

# Background estimation in the search for $H^+ \rightarrow \tau\nu$ in the ATLAS experiment

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# World of subatomic particles – the Standard Model

	mass → $\approx 2.3 \text{ MeV}/c^2$ charge → $2/3$ spin → $1/2$ <b>u</b> up	mass → $\approx 1.275 \text{ GeV}/c^2$ charge → $2/3$ spin → $1/2$ <b>c</b> charm	mass → $\approx 173.07 \text{ GeV}/c^2$ charge → $2/3$ spin → $1/2$ <b>t</b> top	mass → 0 charge → 0 spin → 1 <b>g</b> gluon	mass → $\approx 126 \text{ GeV}/c^2$ charge → 0 spin → 0 <b>H</b> Higgs boson
<b>QUARKS</b>	mass → $\approx 4.8 \text{ MeV}/c^2$ charge → $-1/3$ spin → $1/2$ <b>d</b> down	mass → $\approx 95 \text{ MeV}/c^2$ charge → $-1/3$ spin → $1/2$ <b>s</b> strange	mass → $\approx 4.18 \text{ GeV}/c^2$ charge → $-1/3$ spin → $1/2$ <b>b</b> bottom	mass → 0 charge → 0 spin → 1 <b><math>\gamma</math></b> photon	
	mass → $0.511 \text{ MeV}/c^2$ charge → -1 spin → $1/2$ <b>e</b> electron	mass → $105.7 \text{ MeV}/c^2$ charge → -1 spin → $1/2$ <b><math>\mu</math></b> muon	mass → $1.777 \text{ GeV}/c^2$ charge → -1 spin → $1/2$ <b><math>\tau</math></b> tau	mass → $91.2 \text{ GeV}/c^2$ charge → 0 spin → 1 <b>Z</b> Z boson	
<b>LEPTONS</b>	mass → $< 2.2 \text{ eV}/c^2$ charge → 0 spin → $1/2$ <b><math>\nu_e</math></b> electron neutrino	mass → $< 0.17 \text{ MeV}/c^2$ charge → 0 spin → $1/2$ <b><math>\nu_\mu</math></b> muon neutrino	mass → $< 15.5 \text{ MeV}/c^2$ charge → 0 spin → $1/2$ <b><math>\nu_\tau</math></b> tau neutrino	mass → $80.4 \text{ GeV}/c^2$ charge → $\pm 1$ spin → 1 <b>W</b> W boson	<b>GAUGE BOSONS</b>

# Beyond the SM –

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## The search for $H^+ \rightarrow \tau\nu$

Valid problems of the SM:

- ▶ does not include gravity,
- ▶ dark matter and dark energy,
- ▶ neutrino masses etc.

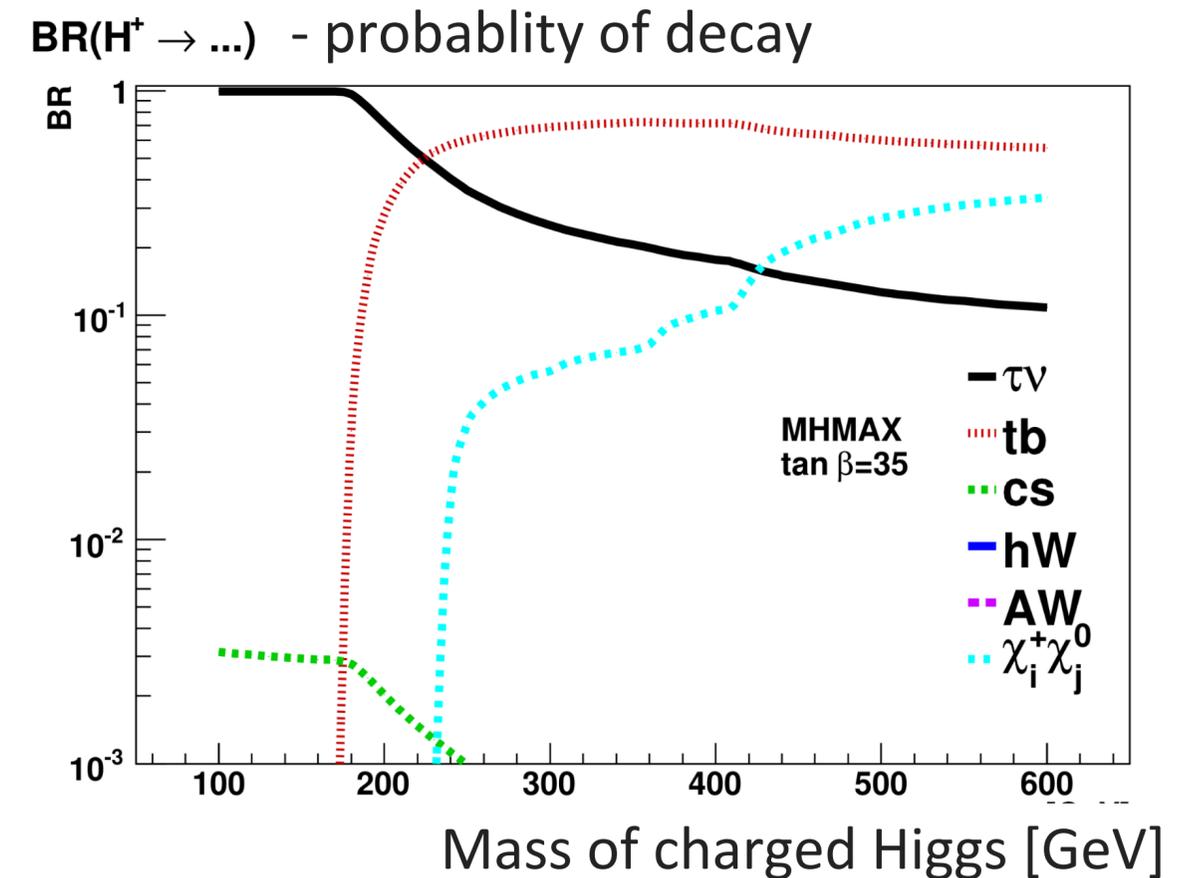
# Beyond the SM – The search for $H^+ \rightarrow \tau\nu$

New fundamental theory needed



Minimal Supersymmetric Standard Model (MSSM) predicting additional particles like charged Higgs boson

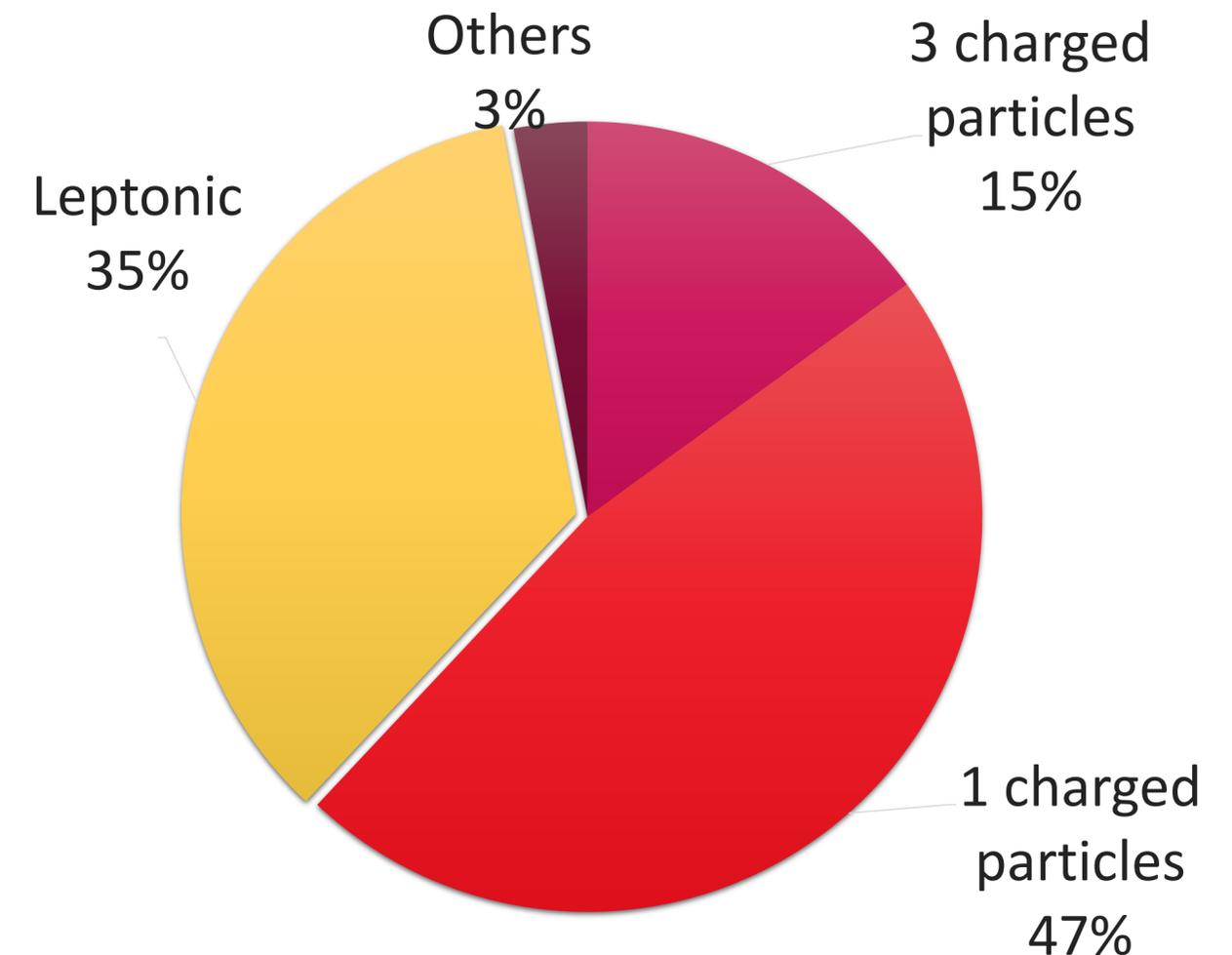
- ▶ To prove (or disprove) a MSSM theory we're looking for  $H^+ \rightarrow \tau\nu$  process in the ATLAS data collected in 2015-2018



# Tau lepton in ATLAS analyses

- Tau lepton is a short-lived particle, hence we can observe only its decay products
- It is the heaviest from leptons, so it can decay into mostly **hadrons** (particles made of quarks) or **leptons**
- **Hadronically decaying taus are difficult to identify in data**

## Tau decay



# Background estimation - introduction

In particle physics analysis:

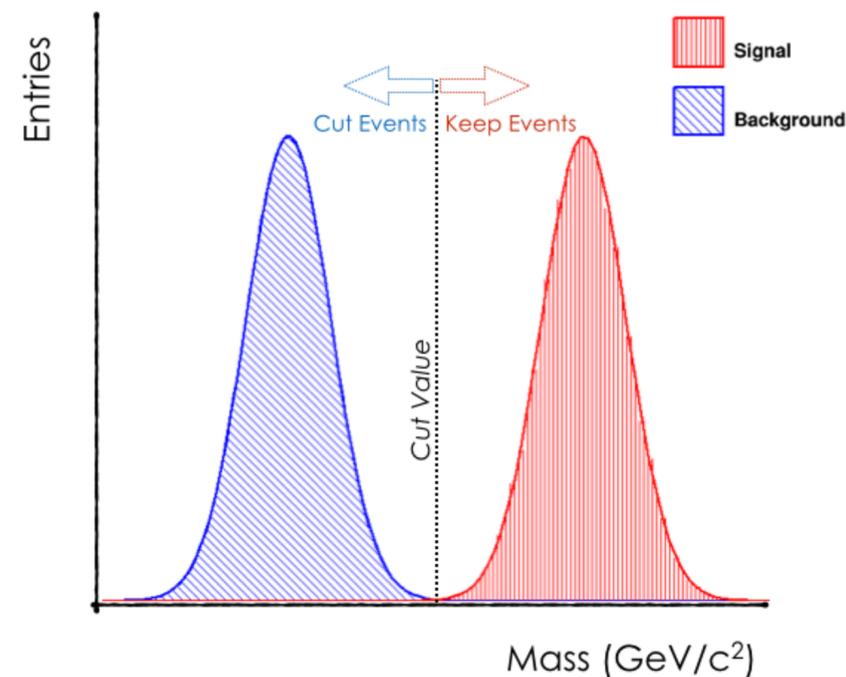
➤ **”signal”** - whatever particular process we are interested in

➤ **„background”** - processes which might look a little bit (or a lot) like the signal process we care about, but are not

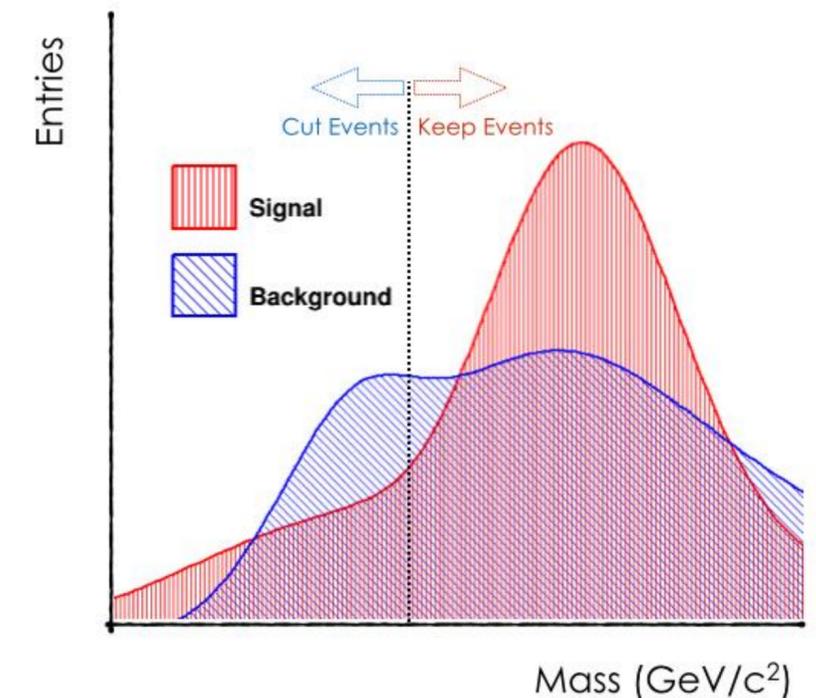


**Everything else**

**We want to separate them by using cuts**  
– set of restrictions on the particle parameters that would give us the biggest signal part possible (and lowest background)



Perfect world

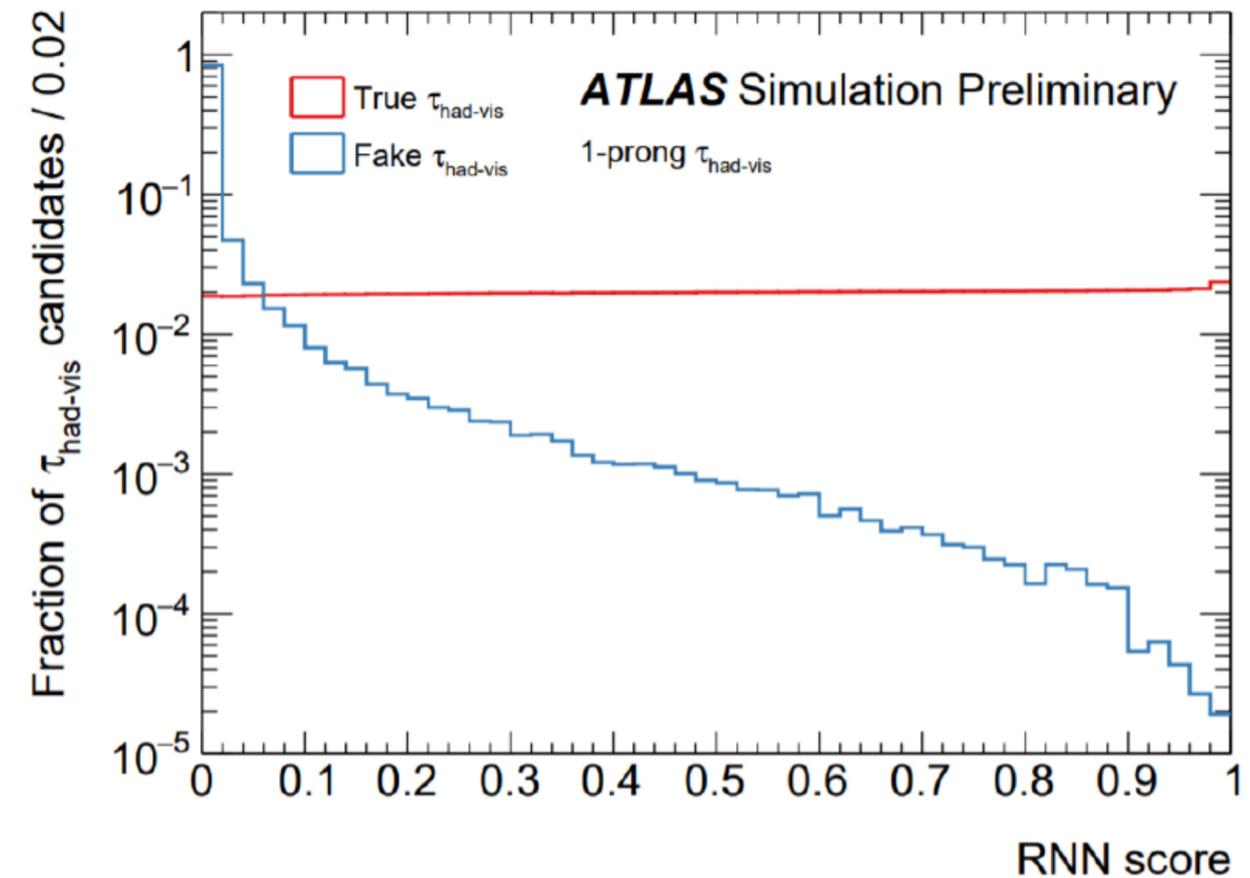


Reality

# Background estimation

To estimate  $H^+ \rightarrow \tau\nu$  background:

- I. **Monte Carlo generators** are used to predict which SM physical processes we can observe during collision ( $t\bar{t}$ ,  $Wt$ ,  $V$ +jets events)
- II. **Machine learning** (Recurrent Neural Networks) for proper tau identification



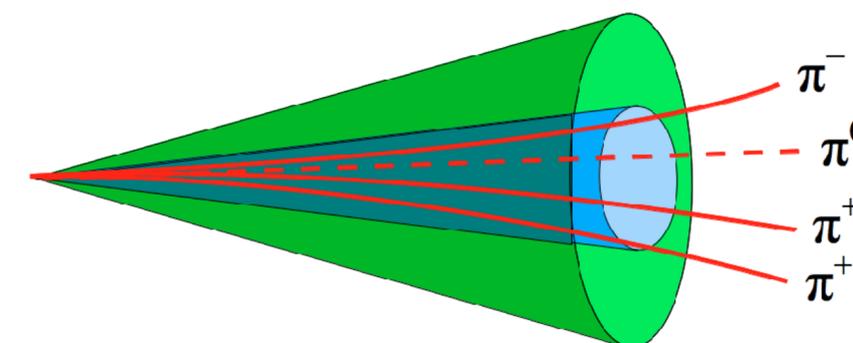
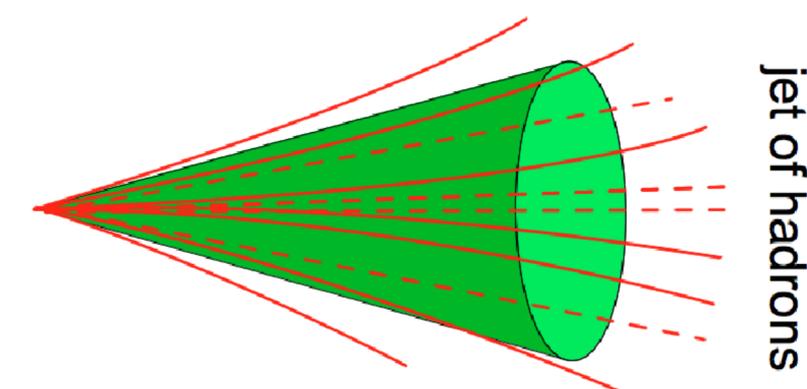
# Background estimation

To estimate  $H^+ \rightarrow \tau\nu$  background:

- I. Monte Carlo simulation which
- II. Monte Carlo simulation for proper tau identification

**However, for analysis with tau leptons in final states it is not enough**

**Jet - narrow cone of hadrons and other particles produced during collisions of gluons and quarks**



- Main source of background comes from the jets misidentified as tau, these are poorly modelled by the Monte Carlo method

# Fake Factor Method

- Data-driven method of the fake-tau background determination
- The fake- $\tau$  background in SR with the nominal object selection ( $\tau$  ID) is estimated from events with a reversed selection (**anti- $\tau$  ID**), using **fake factors (FF)** for certain control region (CR) defined as:

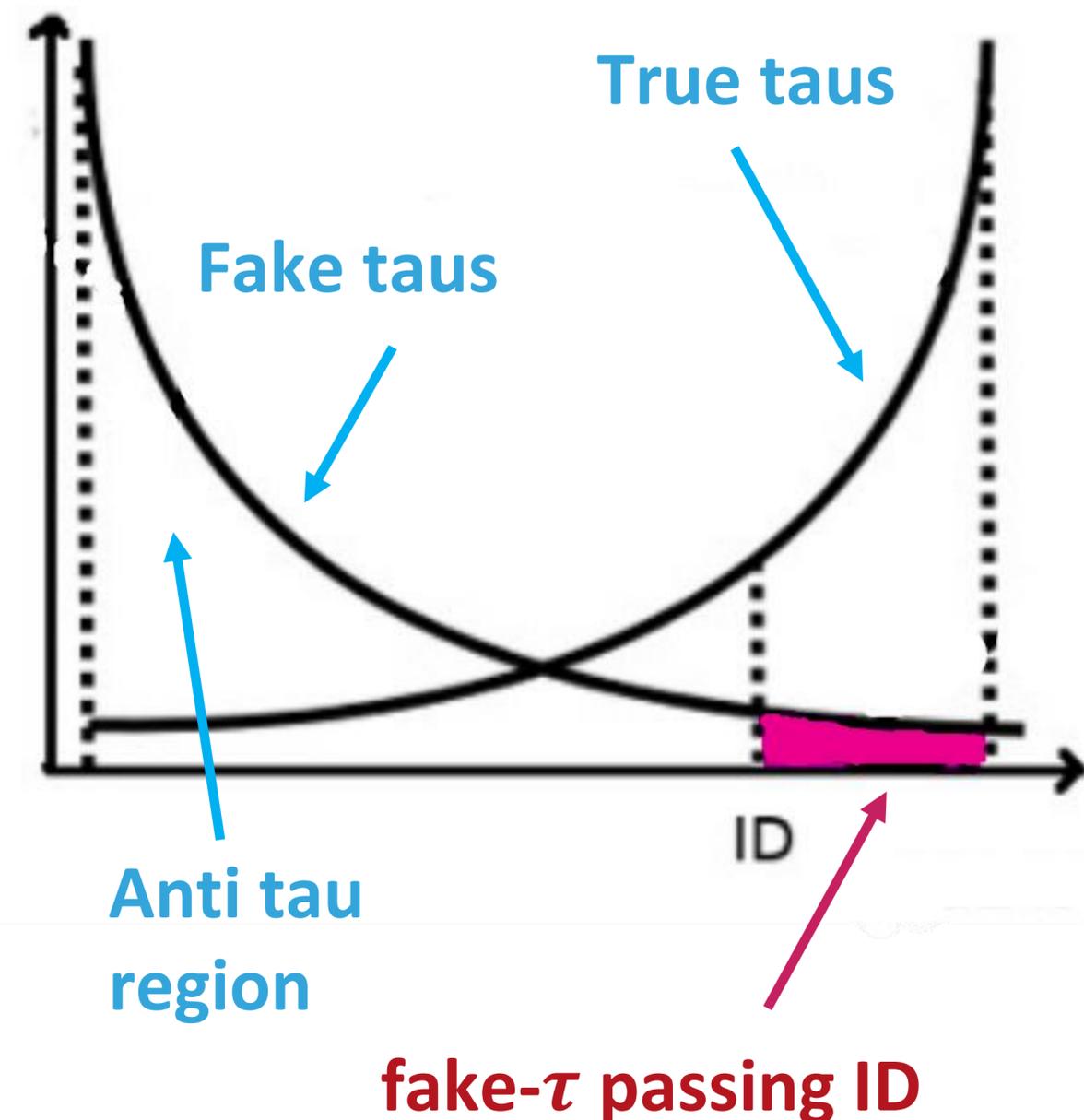
$$FF = \frac{N_{\tau-id}^{CR}}{N_{anti\tau-id}^{CR}}$$

Events that  
pass  $\tau$   
identification

Events that fail  
the nominal  $\tau$   
identification

- Total number of background events from jets ( $N_{fakes}^{\tau}$ ) in signal region:

$$N_{fakes}^{\tau} = N_{anti-\tau}^{SR} \times FF$$



# Fake Factor Method in $H^{\pm} \rightarrow \tau\nu$ analysis

- **Complication:** some true taus sneak into the CR and the “SR but fail ID” regions
- **Solution - subtracting them with the use of MC:**

$$FF = \frac{N_{\tau}^{\text{CR}}(\text{data}) - N_{\tau}^{\text{CR}}(\text{MC, true} - \tau)}{N_{\text{anti-}\tau}^{\text{CR}}(\text{data}) - N_{\text{anti-}\tau}^{\text{CR}}(\text{MC, true} - \tau)}$$

- Parametrization of the FFs :  $p_T^{\tau}$  and number of charged tracks (1-track or 3-track  $\tau$ )
- In the  $H^{\pm} \rightarrow \tau\nu$  analysis FFs are extracted in two CRs:
  - I. **Multi-jets CR**  $\rightarrow$  gluon-initiated jets
  - II. **W+jets CR**  $\rightarrow$  quark-initiated jet

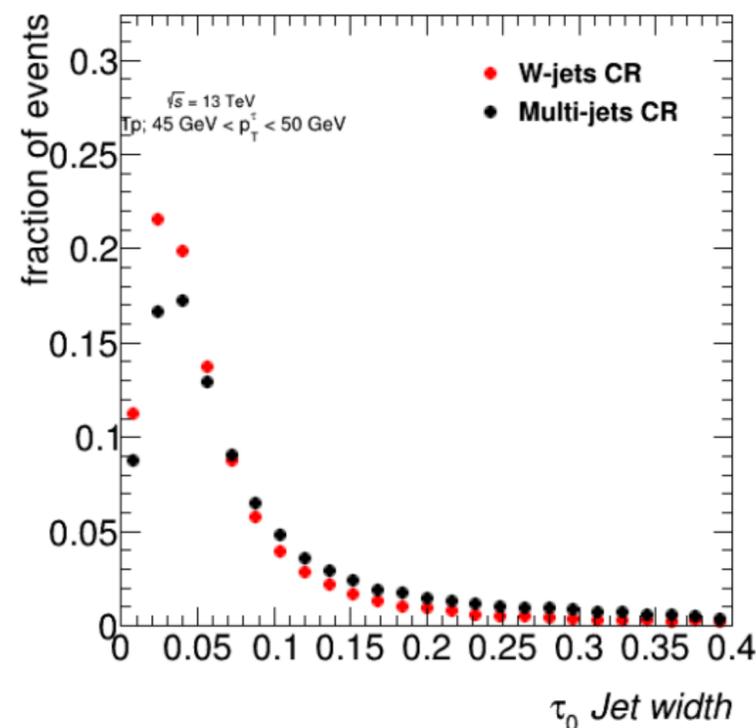
FF determination Data	FF application Data
CR Pass ID	SR Pass ID
CR Fail ID	SR but Fail ID

# Quark/gluon composition determination

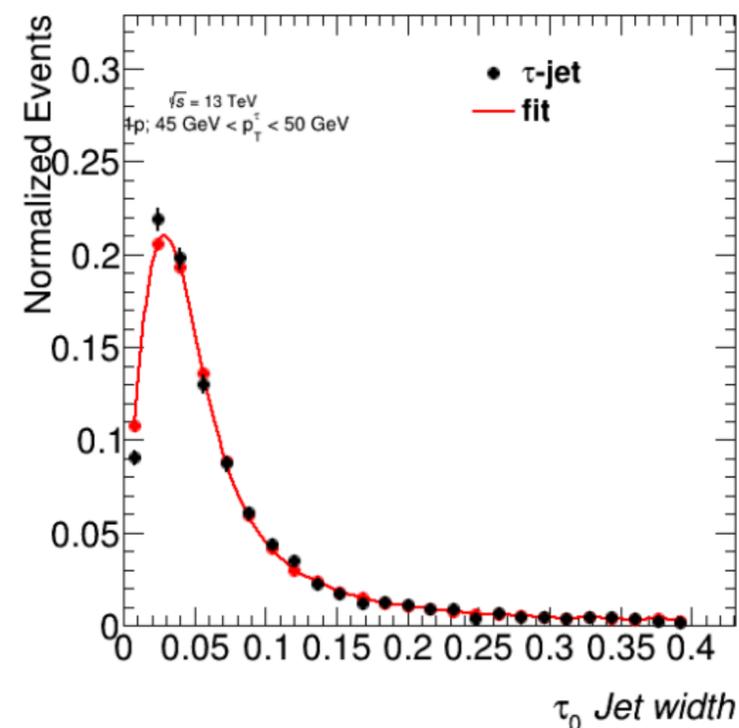
- To obtain proper jets composition in a signal region, FFs from two CRs are combined using the formula:

$$FF_{SR}^{\text{comb}}(i) = \alpha_{MJ}(i) \times FF_{MJ}(i) + (1 - \alpha_{MJ}(i)) \times FF_{W+\text{jets}}(i)$$

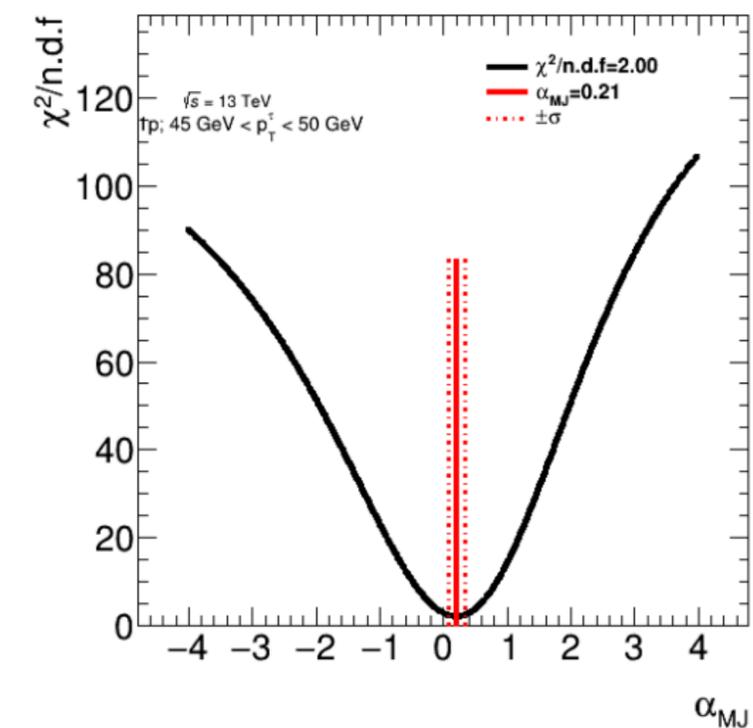
- $\alpha_{MJ}$  parametr is calculated using a template-fit method:



Creates 2 templates for CRs  
(distributions for  
discriminant variable)

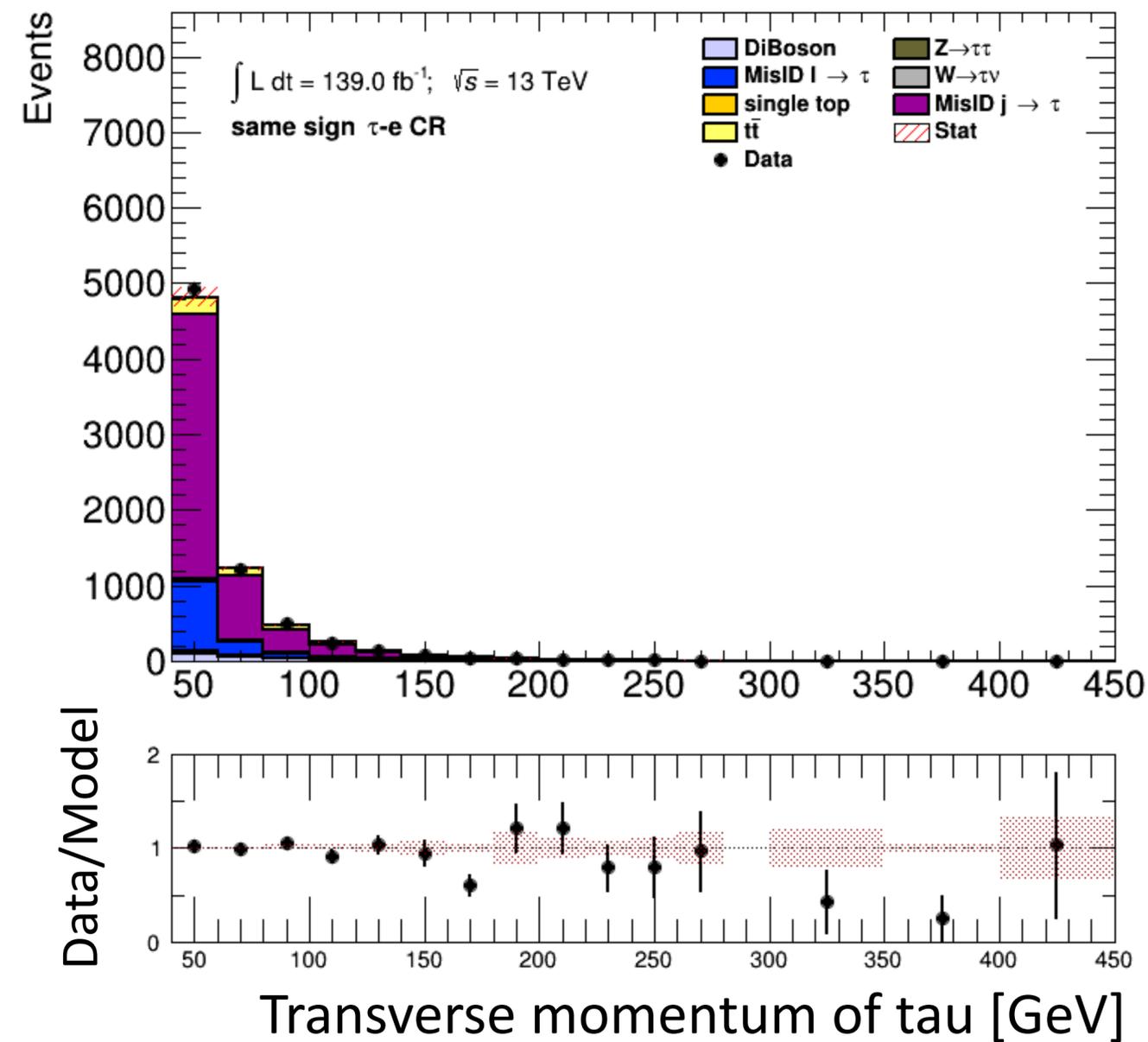
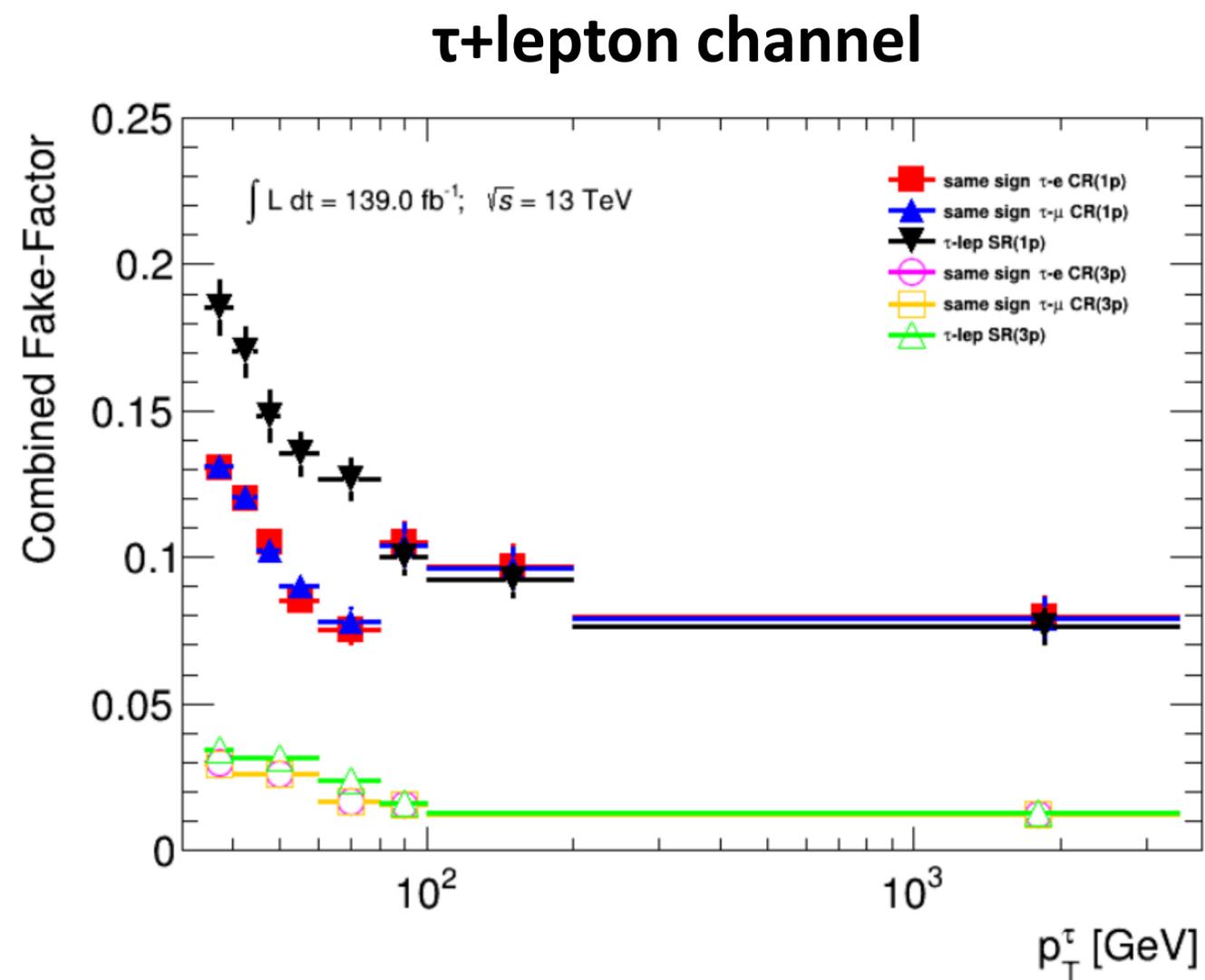


Then fits it to  
the SR  
template



$\chi^2$  statistics  
distribution

# Validation of the FF method



# Full systematics of Fake Factor method

- I. the statistical uncertainties in the event yields entering the computation of FFs (each bin of their parameterization and for each control region)
- II. the level of contamination of true  $\tau$  candidates fulfilling the anti- $\tau$  selection (varied by 50%)
- III. MC modeling of true tau identification - scale factors error of 5% is implemented
- IV. the statistical uncertainty of the best-fit value of  $\alpha_{MJ}$
- V. for the  $\Upsilon$  distribution (used in RNN-based  $\tau$  ID) the uncertainty of the inverse transform sampling method (Smirnov transformation for  $\Upsilon$ )
- VI. the modelling of heavy-flavour jets mimicking  $\tau$  candidates

# Systematics of Fake Factor method in $H^\pm \rightarrow \tau\nu$ analysis (2015-2016 data)

	$\tau_{\text{had-vis+jets}}$	$\tau_{\text{had-vis+lepton}}$
Source of uncertainty	Effect on yield	Effect on yield
I. Fake factors: jet composition	1.6%	0.6%
II. Fake factors: statistical uncertainties	1.6%	1.7%
III. Fake factors: prompt $\tau_{\text{had-vis}}$ in the anti- $\tau$ CR	+5.6% -8.3%	+4.8% -7.2%
IV. Fake factors: $\alpha_{\text{MJ}}$ uncertainty	7%	6.2%
V. Fake factors: Smirnov transform.	0%	0%
VI. Fake factors: heavy flavor jet fraction.	5%	5%

Paper: [JHEP 09\(2018\)139](#)

- Estimation of the fake jets  $\rightarrow \tau$  background is the main source of systematic uncertainties in the low- and intermediate-mass  $H^\pm$  search and the second major source (after the signal modelling) for large  $H^\pm$  masses



# Summary

- ▶ One of the key roles of the ATLAS experiment is to search for physics beyond the Standard Model
- ▶ Data-driven Fake Factor method for fake  $\tau$  leptons background estimation is an important part of the  $H^\pm \rightarrow \tau\nu$  analysis. Optimization of the method and detailed uncertainties estimation are essential as a major source of systematic uncertainties
- ▶ Currently getting ready to publish results with full Run 2 (2015-2018) data with improved analysis strategy

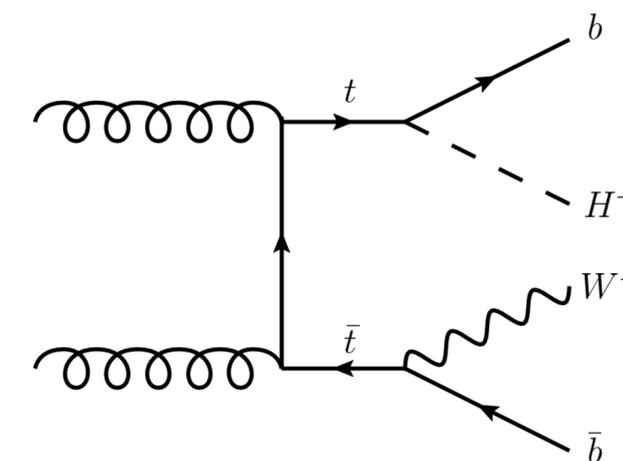
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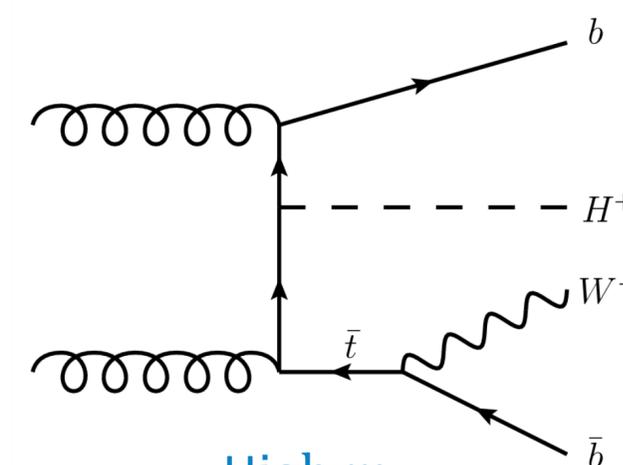
# BACKUP

# $H^\pm \rightarrow \tau\nu$ search

- ▶  $H^\pm$  bosons are predicted in different extensions of the SM that add a second doublet or triplets to the Higgs scalar sector (2HDM, NMSSM, Triplet, etc.)
- ▶ Investigated signal mass range:  $80 \text{ GeV} \leq m_{H^\pm} \leq 3000 \text{ GeV}$
- ▶  $H^\pm$  production process depends on the mass range but is typically associated with  $t$  and  $b$  quarks
- ▶ Analysis split into two sub-channels based on the decay mode of the associated  $t$ :
  - ▶  **$\tau$ +jets**: most sensitive to high  $H^\pm$  mass
  - ▶  **$\tau$ +lepton**: most sensitive to low and intermediate  $H^\pm$  mass
- ▶ The search using 2015+2016 data was published in: [JHEP 09\(2018\)139](#)



Low & intermediate  $m_{H^\pm}$



High  $m_{H^\pm}$

# Analysis Overview

- ▶ Search for charged  $H^\pm$  decaying to  $\tau\nu$  using  $139 \text{ fb}^{-1}$  of  $pp$  data at  $\sqrt{s} = 13 \text{ TeV}$  collected by the ATLAS experiment at the LHC during full Run-2 period
- ▶ **Recurrent Neural Networks (RNN) based  $\tau$  identification** → identification efficiency improved in the range 75–100% compared to one based on Boosted Decision Trees
- ▶ Parametrized Neural Networks (PNN) are trained to separate signal from all backgrounds (more in Zuzanna's talk)
- ▶ Backgrounds with prompt hadronic  $\tau$  :
  - $t\bar{t}$ ,  $V$ +jets and diboson events: estimated with MC
- ▶ Backgrounds with fake  $\tau$ :
  - fake lepton →  $\tau$  : estimated with MC
  - **fake jets →  $\tau$  : estimated with data-driven Fake Factor (FF) method**

# Event Selection & MVA Strategy

## $\tau$ +jets channel

- ▶ 1 medium  $\tau$  candidate with  $p_T^\tau > 40$  GeV
- ▶ no loose leptons ( $e$  or  $\mu$ ) with  $p_T > 20$  GeV
- ▶  $\geq 3$  jets with  $p_T > 25$  GeV, of which at least one is  $b$ -tagged
- ▶  $E_T^{\text{miss}} > 150$  GeV
- ▶  $m_T > 50$  GeV

$$m_T = \sqrt{2p_T^\tau E_T^{\text{miss}}(1 - \cos \Delta\phi_{\tau, \text{miss}})},$$

$\Delta\phi_{\tau, \text{miss}}$  - the azimuthal angle between the  $\tau$  candidate and the direction of  $E_T^{\text{miss}}$

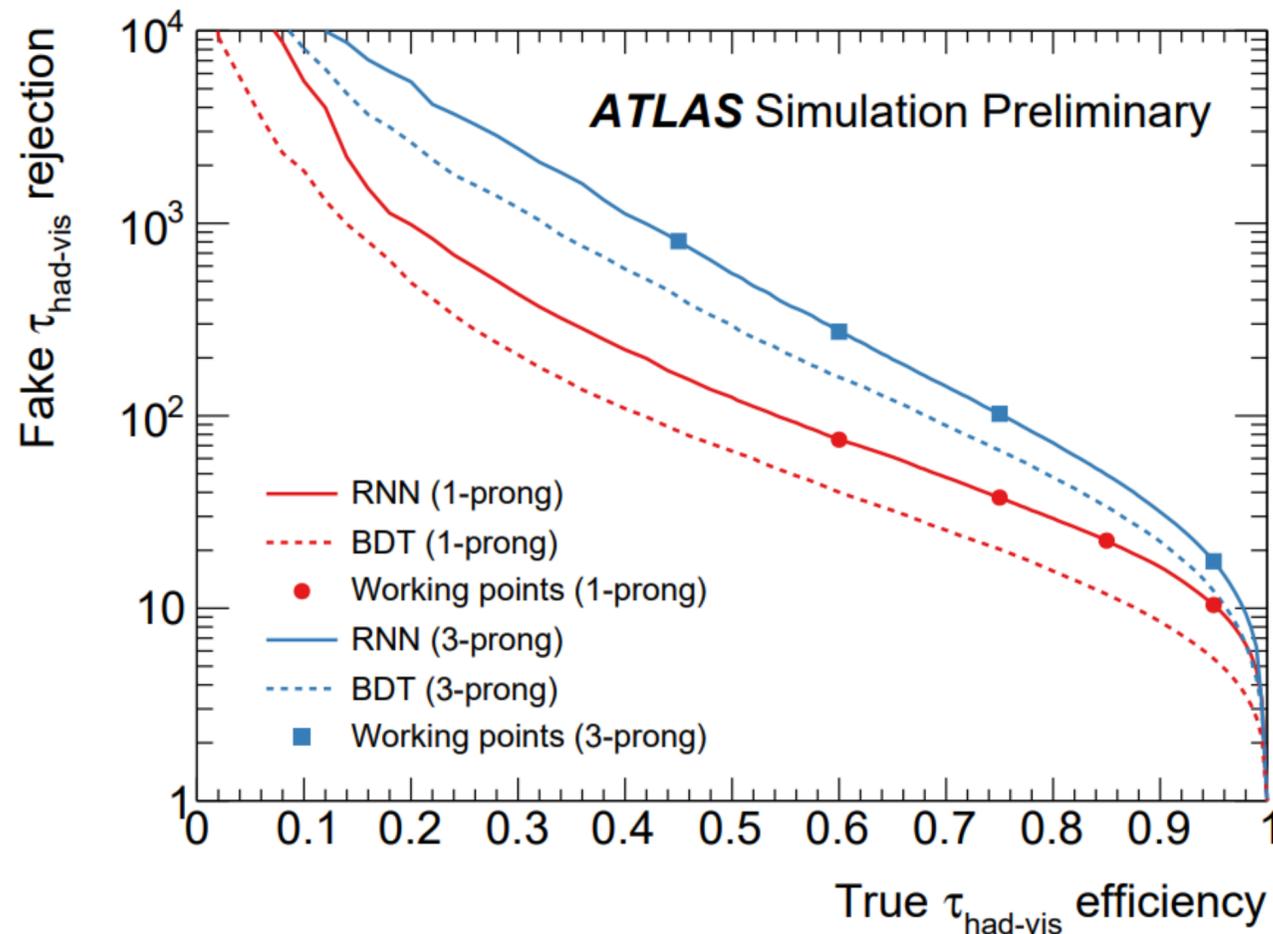
## $\tau$ +lepton channel

- ▶ Exactly 1  $\ell$  ( $e$  or  $\mu$ ), with  $p_T^\ell > 30$  GeV
- ▶ Exactly 1 medium  $\tau$  with no additional loose or tighter  $\tau$  leptons in the event
- ▶  $\ell$  and  $\tau$  with opposite signs
- ▶  $E_T^{\text{miss}} > 50$  GeV
- ▶  $\geq 1$   $b$ -tagged jets

- **New multivariate analysis technique, i.e. Parametrized Neural Networks (PNN), previously Boosted Decision Trees (BDT) used**
- PNN score used as a final discriminant in a likelihood fit

# Tau identification using RNN

A multivariate discriminant is a common tool for background estimation, which enables to differentiate correctly reconstructed tau leptons from the other objects. Right now a recurrent neural network (RNN) is used.



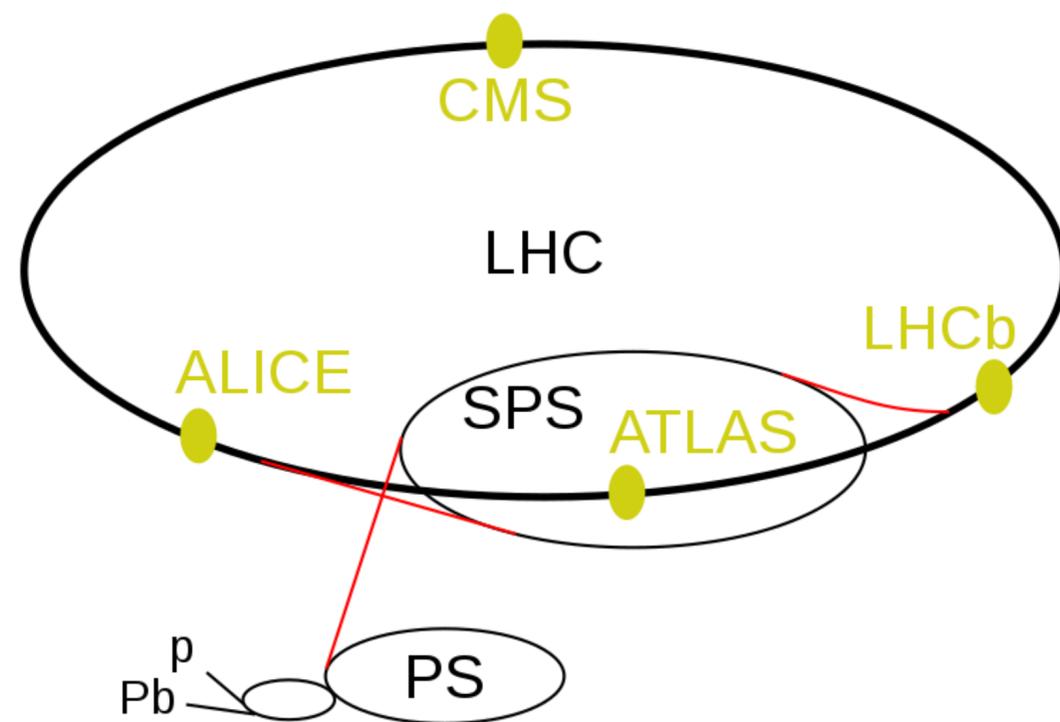
	Observable	1-prong	3-prong
Track inputs	$p_T^{\text{seed jet}}$	•	•
	$p_T^{\text{track}}$	•	•
	$\Delta\eta^{\text{track}}$	•	•
	$\Delta\phi^{\text{track}}$	•	•
	$ d_0^{\text{track}} $	•	•
	$ z_0^{\text{track}} \sin \theta $	•	•
	$N_{\text{IBL hits}}$	•	•
	$N_{\text{Pixel hits}}$	•	•
Cluster inputs	$p_T^{\text{jet seed}}$	•	•
	$E_T^{\text{cluster}}$	•	•
	$\Delta\eta^{\text{cluster}}$	•	•
	$\Delta\phi^{\text{cluster}}$	•	•
	$\lambda_{\text{cluster}}$	•	•
	$\langle \lambda_{\text{cluster}}^2 \rangle$	•	•
	$\langle r_{\text{cluster}}^2 \rangle$	•	•
High-level inputs	$p_T^{\text{uncalibrated}}$	•	•
	$f_{\text{cent}}$	•	•
	$f_{\text{leadtrack}}^{-1}$	•	•
	$\Delta R_{\text{max}}$	•	•
	$ S_{\text{leadtrack}} $	•	•
	$S_T^{\text{flight}}$		•
	$f_{\text{iso}}^{\text{track}}$	•	•
	$f_{\text{EM}}^{\text{track}}$	•	•
	$p_T^{\text{EM+track}} / p_T$	•	•
	$m^{\text{EM+track}}$	•	•
	$m^{\text{track}}$		•

# Definitions of control regions used in FF method

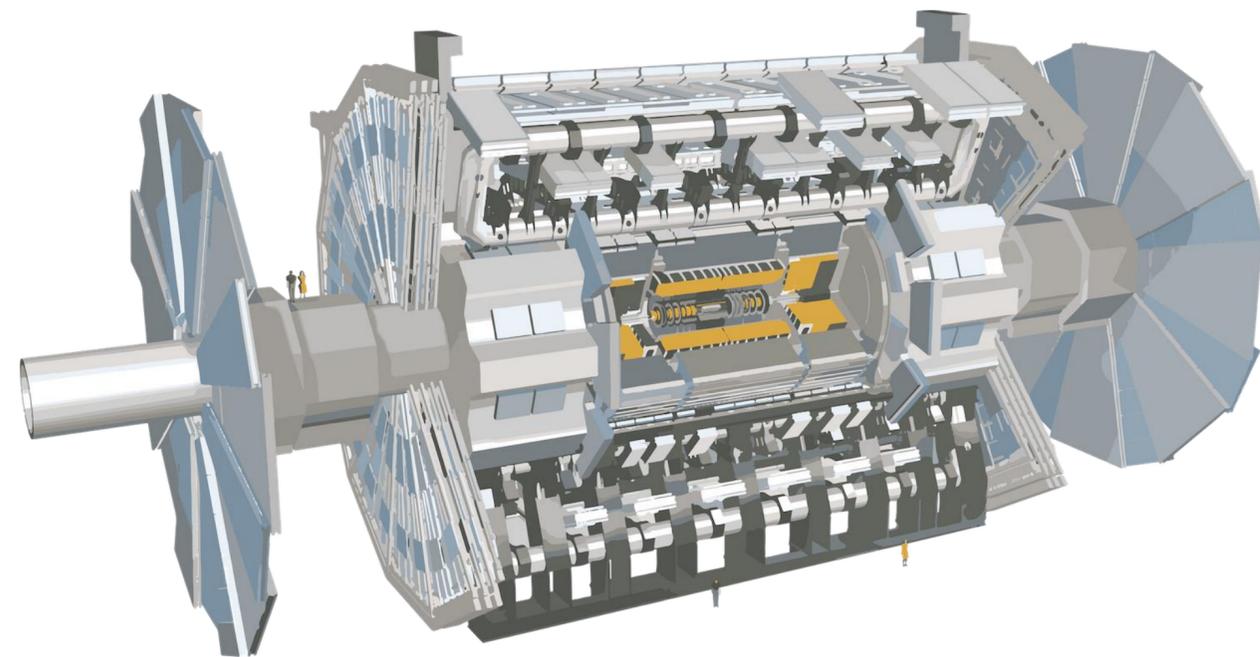
Multi-jets CR	W+jets CR
at least one reconstructed $\tau_{\text{had}}$ candidate with $p_{\text{T}}^{\tau} > 30 \text{ GeV}$ number of jets $\geq 1$ <i>b</i> -jet veto, electron and muon veto $E_{\text{T}}^{\text{miss}} < 80 \text{ GeV}$ $m_{\text{T}}(\tau, E_{\text{T}}^{\text{miss}}) > 50 \text{ GeV}$ transformed $\tau_{\text{had}}$ BDT score $> 0.02$	at least one reconstructed $\tau_{\text{had}}$ candidate with $p_{\text{T}}^{\tau} > 30 \text{ GeV}$ one lepton (electron or muon) <i>b</i> -jet veto $p_{\text{T}}$ of electron or muon $> 30 \text{ GeV}$ $60 < m_{\text{T}}(l, E_{\text{T}}^{\text{miss}}) < 160 \text{ GeV}$ transformed $\tau_{\text{had}}$ BDT score $> 0.02$

# The ATLAS experiment

ATLAS is the largest particle detector experiment at the Large Hadron Collider (LHC), a particle accelerator at CERN in Geneva.



CERN accelerator complex



The ATLAS detector

