

Exploring Non-SUSY New Physics in Polarised Møller Scattering

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Outline

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Introduction

Left-right asymmetry, A_{LR} , very clean probe of new physics.

$$A_{LR} \equiv \frac{d\sigma_L - d\sigma_R}{d\sigma_L + d\sigma_R}$$

Current measurement [SLAC E158 05]:

$$A_{LR} = 131 \pm 14 \text{ (stat)} \pm 10 \text{ (syst) ppb}$$

When the 12 GeV upgrade at JLAB completed, precision 1 ppb measurement will be possible.

- Sensitive to the existence of extra Z' gauge bosons and doubly charged scalars at several TeV range.
- Provide complementary information to LHC direct measurements.

Results from PAMELA and FERMI points to possibility of a low-scale hidden sector.

- Can probe GeV range hidden sector new physics.

Effective Operators

A systematic and model-independent way to parametrise new physics at scale $\Lambda > v$:

$$\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_6$$

$$\begin{aligned} -2\mathcal{L}_6 = & \frac{c_{LL}}{\Lambda^2} (\bar{L}_a \gamma^\mu L_a) (\bar{L}_b \gamma_\mu L_b) + \frac{c_{LR}}{\Lambda^2} (\bar{L}_a \gamma^\mu L_a) (\bar{e}_R \gamma_\mu e_R) + \\ & \frac{c_{RR}}{\Lambda^2} (\bar{e}_R \gamma^\mu e_R) (\bar{e}_R \gamma_\mu e_R) + \frac{d_{LL}}{\Lambda^2} (\bar{L}_a \gamma^\mu L_b) (\bar{L}_b \gamma_\mu L_a) \\ & + h.c. \end{aligned}$$

Relevant operators contributing to A_{LR} :

$$\mathcal{O}_1 = \bar{e} \gamma^\mu \hat{L} e \bar{e} \gamma_\mu \hat{L} e, \quad \mathcal{O}_2 = \bar{e} \gamma^\mu \hat{R} e \bar{e} \gamma_\mu \hat{R} e$$

\hat{L}, \hat{R} chiral projectors.

For $\Lambda^2 \gg s \gg m_e^2$ valid for 12 GeV experiments:

$$\begin{aligned} A_{LR} &= A_{LR}^{SM} + \delta A_{LR} \\ &= \frac{4G_\mu s}{\sqrt{2}\pi\alpha} \frac{y(1-y)}{1+y^4+(1-y)^4} \\ &\quad \times \left[\left(\frac{1}{4} - \sin^2 \theta_W^{\overline{MS}} \right) + \frac{c'_{LL} - c_{RR}}{4\sqrt{2}G_\mu\Lambda^2} \right] \end{aligned}$$

$$y = Q^2/s, \quad s = 2m_e E_{beam}, \quad c'_{LL} = c_{LL} + d_{LL}.$$

Useful variable:

$$\delta_{LR} \equiv \frac{\delta A_{LR}}{A_{LR}^{SM}} = \frac{c'_{LL} - c_{RR}}{\sqrt{2}G_\mu(1 - 4\sin^2 \theta_W^{\overline{MS}})\Lambda^2}$$

New Physics Scenarios

High scale new physics, couples directly to SM:

- GUT models
- Left-right symmetric models
- Bottom-up $U(1)_X$ models
- Models with doubly charged scalars.

New physics can be low scale, couples to SM only through mixings:

- Hidden sector/Shadow Z' models

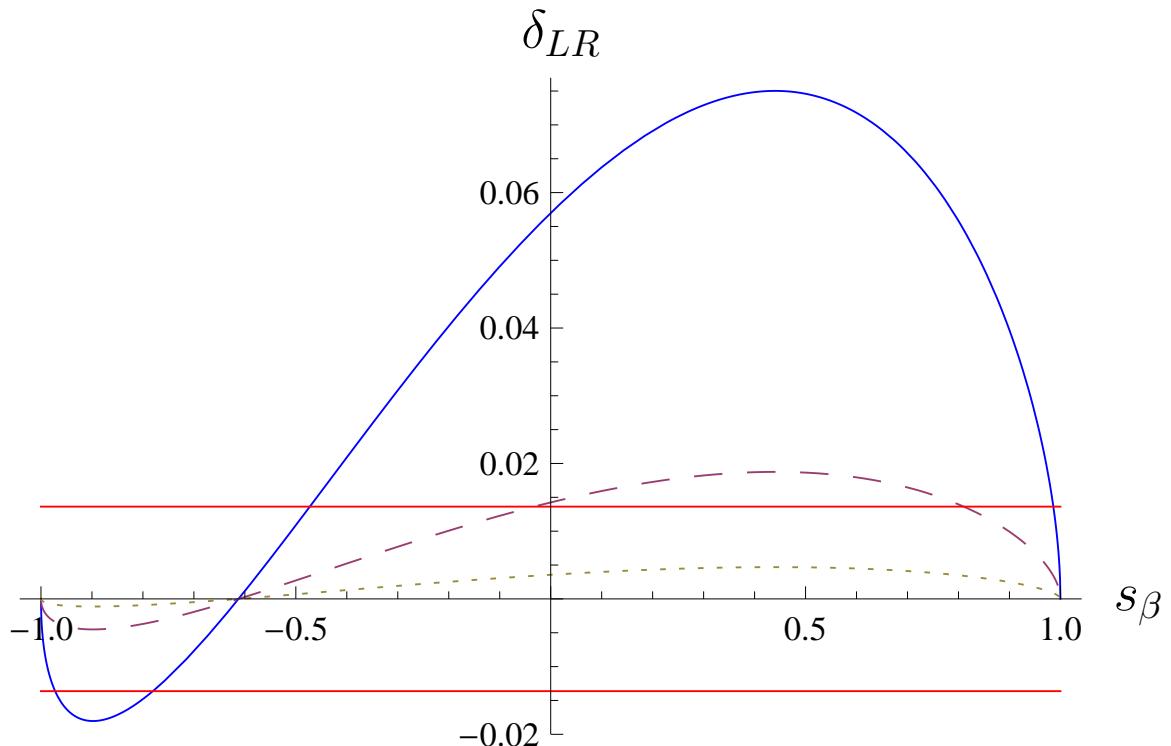
GUT Models

$$E_6 \rightarrow SO(10) \times U(1)_\psi \rightarrow SU(5) \times U(1)_\chi \times U(1)_\psi \rightarrow SM \times U(1)_\beta$$

$$X^\mu(\beta) = B_\psi^\mu \sin \beta + B_\chi^\mu \cos \beta, \quad c'_{LL} - c_{RR} = \frac{4\pi\alpha}{3c_W^2} c_\beta \left(c_\beta + \sqrt{\frac{5}{3}} s_\beta \right)$$

E158 95% CL limits: $-0.337 < \delta_{LR} < 0.122$.

JLAB 95% CL limits: $|\delta_{LR}| < 0.0136$ (projected assuming 1 ppb total 1σ error).



Solid, dashed and dotted lines denote $\Lambda = 1, 2, 4$ TeV.

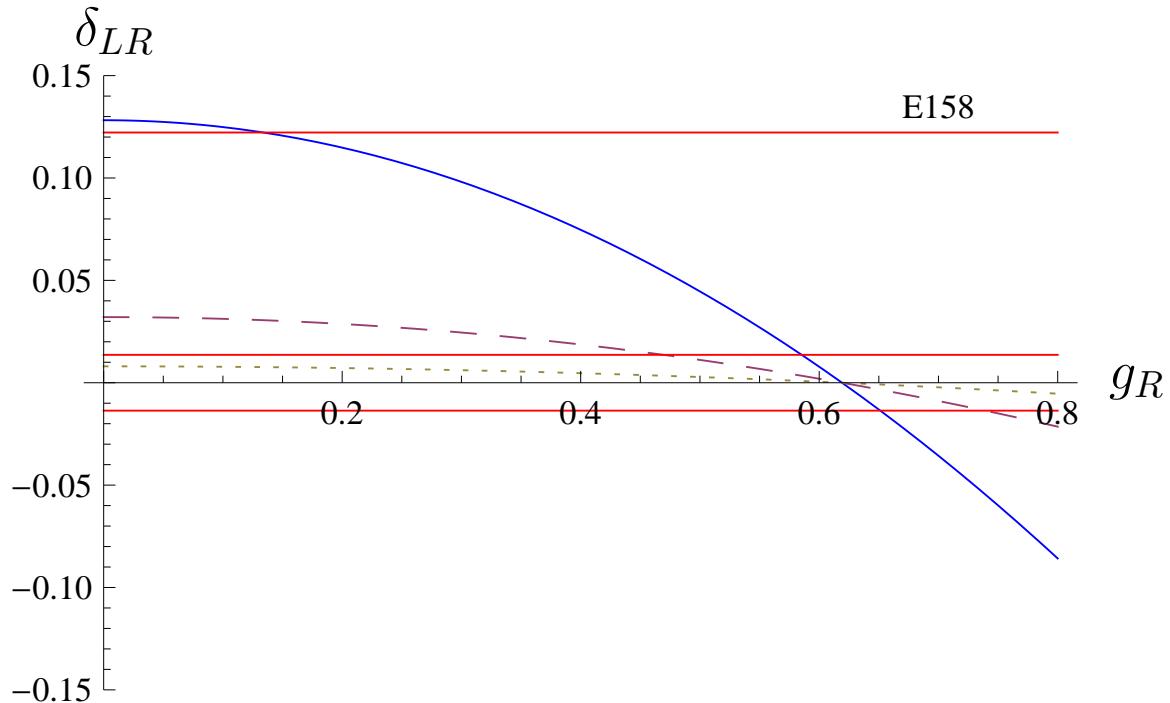
Horizontal lines projected JLAB upper and lower limits.

LR Symmetric Models

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L} \rightarrow SU(2)_L \times U(1)_R \times U(1)_{B-L}$$

$$c'_{LL} - c_{RR} = \frac{1}{4} (3g_Y^2 - g_R^2)$$

Note $\delta_{LR} \neq 0$ even if $g_R = g_L$.



Solid, dashed and dotted lines denote $\Lambda = 1, 2, 4$ TeV.

Narrow band projected JLAB limits.

Bottom-up $U(1)_X$ Models

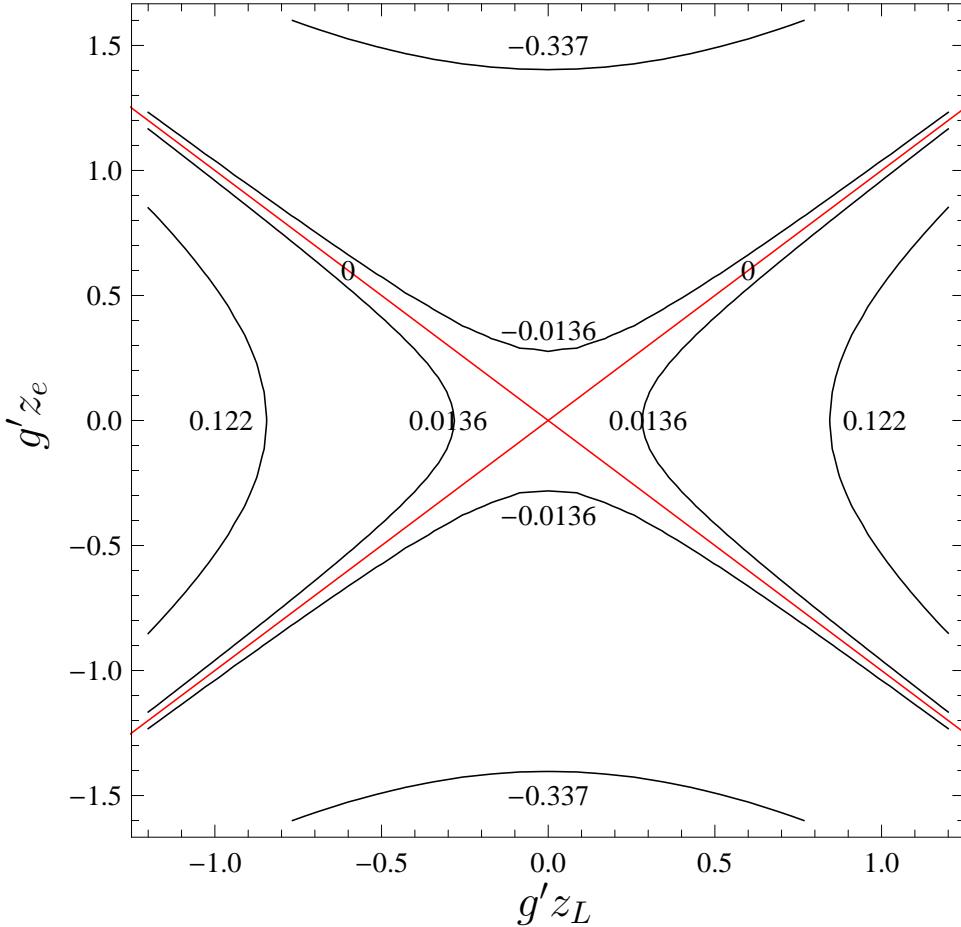
Extra $U(1)$ factor unrelated to larger broken down symmetry.

SM charges under $U(1)_X$ determined by **anomaly cancellation**:

$$\begin{aligned} z_Q &= -\frac{1}{3}z_L, & z_u &= \frac{2}{3}z_L - z_e, & z_d &= -\frac{4}{3}z_L + z_e \\ z_N &= 2z_L - z_e, & z_H &= z_L - z_e \end{aligned}$$

General case: $z_L \neq z_e$, $c'_{LL} - c_{RR} = g'^2(z_L^2 - z_e^2)$

$U(1)_X = U(1)_{B-L}$: $z_L = z_e \Rightarrow \delta A_{LR} = 0$



Contours of constant δ_{LR} for $\Lambda = 1$ TeV (other values of Λ obtained simply by a rescaling factor Λ/TeV).

Doubly Charged Scalars

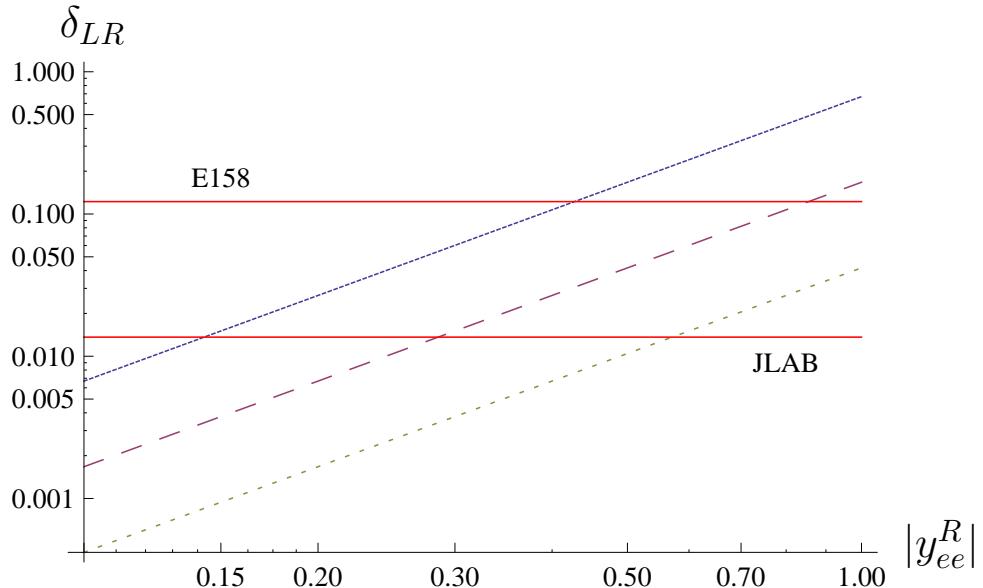
$P^{\pm\pm}$ motivated by Type II seesaw neutrino mass generation mechanisms.

$P^{\pm\pm}$ $SU(2)_L$ singlet:

$$\frac{|y_{ee}^R|^2}{4M_P^2}(\bar{e}\gamma^\mu \hat{R}e)(\bar{e}\gamma_\mu \hat{R}e) \Rightarrow \delta_{LR} = \frac{|y_{ee}^R|^2}{2\sqrt{2}(1-4s_W^2)G_\mu M_P^2}$$

$P^{\pm\pm}$ $SU(2)_L$ doblet:

- Like singlet case but with LL-type Wilson coefficients.
- $\delta_{LR} < 0!$



Solid, dashed and dotted lines denote $M_P = 1, 2, 4$ TeV.

Hidden Sector/Shadow Z' Models

Extra $U(1)$ hidden sector gauge symmetry, which communicate with the SM only through scalar mixing and gauge kinetic mixing [**Kumar and Wells 06, CNW 06**].

Gauge kinetic mixing:

$$\frac{\epsilon}{2} X_{\mu\nu} B^{\mu\nu} \Rightarrow \begin{pmatrix} X \\ B \end{pmatrix} = \begin{pmatrix} c_\epsilon & 0 \\ -s_\epsilon & 1 \end{pmatrix} \begin{pmatrix} X' \\ B' \end{pmatrix}, \quad s_\epsilon = \frac{\epsilon}{\sqrt{1-\epsilon^2}}$$

$$\begin{pmatrix} B' \\ W_3 \\ X' \end{pmatrix} = \begin{pmatrix} c_W & -s_W & 0 \\ s_W & c_W & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_\eta & -s_\eta \\ 0 & s_\eta & c_\eta \end{pmatrix} \begin{pmatrix} \gamma \\ Z \\ Z' \end{pmatrix}$$

$$\tan 2\eta = \frac{2s_W s_\epsilon}{(c_W M_X/M_W)^2 + s_W^2 s_\epsilon^2 - 1}$$

Gauge-fermion couplings:

$$Z^\mu \bar{f} f : \quad i\gamma^\mu \frac{g_L}{c_W} \left\{ \left[c_\eta (T_{f_L}^3 - s_W^2 Q_{f_L}) - \frac{1}{2} s_\eta s_W s_\epsilon Y_{f_L} \right] \hat{L} + (L \leftrightarrow R) \right\}$$

$$Z'^\mu \bar{f} f : \quad -i\gamma^\mu \frac{g_L}{c_W} \left\{ \left[s_\eta (T_{f_L}^3 - s_W^2 Q_{f_L}) + \frac{1}{2} c_\eta s_W s_\epsilon Y_{f_L} \right] \hat{L} + (L \leftrightarrow R) \right\}$$

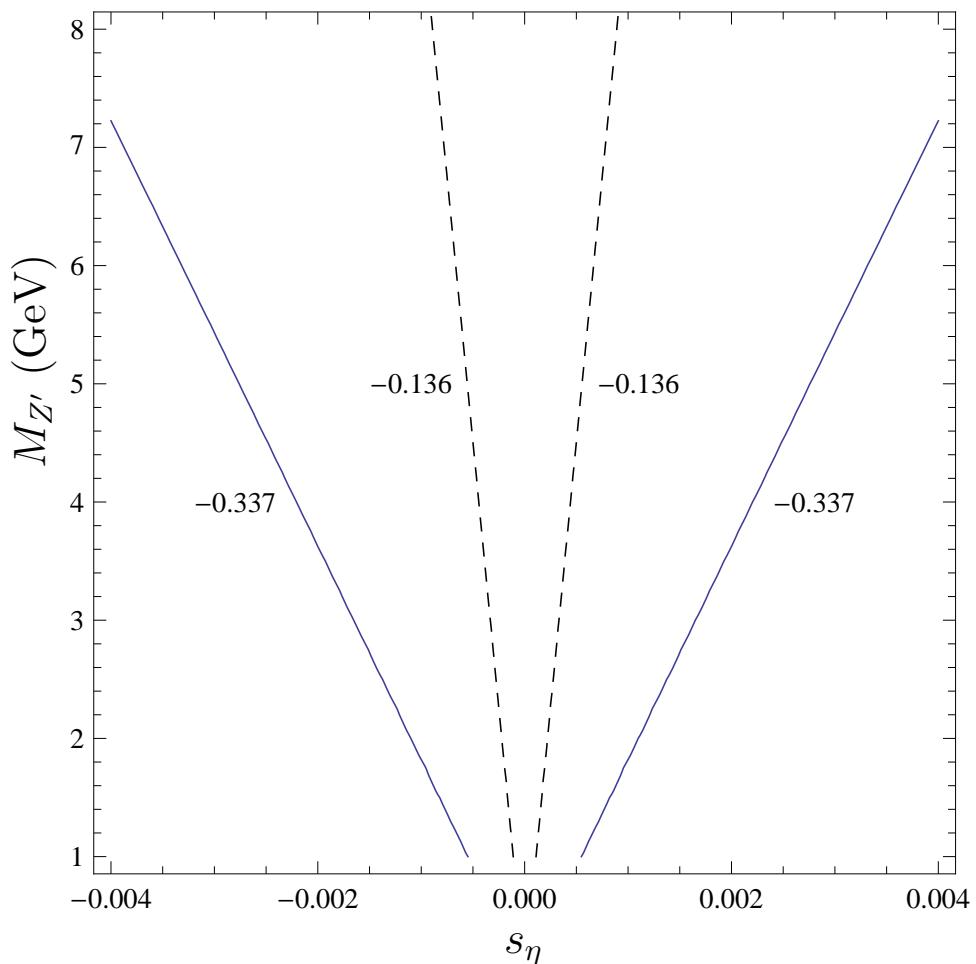
Low scale limit ($M_{Z'} \sim M_X \ll M_W$, $\eta \sim \epsilon$):

$$\delta_{LR} = \left(s_\eta^2 - 15.80 c_\eta^2 s_\epsilon^2 + 31.85 c_\eta s_\eta s_\epsilon \right) \frac{M_Z^2}{M_{Z'}^2}$$

Dominated by γ - Z' interference.

Light Shadow Z' ($M_{Z'} \ll M_W$)

$$s_\epsilon = -\frac{s_\eta}{s_W c_\eta} (1 - c_W^2 r^2) , \quad r = \frac{M_{Z'}}{M_W}$$



Contours of constant δ_{LR} . Solid and dashed contours denote E158 and projected JLAB lower limits.

NO PARAMETER SPACE EXISTS FOR $\delta_{LR} > 0!$.

Summary

Complementary probe of new physics to LHC at several TeV order

Given the mass of new states measured at the LHC:

- Clean determination of Z' gauge couplings to electrons and mixing parameters.
- Can shed light on the origin of the extra Z' .
- Clean determination of Yukawa couplings of doubly charged scalars to electrons.
- Sign of δ_{LR} can determine type of doubly charged scalars.

Probe low scale hidden sector

Provide valuable inputs to dark matter model building employing a hidden/shadow sector:

- Constrain mixing between SM Z and shadow Z' .
- Best measure of shadow Z' couplings to electrons.