution is important for small  $\beta$  (large diffractive mass) - contibutes to  $F_2^D$  up to 20-30%.