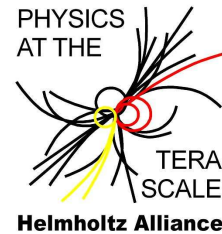


Squark and gluino production at the LHC

Anna Kulesza **RWTH**AACHEN



AK and L. Motyka, Phys. Rev. Let. **102**, 111802 (2009)

AK and L. Motyka, arXiv:0905.4749 [hep-ph]

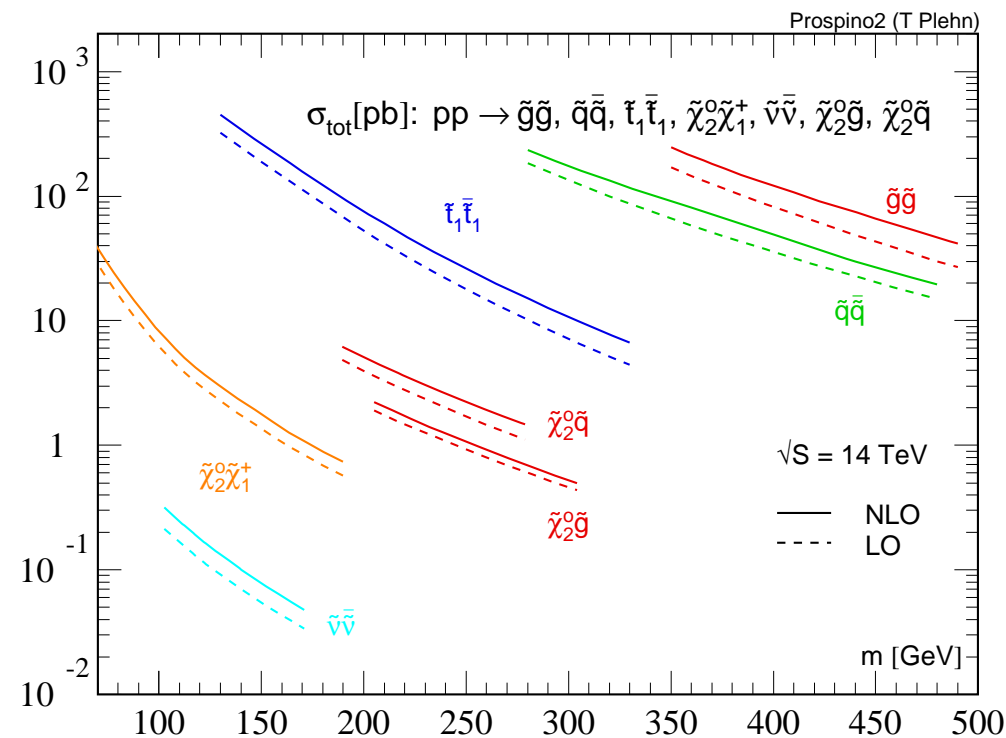
W. Beenakker, S. Brensing, M. Krämer, AK, E. Laenen and I. Niessen, in preparation

The 2009 Europhysics Conference on High Energy Physics

16-22 July 2009, Kraków, Poland

SUSY particle pair-production at the LHC

- MSSM: minimal content of SUSY particles + R -parity conservation
- At the LHC dominant sparticle production channels involve squarks (\tilde{q}) and gluinos (\tilde{g}) in the final state ($\tilde{q}\tilde{q}^*$, $\tilde{q}\tilde{q}$, $\tilde{q}\tilde{g}$, $\tilde{g}\tilde{g}$ pairs)

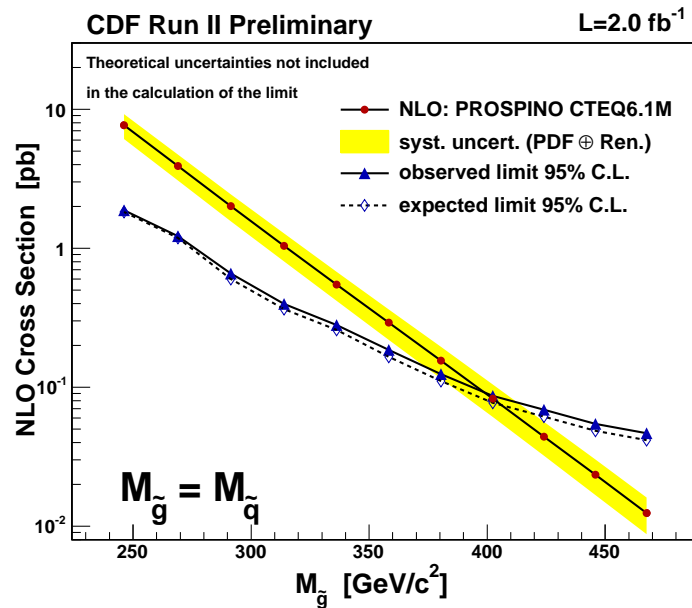


[Plehn, Prospino2]

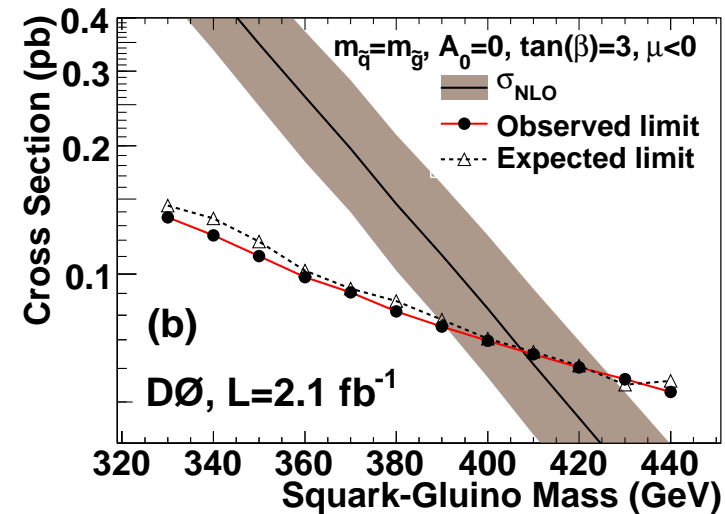
- Production cross sections large \Rightarrow “easy” SUSY discovery

Total cross sections

- Total cross sections: crucial for exclusion limits / useful for mass determination in case of discovery

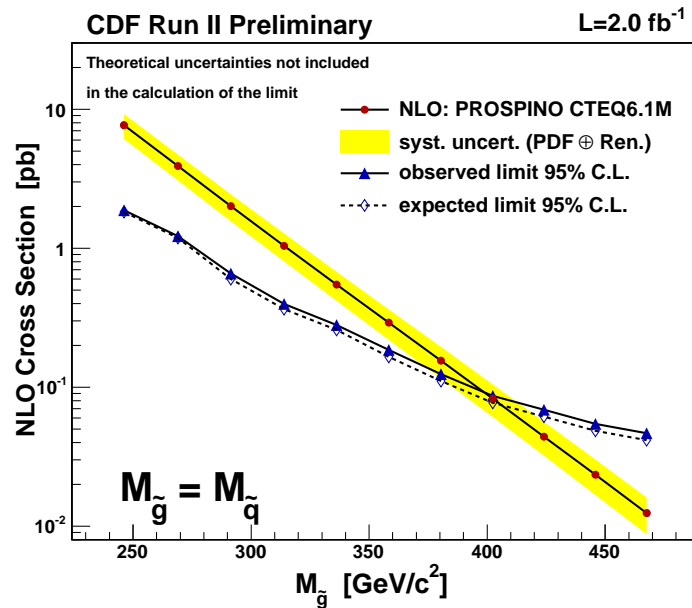


$(A_0 = 0, \text{sgn}(\mu) = -1, \tan \beta = 5)$

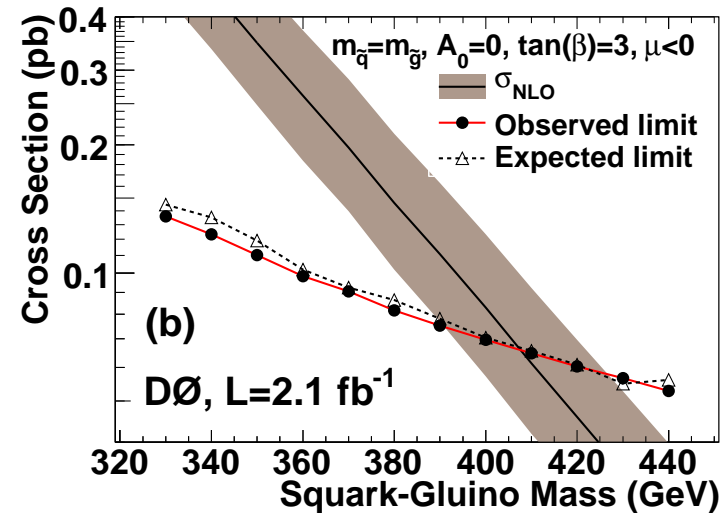


Total cross sections

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$$(A_0 = 0, \text{sgn}(\mu) = -1, \tan \beta = 5)$$



- Important to know total cross sections with high precision

Theoretical status

- Leading order = $\mathcal{O}(\alpha_s^2)$ [*Kane, Leveille'82*][*Harrison, Llewellyn Smith'84*][*Dawson, Eichten, Quigg'85*]
- Higher-order corrections to $\mathcal{O}(\alpha_s^2)$ processes
 - NLO SUSY-QCD corrections $\rightarrow \mathcal{O}(\alpha_s^3)$ [*Beenakker, Höpker, Spira, Zerwas'96*]
 - **This talk: NLL threshold resummed corrections**
 - For $\tilde{q}\tilde{q}$ production:
 - dominant NNLO contributions (NNLL-NNLO, Coulomb, scale dependence) $\rightarrow \mathcal{O}(\alpha_s^4)$ [*Langenfeld, Moch'09*]
 - resummation of Coulomb and threshold corrections \rightarrow see Ch. Schwinn's talk
 - EW corrections $\rightarrow \mathcal{O}(\alpha_s^2\alpha)$ [*Hollik, Kollar, Trenkel'07*][*Hollik, Mirabella'08*][*Hollik, Mirabella, Trenkel'08*][*Beccaria et al.'08*]

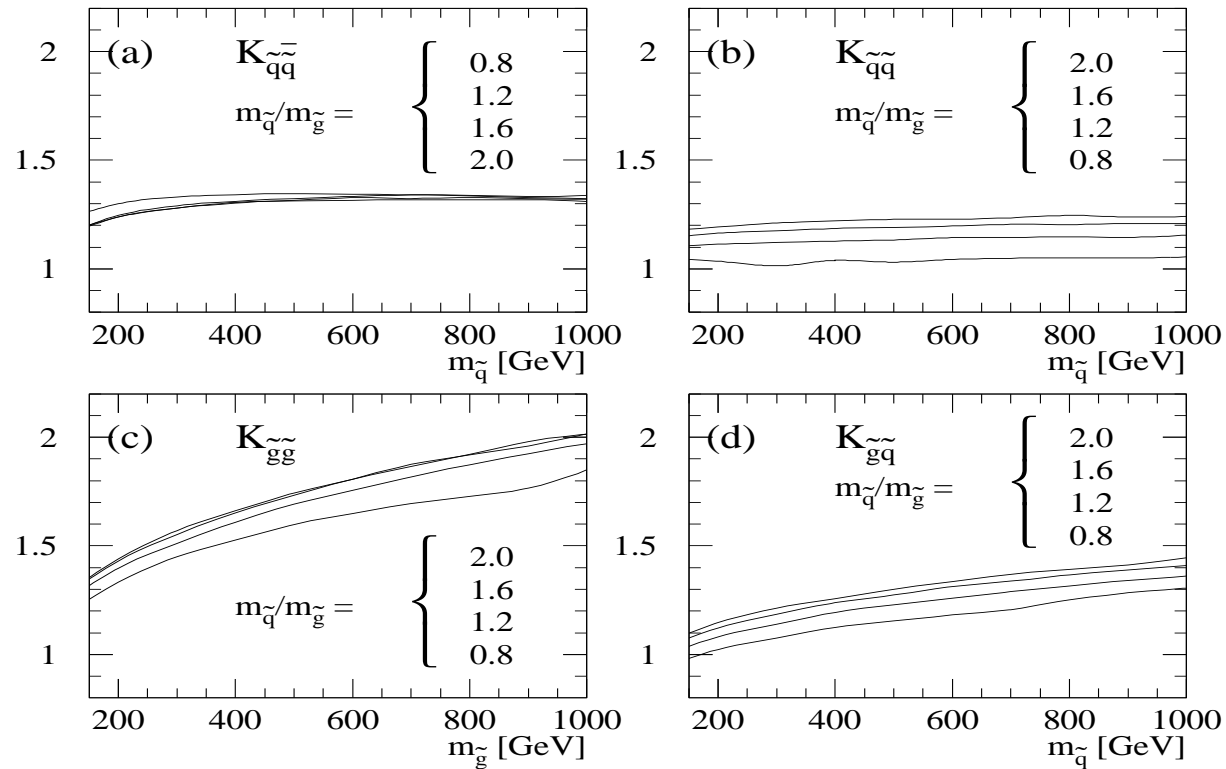
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- Tree-level EW effects of $\mathcal{O}(\alpha\alpha_s)$ and $\mathcal{O}(\alpha^2)$
 - QCD-EW interference and photon-induced contributions, tree-level EW
[*Bornhauser et al.'07*] [*Alan, Cankocak, Demir'07*][*Hollik, Kollar, Trenkel'07*][*Hollik, Mirabella'08*][*Hollik, Mirabella, Trenkel'08*][*Bozzi, Fuks, Klasen'05*]

Coloured sparticle production at NLO (SUSY-QCD)

[Beenakker, Höpker, Spira, Zerwas'96]

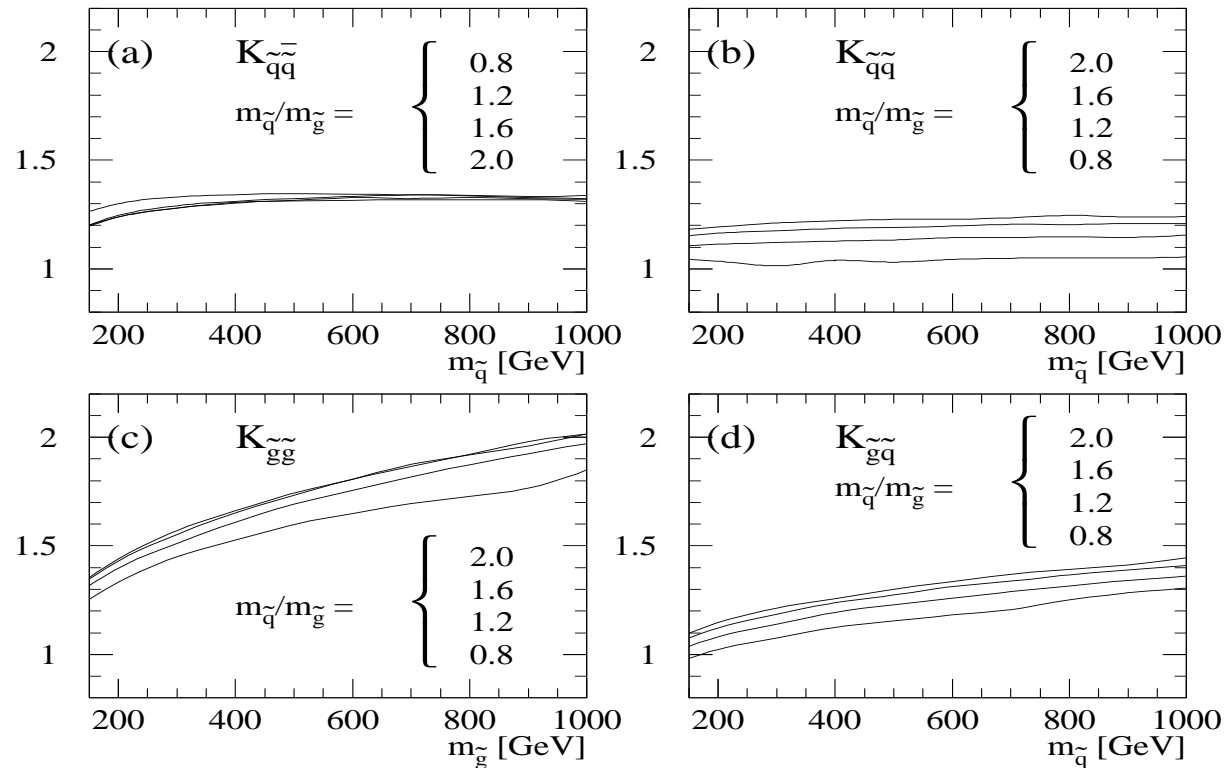
$$K_{ij} = \sigma_{ij}^{\text{NLO}} / \sigma_{ij}^{\text{LO}}$$



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$$K_{ij} = \sigma_{ij}^{\text{NLO}} / \sigma_{ij}^{\text{LO}}$$



⇒ Increase of the cross sections due to NLO SUSY-QCD corrections

Note: assume all squarks (\tilde{q}_L, \tilde{q}_R) mass degenerate; no final state stops ⇒

[Beenakker, Krämer, Plehn, Spira, Zerwas'98]

Higher-order soft gluon effects

- Large masses of squarks and gluons

⇒ often production close to threshold $\hat{s} \sim 4m^2$ ($m = \frac{m_{\tilde{q}} + m_{\tilde{g}}}{2}$ for $\tilde{q}\tilde{g}$ production)

⇒ real emission forced to be predominantly soft

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- At n -th order in α_s (wrt. LO) soft gluon contributions of the form

$$\alpha_s^n \log^m(\beta^2), \quad m \leq 2n \qquad \beta^2 = 1 - \frac{4m^2}{\hat{s}}$$

- the closer to threshold the more important the logarithmic terms

→ **expect large contributions from soft gluon radiation**

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In the limit of $\beta \rightarrow 0$ convergence of fixed-order expansion spoiled

Resummation

reorganization of the perturbative series = resummation

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From

$$\hat{\sigma} \sim c_{00} +$$

$$+ \alpha_s \left(c_{12} \log^2(\beta^2) + c_{11} \log(\beta^2) + c_{10} \right) \leftarrow \text{NLO}$$
$$+ \alpha_s^2 \left(c_{24} \log^4(\beta^2) + c_{23} \log^3(\beta^2) + c_{22} \log^2(\beta^2) + \dots \right) \leftarrow \text{NNLO}$$

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to

$$\hat{\sigma} \sim \mathcal{C}(\alpha_s) \exp \left(\underbrace{Lg_1(\alpha_s L)}_{LL} + \underbrace{g_2(\alpha_s L)}_{NLL} + \alpha_s \underbrace{g_3(\alpha_s L)}_{NNLL} + \dots \right)$$

$L = \log(\dots)$

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- Threshold corrections exponentiate in the space of Mellin moments taken wrt.

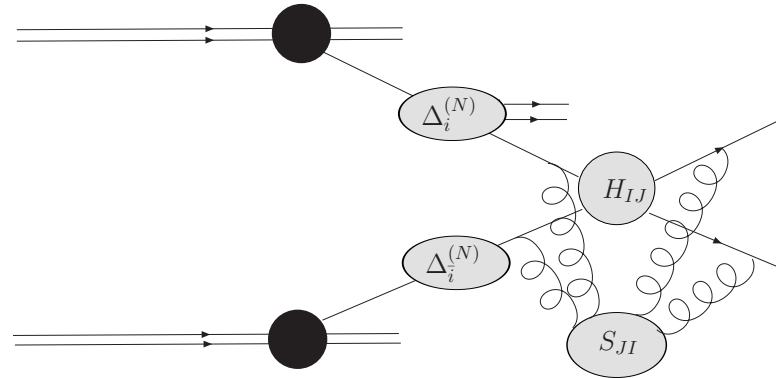
$$z \equiv \frac{4m^2}{\hat{s}} = 1 - \beta^2 \quad [\text{Sterman}'87][\text{Catani, Trentadue}'89]$$

$$f^{(N)} = \int_0^1 dz z^{N-1} f(z) \quad \Rightarrow \quad \log(\beta^2) = \log(1 - z) \longleftrightarrow \log(N)$$

Resummation for $2 \rightarrow 2$ with colour and masses

$2 \rightarrow 2$ with colour flow

[Kidonakis, Sterman'96-97][Kidonakis, Oderda, Sterman'98][Bonciani et al.'03]



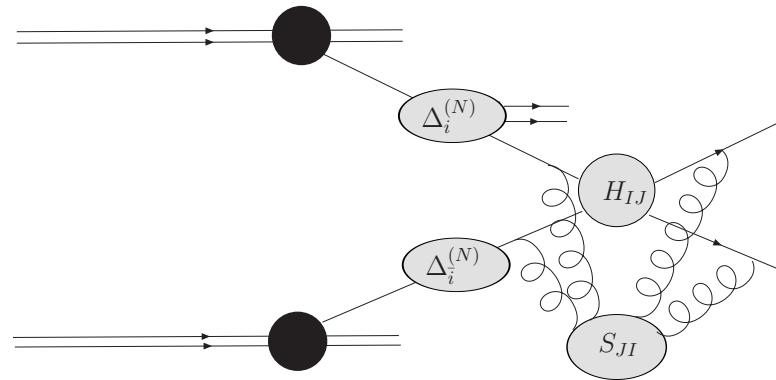
Up to NLL:

$$\hat{\sigma}_{ij \rightarrow kl}^{(N)} = \underbrace{H_{ij \rightarrow kl, IJ}^{(N)}}_{\text{hard function}} \times \underbrace{\Delta_i^{(N)} \Delta_j^{(N)}}_{\substack{\text{soft-collinear radiation} \\ \text{universal factors; KNOWN}}} \times \underbrace{S_{ij \rightarrow kl, JI}}_{\substack{\text{soft wide-angle emission} \\ \text{process-dependent}}}$$

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S_{JI} from solving renormalization group equation

$$\left(\mu \frac{\partial}{\partial \mu} + \beta(g) \frac{\partial}{\partial g} \right) S_{JI}^{(N)} = -\Gamma_{JK}^\dagger S_{KI}^{(N)} - S_{JL}^{(N)} \Gamma_{LI}$$

Anomalous dimensions

- Need 1-loop anomalous dimension matrices in order to resum up to NLL
 - massless $2 \rightarrow 2$ QCD processes [*Kidonakis, Oderda, Sterman'98*][*Bonciani et al.'03*]
 - massive case: heavy quark $Q\bar{Q}$ production [*Kidonakis, Sterman'96*][*Bonciani et al.'98*]
- Calculation of 1-loop soft anomalous dimension matrices Γ_{JI} for $2 \rightarrow 2$ processes with nontrivial colour structure and massive particles in the final state

$$\begin{aligned}\tilde{q}\tilde{q} & \mathbf{3} \otimes \bar{\mathbf{3}} = & \mathbf{1} \oplus \mathbf{8} \\ \tilde{q}\tilde{q} & \mathbf{3} \otimes \mathbf{3} = & \bar{\mathbf{3}} \oplus \mathbf{6} \\ \tilde{q}\tilde{g} & \mathbf{3} \otimes \mathbf{8} = & \mathbf{3} \oplus \bar{\mathbf{6}} \oplus \mathbf{15} \\ \tilde{g}\tilde{g} & \mathbf{8} \otimes \mathbf{8} = & \mathbf{1} \oplus \mathbf{8} \oplus \mathbf{8} \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}\end{aligned}$$

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- Threshold limit $\hat{s} \rightarrow 4m^2$
 - Γ_{JI} matrices calculated in the s-channel colour basis become diagonal
 - Coefficients governing the soft radiation factors correspond to the values of the quadratic Casimir operators for the SU(3) representations

soft gluon radiation from the total colour charge of the heavy particle pair produced at threshold

Resummation-improved NLL+NLO total cross section

NLL resummed expression has to be **matched** with the full **NLO** result

$$\begin{aligned} \sigma_{h_1 h_2 \rightarrow kl}^{(\text{match})}(\rho, m^2, \{\mu^2\}) &= \sum_{i,j=q,\bar{q},g} \int_{C_{MP}-i\infty}^{C_{MP}+i\infty} \frac{dN}{2\pi i} \rho^{-N} f_{i/h_1}^{(N+1)}(\mu_F^2) f_{j/h_2}^{(N+1)}(\mu_F^2) \\ &\times \left[\hat{\sigma}_{ij \rightarrow kl, N}^{(\text{res})}(m^2, \{\mu^2\}) - \hat{\sigma}_{ij \rightarrow kl, N}^{(\text{res})}(m^2, \{\mu^2\}) \Big|_{\text{NLO}} \right] \\ &+ \sigma_{h_1 h_2 \rightarrow kl}^{\text{NLO}}(\rho, m^2, \{\mu^2\}), \end{aligned}$$

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- Inverse Mellin transform evaluated using a contour in the complex N space according to 'Minimal Prescription' [*Catani, Mangano, Nason Trentadue'96*]
- NLO cross sections evaluated with publicly available code PROSPINO

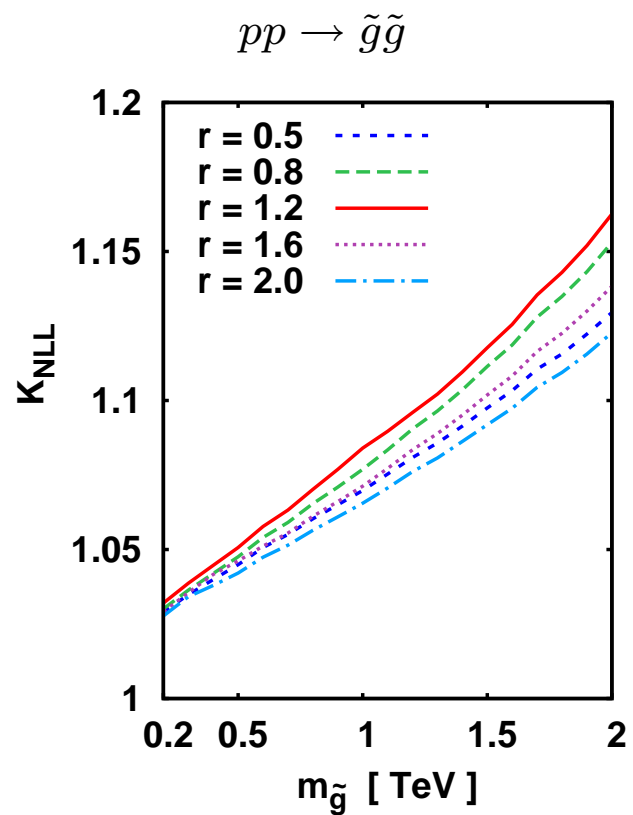
[*Beenakker, Hoepker, Krämer, Plehn, Spira, Zerwas*]

[*Plehn, <http://www.ph.ed.ac.uk/~tplehn/prospino/>*]

The NLL K-factors at the LHC

[AK, Motyka'08]

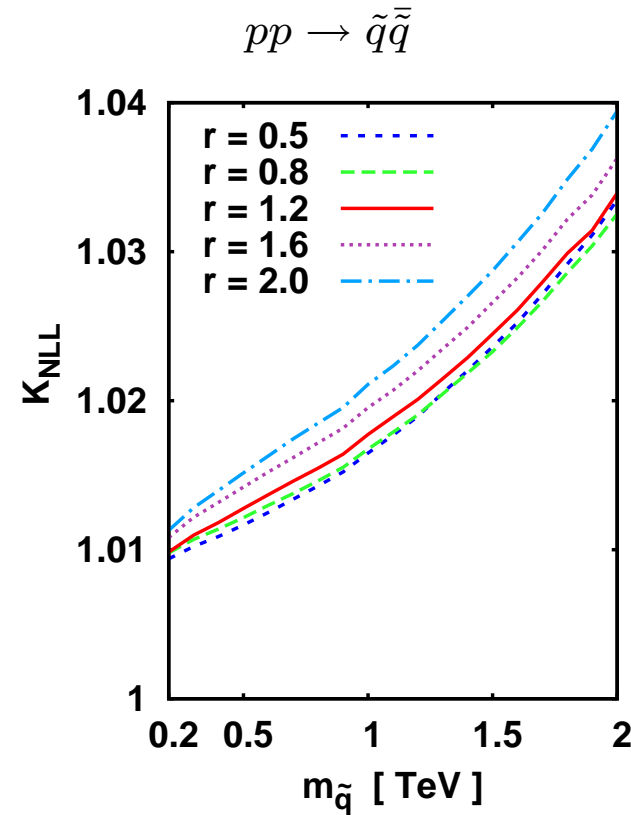
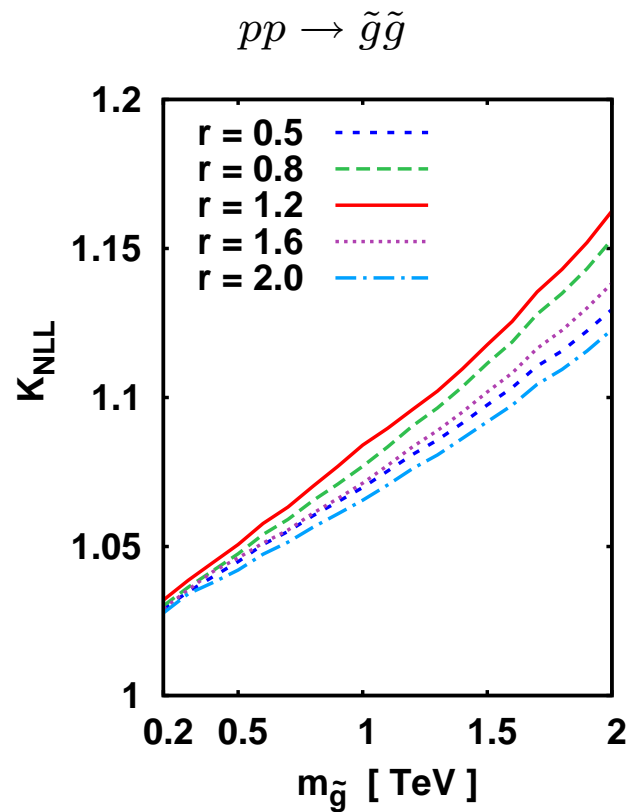
$$K^{\text{NLL}} = \frac{\sigma^{\text{match}}}{\sigma^{\text{NLO}}}$$



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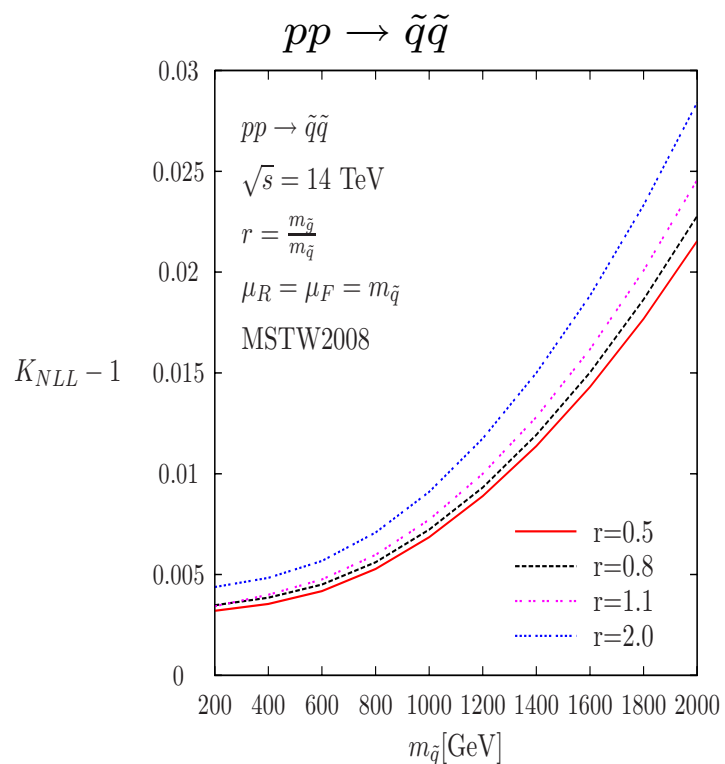


$$r = \frac{m_{\tilde{g}}}{m_{\tilde{q}}}, \mu_F = \mu_R = m, \text{CTEQ6M pdfs}$$

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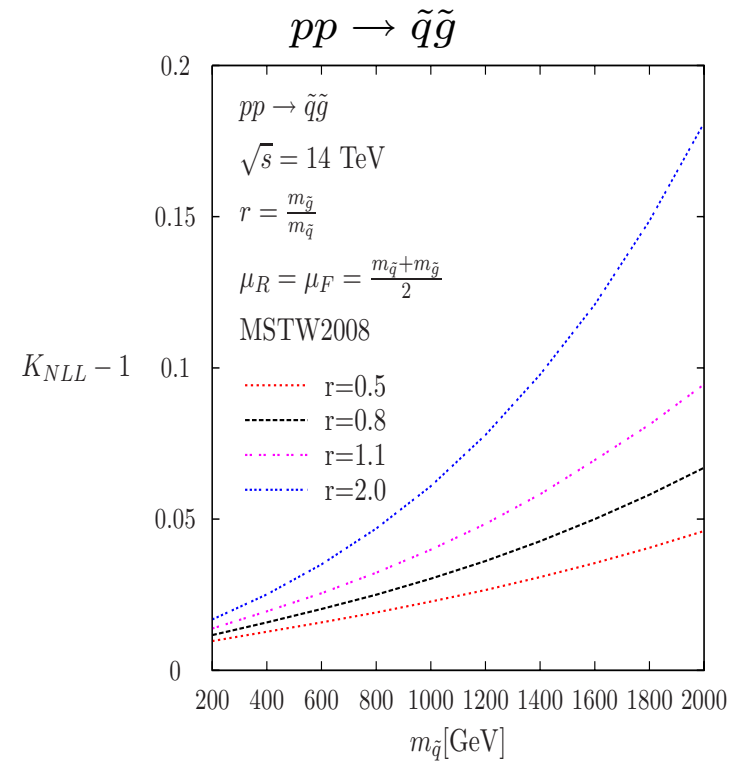
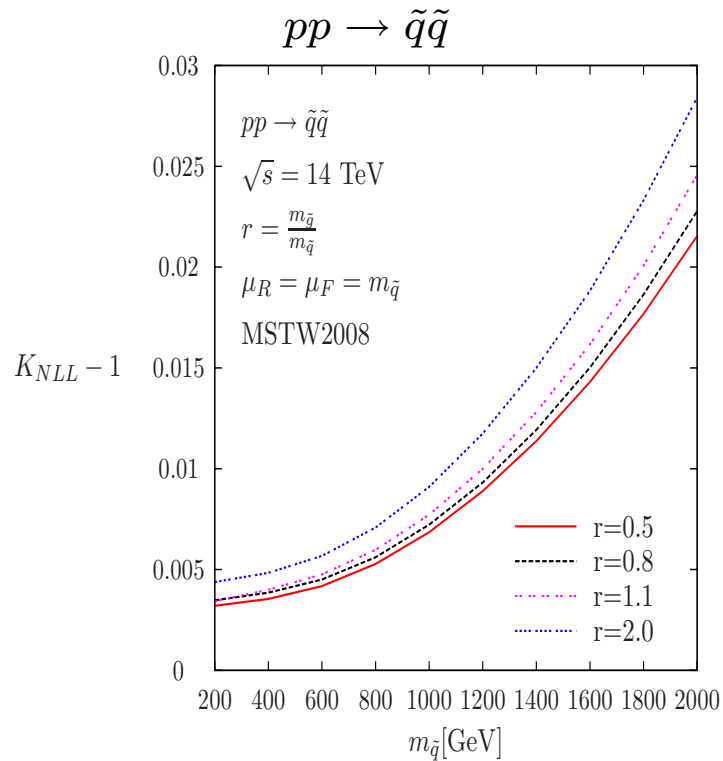


[PRELIMINARY]

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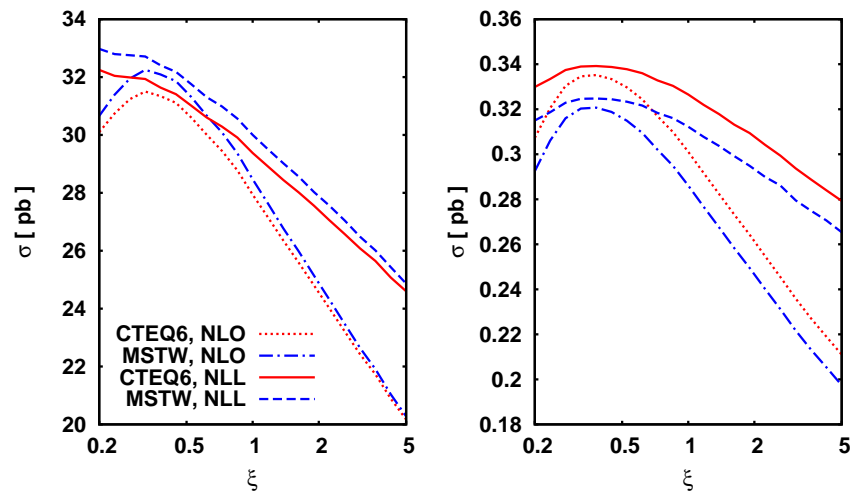
[PRELIMINARY]

The scale dependence

[AK, Motyka'09]

$$\xi = \mu/m, \mu = \mu_F = \mu_R, r = \frac{m_{\tilde{g}}}{m_{\tilde{q}}} = 1.2$$

$pp \rightarrow \tilde{g}\tilde{g}$



$m_{\tilde{g}} = 500$ GeV

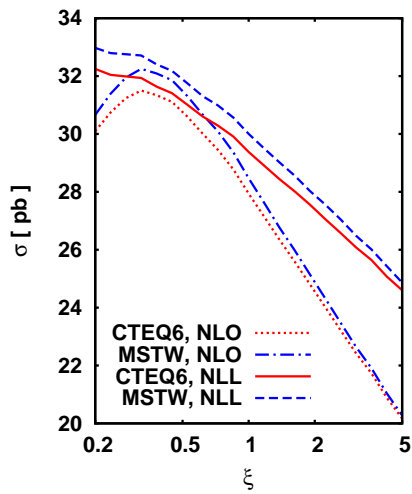
$m_{\tilde{g}} = 1$ TeV

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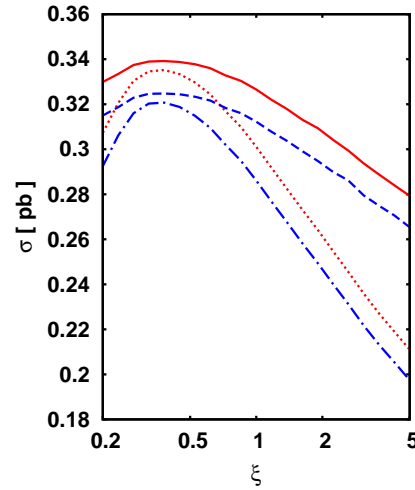
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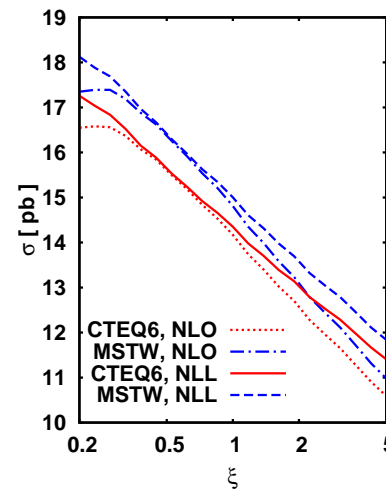


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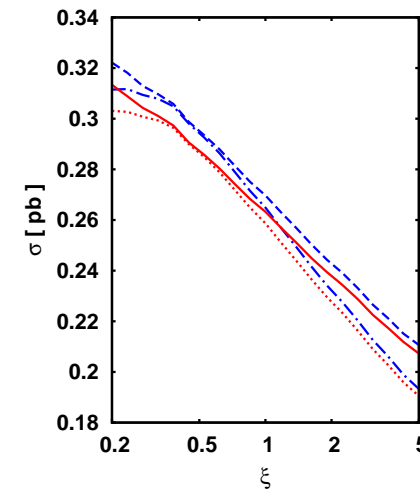


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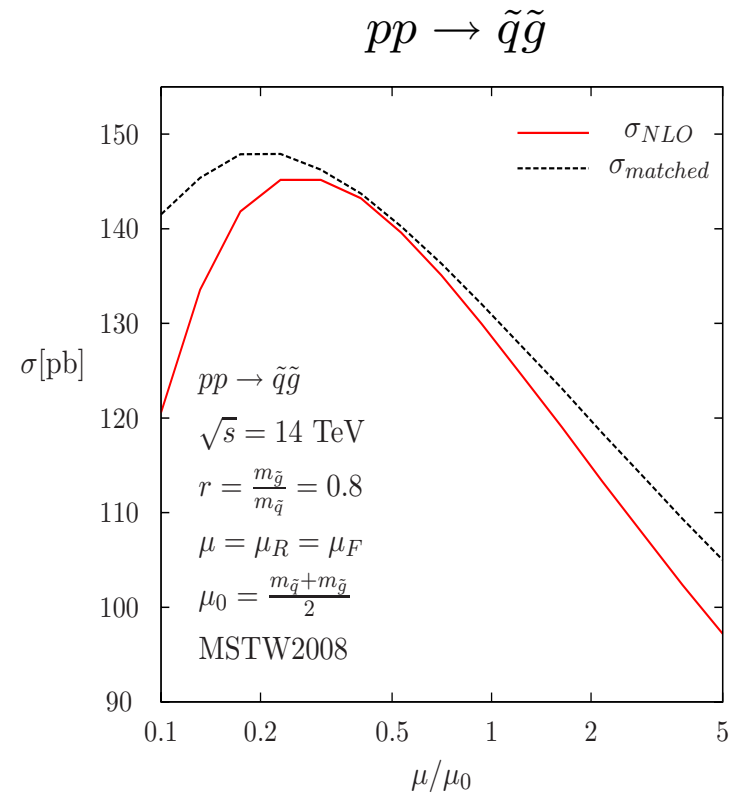
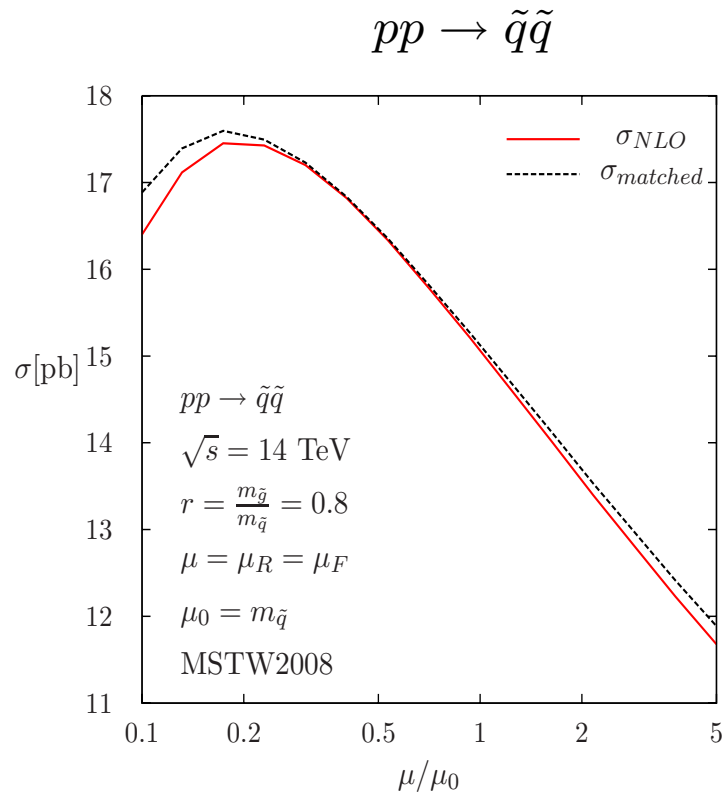


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The scale dependence cntd.

[Beenakker, Brensing, Krämer, AK, Laenen, Niessen]

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[PRELIMINARY]

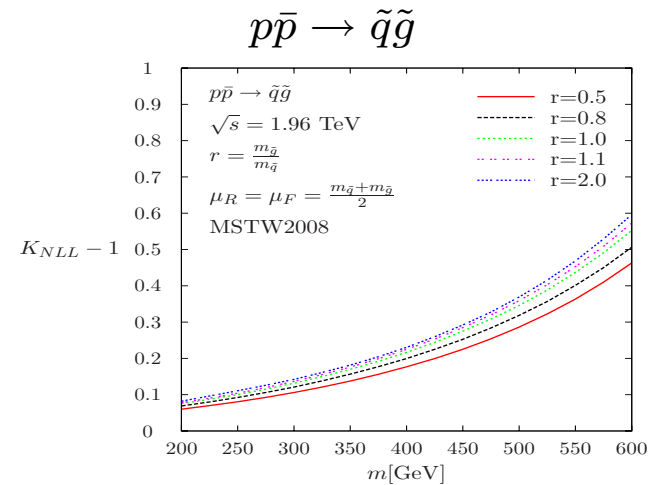
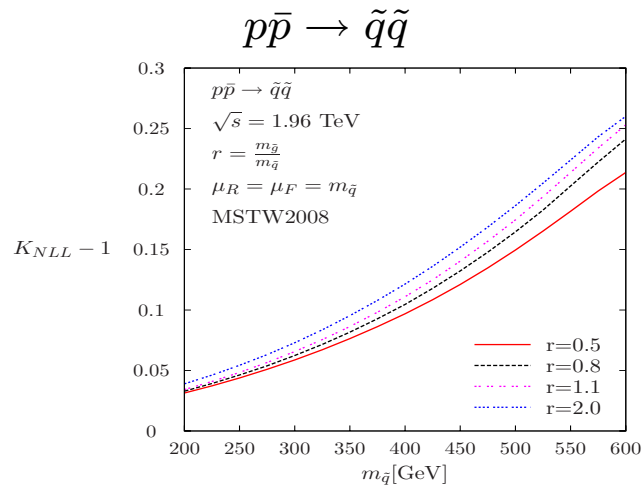
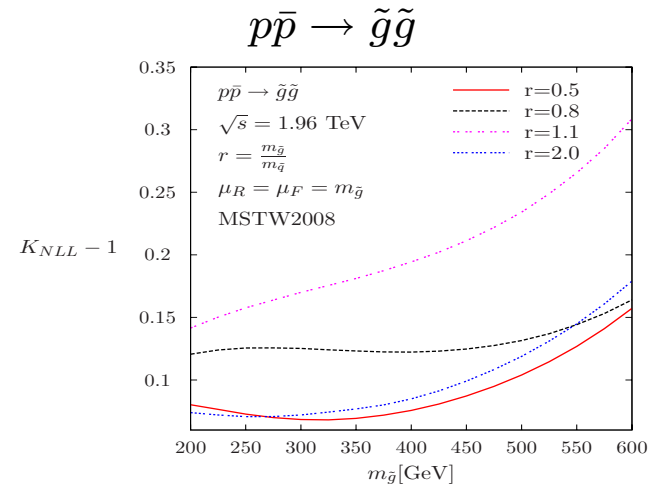
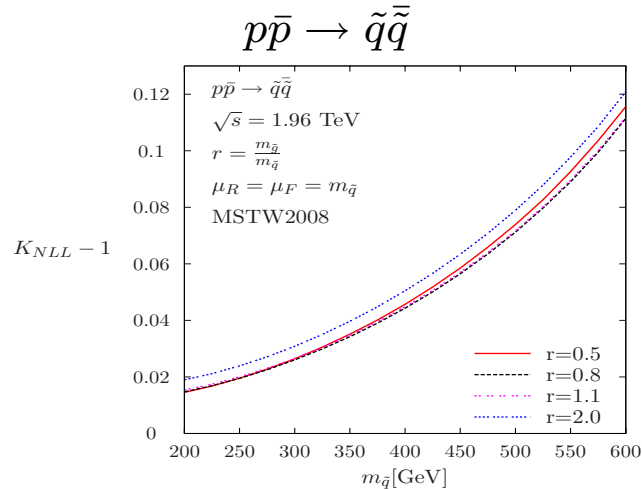
The NLL K-factors at the Tevatron

[Beenakker, Brensing, Krämer, AK, Laenen, Niessen]

Results for resummed total cross sections available for all four production processes

$$K_{\text{NLL}} - 1 = \frac{\sigma^{\text{match}} - \sigma^{\text{NLO}}}{\sigma^{\text{NLO}}}$$

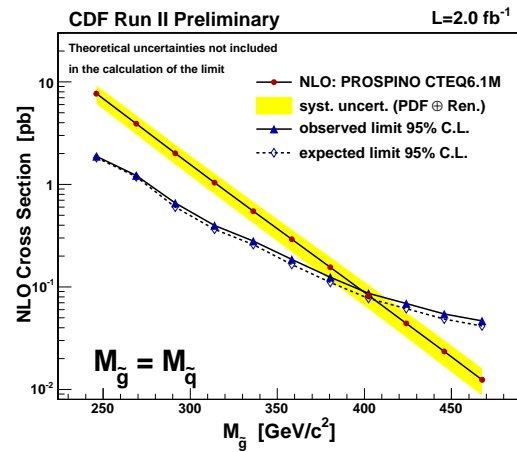
[PRELIMINARY]



Limits on squark and gluino masses from the Tevatron

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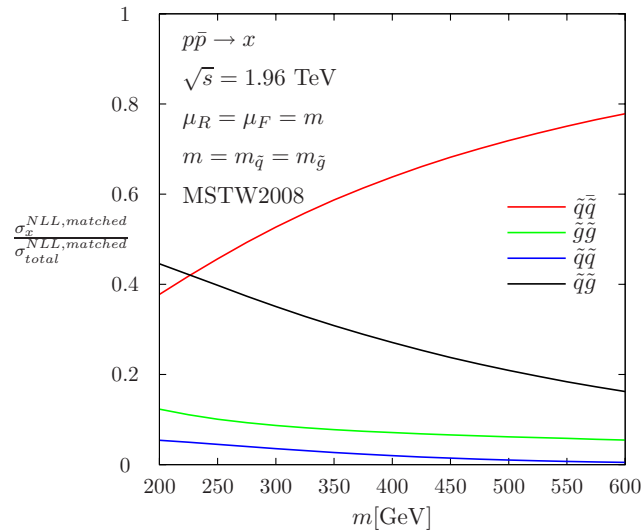
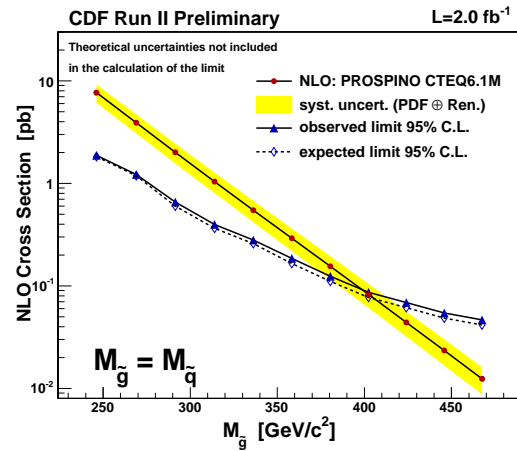
Impact on limits on squark and gluino masses under investigation



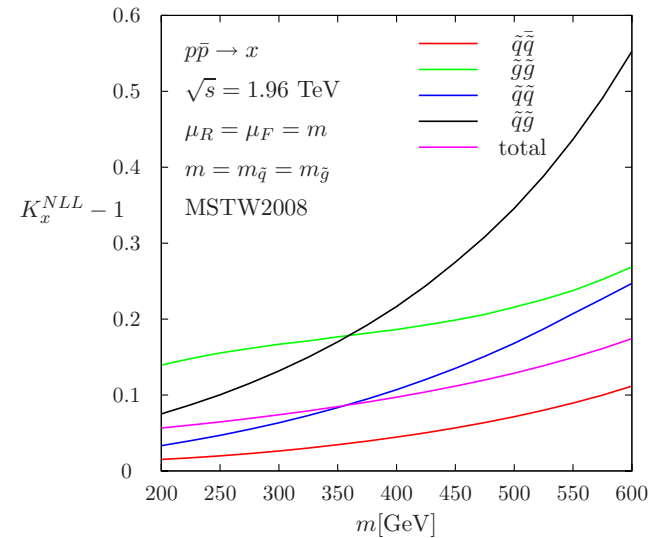
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- NLL resummed (matched with NLO) predictions for total cross sections: corrections due to soft gluon emissions above NLO

Summary

- If SUSY realized in Nature, \tilde{q} and \tilde{g} production processes will be among the most dominant mechanisms of sparticle production at the LHC
- NLL resummed (matched with NLO) predictions for total cross sections: corrections due to soft gluon emissions above NLO
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 - Corrections due to NLL resummation highest for $\tilde{g}\tilde{g}$ production
 - $\mathcal{O}(10\%)$ correction for at $1 \text{ TeV} \lesssim m_{\tilde{g}} \lesssim 2 \text{ TeV}$
 - reduction of sensitivity to scale choices
- Tevatron: resummed corrections to all squark and gluino production processes calculated