# An Overview of Diffractive Photon+Jet Production at the ATLAS Detector 

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## Introduction: The Standard Model

- The Standard Model (SM) of particle physics is the most successful theory to date in describing the fundamental particles and their interactions.
- Needs to be tested as strictly as possible to be verified or find avenues to new physics.
- For this purpose the most powerful particle accelerator in the world, the Large Hadron
 Collider LHC, was built.


## Introduction: LHC and the ATLAS detector

- The Large Hadron Collider (LHC) is a particle accelerator that can accelerate proton beams to up to 7 TeV .
- The experiments are built in the interaction points, where the beams cross.

- The ATLAS detector has a multilayered structure, with each layer focusing in measuring different physical properties.
- Its one of the two multipurpose detectors at the LHC.


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- in QM the bound states for a spherically symmetric potential fall into families with increasing angular momentum and energy,
- these bound states appear as poles of the partial wave amplitude with a given integer angular momentum,
- idea: continue these amplitudes to complex values of angular momentum,
- for 'well behaved' potentials (like the Yukawa one) the poles lie on a straight line, called the Regge trajectory:
$\alpha_{R}(t)=\alpha_{R}(0)+\alpha_{R}^{\prime}(0) \cdot t$, where
$\alpha_{R}(0)$ is called the intercept
$\alpha_{R}^{\prime}(0)$ - the slope,
- object exchanged in the $t$ channel between two hadrons is not a single particle, but all particles lying on the Regge trajectory.


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- All known resonances lie on trajectories with an intercept smaller than $1 \rightarrow$ the total cross section should decrease with increasing collision energy.
- Which was the case at low energies, but clearly not true at higher ones.




## Regge Theory and Pomeron Concept

Regge theory studies the analytical properties of scattering. By expanding the partial wave equation to imaginary angular momentum is possible to build (Regge) trajectories which contain all bound states. For the case of Yukawa potential.

$$
\alpha_{R}(t)=\alpha_{R}(0)+\alpha_{R}^{\prime}(0) t
$$

- The Reggeon object, which is what is exchanged, is a mixture of all the resonances in the trajectory. It predicts a descending cross-section with increase of centre-of-mass energy.
- Pomeron trajectory is introduced, which produces a cross-section that increases logarithmically with energy.


Figure: Selected Regge trajectories. On the dashed line Pomeron trajectory, with $\alpha(0) \sim 1.1$. Extracted from [3]

## Diffraction at the LHC

- Diffraction in high energy physics is referred to as events governed by the mechanism of colourless exchange of vacuum quantum numbers:
- photon in case of electromagnetic and
- Pomeron for strong interactions.
- The Pomeron trajectory does not hold any known resonance:
- its actual structure is as yet unknown,
- the simplest possibility being a two-gluon glueball (+h.o. contributions).

"Regge Pomeron"


two-gluon glueball

h.o. corrections $\rightarrow$ internal structure
- Main signatures are the presence of a large rapidity gap devoid of particles, that can be destroyed by further interactions, and protons scattered at very small angles ( $\mu \mathrm{rad}$ )
- Generally soft, low transverse momentum $p_{T}$ transfer, makes them intractable by perturbation methods. Need for effective theories.

Measuring forwards protons: The AFP, Atlas Forwards Proton detector


- Silicon Tracker (SiT): A set of four planes in each Roman Pot (RP) station.
- $50 \times 250 \mu \mathrm{~m}$ pixel size.
- Planes tilted $14^{\circ}$ to improve resolution.
- Resolution: $\sigma_{x}=6 \mu \mathrm{~m}, \sigma_{y}=30 \mu \mathrm{~m}$.
- Time-of-flight (ToF): Designed to measure the primary vertex z-coordinate.
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## Measurement Principle: SiT

The proton trajectory depends on:

- The energy loss on the interaction $\xi=1-\frac{E_{\text {proton }}}{E_{\text {beam }}}$.
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Figure: AFP measures displacement, which is related to mass of the central system.

Acceptance of the detector limited by collimator apertures and beam-detector distance.

## Backgrounds and Time of Flight detector

- Cross sections are expected to be low, while production of single diffraction protons is around $10 \%$ of the total $\rightarrow$ High probability coincidental fake double tag events.
- Need to operate at regimes with low number of interaction per bunch crossing (low $\mu$ ) $\rightarrow$ Need longer times to accumulate enough events.
- ToF times the arrival of the protons, which allows to reconstruct the longitudinal position of the event. Comparing with information of the central detectors permits to
 reject background events.




## Double Pomeron Exchange Photon + Jet Production

- Analogically to the proton internal structure, Pomeron might be a quite complicated object:
- gluonic structure can be probed by looking at properties of events coming from gluon-gluon interactions, e.g. studies of Double Pomeron Exchange Jet Production (top diagram),
- quark structure can be studied using processes like DPE $\gamma+$ jet (see middle diagram) or DPE $W$ production (bottom diag.).
- DPE $\gamma+$ jet production signature:
- both protons exchange a Pomeron, remaining intact,
- one Pomeron emits a gluon and takes a quark from the other, generating a photon and a jet,
- intact protons might be later destroyed by further soft interactions; this is modelled by a gap survival probability, which is expected to be of about 0.03 for DPE processes at the LHC energies.
- Quark composition can be studied by measuring the ratio of DPE $\gamma+$ jet to DPE JJ:






## Pomeron Quark Content: Observables

- Top left: ratio of DPE $\gamma+$ jet to DPE JJ as a function of $p_{T}$. Color lines represent various assumption of $d / u$ ratio in the Pomeron.
- Bottom: distribution of diffractive $\operatorname{mass} M=\sqrt{\xi_{1} \xi_{2} s} . \xi_{1,2}$ denotes energy lost by protons for various assumption of $d / s$ (left) and $d / u$ (right) ratio in the Pomeron.





## Datasets for SD and DPE $\gamma+$ jet

Run 2 (tables taken from SD JJ analysis):
Table 5.1: An overview of 2017 low- $\mu$ runs with integrated luminosity from LBs passing the GRL requirements separate for protons on the ATLAS A and C sides.

| ATLAS <br> Run Number | $\begin{aligned} \text { LHC } \\ \hline \text { er } \quad \text { Fill } \\ \hline \end{aligned}$ | $\begin{gathered} \text { Pile-up } \\ \mu \\ \hline \end{gathered}$ | Int. Luminosity $\left[\mathrm{nb}^{-1}\right]$for protons on side $A$ |  | Int. Luminosity $\left[\right.$ nb $\left.^{-1}\right]$for protons on side C |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 331020336505 | 6019 | $\sim 1.0$ |  | 56.866 |  | 510.841 |
|  | 6238 | $\sim 0.04$ |  | 44.751 |  | 60.2411 |
| 341294 | 6404 | $\sim 2.0$ |  | 709.542 |  | 709.542 |
| 341312 | 6405 | $\sim 2.0$ |  | 18245.492 |  | 18234.639 |
| 341419 | 6411 | $\sim 2.0$ |  | 31636.072 |  | 31593.050 |
| 341534 | 6413 | $\sim 2.0$ |  | 47663.701 |  | 52680.387 |
| 341615 | 6349 | $\sim 2.0$ |  | 31772.631 |  | 31772.631 |
| 341649 | 6417 | $\sim \sim$ | 3325.167 |  |  | 6449.680 3325.167 |
|  |  | $\sim 1.0$ |  |  |  | 3325.167 |
| Sample | Side |  | Consecutive Cut |  |  |  |
|  |  | All | 1 Vertex | GRL | 2 Jets | 1 Proton |
| SD MC | A | 239499 | 136390 | 136390 | 4825 | 2061 |
|  | C |  |  |  |  | 2654 |
| 331020 | A | 1952990 | 357983 | 35404 | 23810 | 4447 |
|  | C |  |  | 321605 | 217344 | 86869 |
| 336505 | A | 41908 | 30857 | 22712 | 8536 | 658 |
|  | C |  |  | 30784 | 11602 | 4262 |
| 341294 | A | 116385 | 7124 | 6977 | 4651 | 1122 |
|  | C |  |  | 6977 | 4651 | 1214 |
| 341312 | A | 940071 | 70859 | 70717 | 47080 | 3111 |
|  | C |  |  | 70710 | 47076 | 19056 |
| 341419 | A | 1649173 | 123268 | 123124 | 81825 | 5436 |
|  | C |  |  | 123003 | 81745 | 32564 |
| 341534 | A | 2864217 | 246297 | 212460 | 143152 | 10453 |
|  | C |  |  | 244662 | 164771 | 63254 |
| 341615 | A | 1650721 | 130234 | 130114 | 86375 | 6217 |
|  | C |  |  | 130114 | 86375 | 34642 |
| $\begin{gathered} 341649 \\ \mu \sim 1.0 \end{gathered}$ | A | 236404 | 44633 | 39653 | 26460 | 2011 |
|  | C |  |  | 39653 | 26460 | 10533 |
| $\begin{aligned} & 341649 \\ & \mu \sim 2.0 \end{aligned}$ | A | 354020 | 27503 | 26387 | 17611 | 1315 |
|  | C |  |  | 26124 | 17439 | 6943 |

## Run 3:

- $\mu \sim 1$ :

- $\mu \sim 0.2$ :
- $455818 \approx 35 \mathrm{nb}^{-1}$,
- $\mu \sim 0.05$ :
- $428770 \approx 34 \mathrm{nb}^{-1}$,
- $455818 \approx 20 \mathrm{nb}^{-1}$,
- $455838 \approx 43 \mathrm{nb}^{-1}$,
- $\mu \sim 0.02$ (higher $\xi_{\text {min }}$ ):
- $435229 \approx 155 \mathrm{nb}^{-1}$,
- $435333 \approx 15 \mathrm{nb}^{-1}$,
- $\mu \sim 0.005$ :
- $427929 \approx 0.46 \mathrm{nb}^{-1}$,
- $\mu \sim 0.005$ (low-B):
- $460348 \approx 1.75 \mathrm{nb}^{-1}$.


## A NEAR, Cluster Reco. Efficiency (2022)



## Summary

- Diffraction in high energy physics, characterized by large rapidity gaps and presence of forward protons, can be studied using data collected by ATLAS Roman Pots.
- Double Pomeron Exchange $\gamma+$ jet production offers a probe in the quark content of the Pomeron; note that by using ratio to DPE JJ, the impact of gap survival will be effectively cancelled out.
- cross-section determination,
- if enough statistics - quark structure of Pomeron.
- Existing Run 3 datasets will be checked for evidence of single diffractive $\gamma+$ jet events. First task is to measure cross section $\rightarrow$ determine gap survival probability.
- Not enough data to see the evidence of DPE $\gamma+$ jet in Run 3:
- $p p$ reference run may be a nice opportunity to make such measurement at $\sqrt{s}=5.36 \mathrm{TeV}$,
- Run $2, \sqrt{s}=13 \mathrm{TeV}$, 2017, $\mu \sim 2$ data-sets to be investigated.
- measurement at $\sqrt{s}=13.6 \mathrm{TeV}$ would require a few days of low- $\mu$ run.


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Backup

## Digression: Why "Diffraction"?



- 'Diffraction' in optics:
- light with wavelength of $\lambda$ is shining on black disc with radius $R_{0}$,
- distant screen - characteristic 'diffractive' pattern:
- large forward peak for scattering angle $\theta=0$,
- series of symmetric minima and maxima, with the first minimum at $\theta_{\text {min }} \simeq \pm \lambda /\left(2 R_{0}\right)$,
- intensity as a function of scattering angle: $\frac{I(\theta)}{I(\theta=0)}=\frac{\left[2 J_{1}(x)\right]^{2}}{x^{2}} \simeq 1-\frac{R_{0}^{2}}{4}(k \theta) 2$
- $J_{1}$ is the Bessel function of the first order,
- $x=k R_{0} \sin \theta \simeq k R_{0} \theta$ with $k=2 \pi / \lambda$.
- diffraction pattern is related to the size of the target and to the wavelength of the light beam.


Differential cross-section for $p p \rightarrow p p$ :
$\frac{\frac{d \sigma}{d t}(t)}{\frac{d \sigma}{d t}(t=0)} \simeq e^{-b|t|} \simeq 1-b(P \theta)^{2}$

- $|t| \simeq(P \theta)^{2}$ - absolute value of the squared four-momentum transfer,
- $P$ is the incident proton momentum,
- $\theta$ is the scattering angle,
- $b=R^{2} / 4$, where $R$ is related to the target size,

Data: a dip followed by a secondary maximum.
Similar $t$ distributions observed for other reactions $\rightarrow$ diffractive processes.

## The ATLAS Forward Proton Detector



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