

**EPS-HEP 2009 Kraków, 17.07. 2009**

**EVOLUTION of the Universe**

**to the present Inert Phase**

**(Dark)**

**2HDMs**

**Today: Inert Model**

**In the Past?**

Maria Krawczyk

University of Warsaw

I. Ginzburg, K. Kanishev

(Novosibirsk University),

D.Sokołowska

(University of Warsaw)

# THE THEORY OF MATTER and STANDARD MODEL(S)

F. Wilczek, LEPFest, Nov.2000 (hep-ph/0101187)

**Theory of Matter =  $SU(2)_{I_{weak}} \times U(1)_{Y_{weak}} \times SU(3)_{color}$**

Theory of Matter refers to the core concepts

- quantum field theory
- gauge symmetry
- spontaneous symmetry breaking
- asymptotic freedom
- the assignments of the lightest quarks and leptons

**Standard Models:** Choose the number of Higgs (scalar) doublets  
SM=1HDM, 2HDM (MSSM), 3HDM ...

Note, that the lightest scalar is often **SM-like**

**NonStandard Models** are based on more radical assumptions.

See eg. *CP Study and the Nonstandard Higgs Workshop 2002-2006*  
(CERN Report hep-ph/0608079)

# Brout-Englert-Higgs mechanism

Spontaneous breaking of EW symmetry

$$SU(2) \times U(1) \rightarrow U(1)_{\text{QED}}$$

## Standard Model

Doublet of  $SU(2)$ :  $\Phi = (\varphi^+, v + H + i\zeta)^T$

Masses for  $W^{+/-}$ ,  $Z$  ( $\rho = 1$ ), no mass for photon

Fermion masses via Yukawa interaction

Higgs particle  $H_{\text{SM}}$  - spin 0, neutral, CP even

mass  $\leftrightarrow$  selfinteraction unknown

couplings to  $WW/ZZ$ , Yukawa couplings to fermions

# Brout-Englert-Higgs mechanism

Spontaneous breaking of EW symmetry

$SU(2) \times U(1) \rightarrow ?$

## Two Higgs Doublet Models

Two doublets of  $SU(2)$  ( $Y=1$ ,  $\rho=1$ ) -  $\Phi_1, \Phi_2$

Masses for  $W^{+-}$ ,  $Z$  ( $\rho=1$ ), no mass for photon?

Fermion masses via Yukawa interaction –

various models – Model I, II, III, IV...

5 scalars:  $H^+$  and  $H^-$  and neutrals:

- CP conservation: CP-even  $h, H$  & CP-odd  $A$
- CP violation:  $h_1, h_2, h_3$  with indefinite CP parity\*

Sum rules (relative couplings to SM  $\chi$ )

# 2HDM Potential

(Lee'73, Haber, Gunion, Weinberg, Branco, Rebelo, Lavoura, Ferreira, Barroso, Santos, Silva, Diaz-Cruz, Grimus, Ecker, Ivanov, Ginzburg, Osland, Nishi, Nachtamnn, Akeroyd, Kanemura, Kalinowski, Grzadkowski, Hollik, Rosiek)

$$\begin{aligned} V = & \lambda_1(\Phi_1^\dagger\Phi_1)^2 + \lambda_2(\Phi_2^\dagger\Phi_2)^2 + \lambda_3(\Phi_1^\dagger\Phi_1)(\Phi_2^\dagger\Phi_2) \\ & + \lambda_4(\Phi_1^\dagger\Phi_2)(\Phi_2^\dagger\Phi_1) + [\lambda_5(\Phi_1^\dagger\Phi_2)^2 + \text{h.c.}] \\ & + [(\lambda_6(\Phi_1^\dagger\Phi_1) + \lambda_7(\Phi_2^\dagger\Phi_2))(\Phi_1^\dagger\Phi_2) + \text{h.c.}] \\ & - m_{11}^2(\Phi_1^\dagger\Phi_1) - m_{22}^2(\Phi_2^\dagger\Phi_2) - [m_{12}^2(\Phi_1^\dagger\Phi_2) + \text{h.c.}] \end{aligned}$$

$Z_2$  symmetry transformation:  $\Phi_1 \rightarrow \Phi_1$   $\Phi_2 \rightarrow -\Phi_2$

Hard  $Z_2$  symmetry violation:  $\lambda_6, \lambda_7$  terms

Soft  $Z_2$  symmetry violation:  $m_{12}^2$  term

Explicit  $Z_2$  symmetry in  $V$ :  $\lambda_6, \lambda_7, m_{12}^2 = 0$

# $Z_2$ symmetry: $\Phi_1 \rightarrow \Phi_1$ $\Phi_2 \rightarrow -\Phi_2$

- If  $Z_2$  symmetry holds in the Lagrangian  $\rightarrow$ 
  - no CP violation in the scalar sector *Branco, Rebelo*
- Softly broken  $\rightarrow$  *Weinberg, Paschos*
  - CP violation possible, tree-level FCNC absent
- Hard breaking  $\rightarrow$ 
  - CP violation possible, even without CP mixing (\*)
  - tree-level FCNC danger

*Ma'78*

$Z_2$  symmetry both in L and in vacuum – Inert Model

doublet  $\Phi_1$  as in SM, with Higgs boson  $h$  SM-like,  
4 scalars from  $\Phi_2$  with  $Z_2$ -odd parity (dark scalars)

The lightest dark scalar candidate for a dark matter



# Possible vacuum states

(parameters in V- real)

The most general vacuum state

$$\langle \phi_1 \rangle = \begin{pmatrix} 0 \\ \frac{1}{\sqrt{2}}v_1 \end{pmatrix}, \quad \langle \phi_2 \rangle = \begin{pmatrix} u \\ \frac{1}{\sqrt{2}}v_2 e^{i\xi} \end{pmatrix}$$

with  $v_1, v_2, u, \xi$  - real, positive

For the **Inert** vacuum

$$u = v_2 = 0, \quad v = v_1$$

For the CP conserving **CPc**

$$u = \xi = 0$$

For the Charge Breaking **ChB**

$$v_2 = \xi = 0$$

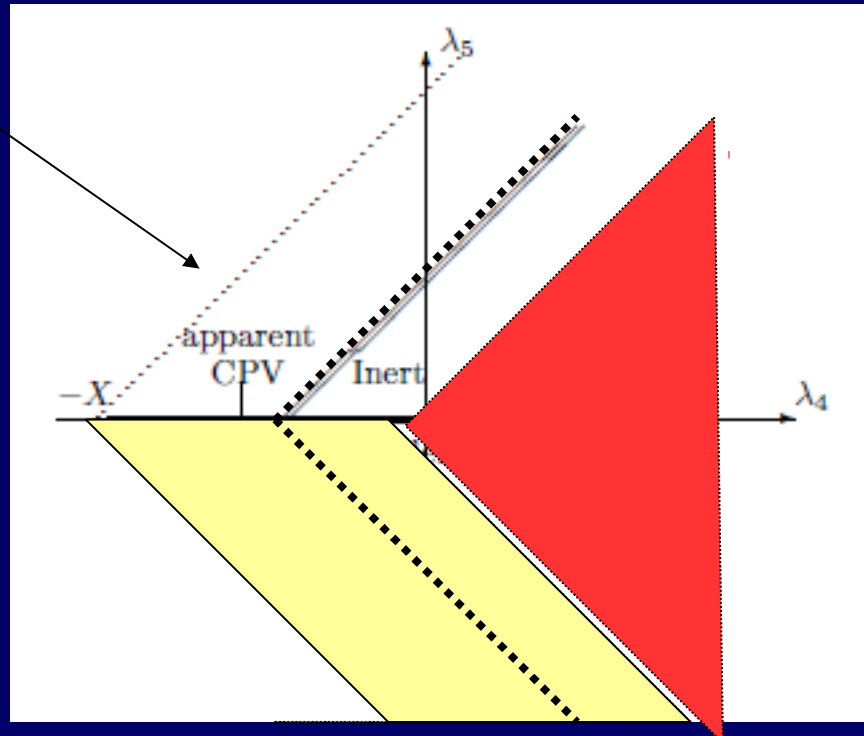
# Various vacua on $\lambda_5 \times \lambda_4$ plane

for  $V$  with an explicit  $Z_2$  symmetry

Positivity constrains on  $V$ :

$$X = \sqrt{\lambda_1 \lambda_2 + \lambda_3} > 0$$

$$\lambda_4 \pm \lambda_5 > -X$$

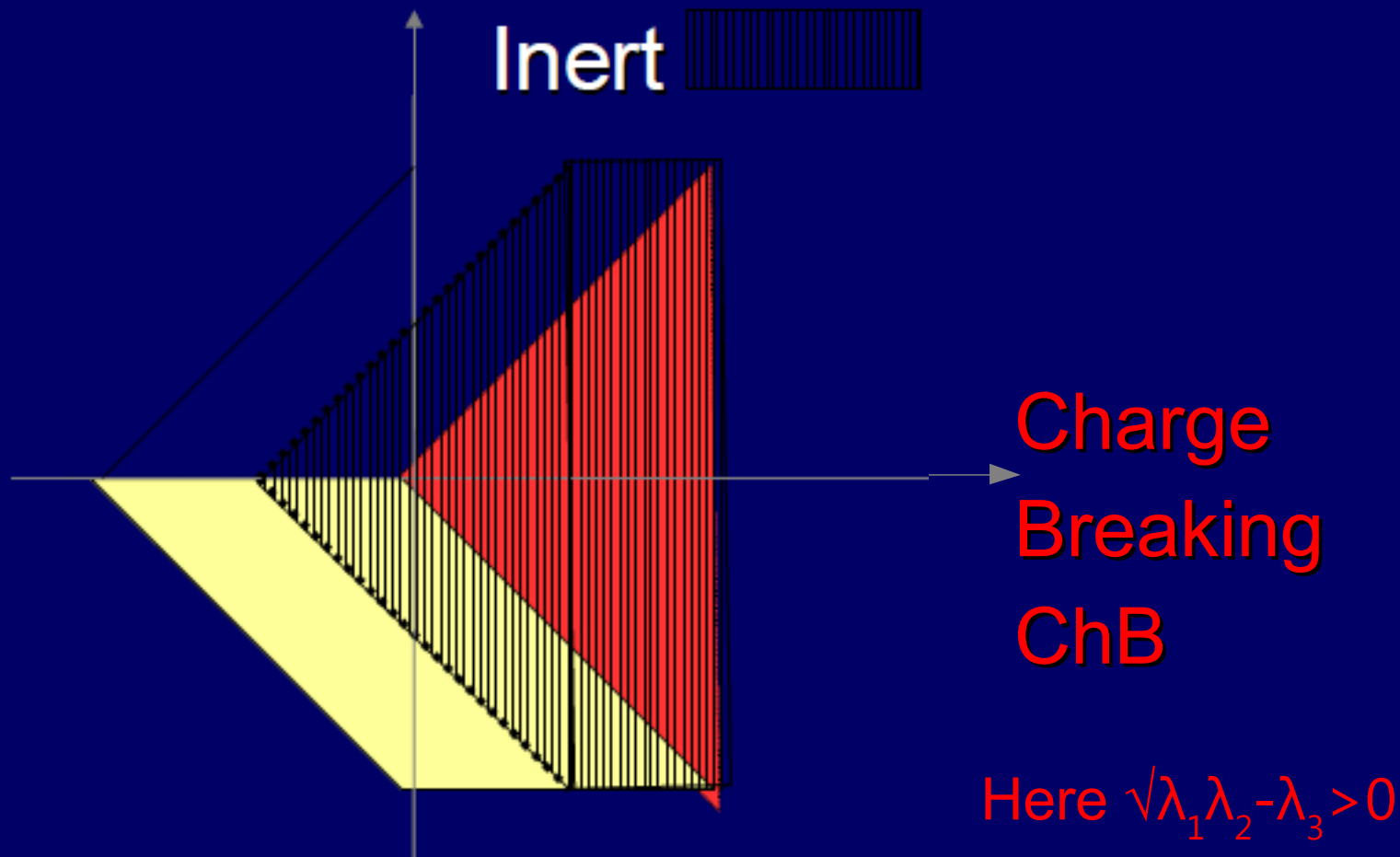


Charge  
Breaking  
ChB

CP conserving – CPc



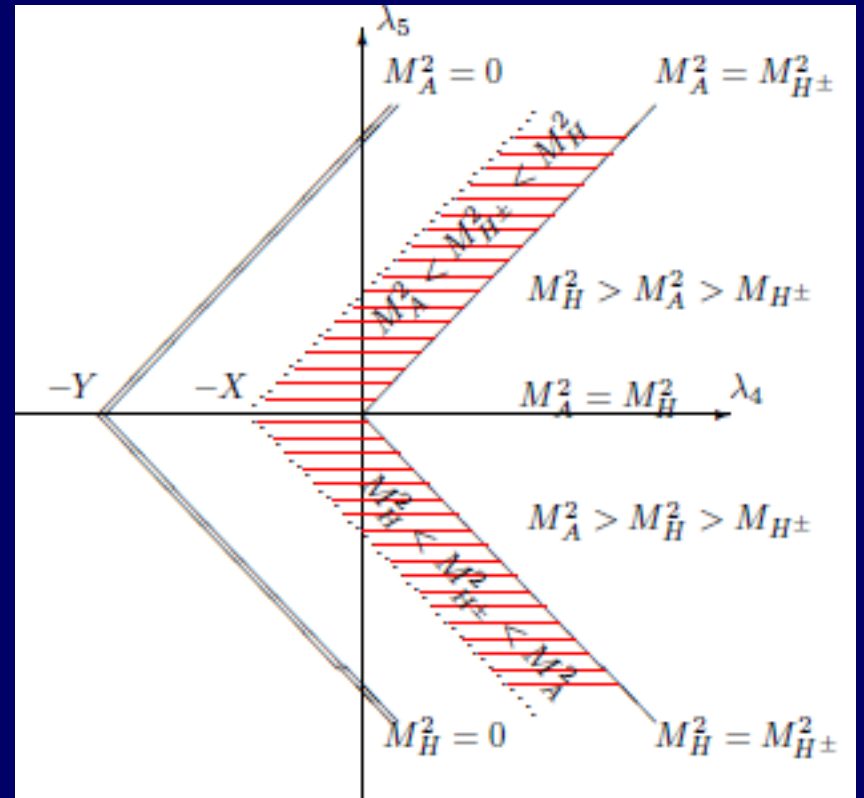
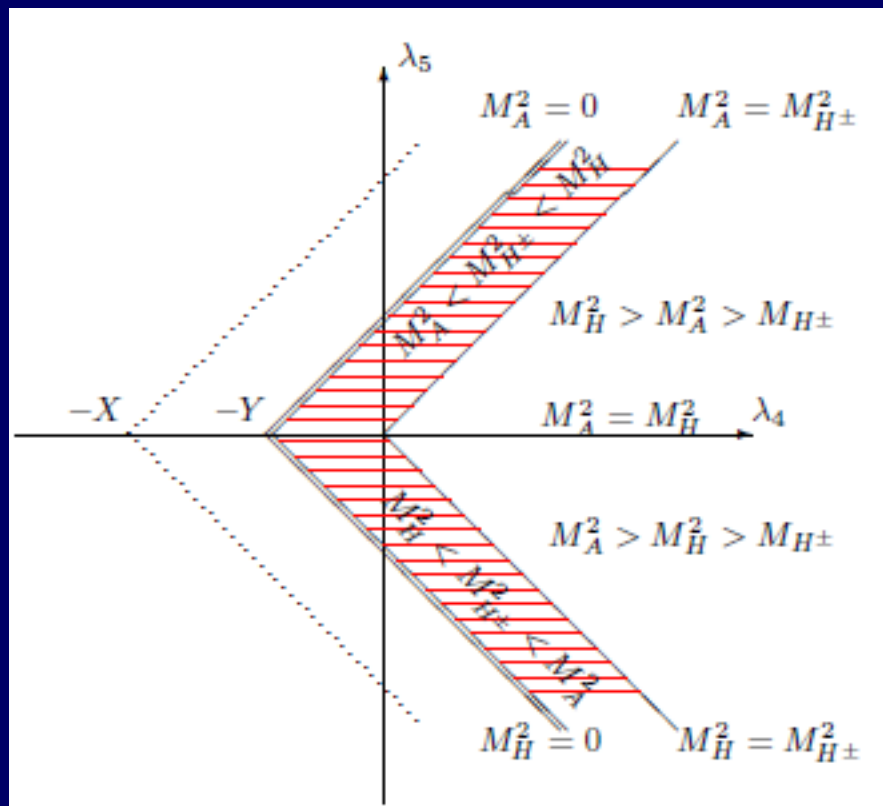
# Various vacua on $\lambda_5 \times \lambda_4$ plane



CP conserving – CPc

# Intert Model – dark scalar masses

using  $X$  (positivity) and  $Y = M_{H^\pm}^2 - 2/v^2$



If  $H$  is the dark matter candidate

$$\rightarrow \lambda_5 < 0$$

# Inert Model (Dark 2HDM) vs data

Ma.. ' 78, Barbieri.. ' 2006

Exact  $Z_2$  symmetry in L and in vacuum  $\rightarrow$

$Z_2$  parity: odd is only  $\Phi_2$

$\Phi_1$

- Nonzero vev only doublet  $\Phi_1$  (Higgs doublet) only it couples to fermions (Model I)
- Higgs boson  $h$  SM-like:  $M_h^2 = m_{11}^2 = \lambda_1 v^2$

$\Phi_2$

- Zero vev for  $\Phi_2$  (scalar doublet) and no Yukawa int. i
- Four scalars with  $Z_2$  - odd parity (dark scalars D)
- The lightest dark scalar -a candidate for dark matter

# Dark scalars

- Masses

$$\begin{aligned}M_{H^+}^2 &= -\frac{m_{22}^2}{2} + \frac{\lambda_3 v^2}{2} \\M_H^2 &= -\frac{m_{22}^2}{2} + \frac{\lambda_3 + \lambda_4 + \lambda_5}{2} v^2 \\M_A^2 &= -\frac{m_{22}^2}{2} + \frac{\lambda_3 + \lambda_4 - \lambda_5}{2} v^2\end{aligned}$$

- D couple to W/Z (eg.  $AZH$ ,  $H^- W^+ H$ ),  
not to  $VV$
- Selfcouplings  $DDDD$  proportional to  $\lambda_2$
- Couplings between Higgs boson  $h$  and  $D$ -  
proportional to  $M_D^2 + m_{22}^2/2$

# Testing Inert Model

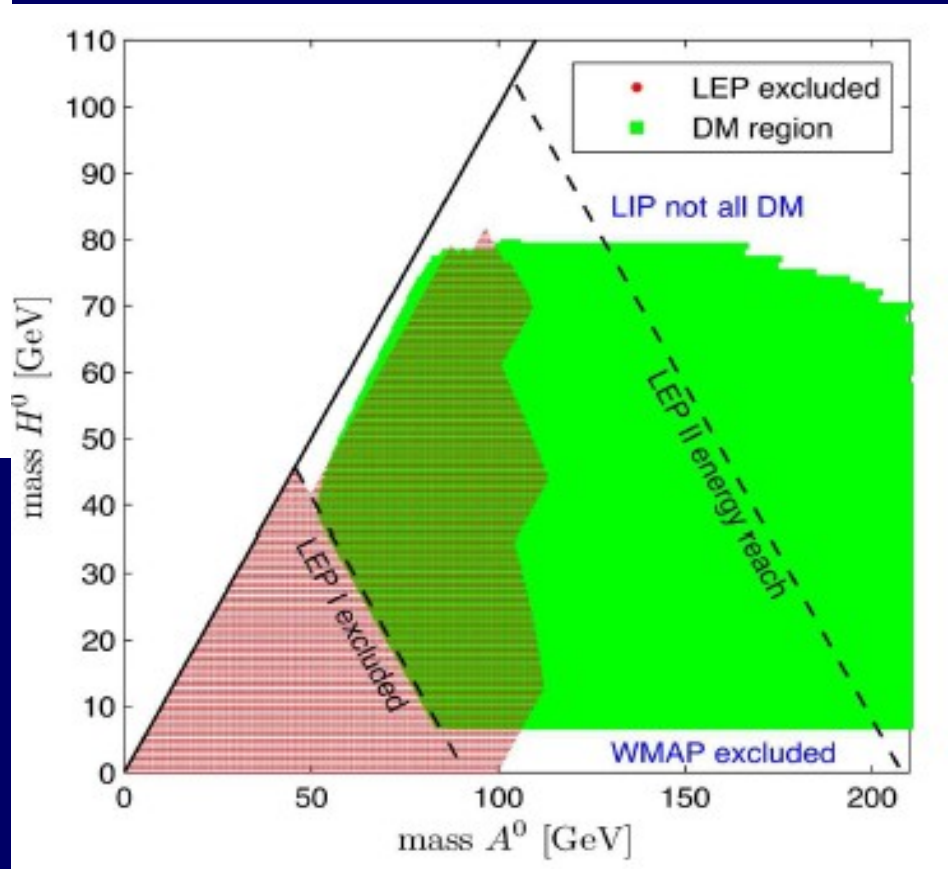
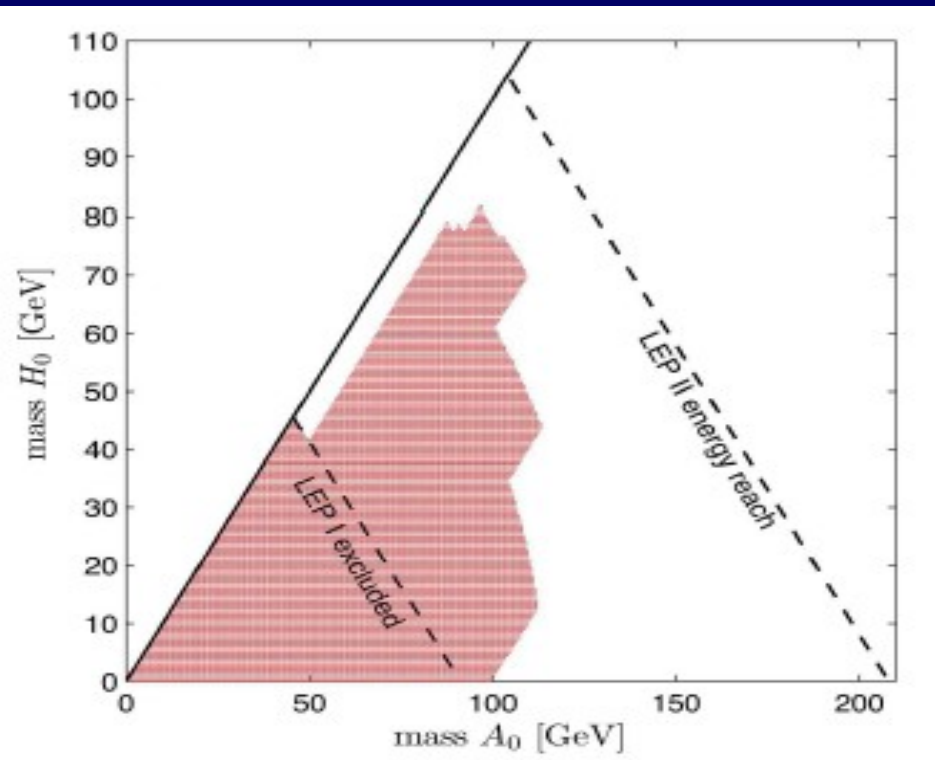
- Since only dark scalars have odd  $Z_2$ -parity
  - they can be produced only in pairs
  - the lightest dark particle – a suitable DM
- The strategy to test such model for light h
  - (Barbieri et al '2006 heavy h)
  - consider properties of SM-like h and look for deviations
  - consider properties of dark scalars, especially DM candidate
- Cao, Ma, Rajasekaren' 2007: colliders signal/constraints

LEP II:  $M_H + M_A > M_Z$ ,  $\Delta(A, H) = 5 - 30$  GeV for  $M_h = 105 - 110$  GeV  
EW precision data:  $(M_{H^+} - M_A)(M_{H^+} - M_H) = M^2$ ,  $M = 120_{-30}^{+20}$  GeV

# Dark 2HDM: LEP II exclusion

Lundstrom et al 0810.3924

$$M_A - M_H = 8 \text{ GeV}$$



**LEP II + WIMP**  
 **$M_h = 200 \text{ GeV}$**

# Evolution of Universe – different vacua in the past ?

We consider 2HDM with an explicit  $Z_2$  symmetry

Useful parametrization with  $k$  and  $\delta$

$$\lambda_2/\lambda_1 = k^4, \quad m_{11}^2 = m^2(1 - \delta), \quad m_{22}^2 = k^2 m^2(1 + \delta)$$

$$\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5, \quad \tilde{\lambda}_{345} = \lambda_3 + \lambda_4 - \lambda_5.$$

$$\Lambda_{345\pm} = \sqrt{\lambda_1\lambda_2} \pm \lambda_{345}, \quad \tilde{\Lambda}_{345\pm} = \sqrt{\lambda_1\lambda_2} \pm \tilde{\lambda}_{345}, \quad \Lambda_{3\pm} = \sqrt{\lambda_1\lambda_2} \pm \lambda_3.$$

Yukawa interaction – **Model I**  $\rightarrow$

all fermions couple only to  $\Phi_1$

# Possible vacua:

ChB

$$v_1^2 = \frac{m^2 k^2}{2} \left( \frac{1}{\Lambda_{3+}} - \frac{\delta}{\lambda_{3-}} \right), \quad v_2 = 0, \quad u^2 = \frac{m^2}{2} \left( \frac{1}{\Lambda_{3+}} + \frac{\delta}{\lambda_{3-}} \right)$$

$$\mathcal{E}_{ch} = -\frac{m^4 k^2}{8} \left( \frac{1}{\Lambda_{3+}} + \frac{\delta^2}{\lambda_{3-}} \right)$$

Inert

$$v_2 = 0, \quad v^2 = v_1^2 = \frac{m^2(1-\delta)}{\lambda_1}, \quad \mathcal{E}_A = -\frac{m^4(1-\delta)^2}{8\lambda_1}$$

B

$$v_1 = 0, \quad v^2 = v_2^2 = \frac{m^2(1+\delta)}{k^2\lambda_1}, \quad \mathcal{E}_B = -\frac{m^4(1+\delta)^2}{8\lambda_1}$$

N(CP<sub>C</sub>)

$$v_1^2 = \frac{m^2 k^2}{2} \left( \frac{1}{\Lambda_{345+}} - \frac{\delta}{\Lambda_{345-}} \right), \quad v_2^2 = \frac{m^2}{2} \left( \frac{1}{\Lambda_{345+}} + \frac{\delta}{\Lambda_{345-}} \right),$$

$$\mathcal{E}_c = -\frac{k^2 m^4}{4} \left( \frac{1}{\Lambda_{345+}} + \frac{\delta^2}{\Lambda_{345-}} \right)$$

Depending on value of  $\delta \rightarrow$  a true vacuum, ie.  
the one with the minimal energy



# Thermal evolution of parameters

Matsubara method ( $T \gg m^2$ ) –

only quadratic parameters change with  $T$

$$m_{11}^2(T) = m_{11}^2(0) - 2c_1 m^2 w, \quad m_{22}^2(T) = m_{22}^2(0) - 2c_2 m^2 w$$

$$c_1 = \frac{3\lambda_1 + 2\lambda_3 + \lambda_4}{2}, \quad c_2 = \frac{3\lambda_2 + 2\lambda_3 + \lambda_4}{2}, \quad w = \frac{T^2}{12m^2}$$

$$m^2(T) = m^2(1 - (c_2 + c_1)w)$$

$$\delta(T) = \frac{m^2}{m^2(T)} \left( \delta - \frac{c_2 - c_1}{c_2 + c_1} \right) + \frac{c_2 - c_1}{c_2 + c_1}$$

$\delta$  evolves possibly passing different phases

# Phase transitions from the EW phase

Sector I:  $\Lambda_{3-} > 0, \lambda_4 + \lambda_5 > 0$

Two second order phase transitions:

EW  $\xrightarrow{II}$  Phase B  $\xrightarrow{II}$  Charged phase  $\xrightarrow{II}$  Inert phase

Sector II:  $\Lambda_{345-} > 0, \lambda_4 + \lambda_5 < 0$

Two second order phase transitions:

EW  $\xrightarrow{II}$  Phase B  $\xrightarrow{II}$  Phase N  $\xrightarrow{II}$  Inert phase

Sector III: Complementary to sectors I and II

First order phase transition:

EW  $\xrightarrow{II}$  Phase B  $\xrightarrow{I}$  Inert phase

# Conclusions

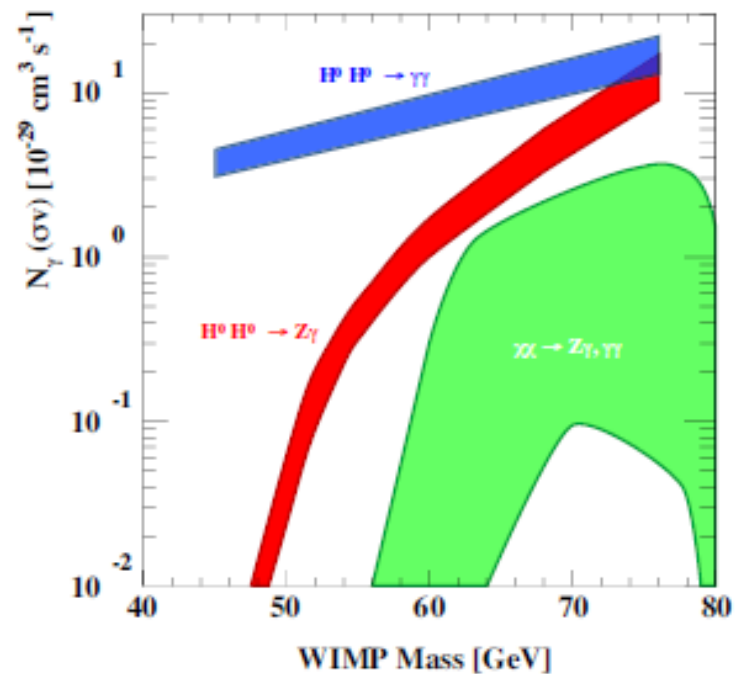
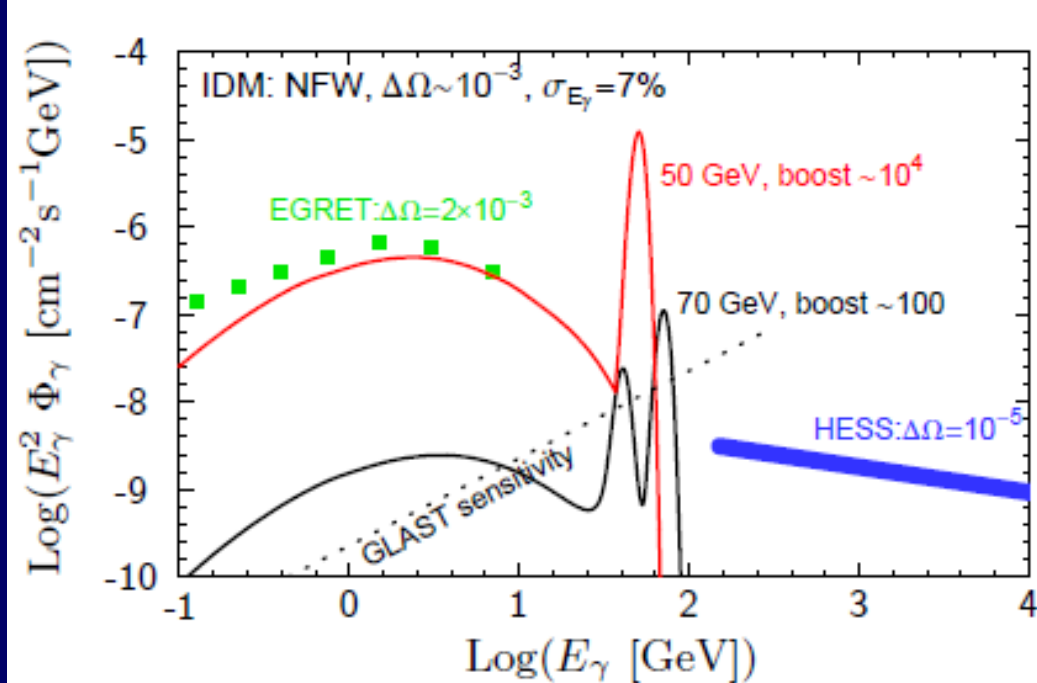
- Rich content of 2HDMs
- Intert Model in agreement with present data – soon tests at FERMI and LHC
- What was in the Past?

- $EW \xrightarrow{||} \text{Phase B} \xrightarrow{|} \text{Inert phase}$
- $EW \xrightarrow{||} \text{Phase B} \xrightarrow{||} \text{Phase N} \xrightarrow{||} \text{Inert phase}$
- $EW \xrightarrow{||} \text{Phase B} \xrightarrow{||} \text{Charged phase} \xrightarrow{||} \text{Inert phase}$
- $EW \xrightarrow{||} \text{Inert phase}$

- Various scenarios
- Can we find clear signals ?

## Significant Gamma Lines from Inert Doublet Model

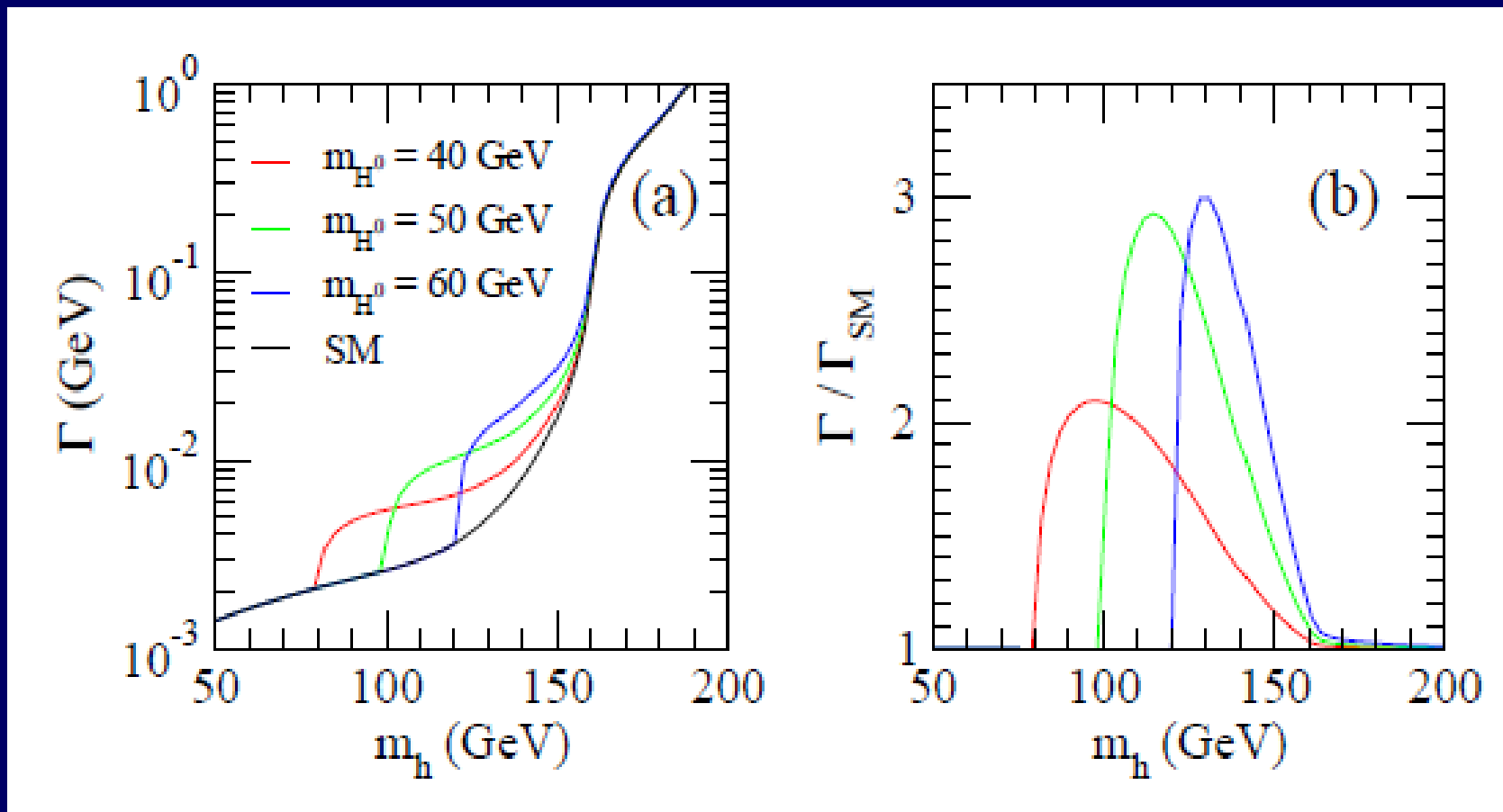
Gustafsson, Lundstrom, Bergstrom, Edsjo' 2007 studied direct annihilation of  $HH$  into  $\gamma\gamma$  and  $Z\gamma$  for  $M_H$  between 40-80 GeV (loop process, energy below WW threshold).



# Conclusion on gamma lines

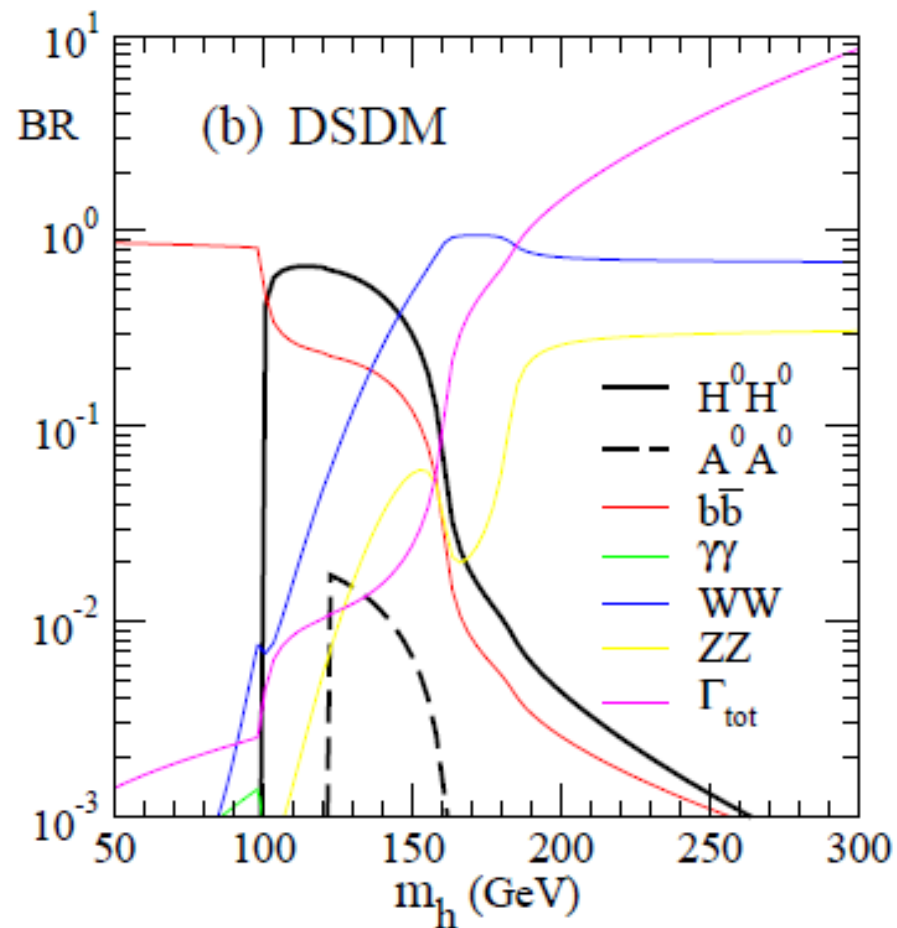
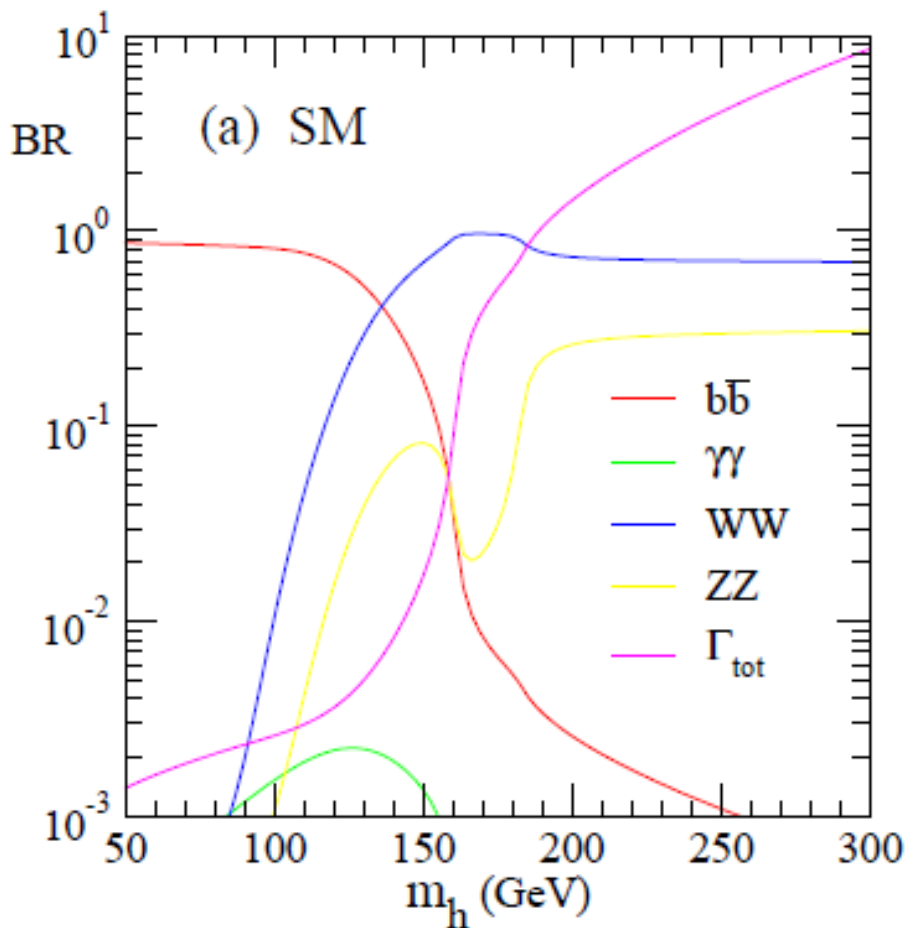
- Gustafsson et al.2007:  
*striking DM line signals -promising features to search with GLAST*  
Mass of:  $H = 40-80$  GeV,  $H^+ = 170$  GeV,  
 $A = 50-70$  GeV,  $h = 500$  and  $120$  GeV
- Honorez, Nezri, Oliver, Tytgat 2006-7:  
*H as a perfect example or archetype of WIMP – within reach of GLAST (FERMI)*  
Here mass of  $h = 120$  GeV, large mass  $H^+$   
close to  $A = 400 - 550$  GeV

# Dark 2HDM – total width of $h$



# Dark 2HDM – additional decays for h

Ma' 2007



For  $M_H = 50$  GeV,  $\Delta(A, H) = 10$  GeV,  $M_{H^\pm} = 170$  GeV,  $m_{22} = 20$  GeV

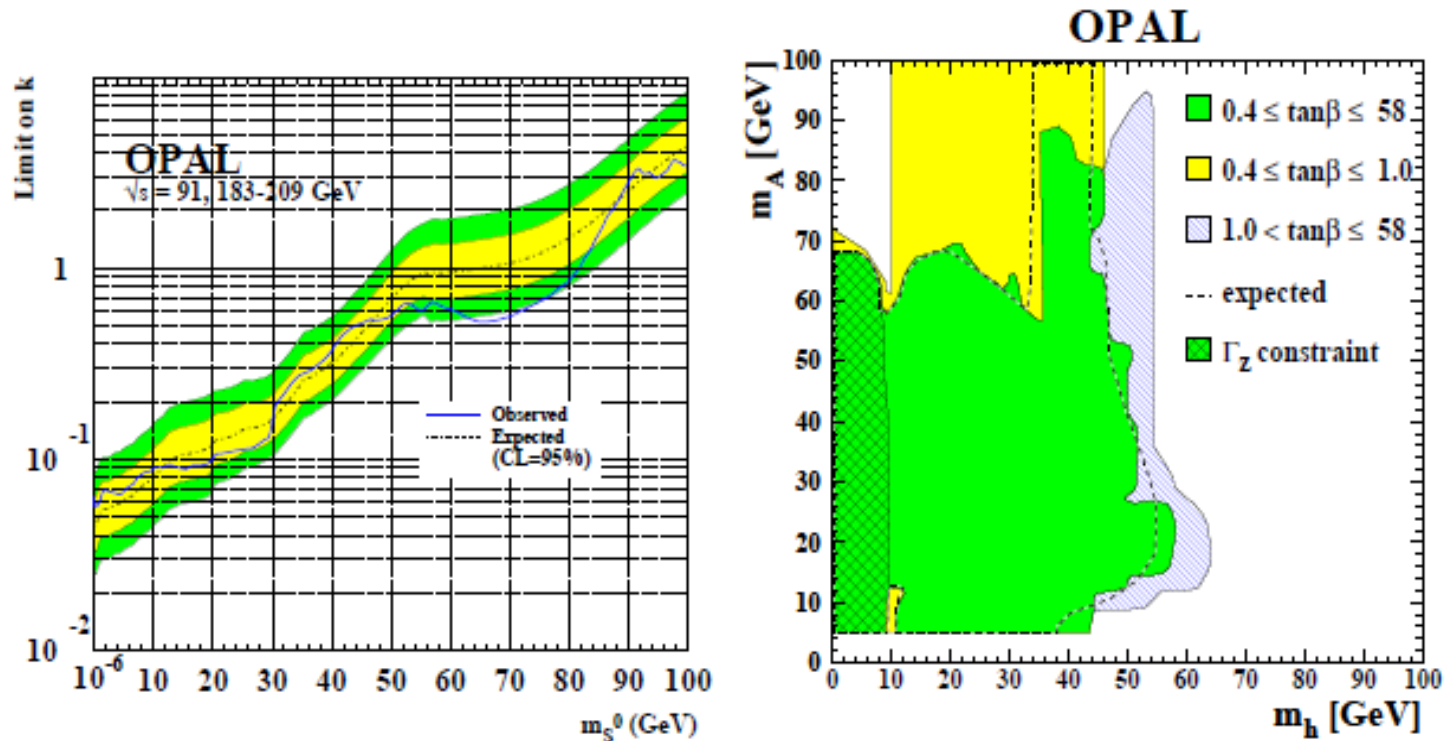
# Backup slides



## LEP: 2HDM with Z2 symmetry

Light  $h$  OR light  $A$  in agreement with current data

$hZZ$ :  $\sin(\beta - \alpha)$  and  $hAZ$ :  $\cos(\beta - \alpha)$



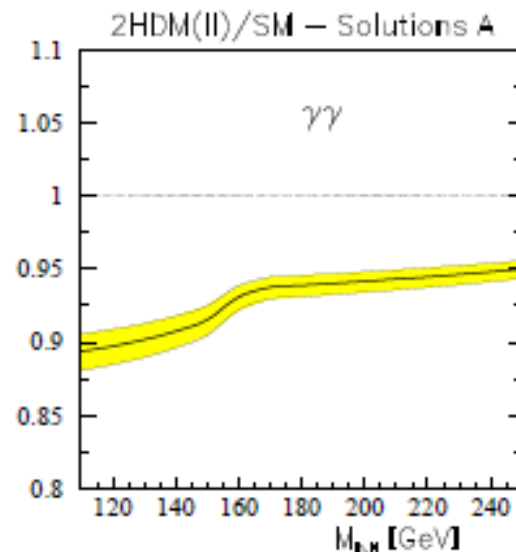
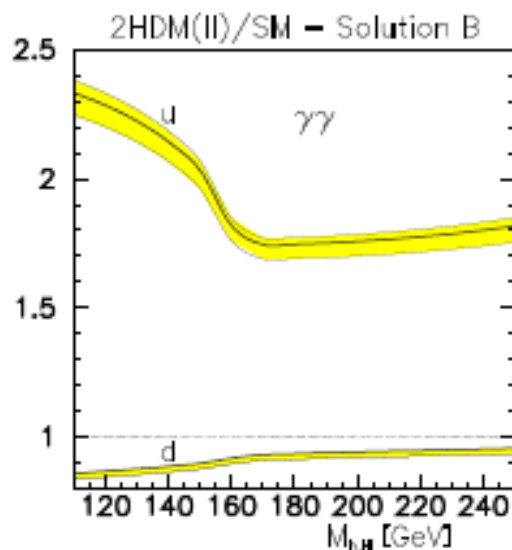
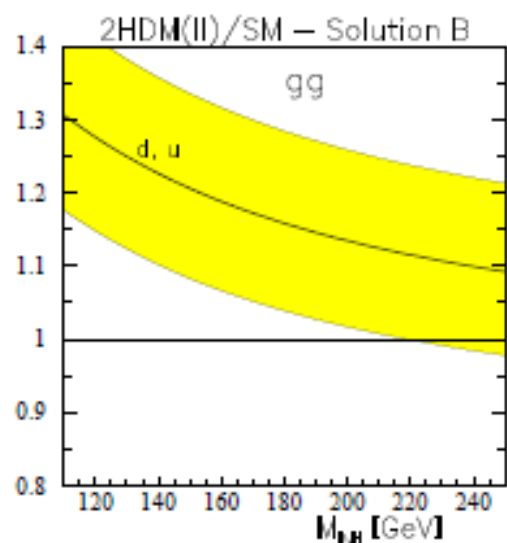
Light scalar  $h \rightarrow$  small  $k = \sin^2(\beta - \alpha)$  !

# Loop couplings $ggh$ , $\gamma\gamma h$ in 2HDM

The coupling between neutral and charged Higgs bosons:  $h_i H^+ H^-$

$$\gamma\gamma h \quad \chi_{H^\pm}^{(i)} = \left(1 - \frac{M_i^2}{2M_{H^\pm}^2}\right) \chi_V^{(i)} + \frac{M_i^2 - \nu v^2}{2M_{H^\pm}^2} \text{Re}(\chi_u^{(i)} + \chi_d^{(i)}).$$

For small  $m_{12}^2$ ... even for  $\chi_V, \chi_u, \chi_d$  equal 1, (SM-like scenario A)  $\rightarrow$  large non-decoupling effects due to heavy  $H^\pm$ . 600 GeV



$ggh$  - solution B „wrong” signs of fermion couplings

# Dark 2HDM: $\gamma\gamma h$

