

Prospects for R-parity Conserving SUSY searches at the LHC



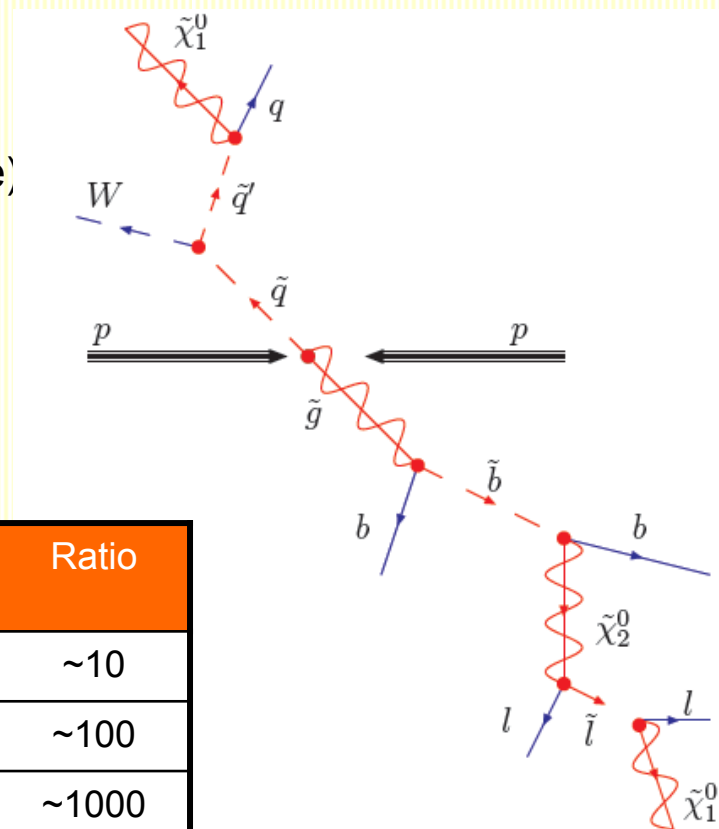
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EPS 2009



On behalf of the CMS and ATLAS collaborations

SUSY at the LHC

- Potentially long decay chains:
 - high- p_T jets
 - leptons (from charginos / neutralinos)
 - missing transverse energy (Lightest SParticle)
 - Main backgrounds:
 - $t\bar{t}$ bar
 - QCD
 - W/Z+jets
- Look for different types of signatures



	Tevatron: σ (pb)	LHC ($\sqrt{s}=14$ TeV): σ (pb)	Ratio
W^\pm (80 GeV)	2600	20000	~ 10
$t\bar{t}$ bar (2x172 GeV)	7	800	~ 100
$\tilde{q}\tilde{q}$ (2x400 GeV)	0.05	60	~ 1000
$\tilde{g}\tilde{g}$ (2x400 GeV)	0.005	100	~ 20000

The different search channels

- 0-lepton channel (jets + E_T^{miss})
 - the least model-dependent SUSY signature search
 - Strong cuts to reduce the QCD background
 - Simple cuts for higher sensitivity
 - Main backgrounds: QCD, ttbar, W, Z
- 1-lepton channel (1 lepton + jets + E_T^{miss})
 - Requiring one lepton is a slightly less generic search but increases the robustness
 - Golden channel at the beginning of data-taking
 - Main backgrounds: ttbar, W

The different search channels

- **Multi-lepton channel (2/3 leptons)**
 - Neutralinos/charginos in the decay chain \rightarrow 2 or 3 leptons in the final state
 - 2 leptons:
 - Opposite sign: non-resonant excess but SM background
 - Same sign: low SM background but lower statistics
 - 3 leptons:
 - 1 jet (avoid cutting on E_T^{miss})
 - E_T^{miss} + Opposite Sign Same Flavour (probe the gaugino direct production)
- **τ channel**
 - Large $\tan\beta$ favours light 3rd generation superpartners
 - Leptonically decaying τ s indistinguishable from leptons (included in other channels)
 - Main backgrounds: $t\bar{t}$, W, Z
- **b channel**
 - The presence of b-jets in SUSY cascades is also enhanced by $\tan\beta$
 - Asking for 2/3 b-jets in the final state (standard b-tagging in ATLAS: $\varepsilon \sim 50\text{-}60\%$)
 - Main background: $t\bar{t}$ (W,Z reduced)

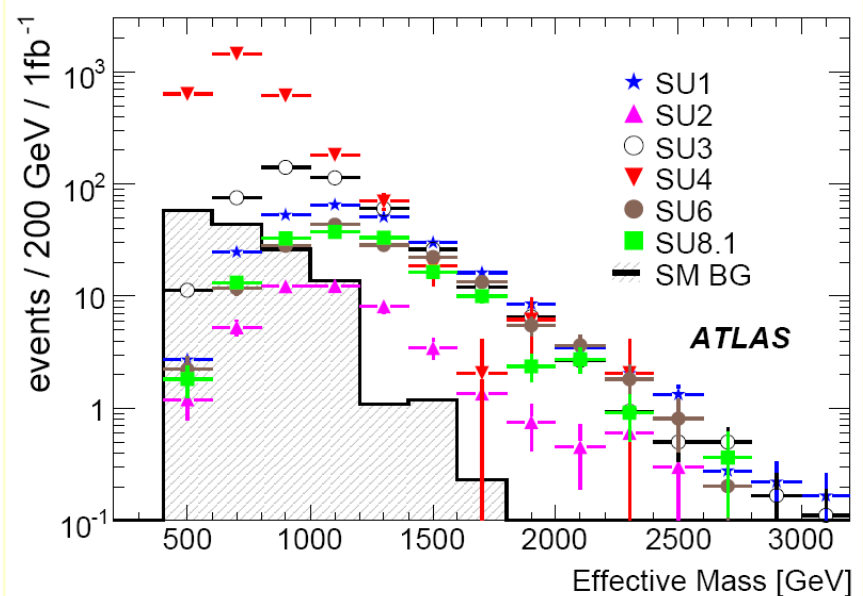
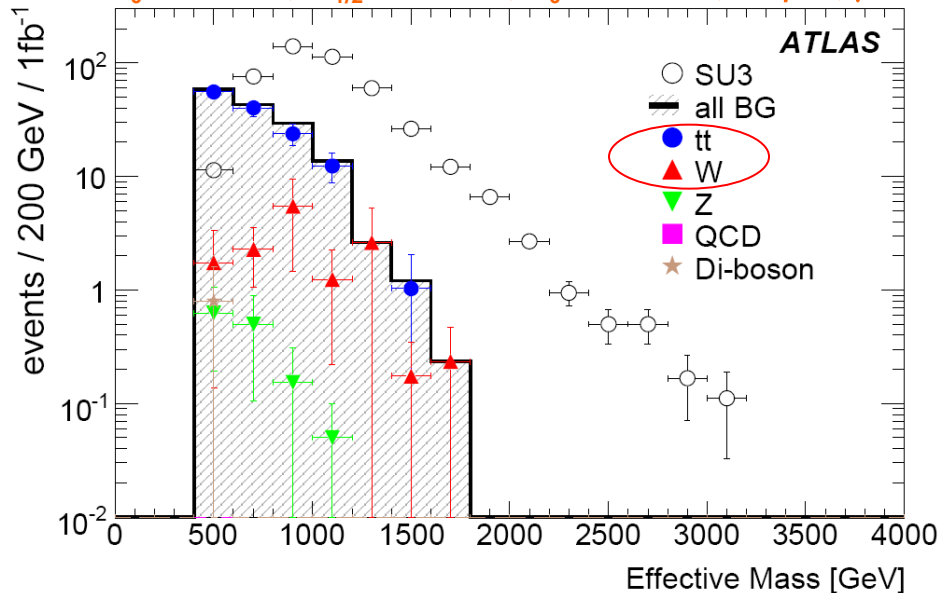
An example: the 1-lepton channel in ATLAS

- One isolated lepton with $p_T > 20$ GeV and no other lepton with $p_T > 10$ GeV
- At least 4 jets with $p_T > 50$ GeV and $p_T^{J1} > 100$ GeV
- $E_T^{miss} > \max(100 \text{ GeV}, 0.2M_{eff})$
- Transverse sphericity $S_T > 0.2$
- Transverse mass $M_T > 100$ GeV

$$M_{eff} = \sum_i p_T^{jet,i} + \sum_j p_T^{lepton,j} + E_T^{miss}$$

$$M_T = \sqrt{2 p_T^{lepton} E_T^{miss} (1 - \cos \varphi(\vec{p}_T^{miss}, \vec{p}_T^{lepton}))}$$

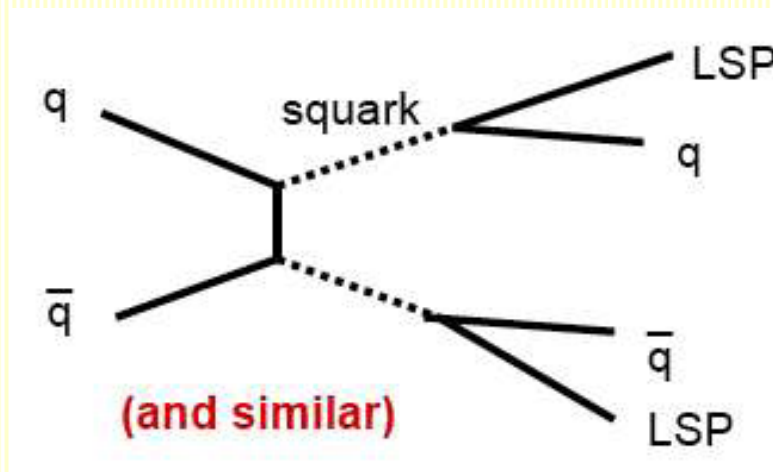
SU3: $m_0=100$ GeV, $m_{1/2}=300$ GeV, $A_0=-300$ GeV, $\tan\beta=6$, $\mu>0$



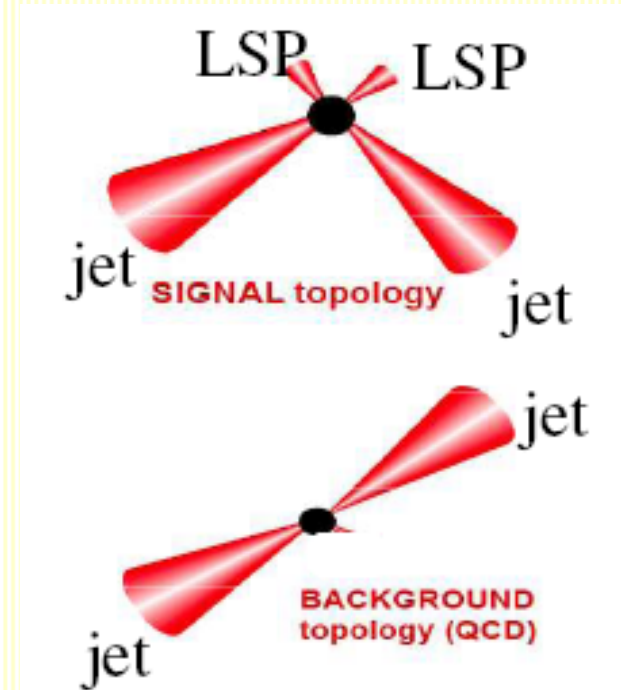
An example: the 0-lepton di-jet channel in CMS

Inspired by L.Randall & D.Tucker-Smith doi:10.1103/PhysRevLett.101.221803

- Probing the squark \rightarrow quark + neutralino channel
($m_{\text{squark}} < m_{\text{gluino}}$)



- Topology: 2 jets + $E_{\text{T}}^{\text{miss}}$



An example: the 0-lepton di-jet channel in CMS

- Introduce α_T (= 0.5 for perfectly measured QCD events):

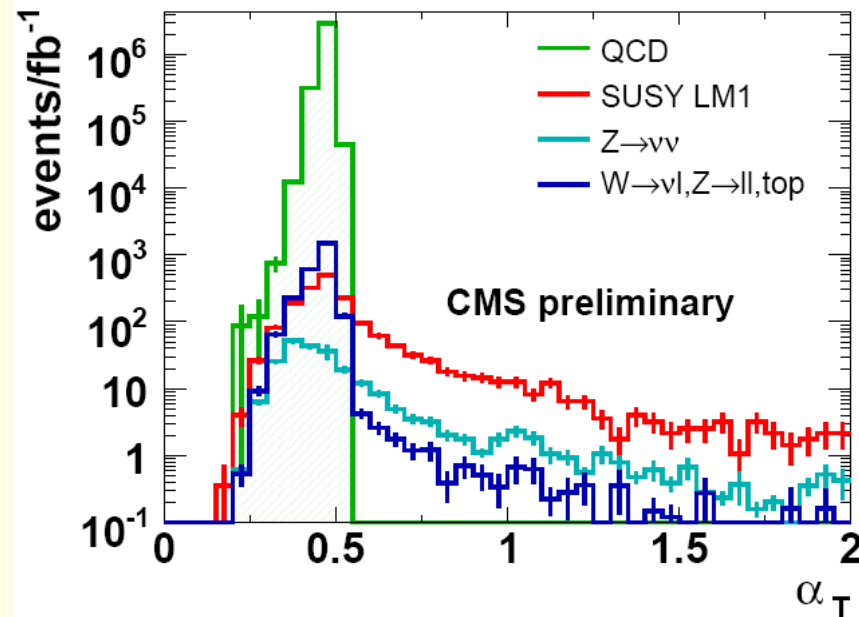
$$\alpha_T = \frac{E_T^{j2}}{\sqrt{2E_T^{j1}E_T^{j2}(1 - \cos \Delta\phi)}} = \frac{\sqrt{E_T^{j2}/E_T^{j1}}}{\sqrt{2(1 - \cos \Delta\phi)}}$$

E_T^{miss} is not the dominant search variable!

- $|\eta^{J1}| < 2.5$
- $E_T^{J1,J2} > 50 \text{ GeV}$
- No e or μ with $p_T > 10 \text{ GeV}$
- No 3rd jet with $p_T > 50 \text{ GeV}$
- $\Delta\phi(-|\Sigma\vec{p}_T|, J_i) > 0.3$ for $i=1,2,3$
- $p_T^{J1} + p_T^{J2} > 500 \text{ GeV}$

An example: the 0-lepton di-jet channel in CMS

- All cuts but α_T applied:

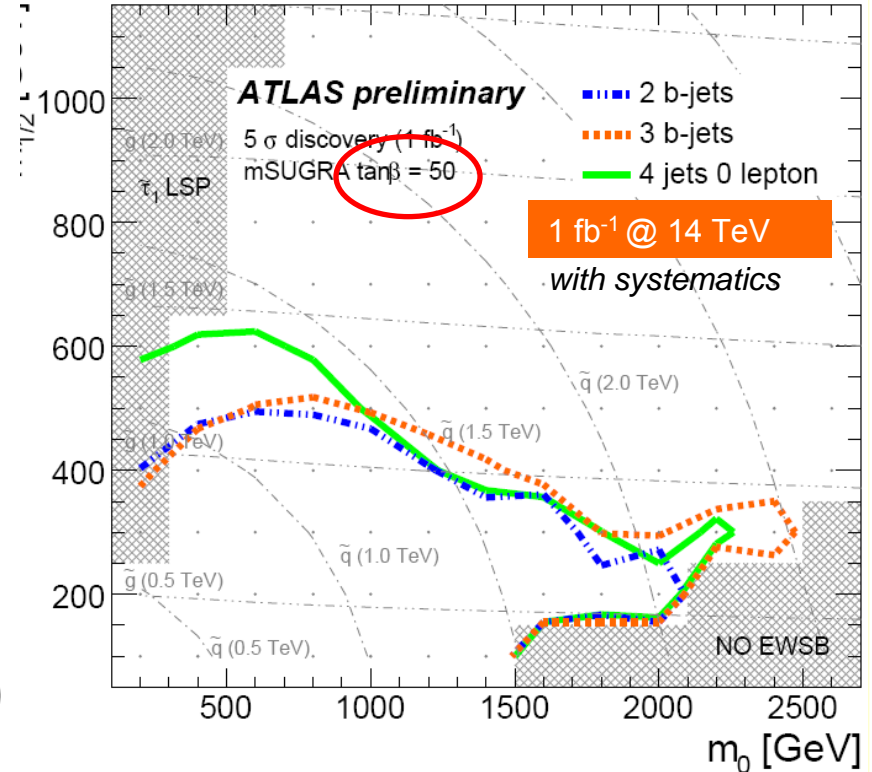
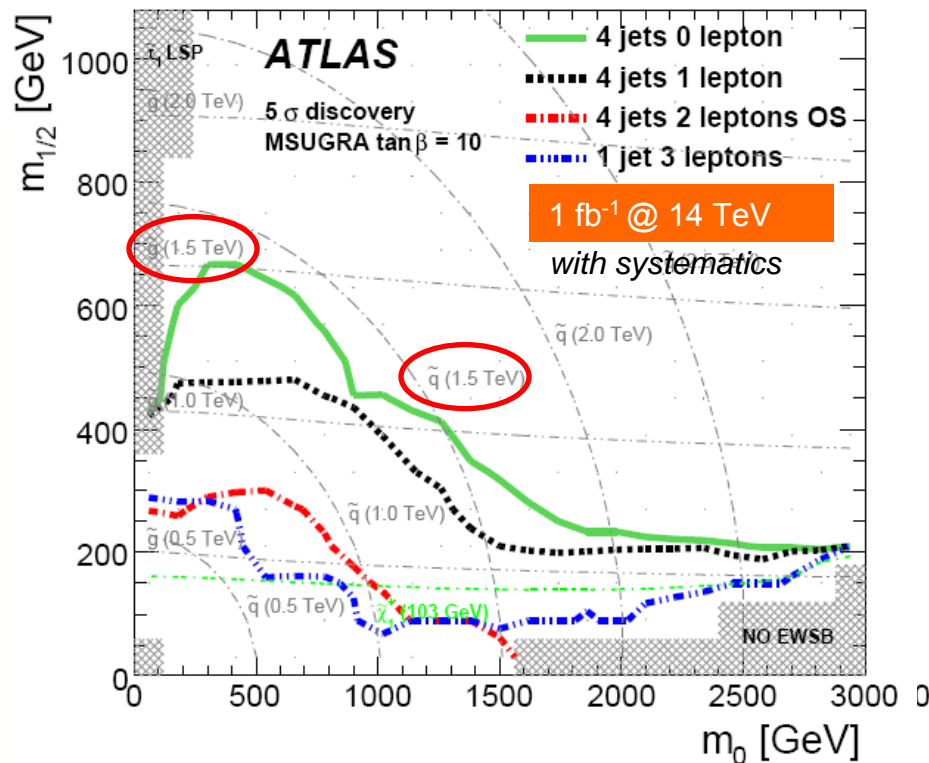


LM1: $m_0=60$ GeV, $m_{1/2}=250$ GeV, $A_0=0$, $\tan\beta=10$, $\mu>0$

- Number of events for 1 fb⁻¹ @ 14 TeV after $\alpha_T > 0.55$:

LM1	QCD	Z+jets	tt, W+jets
439	0	58	19

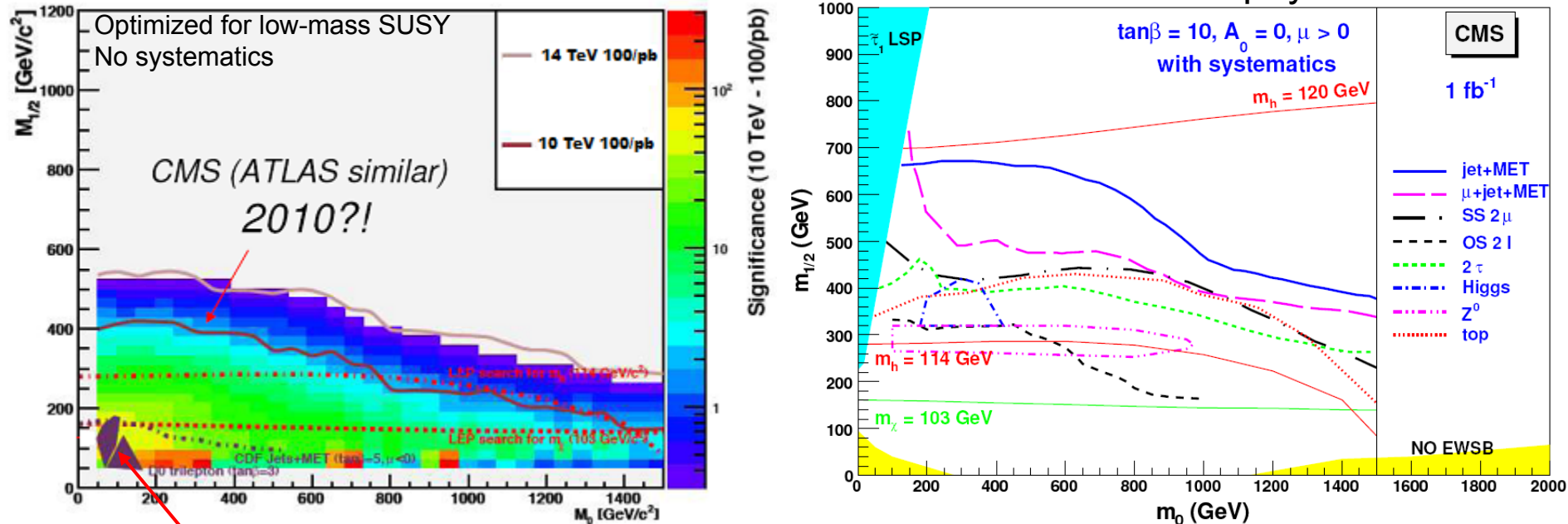
Discovery reach in ATLAS



- The 0-lepton channel has the highest reach
- The 1-lepton channel is more robust against QCD
- The b-jet channel is competitive with the 0-lepton one at high $\tan \beta$ (better reach at high m_0 : 3-body decays of the gluinos involving t and b)

Discovery reach in CMS

CMS physics TDR Vol.II



Tevatron

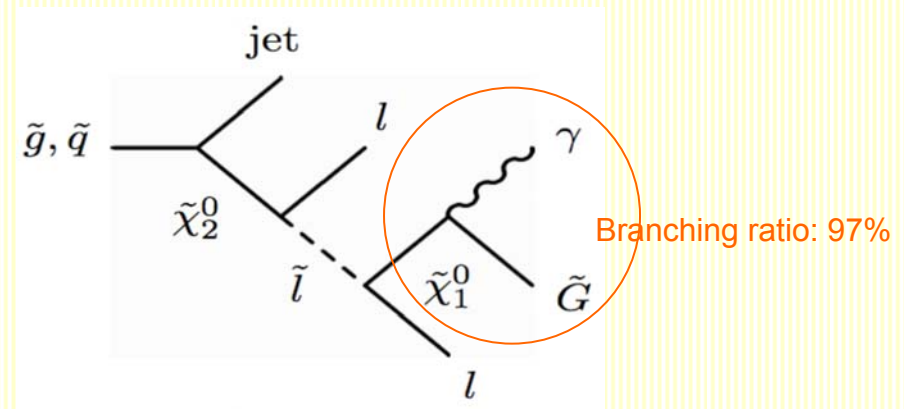
- The most powerful channel in CMS is also the jets + E_T^{miss} channel
- With only 50 pb⁻¹ of understood data @ 10 TeV, the LHC can go significantly beyond the reach of the Tevatron

Gauge Mediated SUSY Breaking in ATLAS

- SUSY broken by gauge interactions through messenger gauge fields
- 6 free parameters:

Λ	SUSY breaking scale
M	Messenger mass scale
$\tan\beta$	Ratio of Higgs VEVs
N	Number of messenger multiplets
$\text{Sign}(\mu)$	Sign of Higgs mass parameter
C_{grav}	Scale factor of Gravitino coupling

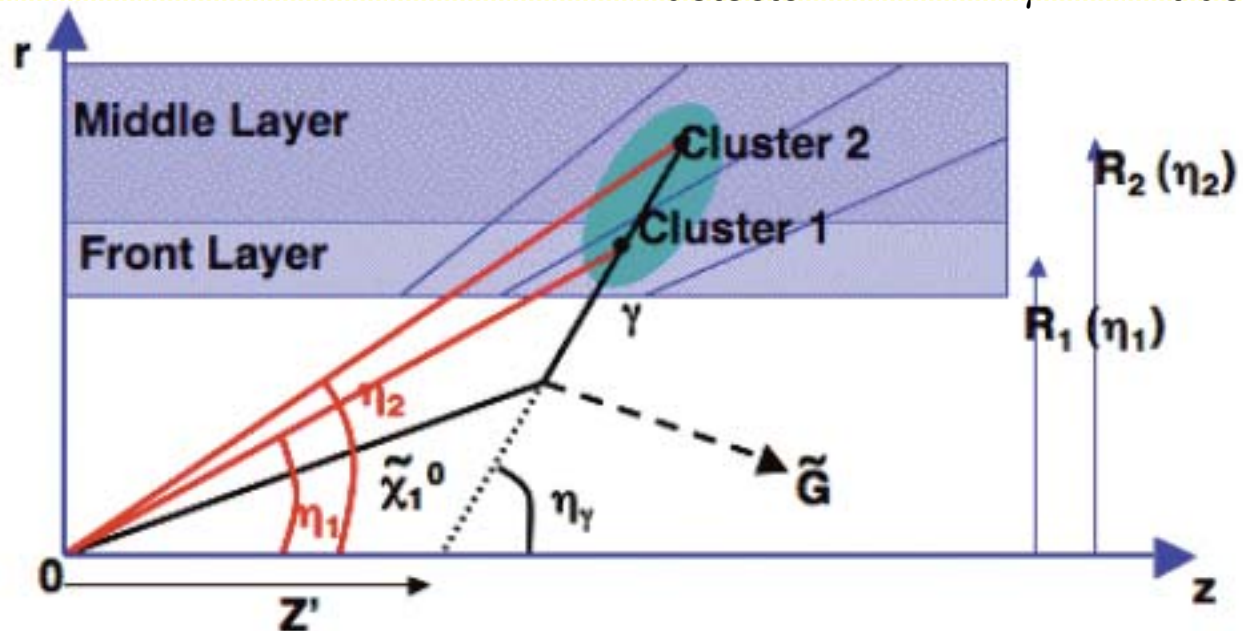
- If $\tan\beta$ is low and $N=1$: $\tilde{\chi}_1^0$ is the NLSP
- Signature:
Jets, 2 leptons, E_T^{miss} , 2 γ



Non-pointing photons in ATLAS

- If the decay length of the $\tilde{\chi}_1^0$ is comparable to the size of ATLAS inner detector, there could be high- p_T *non-pointing* photons:

$$\eta_2 (\ll \eta_{\text{detector}}) \neq \eta_\gamma (\ll \eta_{\text{true}})$$

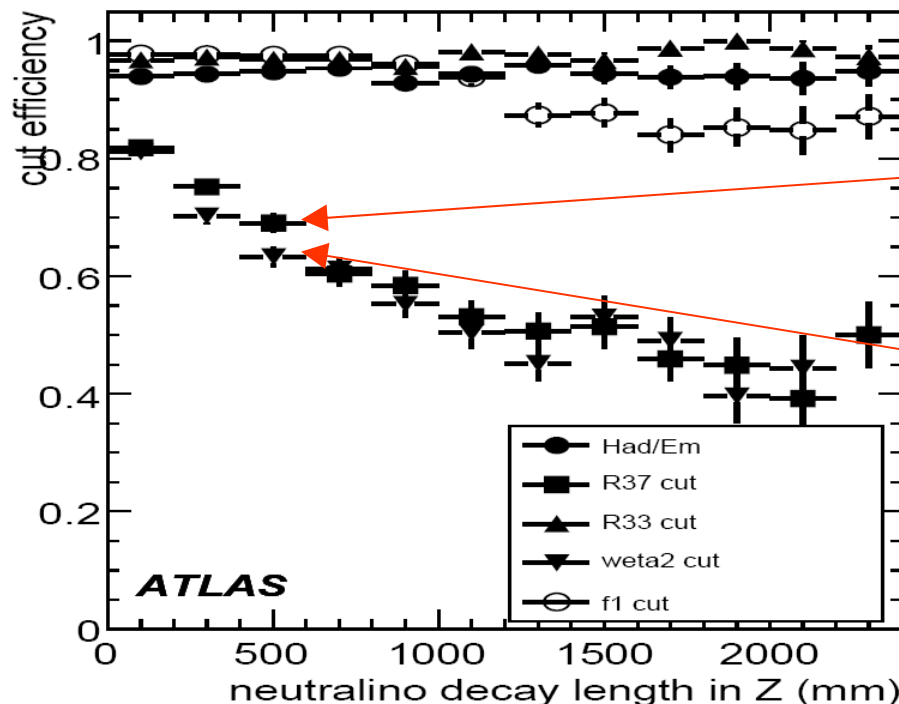


Non-pointing photons
can have a wider
shower profile than
pointing photons

Identifying non-pointing photons in ATLAS

- Idea:

- only use the photon identification selections which are unbiased w.r.t. the neutralino decay length (loosen the ID cuts)
- this increases the fake rate from jets:
 - from $0.19 \pm 0.03\%$ to $0.70 \pm 0.07\%$



For example:

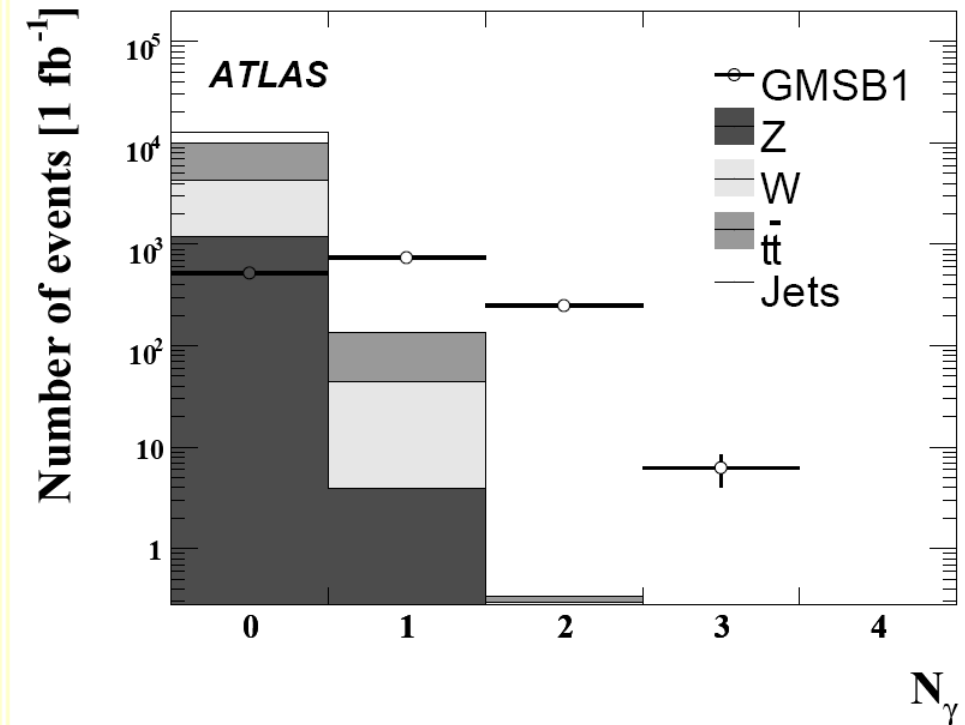
$E_{3 \times 7 \text{ window}} / E_{7 \times 7 \text{ window}}$ in the 2nd sampling layer

Width of the cluster in the 2nd sampling layer

Gauge Mediated SUSY Breaking in ATLAS

Selection:

- $N_{\text{Jets}} \geq 4$ with $p_T > 50$ GeV
- $p_T^{j1} > 100$ GeV
- $E_{T^{\text{miss}}} > \max(100 \text{ GeV}, 0.2M_{\text{eff}})$
- $t\bar{t}$ is the main background

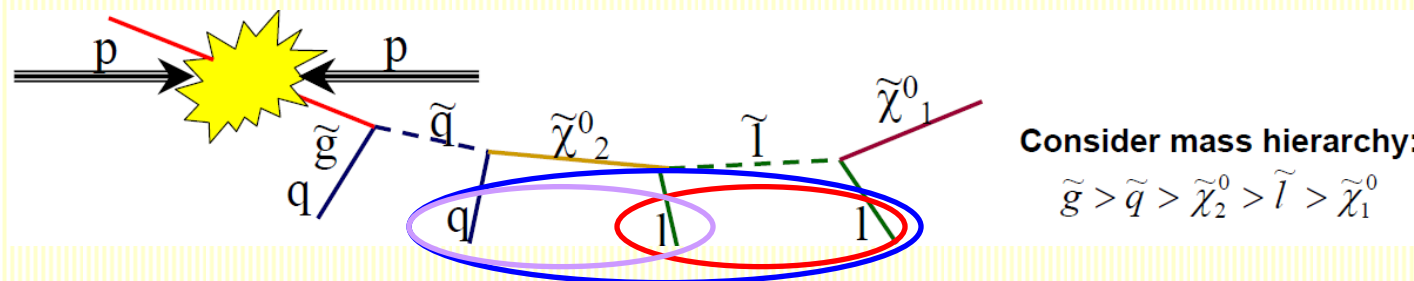


GMSB1: $\Lambda=90$ TeV, $M_m=500$ TeV, $C_G=1.0$, $N_5=1$, $\tan\beta=5$, $\mu>0$

- Techniques are being developed to extract the lifetime of the neutralino using timing and directional information from LAr calorimeter
→ talk by Andrea Giammanco in tomorrow's Higgs and New Physics session

Exclusive measurements

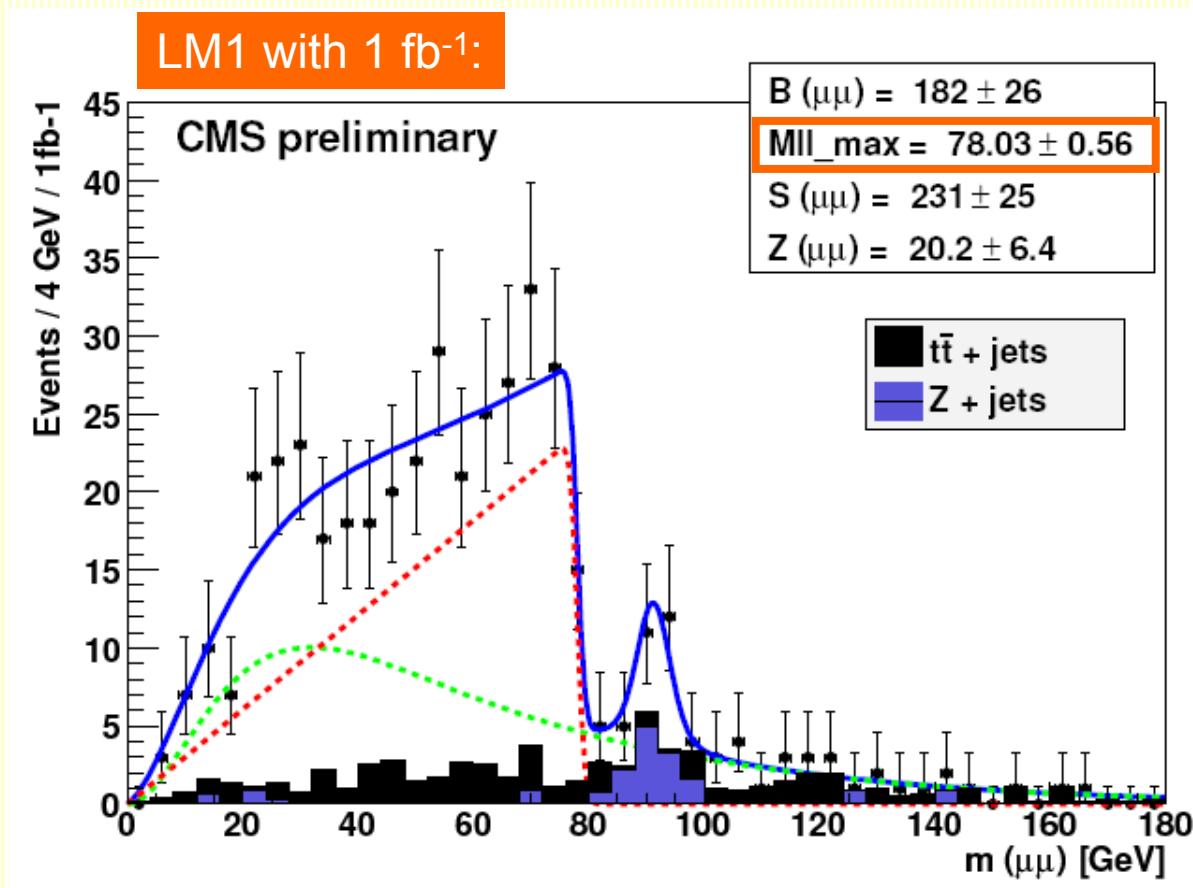
- When SUSY is discovered (if SUSY exists!), the experimental emphasis will be put on measuring the sparticle mass spectrum and deriving the parameters of the model
- The strategy is to exploit kinematics of long decay chains
- The first chain likely to be reconstructed:



- Due to the LSPs, the decay chain cannot be completely reconstructed
- Edges, rather than mass peaks, are measured in the invariant mass distribution of the decay products (e.g. m_{ll} , m_{llq} , m_{lq} , etc)

Exclusive measurements in CMS

Measure the different endpoints, for example $m(\mu\mu)^{\max}$:



— Fit
 Signal
 Flavour-symmetric background

$$\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_{L,R} \ell \rightarrow \ell \ell \tilde{\chi}_1^0$$

$$m_{\ell\ell}^{\max} = m_{\tilde{\chi}_2^0} \sqrt{1 - \frac{m_{\ell_R}^2}{m_{\tilde{\chi}_2^0}^2}} \sqrt{1 - \frac{m_{\tilde{\chi}_1^0}^2}{m_{\ell_R}^2}}$$

Exclusive measurements in ATLAS

Even with 1 fb^{-1} some measurements are possible:

Example point with 1 fb^{-1} :

Observable	SU3 m_{meas} [GeV]	SU3 m_{MC} [GeV]
$m_{\tilde{\chi}_1^0}$	$88 \pm 60 \mp 2$	118
$m_{\tilde{\chi}_2^0}$	$189 \pm 60 \mp 2$	219
$m_{\tilde{q}}$	$614 \pm 91 \pm 11$	634
$m_{\tilde{\ell}}$	$122 \pm 61 \mp 2$	155
Observable	SU3 Δm_{meas} [GeV]	SU3 Δm_{MC} [GeV]
$m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$	$100.6 \pm 1.9 \mp 0.0$	100.7
$m_{\tilde{q}} - m_{\tilde{\chi}_1^0}$	$526 \pm 34 \pm 13$	516.0
$m_{\tilde{\ell}} - m_{\tilde{\chi}_1^0}$	$34.2 \pm 3.8 \mp 0.1$	37.6

Error: MIGRAD + jet energy scale

However:

- The statistical errors are big
- The mass differences are better measured than the absolute masses

Similar findings in CMS

Exclusive measurements in ATLAS

Ultimate goal:

Determine SUSY parameters from endpoint measurements

- *Fittino* package to determine the mSUGRA parameters
- First results:
 - $\text{Sign}(\mu) > 0$ is favoured but $\text{sign}(\mu) < 0$ is not ruled out ($\chi^2=12.6$ vs $\chi^2=15.4$)
 - M_0 and $M_{1/2}$ are well constrained
 - A_0 and $\tan\beta$ are more problematic: there is no information from the Higgs sector at low integrated luminosity

Similar findings in CMS

Example point with 1 fb⁻¹:

Parameter	SU3 value	fitted value	exp. unc.
$\text{sign}(\mu) = +1$			
$\tan\beta$	6	7.4	4.6
M_0	100 GeV	98.5 GeV	± 9.3 GeV
$M_{1/2}$	300 GeV	317.7 GeV	± 6.9 GeV
A_0	-300 GeV	445 GeV	± 408 GeV
$\text{sign}(\mu) = -1$			
$\tan\beta$		13.9	± 2.8
M_0		104 GeV	± 18 GeV
$M_{1/2}$		309.6 GeV	± 5.9 GeV
A_0		489 GeV	± 189 GeV

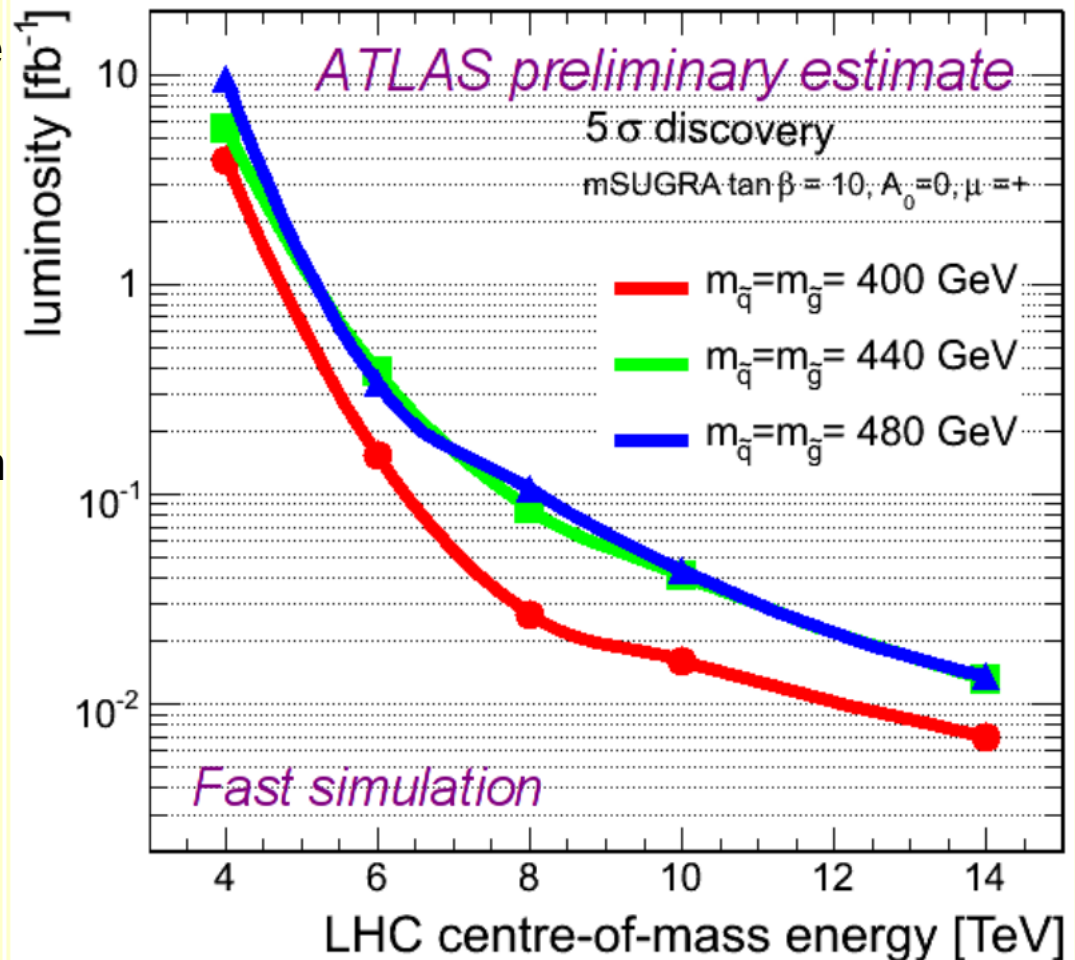
Conclusion

- Looking for excesses in inclusive search channels is the first step towards finding SUSY at the LHC
- Strategies are being developed for different scenarios, e.g. the non-pointing photons in GMSB
- The LHC should be able to discover squarks and gluinos with masses up to 1-1.5 TeV (if systematics are as expected for 1fb^{-1})
- Backgrounds and detectors need to be understood
- After a discovery, the next step would be to select specific SUSY decay chains to measure the properties of new particles
- Some masses can already be measured with 1fb^{-1} for low-mass SUSY scenarios

Backup slides

Different centre-of-mass energies

- Discovery sensitivity possible for SUSY, in an mSUGRA model with equal mass squarks and gluinos.
- The lepton plus jets plus missing-ET channel is employed, as this should be understood more rapidly than the statistically more powerful inclusive jets+missing-Et channel.
- The current published Tevatron limits in this model are around 400 GeV



The SUSY points in this talk

- mSUGRA:

Point	m_0 [GeV]	$m_{1/2}$ [GeV]	A_0 [GeV]	$\tan\beta$	$\text{sign}(\mu)$
Coannihilation (SU1)	70	350	0	10	+
Focus Point (SU2)	3550	300	0	10	+
Bulk (SU3)	100	300	-300	6	+
Low Mass (SU4)	200	160	-400	10	+
Funnel (SU6)	320	375	0	50	+
Coannihilation (SU8.1)	210	360	0	40	+

Label	σ^{LO} (pb)
SU1	8.15
SU2	5.17
SU3	20.85
SU4	294.46
SU6	4.47
SU8.1	6.48
SU9	2.46

Particle	SU1	SU2	SU3	SU4	SU6	SU8.1	SU9
\tilde{d}_L	764.90	3564.13	636.27	419.84	870.79	801.16	956.07
\tilde{u}_L	760.42	3563.24	631.51	412.25	866.84	797.09	952.47
\tilde{b}_1	697.90	2924.80	575.23	358.49	716.83	690.31	868.06
\tilde{t}_1	572.96	2131.11	424.12	206.04	641.61	603.65	725.03
\tilde{d}_R	733.53	3576.13	610.69	406.22	840.21	771.91	920.83
\tilde{u}_R	735.41	3574.18	611.81	404.92	842.16	773.69	923.49
\tilde{b}_2	722.87	3500.55	610.73	399.18	779.42	743.09	910.76
\tilde{t}_2	749.46	2935.36	650.50	445.00	797.99	766.21	911.20
\tilde{e}_L	255.13	3547.50	230.45	231.94	411.89	325.44	417.21
$\tilde{\nu}_e$	238.31	3546.32	216.96	217.92	401.89	315.29	407.91
$\tilde{\tau}_1$	146.50	3519.62	149.99	200.50	181.31	151.90	320.22
$\tilde{\nu}_\tau$	237.56	3532.27	216.29	215.53	358.26	296.98	401.08
\tilde{e}_R	154.06	3547.46	155.45	212.88	351.10	253.35	340.86
$\tilde{\tau}_2$	256.98	3533.69	232.17	236.04	392.58	331.34	416.43
\tilde{g}	832.33	856.59	717.46	413.37	894.70	856.45	999.30
$\tilde{\chi}_1^0$	136.98	103.35	117.91	59.84	149.57	142.45	173.31
$\tilde{\chi}_2^0$	263.64	160.37	218.60	113.48	287.97	273.95	325.39
$\tilde{\chi}_3^0$	466.44	179.76	463.99	308.94	477.23	463.55	520.62
$\tilde{\chi}_4^0$	483.30	294.90	480.59	327.76	492.23	479.01	536.89
$\tilde{\chi}_1^+$	262.06	149.42	218.33	113.22	288.29	274.30	326.00
$\tilde{\chi}_2^+$	483.62	286.81	480.16	326.59	492.42	479.22	536.81
h^0	115.81	119.01	114.83	113.98	116.85	116.69	114.45
H^0	515.99	3529.74	512.86	370.47	388.92	430.49	632.77
A^0	512.39	3506.62	511.53	368.18	386.47	427.74	628.60
H^\pm	521.90	3530.61	518.15	378.90	401.15	440.23	638.88
t	175.00	175.00	175.00	175.00	175.00	175.00	175.00

Sample	m_0 (GeV)	$m_{1/2}$ (GeV)	A_0	$\tan\beta$	$\text{sign}(\mu)$	σ NLO (pb)	(LO) (pb)	lightest \tilde{q} (GeV)	χ_1^0 (GeV)
LM1	60	250	0	10	+	54.86	(43.28)	410 (\tilde{t}_1)	97

- GMSB:

name	NLO (LO) σ [pb]	Λ [TeV]	M_m [TeV]	C_G	$c\tau$ [mm]	$M_{\tilde{\chi}_1^0}$ [GeV]
GMSB1	7.8 (5.1)	90	500	1.0	1.1	118.8

$N_5 = 1, \tan\beta = 5, \text{sgn}(\mu) = +$