Construction of the ground state of Matrix Theory: Near the Origin

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J. Hoppe, D. Lundholm, M.T., arXiv:0803.1316, arXiv:0809.5270, arXiv:0809.5271

Plan

- 1. Matrix Theory
- 2. The construction of the ground state near the origin
- 3. Summary

Matrix Theory

 - membranes - Dirac 62' quantum membranes → YMQM - Goldstone, Hoppe 82' supermembranes → SYMQM - de Wit, Hoppe, Nicolai 87' -femoto-universe of YM → YMQM - Bjorken 79'
 - small volume YM → Lüscher - 82'
 -toy model - SYMQM - Claudson, Halpern 85'
 - M theory on a light cone in the IMF → SYMQM
 -Banks, Fischler, Shenker, Susskind 97'

Supermembranes

 \blacktriangleright quantum description \rightarrow matrix regularization:

$$H_{reg.} = Tr\left(P_sP_s - \frac{1}{2}[X_s, X_t][X_s, X_t] - iX_s\theta\gamma^s\theta\right)$$

$$\begin{aligned} X_{s} &= x_{sA}T_{A}, \quad P_{s} = p_{sA}T_{A} \quad \theta_{\alpha} = \theta_{\alpha A}T_{A}, \quad T_{A} \in \mathfrak{su}(N) \\ [x_{sA}, p_{tB}] &= i\delta_{st}\delta_{AB}, \quad \{\theta_{\alpha A}, \theta_{\beta B}\} = \delta_{\alpha\beta}\delta_{AB}, \\ \gamma^{s} - 16 \times 16, real \end{aligned}$$

$$s, t = 1, \dots, 9, \quad A, B = 1, \dots, N^2 - 1, \quad \alpha, \beta = 1, \dots, 16$$
$$G_A|_{\mathcal{H}} = 0, \qquad G_A = f_{ABC} \left(x_{sB} p_{sC} - \frac{i}{2} \theta_{\alpha B} \theta_{\alpha C} \right)$$

-de Wit, Hoppe, Nicolai, 87'

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• Does the ground state exist? $Q_{\alpha}\Psi = 0$

$$2H\delta_{\alpha,\beta} = \{Q_{\alpha}, Q_{\beta}\}, \quad Q_{\beta} = (-i\gamma_{\beta\alpha}^{s}\partial_{sA} + \frac{1}{2}\gamma_{\beta\alpha}^{st}f_{ABC}x_{sB}x_{tC})\theta_{\alpha A}$$

- Witten index computations suggest YES Yi 97', Sethi, Stern 98', Moore, Nekrasov, Shatashvili 98'

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- Can it be constructed?
 - asymptotically (large x_{sA}) Bach, Bordemann, Fröhlich, Graf, Halpern, Hasler, Hoppe, Konechny, Lundholm, Plefka, Schwatrz, Suter, Yau - 97'-07' - small x_{sA} ...

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Construction of the ground state near the origin

Taylor expansion

$$\Psi(x,\theta) = \psi^{(0)} + x_{sA}\psi^{(1)}_{sA} + \frac{1}{2}x_{sA}x_{tB}\psi^{(2)}_{sA\ tB} + \dots$$

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• $Q_{eta}\Psi=0$ gives 3 towers (for $\psi^{(3k)}$, $\psi^{(3k+1)}$ i $\psi^{(3k+2)}$)

$$\begin{split} \gamma^{s}_{\beta\alpha}\theta_{\alpha A}\psi^{(1)}_{sA} &= 0,\\ \gamma^{s}_{\beta\alpha}\theta_{\alpha A}\psi^{(2)}_{sA,tB} &= 0,\\ -i\gamma^{s}_{\beta\alpha}\theta_{\alpha A}\psi^{(3)}_{sA,tB,uC} + f_{ABC}\gamma^{tu}_{\beta\alpha}\theta_{\alpha A}\psi^{(0)} &= 0 \quad \text{etc} \end{split}$$

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• what about $\psi^{(0)}$? Ψ is a SO(9) singlet -Hasler, Hoppe 02'

• for
$$SU(N = 2)$$
:

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• existence and uniqueness of $\psi^{(0)}$

- J. Wosiek 05' In the Fock space $(\theta_{\alpha A} \rightarrow f^{\dagger}_{\hat{\alpha}A}, f_{\hat{\alpha}A})$ we look for an algebraic combination of SU(2) invariants $f^{\dagger}_{\hat{\alpha}A}f^{\dagger}_{\hat{\beta}A}$ which are SO(9) invariant.

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• paper-pencil derivation of $\psi^{(0)}$?

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Maciej Trzetrzelewski, Jagiellonian University Construction of the ground state of Matrix Theory: Near the Origin

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- ▶ the Hilbert space $\mathcal{H} = \mathcal{H}_{256,A=1} \otimes \mathcal{H}_{256,A=2} \otimes \mathcal{H}_{256,A=3}$
- ▶ under *SO*(9), 256 = 44 + 84 + 128 $\rightarrow |st\rangle$, $|stu\rangle$, $|s\alpha\rangle$
- $\psi^{(0)}$ could be eg.

 $|su\rangle_1|tu\rangle_2|st\rangle_3, |s\alpha\rangle_1|t\alpha\rangle_2|st\rangle_3, |suv\rangle_1|tuv\rangle_2|st\rangle_3$

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• we need
$$\theta_{\alpha A} | st \rangle_A = ?$$
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 $|su\rangle_{1}|tu\rangle_{2}|st\rangle_{3}, |s\alpha\rangle_{1}|t\alpha\rangle_{2}|st\rangle_{3}, |suv\rangle_{1}|tuv\rangle_{2}|st\rangle_{3}$ we need $\theta_{\alpha A}|st\rangle_{A} =?, \quad \theta_{\alpha A}|stu\rangle_{A} =?, \quad \theta_{\alpha A}|i\alpha\rangle_{A} =?$ $2\theta_{\alpha}|st\rangle = \gamma_{\alpha\beta}^{s}|t\beta\rangle + \gamma_{\alpha\beta}^{t}|s\beta\rangle,$ $\theta_{\alpha}|stu\rangle = \frac{i}{\sqrt{2}}\left(\gamma_{\alpha\beta}^{st}|u\beta\rangle + \gamma_{\alpha\beta}^{us}|t\beta\rangle + \gamma_{\alpha\beta}^{tu}|s\beta\rangle\right).$

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Define

$$|||1\rangle := |su\rangle_1 |tu\rangle_2 |st\rangle_3,$$

 $||| 1\rangle := |suv\rangle_1 |tuv\rangle_2 |st\rangle_3 + |tuv\rangle_1 |st\rangle_2 |suv\rangle_3 + |st\rangle_1 |suv\rangle_2 |tuv\rangle_3$

then it follows that

$$\phi:=ert ert 1
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angle$$

is SU(2) invariant $\rightarrow \psi^{(0)} \sim \phi$ (no 128 content)

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 an independent approach - explicit construction in the Fock space given by

$$\theta_{\alpha} \to \lambda_{\hat{\alpha}}, \lambda_{\hat{\alpha}}^{\dagger}$$

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$$v_1 := b_{AB}b_{AB}, \quad v_2 := b^i_{AB}b^i_{AB}, \quad w_1 = b^i_{AB}b^j_{BC}b^j_{CD}b^j_{DA}, \quad w_2 = b^i_{AB}b^j_{BC}b^j_{CD}b^j_{DA}.$$

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► the basis $r_1 = v_1 v_1 v_1 |0\rangle, \quad r_2 = v_1 v_1 v_2 |0\rangle, \quad r_3 = v_1 v_2 v_2 |0\rangle, \quad r_4 = v_2 v_2 v_2 |0\rangle,$ $r_5 = v_1 w_1 |0\rangle, \quad r_6 = v_2 w_1 |0\rangle, \quad r_7 = v_1 w_2 |0\rangle.$

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 $\chi = 326304r_1 + 488136r_2 + 72612r_3 + 1377r_4 + 114576r_5 - 176528r_6 + 10296r_7,$

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• $\chi \sim \phi!$

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Summary

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- outstanding problems
 - does the ground state exist?
 - what is it?
 - other bound states?
 - large N limit?

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- G₂ deformation, group averaging techniques (J. Hoppe, D. Lundholm, M.T.)
- Taylor expansion $\rightarrow \Psi(0) = \phi$ (two ways)